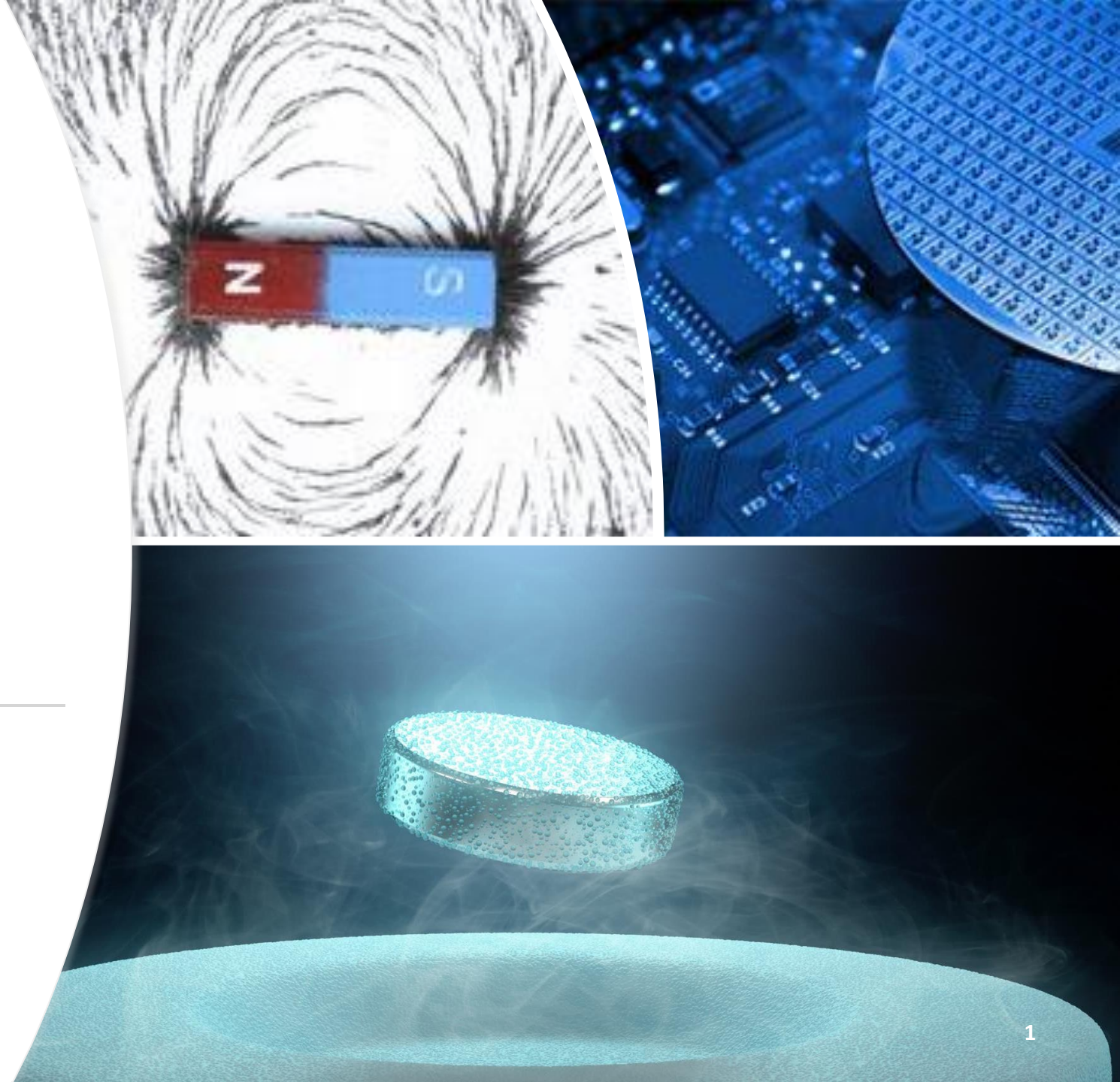




Moderne Experimentalphysik III

Experimentelle Festkörperphysik

M-PHYS-106295, SoSe 2024



Letzte Vorlesung

■ Anmeldung für die Prüfung!!

■ Ab ca. 12:00 Uhr:

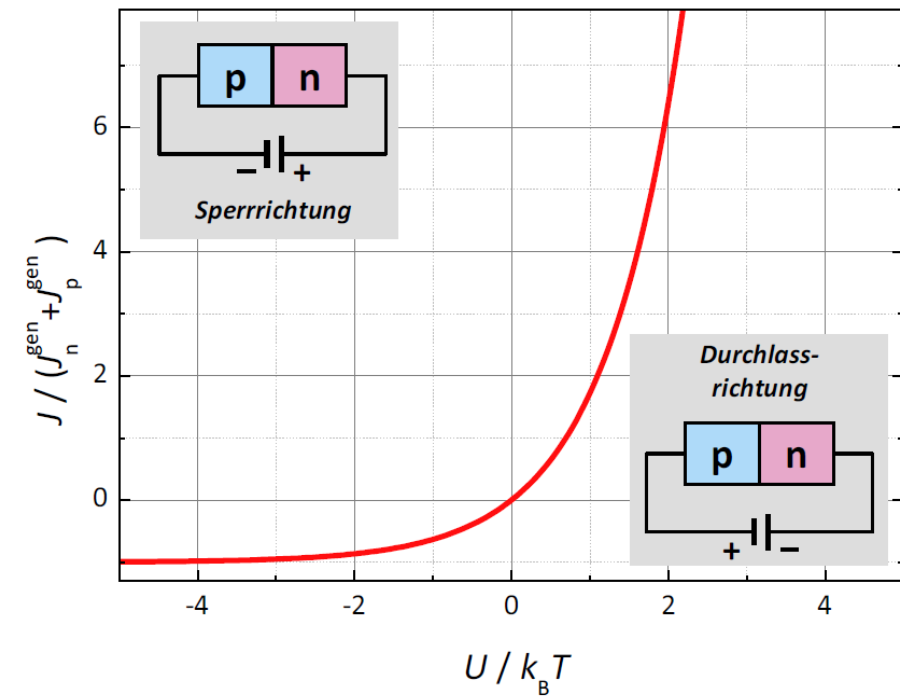
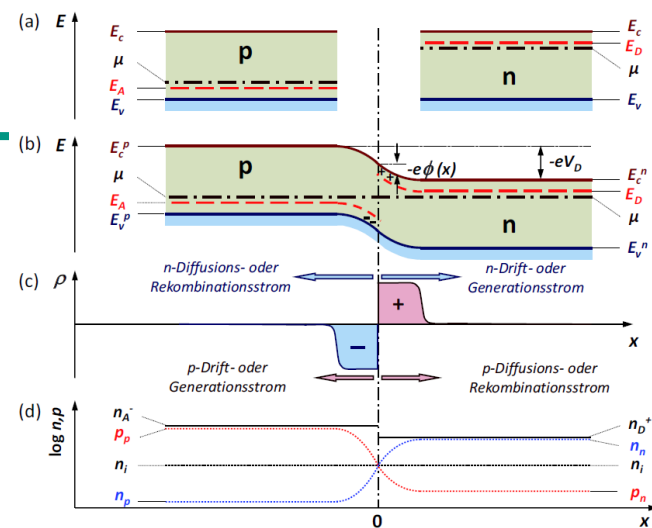
Wiederholung / Labortours

- Halbleiter-Spin Qubits (AG Wernsdorfer)
- SQUID lab (AG Wernsdorfer / AG Pop)
- Rastertunnelmikroskopie-Labor, Spins on Surfaces / Supraleiter (AG Willke)

