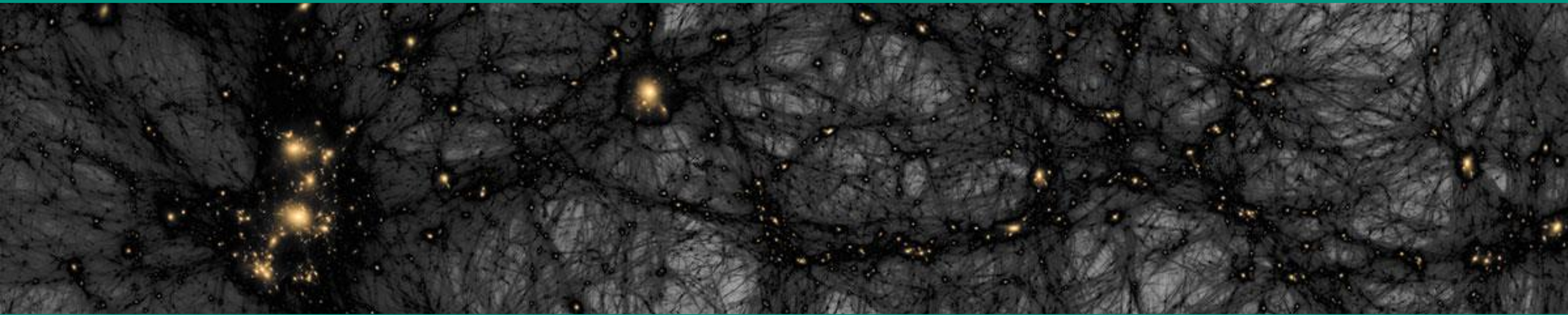


Astroparticle physics I – Dark Matter

WS22/23 Lecture 21

Feb. 15, 2023



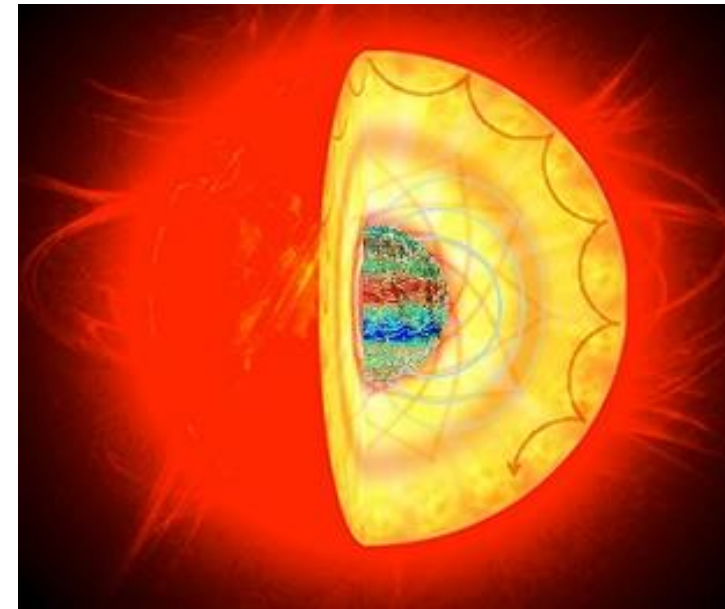
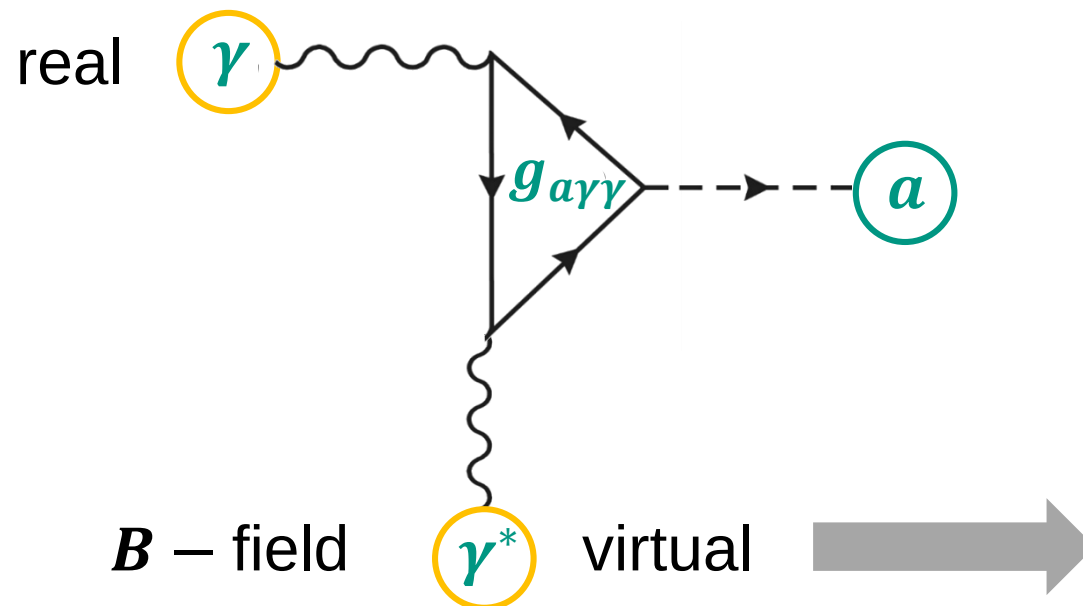
Recap of Lecture 20

