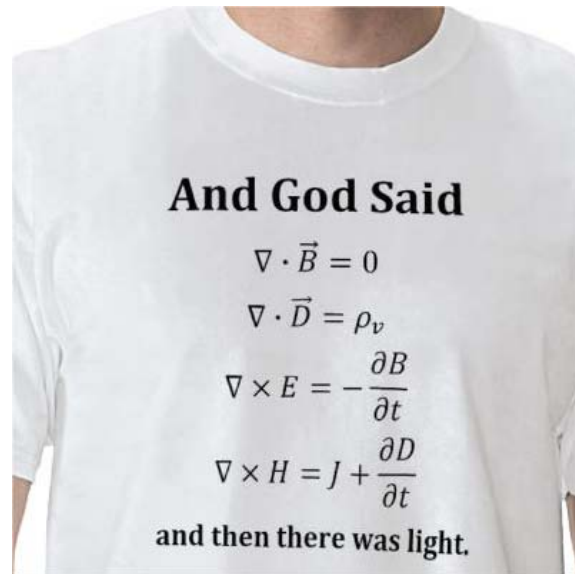


Elektromagnetische Wellen (im Vakuum)

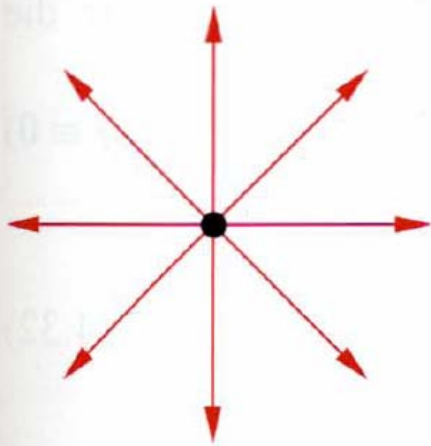
1. Maxwell-Gleichungen

$$\begin{aligned}\operatorname{rot} \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t}, \\ \operatorname{rot} \mathbf{H} &= \mathbf{j} + \frac{\partial \mathbf{D}}{\partial t}, \\ \operatorname{div} \mathbf{D} &= \varrho, \\ \operatorname{div} \mathbf{B} &= 0.\end{aligned}$$

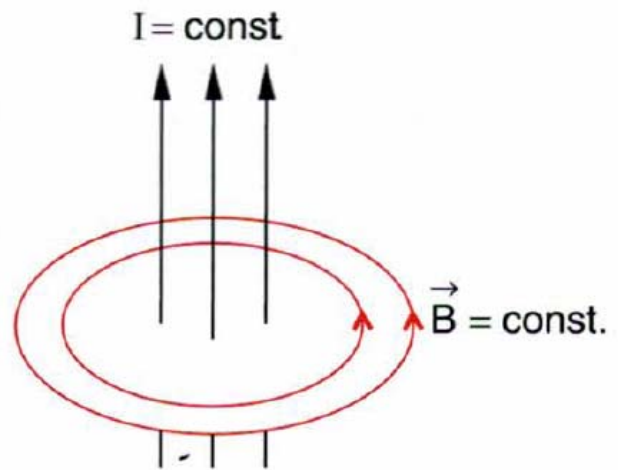
$$\rho = 0 ; \vec{j} = 0$$



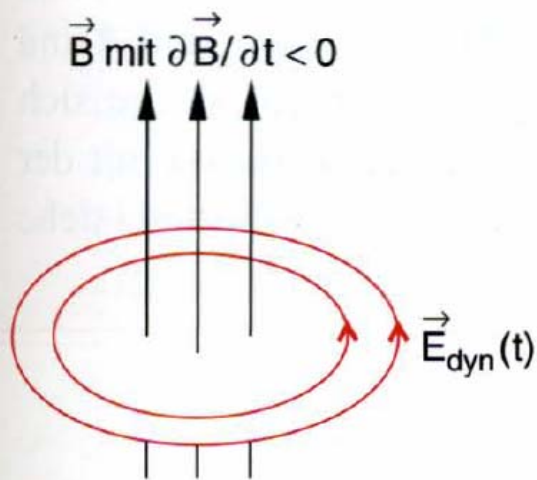
Elektrische Felder werden sowohl von Ladungen als auch von zeitlich sich ändernden Magnetfeldern erzeugt. Magnetische Felder werden von Strömen oder von zeitlich sich ändernden elektrischen Feldern erzeugt.