Wiederholung

pn-Übergang

Solarzelle



Halbleiterheterostrukturen

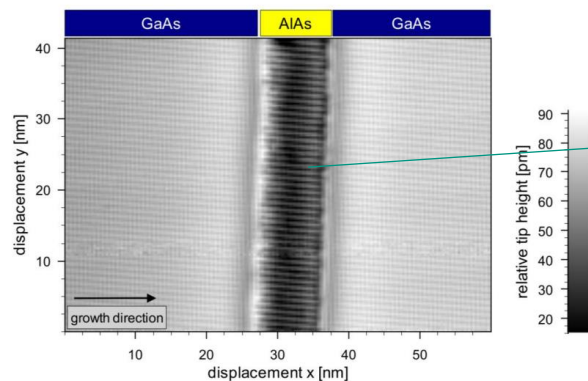
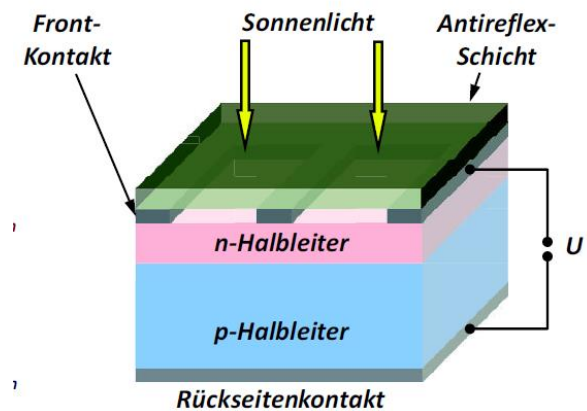
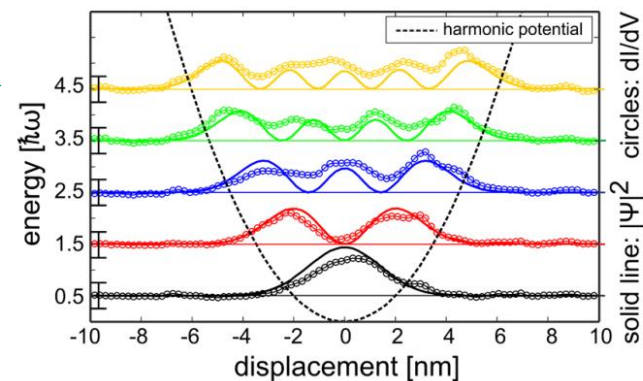


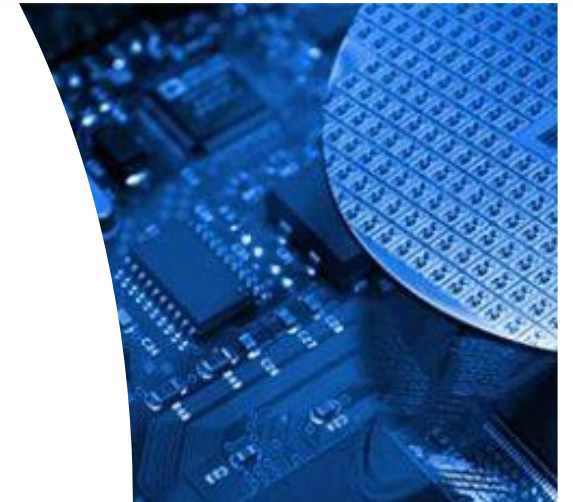
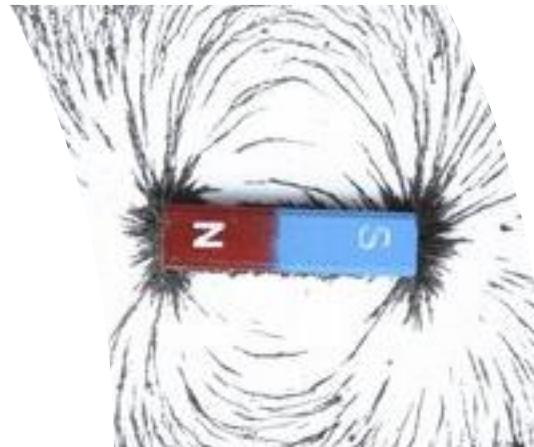
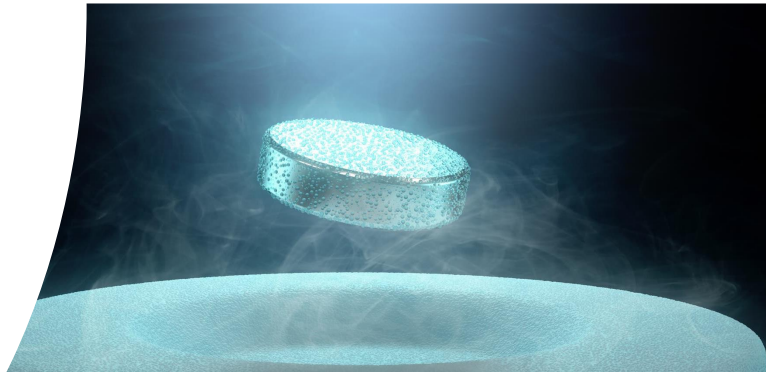
Figure 4.3: Constant current topography of the AlAs layer (marked yellow at the top of the image) embedded in GaAs (marked blue) taken at a setpoint of -2V and 0.1nA.



Wiederholung

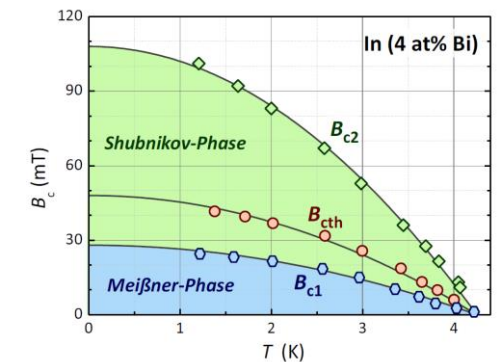
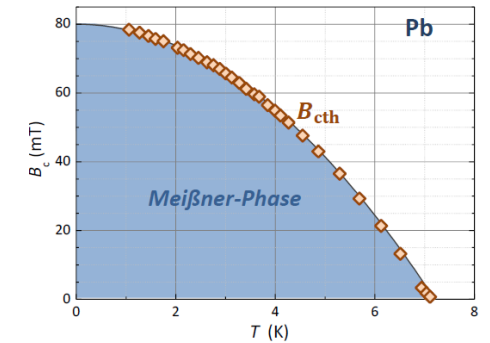
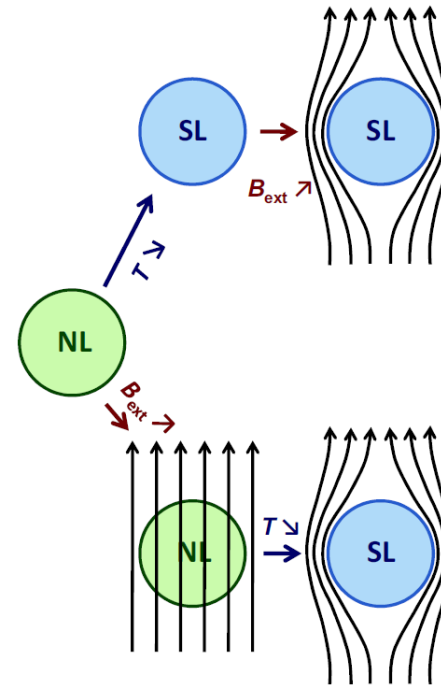
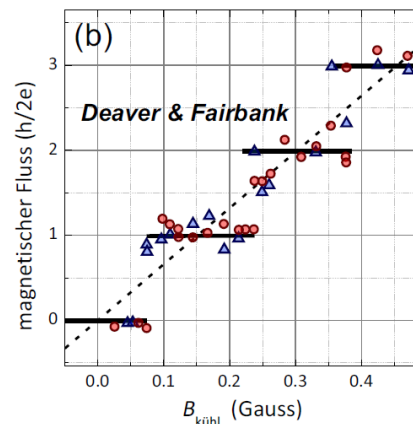
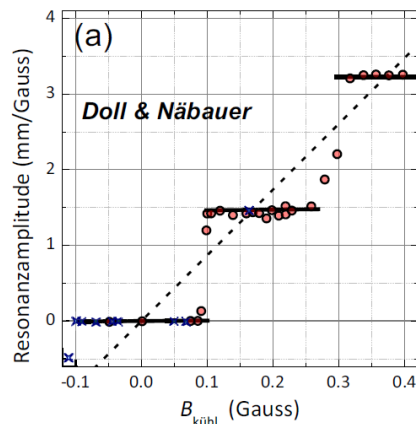
Inhalt