- Axions: a non-thermal *WISP* as a ‘totally different’ *DM* – candidate
 - axions introduced to solve the **strong *CP* – problem**: a light pseudo-scalar
 - **mass m_a** depends on (unknown) **scale f_a** of spontaneous breaking of **$U(1)_{PC}$**
 - broad mass range possible: **$m_a = \mu eV \dots meV$** (latest: **$m_a = 40 \dots 180 \mu eV$**)
 - extension: ***ALPs*** = ***A*xion-*L*ike-*P*articles**, do not solve ***QCD* –/*DM* – problem**
 - axions form a **Bose-Einstein condensate** in the galactic ***DM*** – halo
 - various astrophysical limits: ***CDM* –requirement**, ***SN* – ν –pulse**: **$m_a < 16 meV$**

Generating axions via the Primakoff effect

■ A 'natural' process to generate axions: here - in a **magnetized plasma**

- for axion emission: will occur in magnetized plasma state from **virtual photons** γ^* to transform **real photons** γ (from the plasma) into axions via Primakoff effect
- axions then leave the **solar interior** without further interaction & can be detected

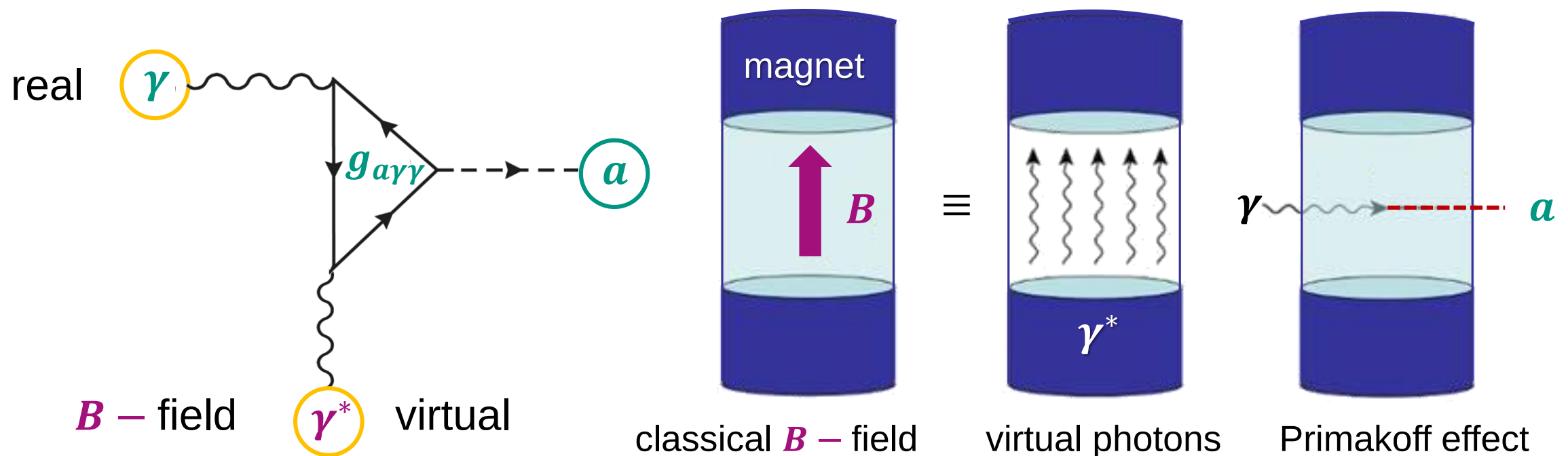


example:
**magnetised
plasma** in the
solar interior

Generating axions via the Primakoff effect

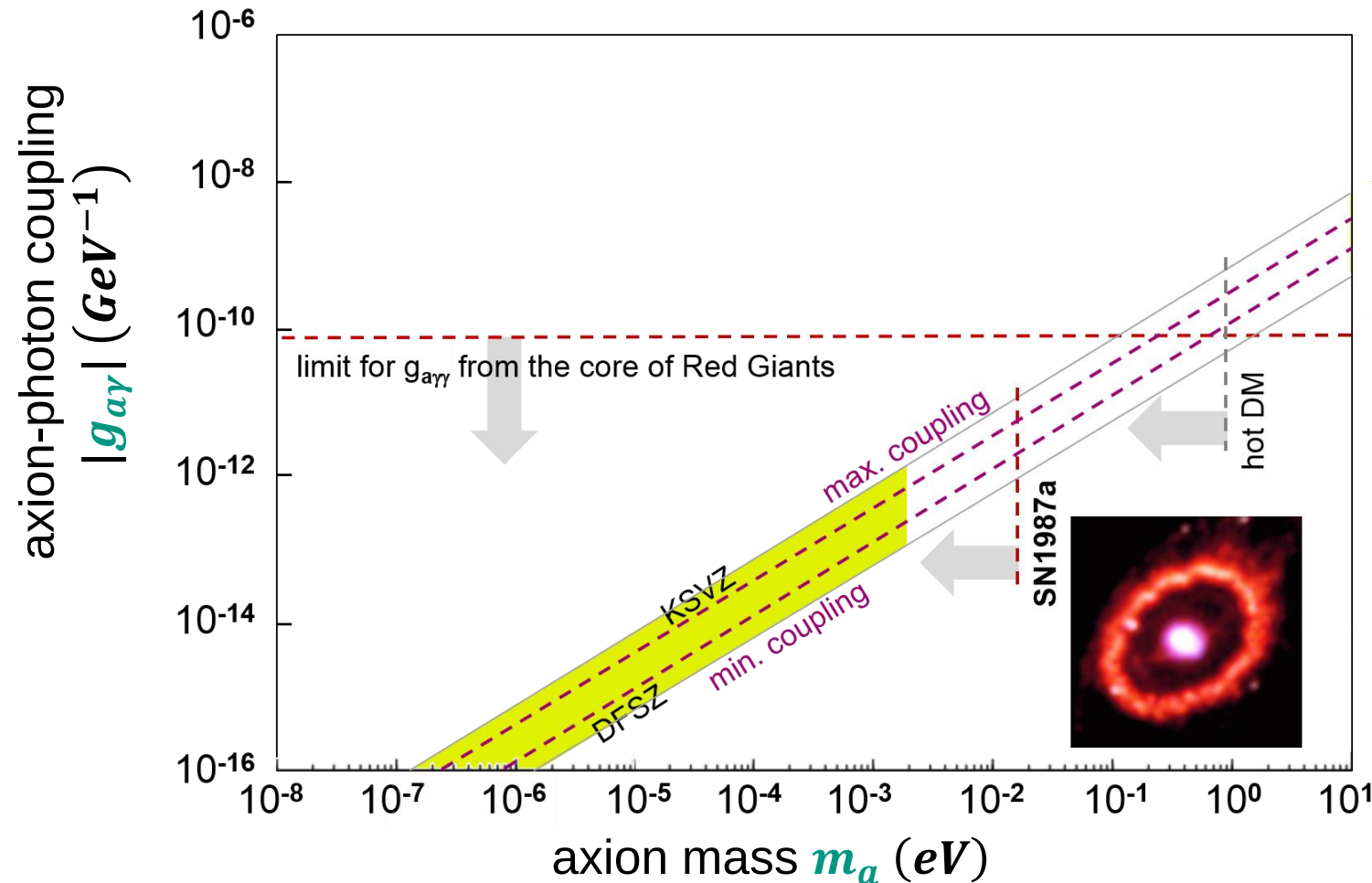
■ A **coherent** process to generate axions: here - we use a **magnet**

- the stronger the B – field (lab: up to 10 T) the more **virtual photons** γ^* we have to transform **real photons** γ into axions