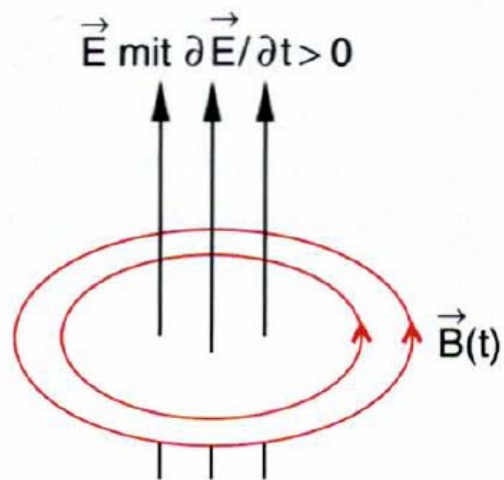
$\vec{E}_{\text{stat}}(q)$



$\text{rot } \vec{B} = \mu_0 \vec{j}$



a) $\frac{\partial \vec{B}}{\partial t} \neq 0; \text{rot } \vec{E} = -\frac{\partial \vec{B}}{\partial t}$



b) $\frac{\partial \vec{E}}{\partial t} \neq 0; \text{rot } \vec{B} = \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}$

2. Die Wellengleichung

$$\Delta \vec{E} = \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\Delta \vec{B} = \epsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$\vec{\nabla} \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}$$

$$\begin{aligned} \vec{\nabla} \times \vec{\nabla} \times \vec{E} &= - \vec{\nabla} \times \frac{\partial \vec{B}}{\partial t} = - \frac{\partial}{\partial t} (\vec{\nabla} \times \vec{B}) \\ &= - \epsilon_0 \mu_0 \frac{\partial}{\partial t} \left(\frac{\partial \vec{E}}{\partial t} \right) = - \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} \end{aligned}$$

$$\begin{aligned} \vec{\nabla} \times \vec{\nabla} \times \vec{E} &= \nabla(\nabla \cdot \vec{E}) - \nabla \cdot (\nabla \vec{E}) \\ &= \underset{0}{\text{grad}(\text{div} \vec{E})} - \text{div}(\text{grad} \vec{E}) \end{aligned}$$

$$\text{div}(\text{grad} \vec{E}) = \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

der Laplace - Operator $\Delta = \text{div}(\text{grad} \dots)$

$$\text{Lichtgeschwindigkeit } c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

$$\Delta \vec{E} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} \quad \approx 3 \cdot 10^8 \frac{\text{m}}{\text{s}}$$

$$\frac{\partial^2 E_x}{\partial x^2} + \frac{\partial^2 E_x}{\partial y^2} + \frac{\partial^2 E_x}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 E_x}{\partial t^2}$$

ähnlich E_y, E_z

identisch für $\vec{B}(\vec{r}, t)$

$$\vec{\nabla} \times \vec{\nabla} \times \vec{B} = \dots$$

3. Ebene elektrische Wellen

$$\frac{\partial \vec{E}}{\partial x} = 0 \quad ; \quad \frac{\partial \vec{E}}{\partial y} = 0$$

$$\text{div } \vec{E} = 0 \quad \longrightarrow \quad \frac{\partial E_z}{\partial z} = 0$$

↓

$$E_z = A = 0 \text{ konstante}$$

$$\vec{E} = \begin{pmatrix} E_x \\ E_y \\ 0 \end{pmatrix}$$

ebene Welle

$$\frac{\partial^2 \vec{E}}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$E_x(z, t) = f_x(z - ct) + g_x(z + ct)$$