- Dielektrische Eigenschaften von Isolatoren
- Halbleiter
- p-n-Übergang
- Niedrigdimensionale Elektronensysteme
- 1D- und 2D-Elektronengas
- Quanten-Hall Effekt
- Magnetische Eigenschaften
- Magnetismus der Leitungselektronen.
- Atomarer Magnetismus
- Magnetische Wechselwirkungen
- Ferro- und Antiferromagnetismus
- Grundbegriffe der Supraleitung
- London-Gleichungen
- Cooper-Paare
- Supraleiter 1. und 2. Art
- Josephson-Effekte



Wiederholung

- Supraleitung führt zu widerstandslosem Stromtransport bei tiefen Temperaturen
- Supraleiter sind perfekte Diamagneten (Meissner-Effekt)
- Typ I und Typ II Supraleiter
- Flussquantisierung



Wiederholung

■ London Gleichungen

1. London-Gleichung

$$\frac{m_s}{q_s^2 n_s} \frac{\partial \mathbf{J}_s}{\partial t} = \mathbf{E}$$

mit $\Lambda = m_s / n_s q_s^2$ = London-Koeffizient

→ beschreibt **widerstandslosen Stromtransport**

2. London-Gleichung

$$\nabla \times (\Lambda \mathbf{J}_s) + \mathbf{b} = \mathbf{0}$$

→ beschreibt **perfekte Feldverdrängung**

Wiederholung

■ Makroskopische Wellenfunktion

Grundhypothese: es existiert eine *makroskopische Wellenfunktion*

$$\psi_s(\mathbf{r}, t) = \psi_0(\mathbf{r}, t) e^{i\theta(\mathbf{r}, t)}$$

$$|\psi(\mathbf{r}, t)|^2 = \psi^*(\mathbf{r}, t) \psi(\mathbf{r}, t) = n_s(\mathbf{r}, t)$$



Strom-Phasen-Beziehung

$$\mathbf{J}_s(\mathbf{r}, t) = \frac{q_s n_s(\mathbf{r}, t) \hbar}{m_s} \left\{ \nabla \theta(\mathbf{r}, t) - \frac{q_s}{\hbar} \mathbf{A}(\mathbf{r}, t) \right\}$$

v



Energie-Phasen-Beziehung

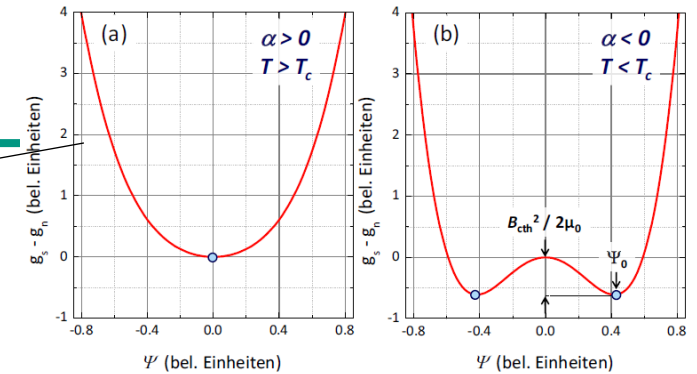
$$\hbar \frac{\partial \theta(\mathbf{r}, t)}{\partial t} = - \left\{ \frac{1}{2n_s} \Lambda J_s^2(\mathbf{r}, t) + q_s \phi_{\text{el}}(\mathbf{r}, t) + \mu(\mathbf{r}, t) \right\}$$

Gesamtenergie



Wiederholung

Ginzburg-Landau Theorie



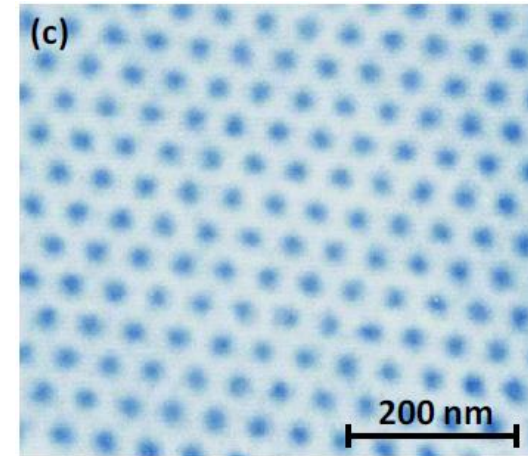
$$g_s = \underbrace{g_n + \alpha|\Psi|^2 + \frac{1}{2}\beta|\Psi|^4}_{(\mu_0 M)^2} + \frac{1}{2\mu_0} (\mathbf{B}_{\text{ext}} - \mathbf{b})^2 + \frac{1}{2m_s} \left| \left(\frac{\hbar}{i} \nabla - q_s \mathbf{A} \right) \Psi \right|^2 + \dots$$

1. Ginzburg-Landau Gleichung

$$0 = \frac{1}{2m_s} \left(\frac{\hbar}{i} \nabla - q_s \mathbf{A} \right)^2 \Psi + \alpha \Psi + \beta |\Psi|^2 \Psi$$

Ginzburg-Landau Kohärenzlänge

$$\xi_{\text{GL}} = \sqrt{-\frac{\hbar^2}{2m_s \alpha}}$$



2. Ginzburg-Landau Gleichung

$$\mathbf{J}_s = \frac{q_s \hbar}{2m_s} \frac{1}{i} (\Psi^* \nabla \Psi - \Psi \nabla \Psi^*) - \frac{q_s^2}{m_s} |\Psi|^2 \mathbf{A}$$

$$\kappa \leq \frac{1}{\sqrt{2}}$$

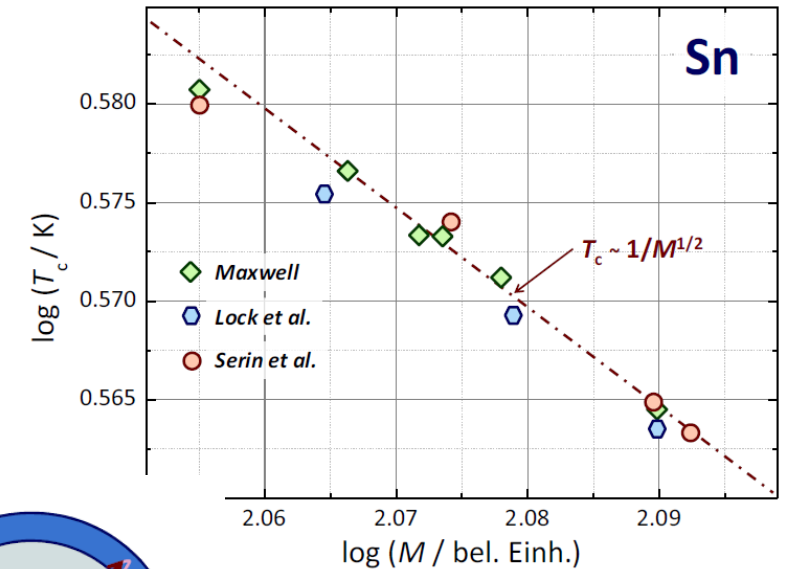
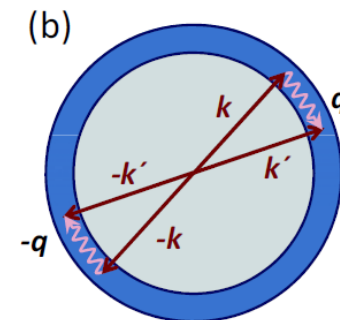
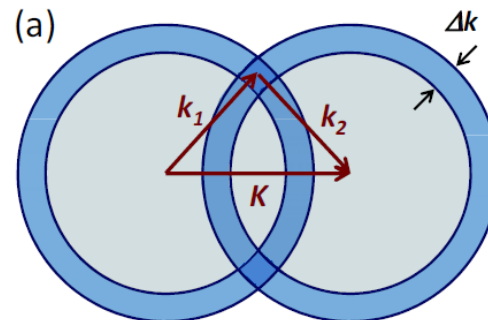
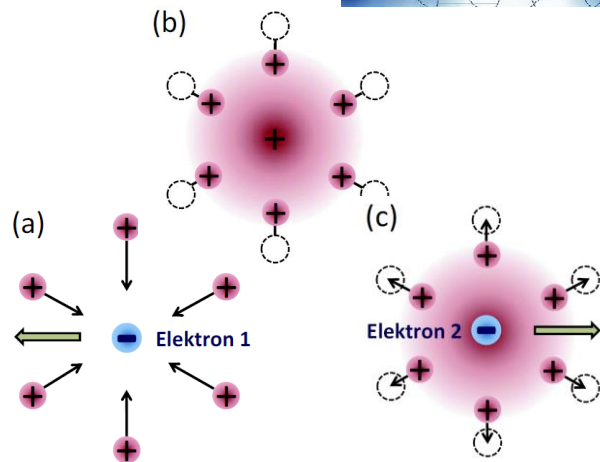
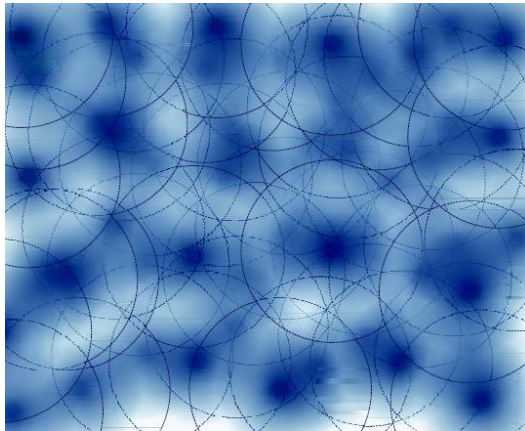
Typ-I Supraleiter