The axion plot: mass m_a vs. coupling $g_{a\gamma\gamma}$

- Experimental limits on axion parameters compared to theoretical estimates



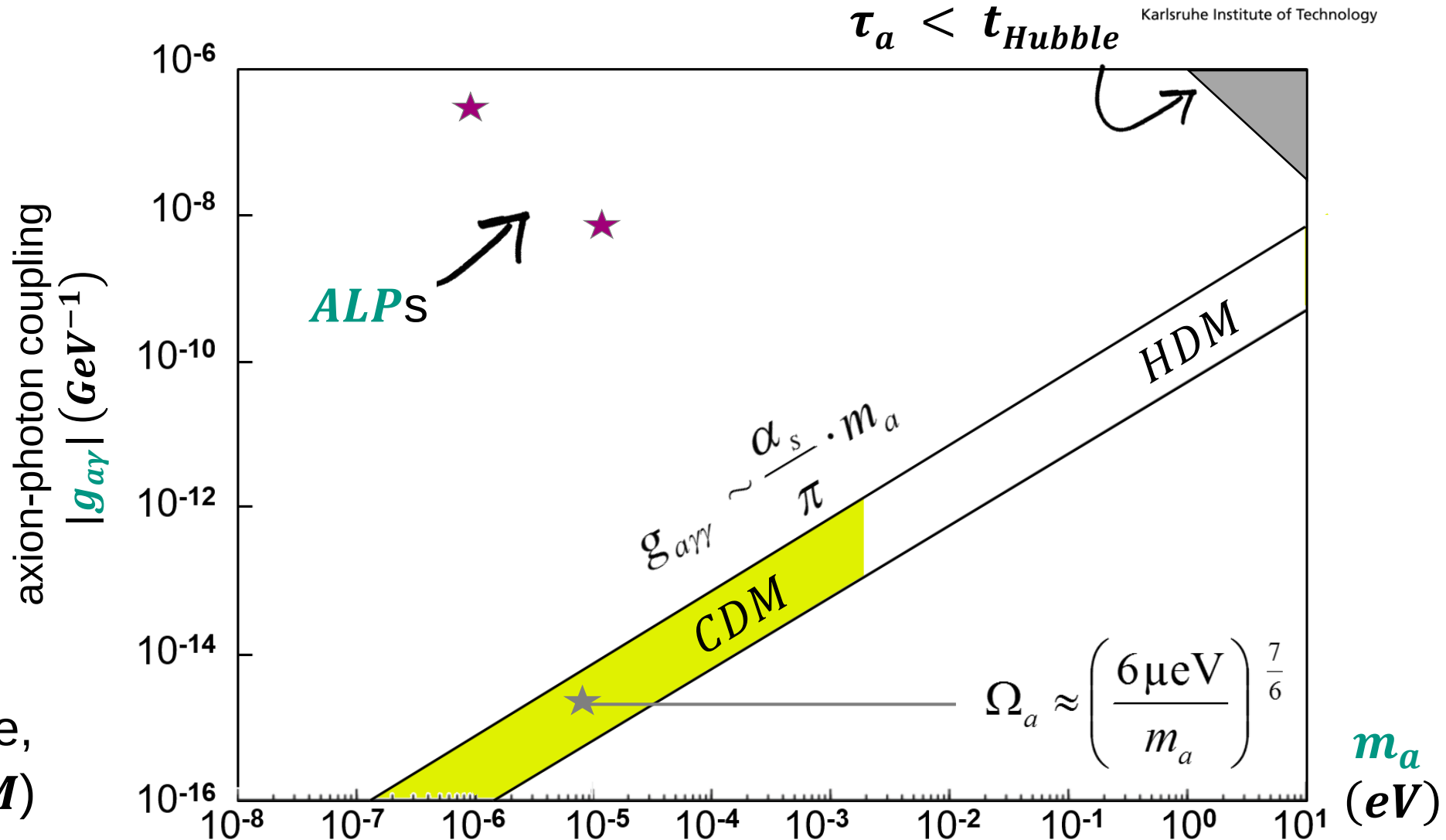
...OK, in principle it's like the *WIMP* plot we know...

mass

Axion plot: parameters for QCD – axions

■ QCD – axions

- expected parameter regions form an allowed band
- beyond:
 $ALPs$ ★
 A xion L ike
 P articles
 (no CP – relevance, little impact on DM)