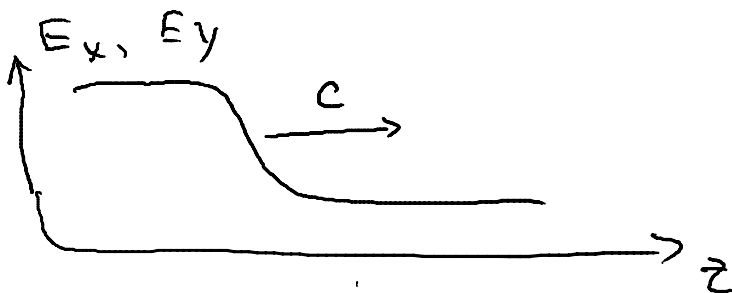
$$E_y(z, t) = f_y(z - ct) + g_y(z + ct)$$

$$z - ct = \text{const} \quad ; \quad \rightarrow \quad \frac{\partial z}{\partial t} - c = 0$$

$$\frac{\partial z}{\partial t} = c$$

$$z + ct = \text{const} \quad ; \quad \rightarrow \quad \frac{\partial z}{\partial t} + c = 0$$

$$\frac{\partial z}{\partial t} = -c$$



transversale Wellen

$$\vec{E} = \begin{pmatrix} \neq 0 \\ \neq 0 \\ 0 \end{pmatrix}$$

4. Periodische Wellen

$\lambda \rightarrow$ Wellenlänge

$$f(z - ct + \lambda) = f(z - ct)$$

$$\vec{E} = \vec{E}_0 \cdot f(z - ct) = \vec{E}_0 \cdot \sin \left[k(z - ct) \right]$$

k - Wellenzahl

$$k \lambda = 2\pi \quad ; \quad k = \frac{2\pi}{\lambda} \quad ; \quad k \equiv k_z$$

$$\vec{E} = \vec{E}_0 \cdot \sin \left(kz - \frac{2\pi c}{\lambda} t \right)$$

$$= \vec{E}_0 \cdot \sin (kz - \omega t)$$

$$c = \frac{\omega}{2\pi} \cdot \lambda \quad \rightarrow \quad v = \frac{\omega}{2\pi} = \frac{c}{\lambda}$$

oder

$$\vec{E} = \vec{E}_0 \cdot \cos (kz - \omega t)$$

oder

$$\vec{E} = \vec{A}_0 \cdot e^{i(kz - \omega t)} + \vec{A}_0^* e^{-i(kz - \omega t)}$$

$$|\vec{A}_0| = \frac{|\vec{E}_0|}{2}$$

Ausbreitungsvektor $\Rightarrow \vec{k} = \begin{pmatrix} k_x \\ k_y \\ k_z \end{pmatrix}$

Wellenvektor

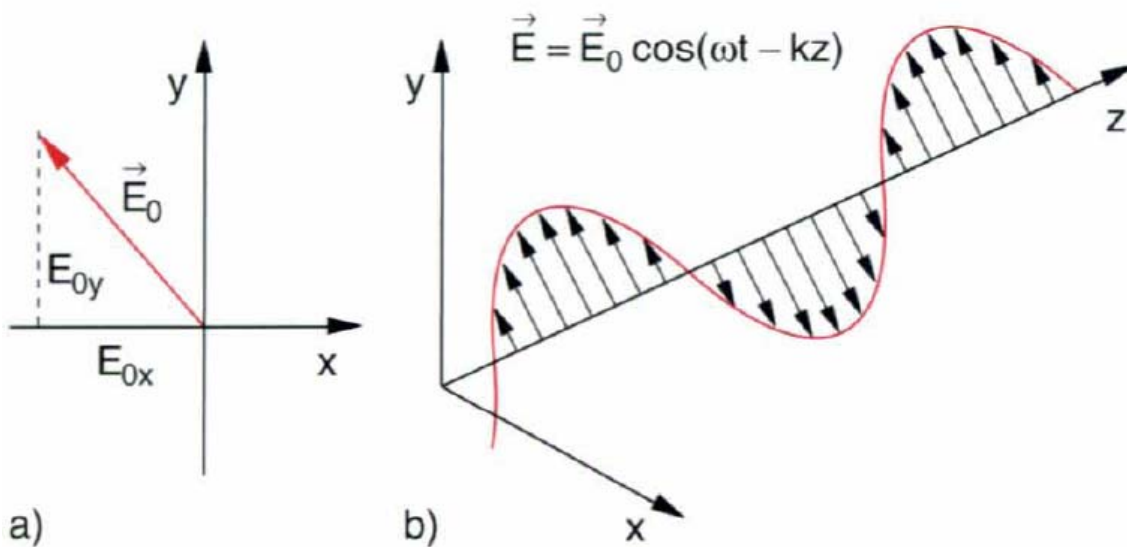
$$|\vec{k}| = k = \frac{2\pi}{\lambda}$$

$$\vec{E} = \vec{A}_0 \cdot e^{i(\vec{k} \cdot \vec{r} - \omega t)} \quad ; \quad \vec{r} = \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