$$\kappa \geq \frac{1}{\sqrt{2}}$$

Typ-II Supraleiter

Wiederholung

■ Mikroskopische Theorie

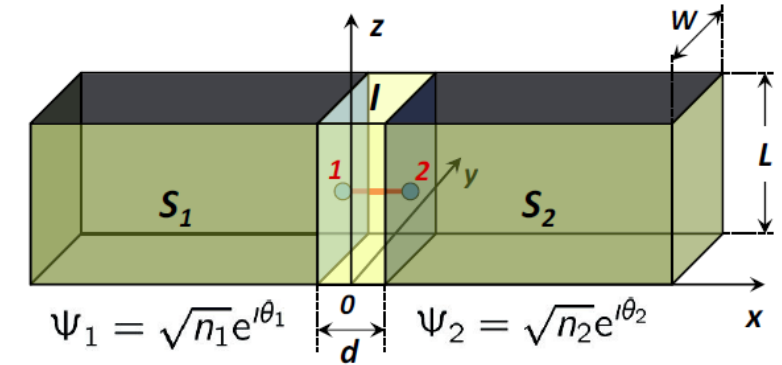


– **Ansatz** für Paarwellenfunktion: **Produkt von ebenen Wellen**

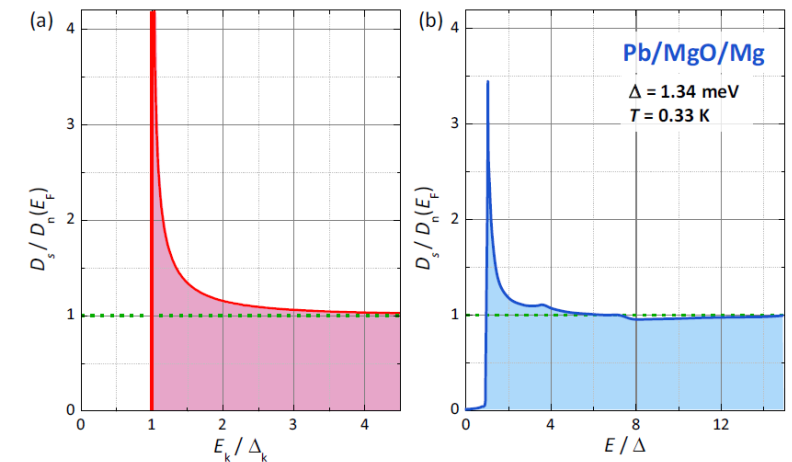
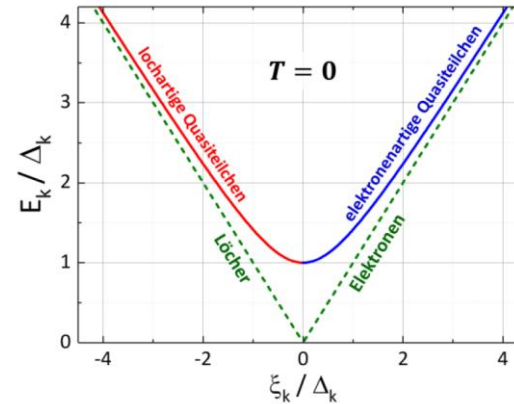
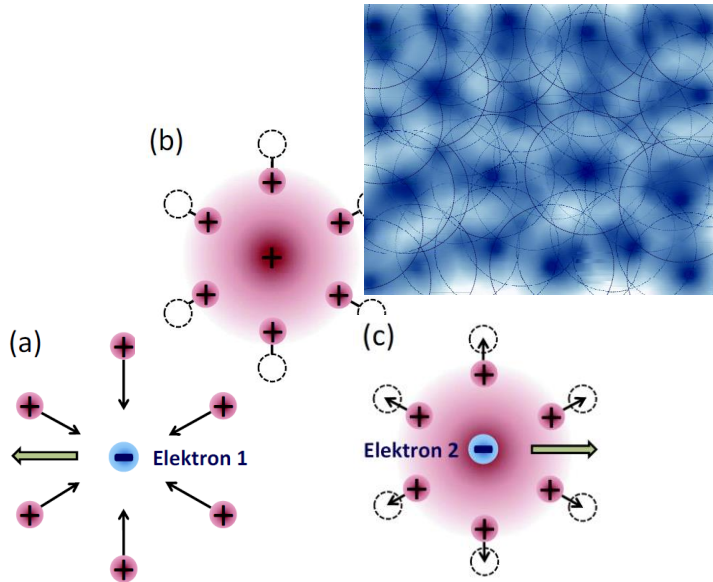
$$\Psi(\mathbf{r}_1, \mathbf{r}_2) = a \exp(i\mathbf{k}_1 \cdot \mathbf{r}_1) \exp(i\mathbf{k}_2 \cdot \mathbf{r}_2) = a \exp(i\mathbf{k} \cdot \mathbf{r})$$

Zusammenfassung

- Mikroskopische Theorie der Supraleitung
- Cooper Paare
- Grundlagen der BCS Theorie
- Energie-Lücke und Quasi-Teilchenanregung
- Josephson Effekt



$$J_s(\varphi) = J_c \sin \varphi$$



Wiederholung

- Atome im magnetischen Feld
- Hund'sche Regeln

– Berechnung der Energieänderungen ΔE_n der atomaren Energien in Störungstheorie 2. Ordnung:

$$\Delta \mathcal{H} = \frac{\mu_B}{\hbar} (L_z + g_s S_z) B_z + \frac{e^2 B_z^2}{8m} \sum_i (x_i^2 + y_i^2)$$

$$\Delta E_n = \langle n | \Delta \mathcal{H} | n \rangle + \sum_{n \neq n'} \frac{|\langle n | \Delta \mathcal{H} | n' \rangle|^2}{E_n - E_{n'}}$$

$$\begin{aligned} \Delta E_n = & \frac{\mu_B B_z}{\hbar} \langle n | L_z + g_s S_z | n \rangle \\ & + \frac{\mu_B^2 B_z^2}{\hbar^2} \sum_{n \neq n'} \frac{|\langle n | L_z + g_s S_z | n' \rangle|^2}{E_n - E_{n'}} \\ & + \frac{e^2 B_z^2}{8m} \langle n | \sum_i (x_i^2 + y_i^2) | n \rangle \end{aligned}$$