Axion plot: parameters for QCD – axions

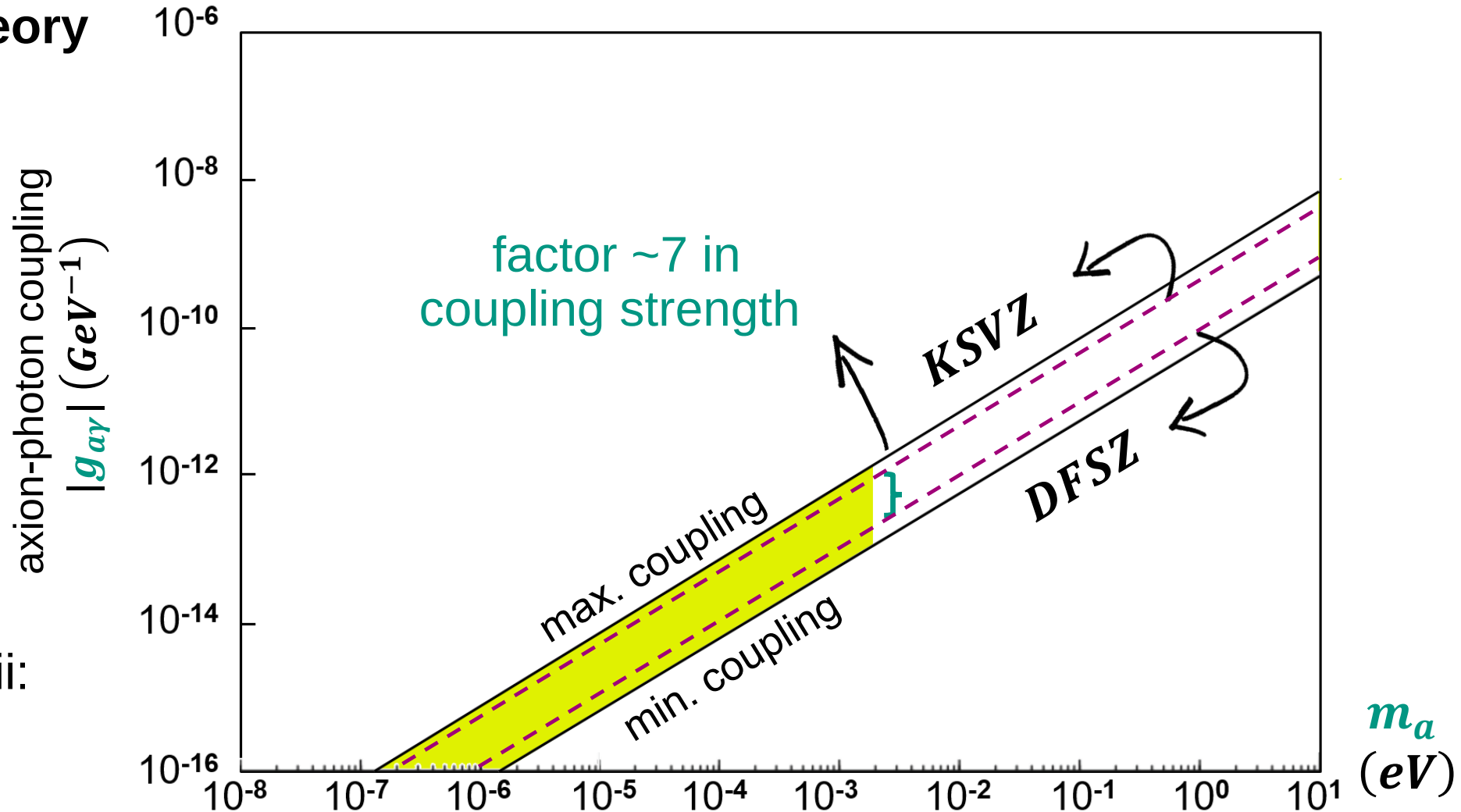
■ QCD – axion theory

- $KSVZ$

Kim-Shifman-
Vainstein-
Zakharov:
stronger axion
coupling

- $DFSZ$

Dine-Fischler-
Srednicki-Zhitnitskii:
weaker axion