5. Polarisation elektromagnetischer Wellen

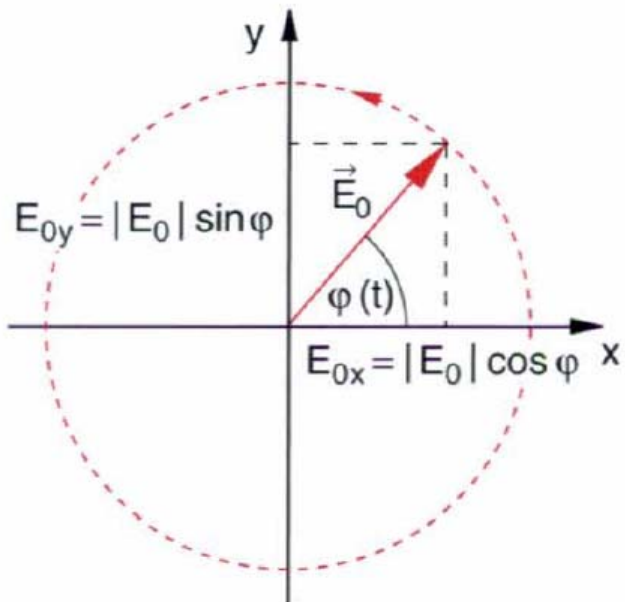
Linear polarisierte Wellen

$$\begin{aligned} \vec{E} &= \vec{E}_0 \cdot \cos(\omega t - kz) \\ E_x &= E_{0x} \cdot \cos(\omega t - kz) \\ E_y &= E_{0y} \cdot \cos(\omega t - kz) \end{aligned}$$

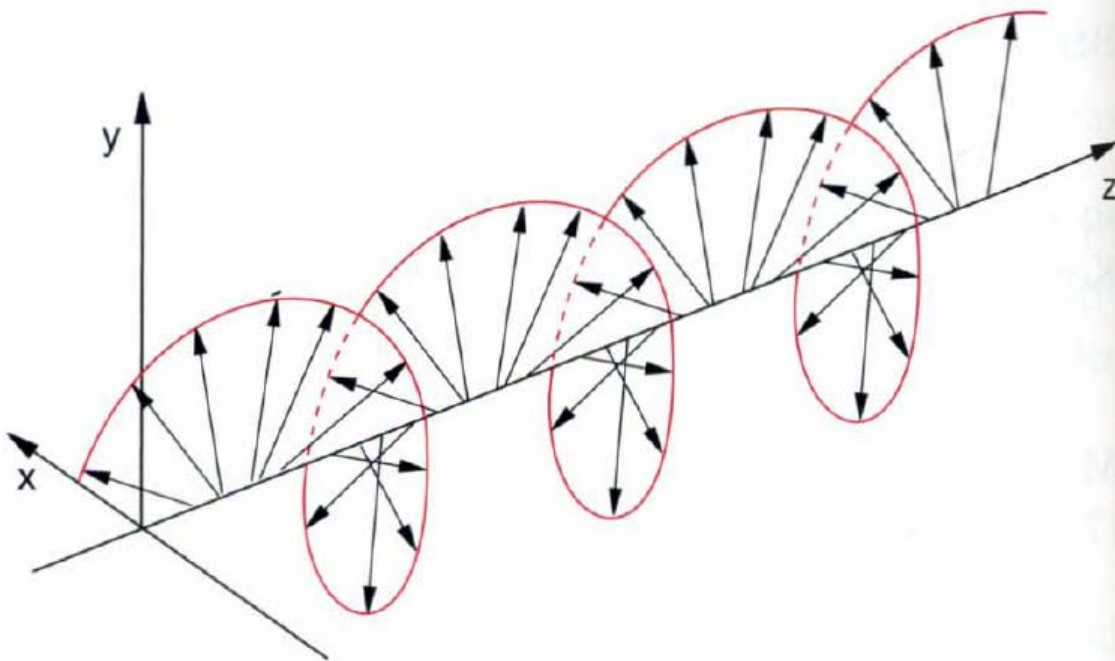


Zirkular polarisierte Wellen

$$\begin{aligned} E_x &= E_{0x} \cdot \cos(\omega t - kz) \\ E_y &= E_{0y} \cdot \cos\left(\omega t - kz - \frac{\pi}{2}\right) = E_{0y} \cdot \sin(\omega t - kz) \end{aligned}$$