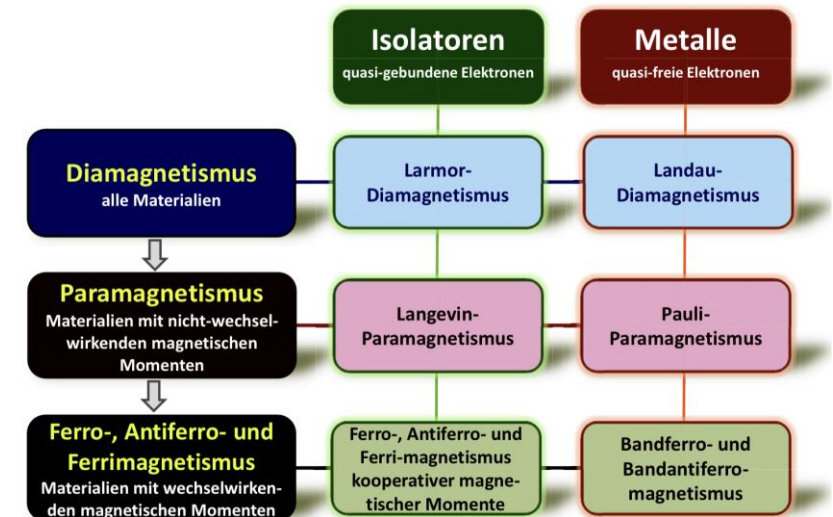
Langevin-Paramagnetismus

Van Vleck-Paramagnetismus

Larmor-Diamagnetismus

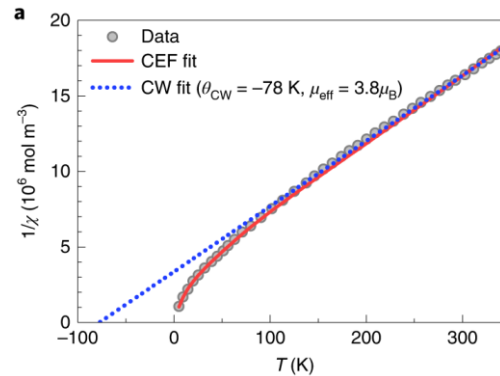
2. Ordn

- I. **Maximierung von S:**
Minimierung der Coulomb-Abstoßung: folglich gleichem Ort sein
- II. **Maximierung von L:**
Reduktion der Coulomb-Energie durch gleichzeitige halbe Füllung liegt $L = 0$ vor (wegen I)
- III. **Kopplung von L und S zu J:**
Minimierung der Spin-Bahn-WW:



Wiederholung

■ Magnetisierungskurven, Brillouin-Funktion

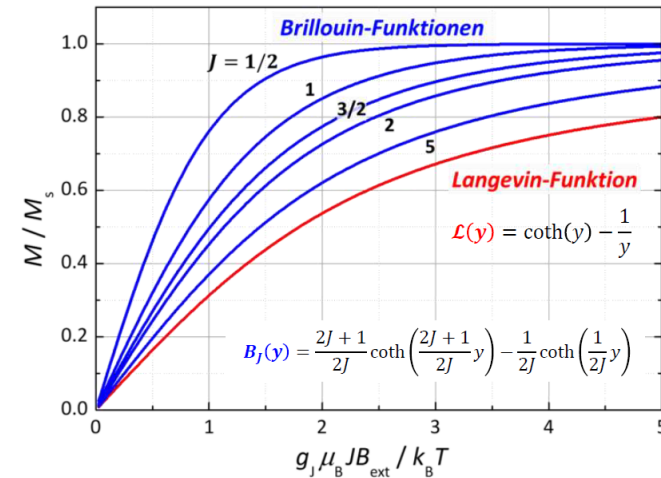


■ Curie-Gesetz

Suszeptibilität

$$\chi = \mu_0 \cdot \left(\frac{\partial M}{\partial B} \right)_{T, V} \approx \frac{\mu_0 \cdot \mu_B^2}{k_B T} = \frac{C}{T}$$

Curie-Gesetz



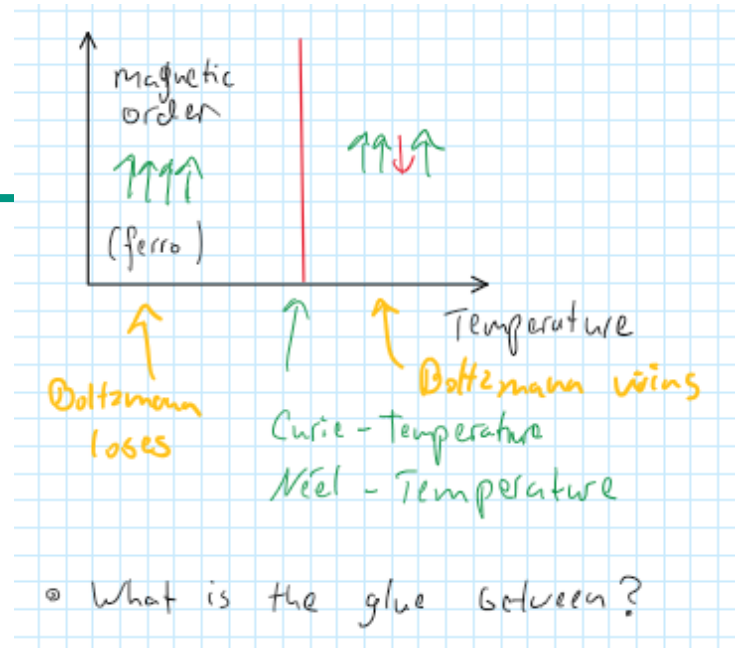
■ Pauli-Paramagnetismus



$$\chi_P = \mu_0 \left(\frac{\partial M}{\partial B_{\text{ext}}} \right)_{V, T} = \frac{3n\mu_0\mu_B^2}{2k_B T_F}$$

Wiederholung

■ Magnetische Ordnung



■ Magnetische Dipol-Wechselwirkung

3.1. Magnetic Dipole Interaction

A diagram showing two magnetic dipoles. Each dipole is represented by a small rectangle with a green top half and a red bottom half, surrounded by a circular loop of arrows indicating the magnetic field lines.

$$E_D = \frac{\mu_0}{4\pi} \cdot \frac{1}{r^3} \left[\vec{m}_1 \cdot \vec{m}_2 - 3(\vec{m}_1 \cdot \hat{r})(\vec{m}_2 \cdot \hat{r}) \right]$$

MDI is anisotropic

■ Austausch-Wechselwirkung

$$E^s(R) = 2E_{1s} + \frac{V + A}{1 + S_{AB}^2}$$

$$E^a(R) = 2E_{1s} + \frac{V - A}{1 - S_{AB}^2}$$

Wiederholung

Landau Diamagnetismus

– für freies Elektronengas (ohne Beweis)

$$\chi_L = -\frac{1}{3} \chi_P, \quad \chi = \chi_L + \chi_P = \frac{n\mu_0\mu_B^2}{k_B T_F}$$