coupling



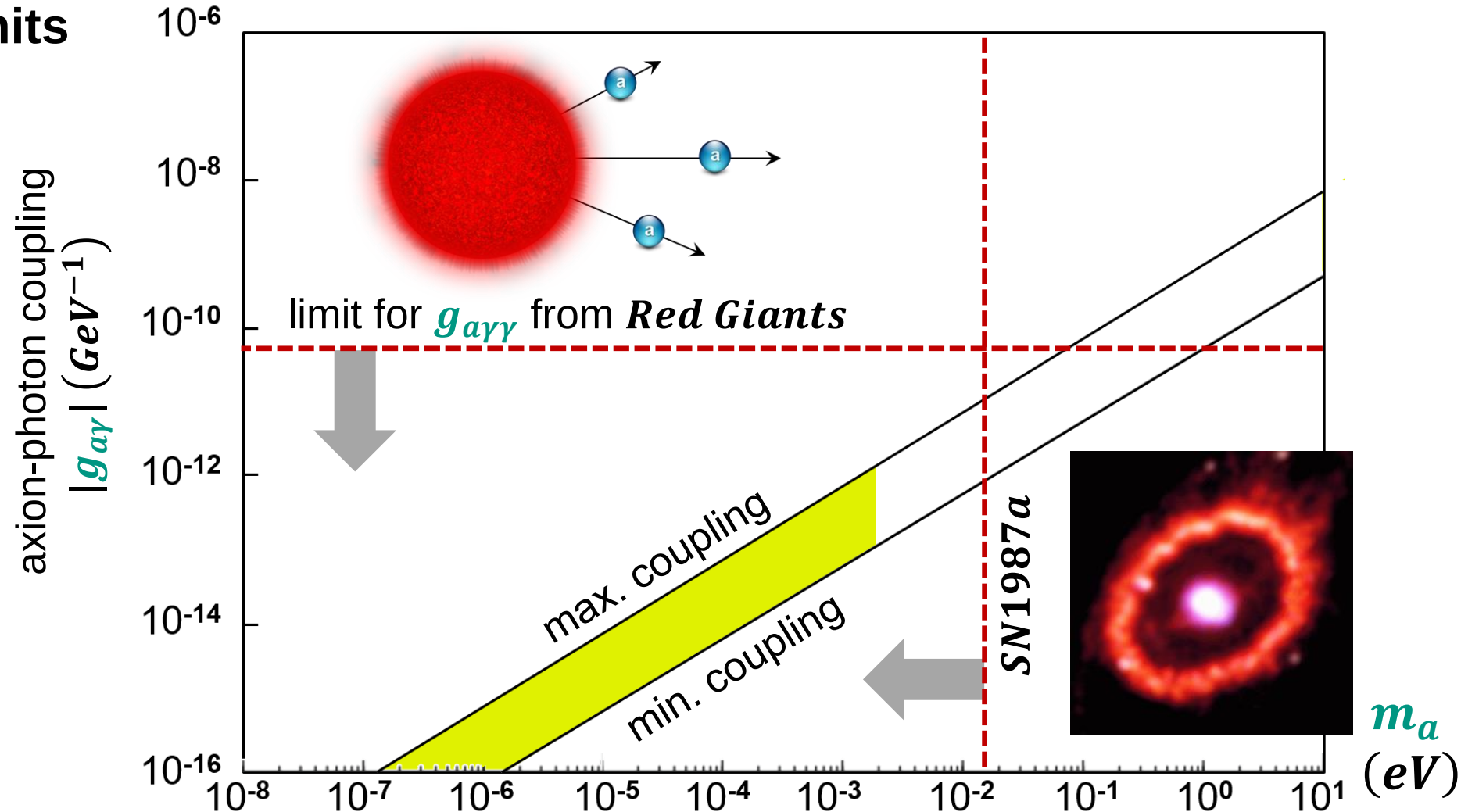
Axion plot: limits from astrophysics

■ QCD – axion limits

- limits due to observed time constants of cooling processes of stars:

Red Giants
Supernovae
White Dwarfs
star clusters

...