a)



b)

Elliptisch polarisierte Wellen:

$$E_{0x} \neq E_{0y}$$

oder $\Delta\varphi \neq 90^\circ$

6. Das Magnetfeld elektromagnetischer Wellen

Für eine in x-Richtung linear polarisierte Welle

$$\mathbf{E} = E_0 \cdot \hat{\mathbf{e}}_x \cdot e^{i(\omega t - kz)}$$

durch Anwendung **rot**

$$(\nabla \times \mathbf{E})_x = 0; \quad (\nabla \times \mathbf{E})_z = 0; \quad (\nabla \times \mathbf{E})_y = \frac{\partial E_x}{\partial z}.$$

$$\frac{\partial \mathbf{B}}{\partial t} = -(\nabla \times \mathbf{E}) \quad \longrightarrow \quad \frac{\partial B_x}{\partial t} = \frac{\partial B_z}{\partial t} = 0$$

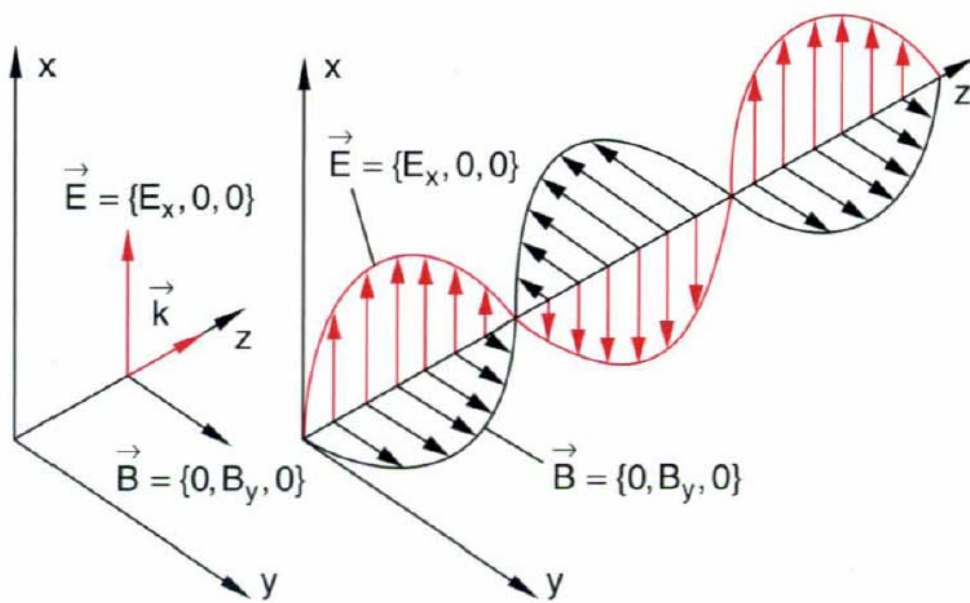
$$-\frac{\partial B_y}{\partial t} = \frac{\partial E_x}{\partial z} = -ikE_x,$$

$$B_y = ikE_0 \int e^{i(\omega t - kz)} dt = \frac{k}{\omega} E_0 e^{i(\omega t - kz)}.$$

$$\omega/k = c \quad \longrightarrow \quad |\mathbf{B}| = \frac{1}{c} |\mathbf{E}|$$

$$\mathbf{E} = \{E_x, 0, 0\} \quad \text{und} \quad \mathbf{B} = \{0, B_y, 0\}$$

$$\mathbf{B} = \frac{1}{\omega} (\mathbf{k} \times \mathbf{E})$$



Das elektromagnetische Frequenzspektrum