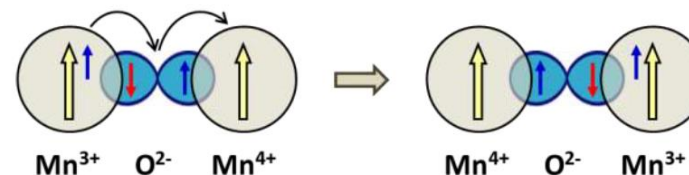
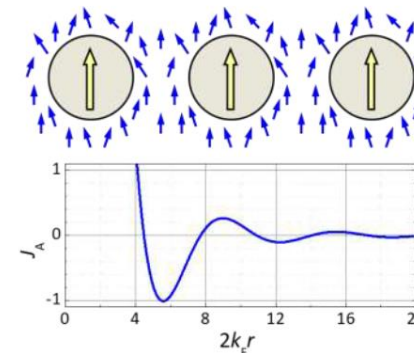
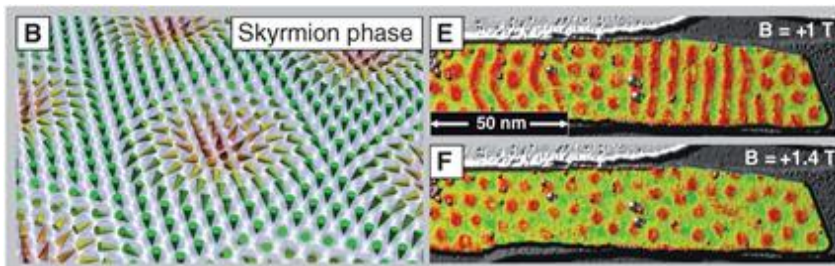
Austauschwechselwirkung

$$\chi_L = -\frac{1}{3} \chi_P \left(\frac{m}{m^*} \right)^2$$

Verschiedene Arten der Austauschwechselwirkung

Stoner Kriterium

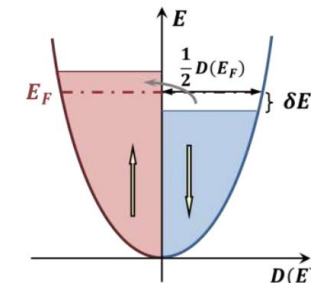
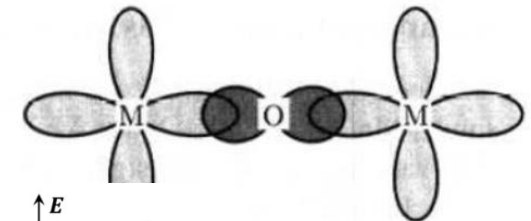
$$\mathcal{H}_{DM} = \mathbf{D}_{ij} \cdot (\mathbf{S}_i \times \mathbf{S}_j)$$



Super Exchange

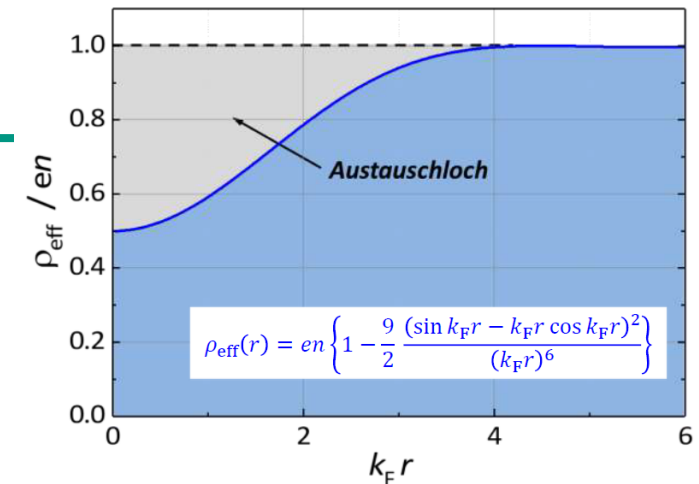
antiferromagnetic

GS	(a)	↑	↑↑	↑
ES1	(b)	↑↑	↑↑	
ES2	(c)	↑↑		↓↓



Wiederholung

■ Austauschloch



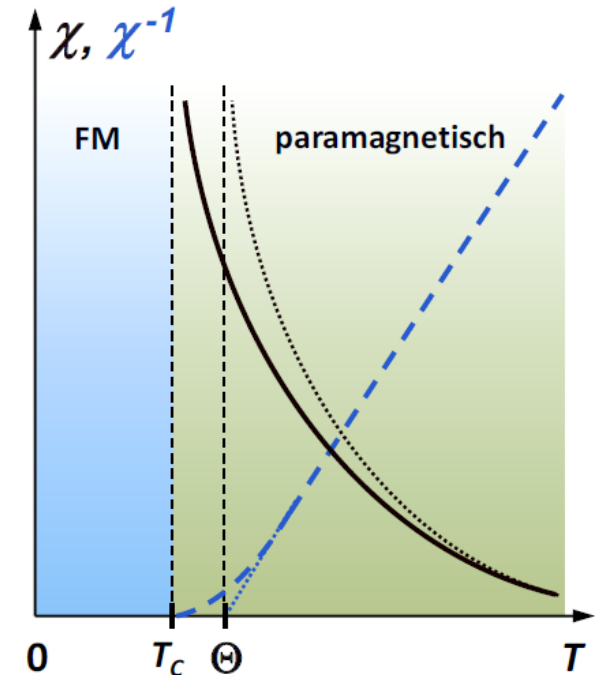
■ Curie-Weiss Gesetz

$$\chi = \mu_0 \left(\frac{\partial M}{\partial B_{\text{ext}}} \right)_{T,V} = \frac{C}{T - T_C}$$

$$C = \frac{n \mu_0 g_J^2 J(J+1) \mu_B^2}{3k_B}$$

$$T_C = \gamma \frac{n \mu_0 g_J^2 J(J+1) \mu_B^2}{3k_B} = \gamma C$$

■ Ferromagnetismus



Wiederholung

■ Eigenschaften von Halbleitern

■ Leitungsband und Valenzband

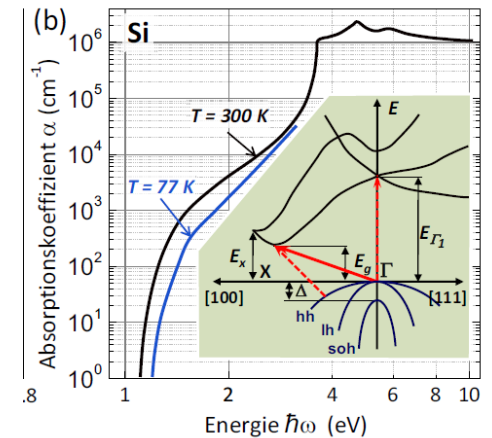
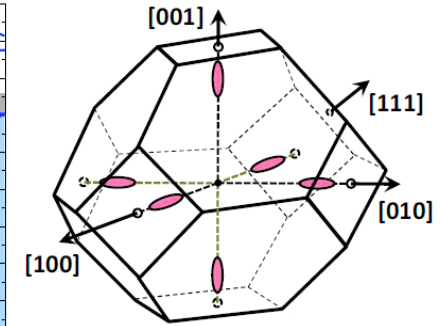
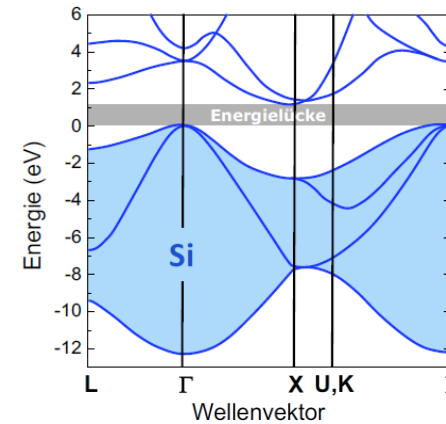
$$E(\mathbf{k}) = E_c + \frac{\hbar^2}{2} \sum_{ij} k_i \left(\frac{1}{m^*} \right)_{ij} k_j \quad (\text{Elektronen})$$

$$E(\mathbf{k}) = E_v + \frac{\hbar^2}{2} \sum_{ij} k_i \left(\frac{1}{m^*} \right)_{ij} k_j \quad (\text{Löcher}).$$

■ Ladungsträgerdichte von intrinsischen Halbleitern

$$\mu = E_v + \frac{1}{2} E_g + \frac{3}{4} k_B T \ln \frac{m_{h,\text{DOS}}^*}{m_{e,\text{DOS}}^*}.$$

$$n_i = \sqrt{n_c \cdot p_v} = \sqrt{n_c^{\text{eff}} p_v^{\text{eff}}} e^{-E_g/2k_B T}$$