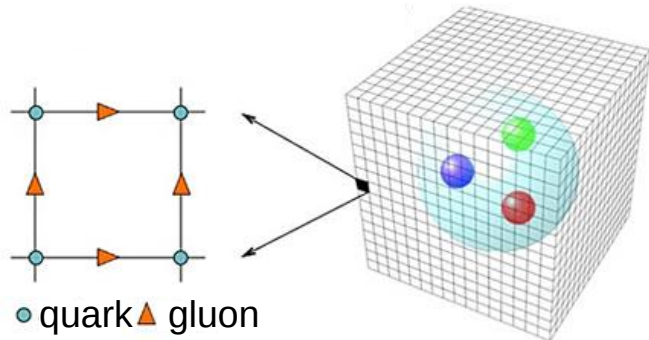


Axion plot: preferred values from lattice– QCD

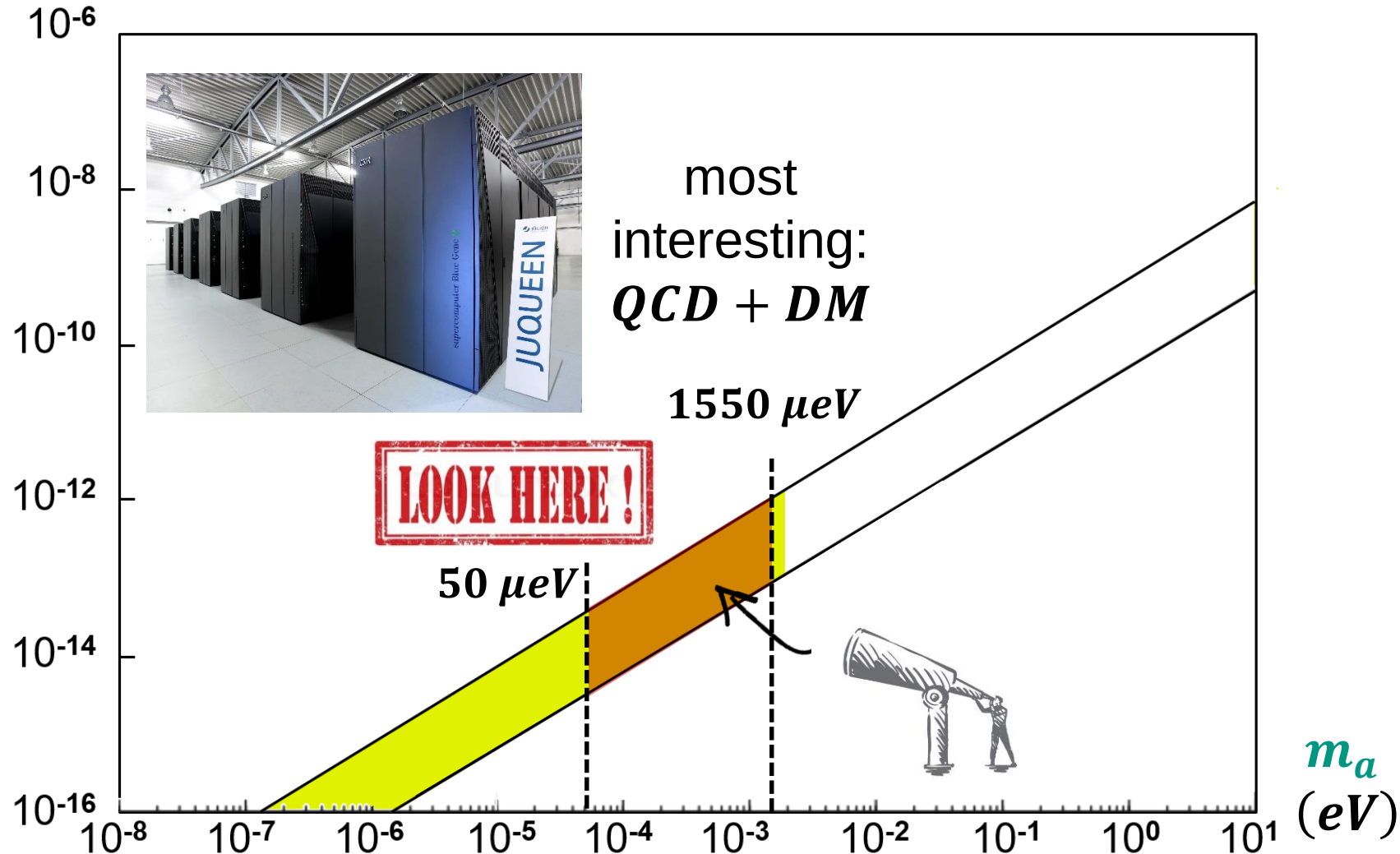
■ QCD – axion limits

- limits from modern lattice QCD-calculations
- 2022 update:

$$m_a = 40 \dots 180 \mu\text{eV}$$



axion-photon coupling
 $|g_{a\gamma}| (\text{GeV}^{-1})$



Axion experiments

■ Detecting axions from

- dark matter halo, sun
- indirectly

Microwave-cavities (haloscopes):

ADMX, HAYSTACK, ...