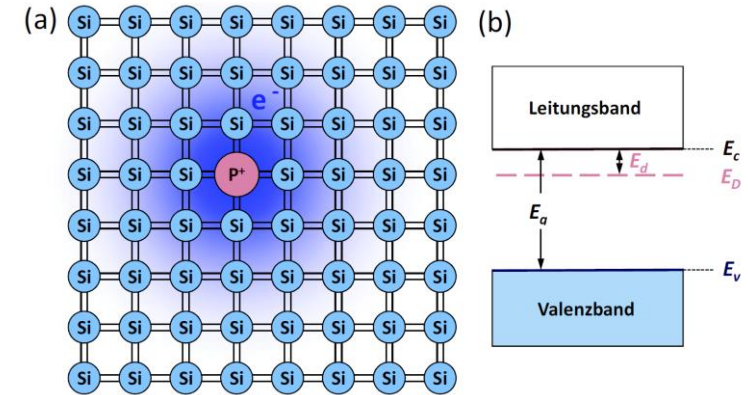


Wiederholung

Eigenschaften von dotierten Halbleitern

Wasserstoffmodell der Dotieratome

Temperaturabhängigkeit



$$r_d = \frac{4\pi\epsilon\epsilon_0\hbar^2}{m_e^*e^2}$$

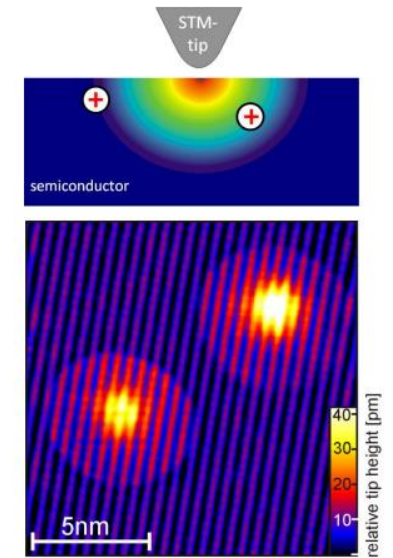
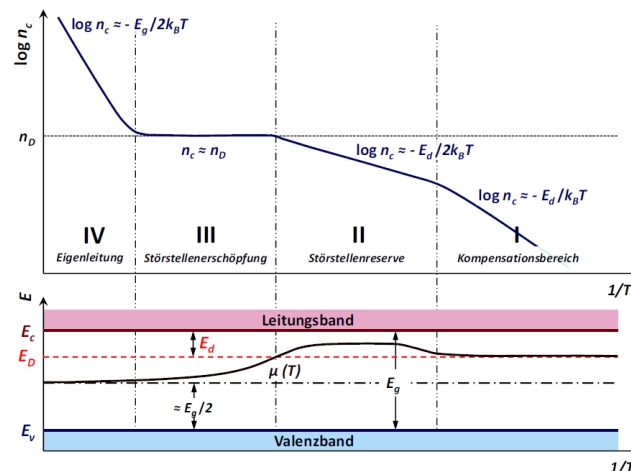
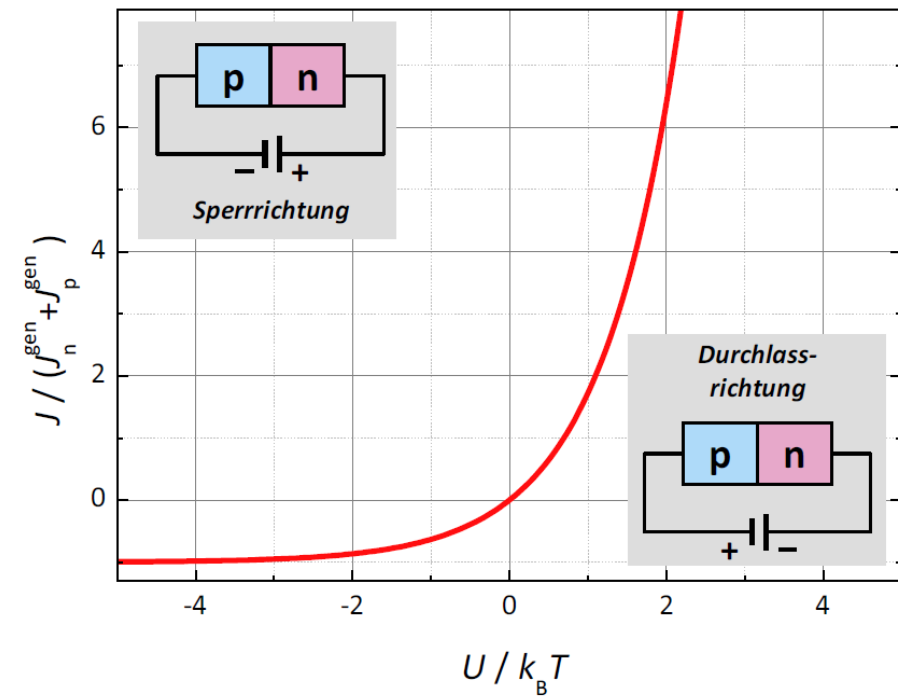
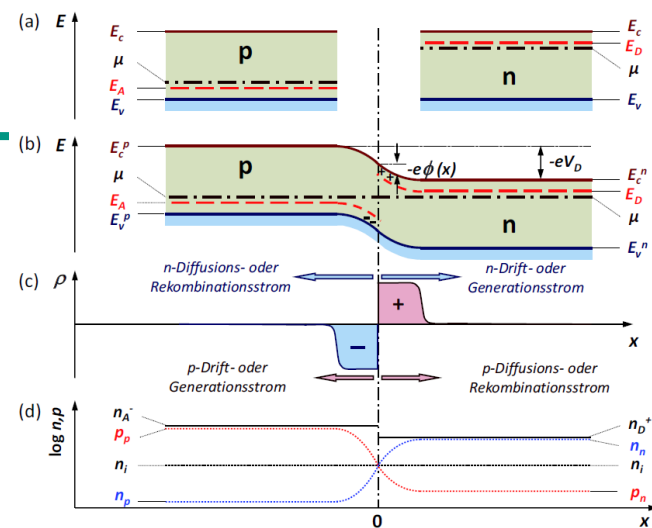


Figure 1.2: The upper image sketches two donors in GaAs and the STM tip. The colored area indicates the space charge region. The lower image shows a constant current topography of two donors. The charge switching is visible by the disk-shape of enhanced topographic height.

Wiederholung

pn-Übergang

Solarzelle



Halbleiterheterostrukturen

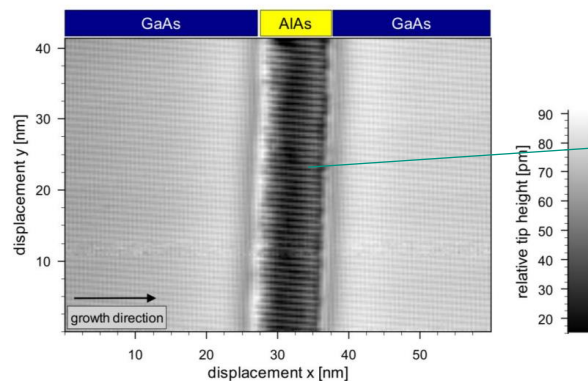
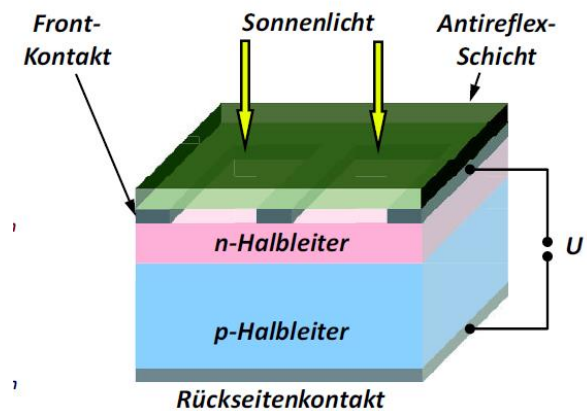
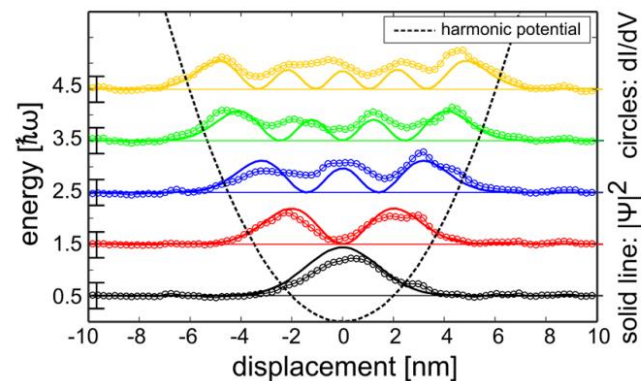
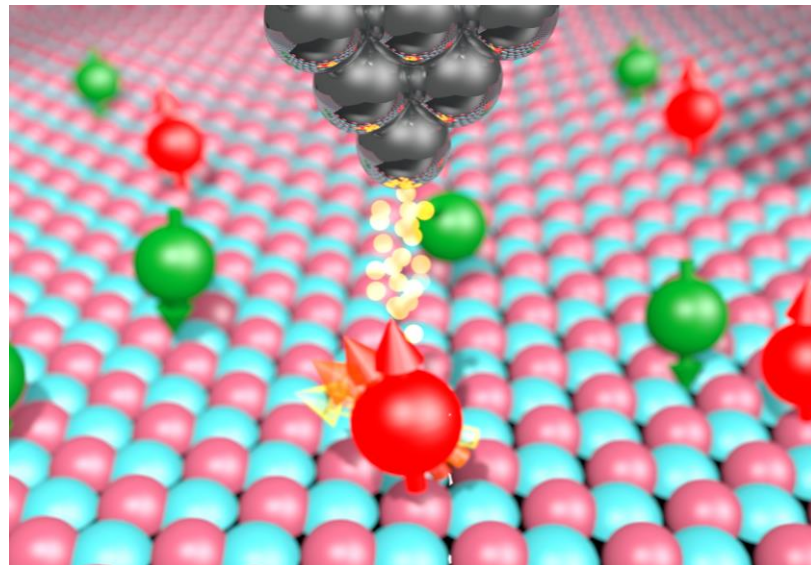


Figure 4.3: Constant current topography of the AlAs layer (marked yellow at the top of the image) embedded in GaAs (marked blue) taken at a setpoint of -2V and 0.1nA.



Werbung

■ Vorlesungen im nächsten Semester:



4021011
4 SWS
Englisch

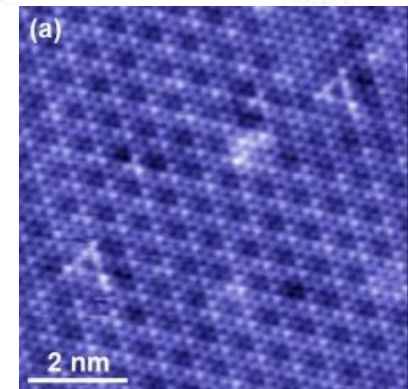
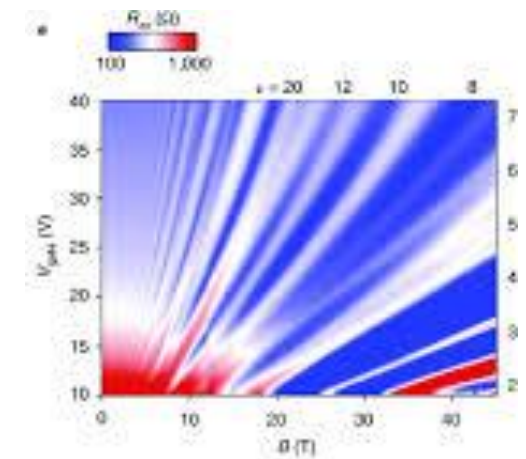
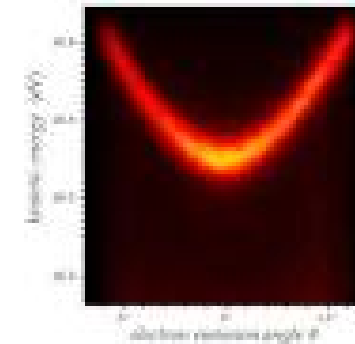
Electronic Properties of Solids I

● Präsenz

Di 09:45-11:15 30.22 Physik-Hörsaal Nr. 4 (Kl. HS B) vom 22.10. bis 11.02.
Fr 09:45-11:15 30.22 Physik-Hörsaal Nr. 4 (Kl. HS B) vom 25.10. bis 14.02.

Vorlesung (V)

Le Tacon, Matthieu
Willke, Philip



4013134
2 SWS

Advanced Seminar: Quantum Science at the Atomic Scale: Advanced Scanning Probe Techniques

● Präsenz

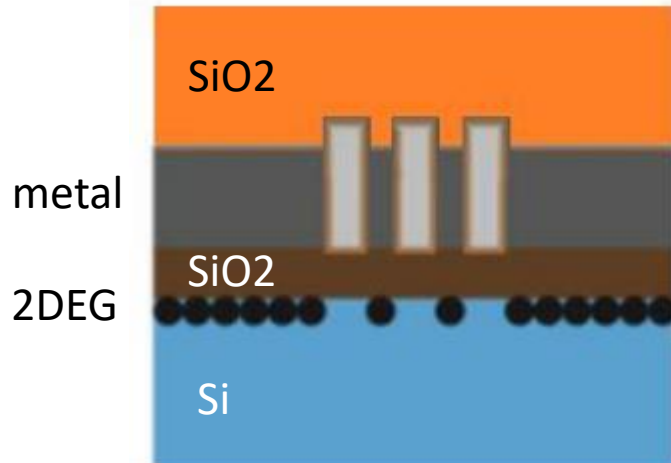
Do 14:00-15:30 30.23 Raum 3/1 vom 24.10. bis 13.02.

Hauptseminar (HS)

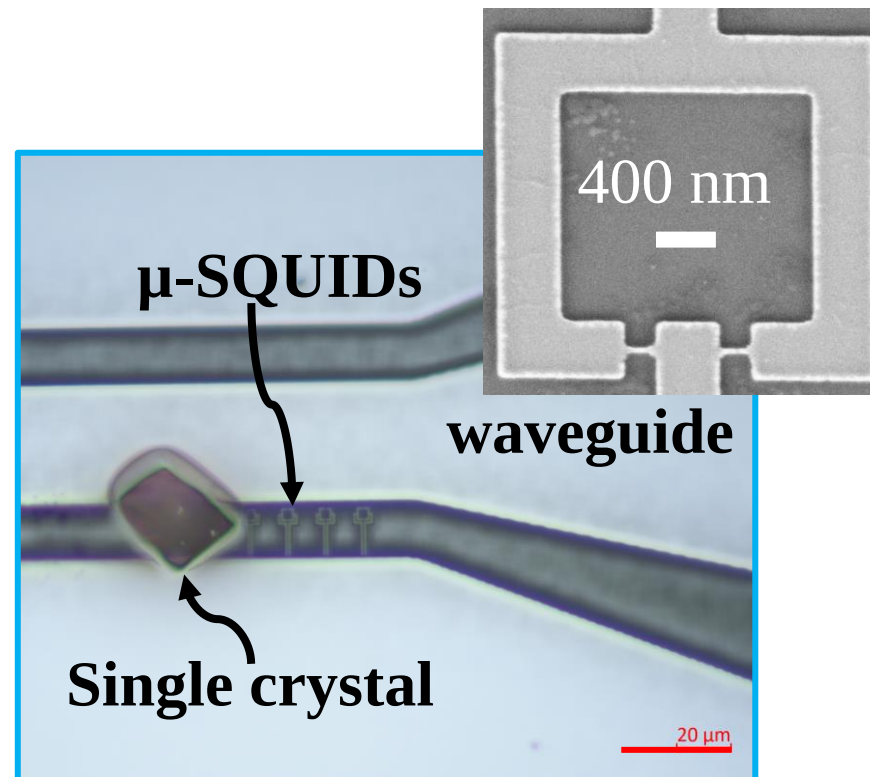
Willke, Philip
Wernsdorfer, Wolfgang
Wulfhekel, Wulf

Laborführungen

Daniel
(Semiconductor+Spins)



Appu
(SC + Spins)



Wantong / Máté
Spins + STM