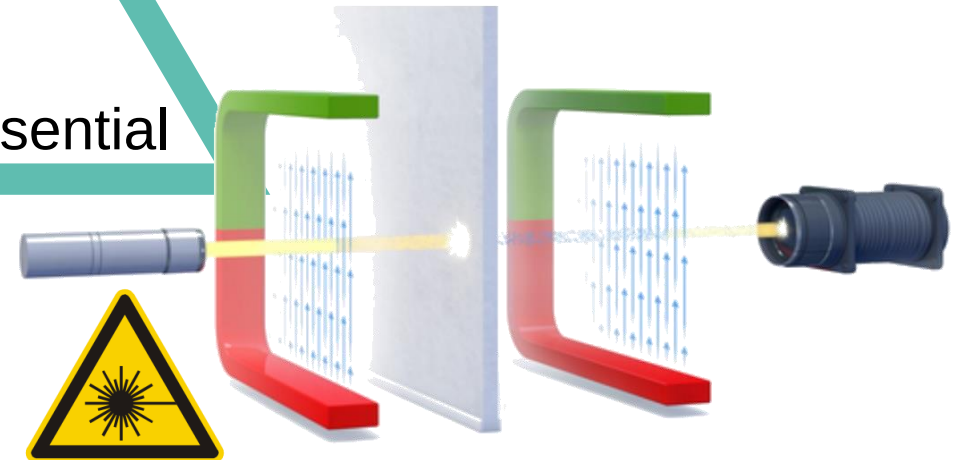
axion- (helio-) telescopes:
CAST, IAXO



Light-Shining-through-Wall (LSW):
ALPS I, II JURA



B – fields are essential

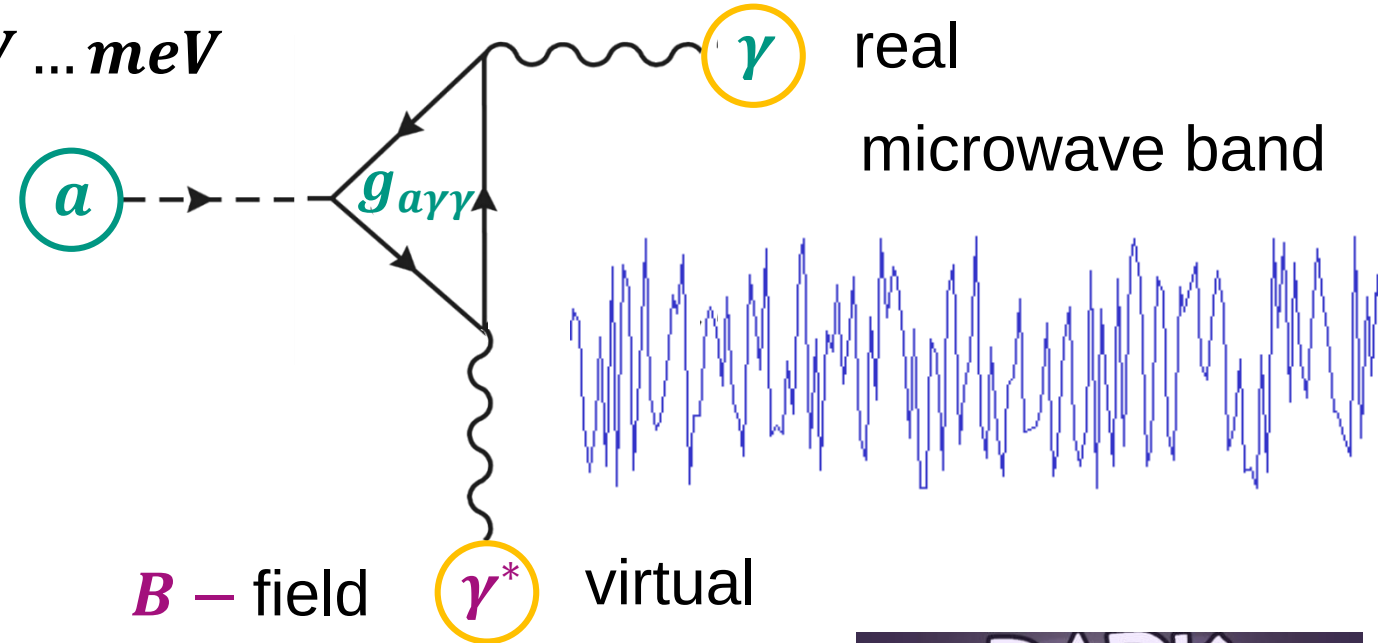


Axion experiments – haloscopes for galactic DM

■ Listening for low-mass axions from the galactic DM – halo

- sensitive mass range: $m_a = \mu eV \dots meV$

Microwave-cavities (haloscopes):
ADMX, HAYSTACK, ...

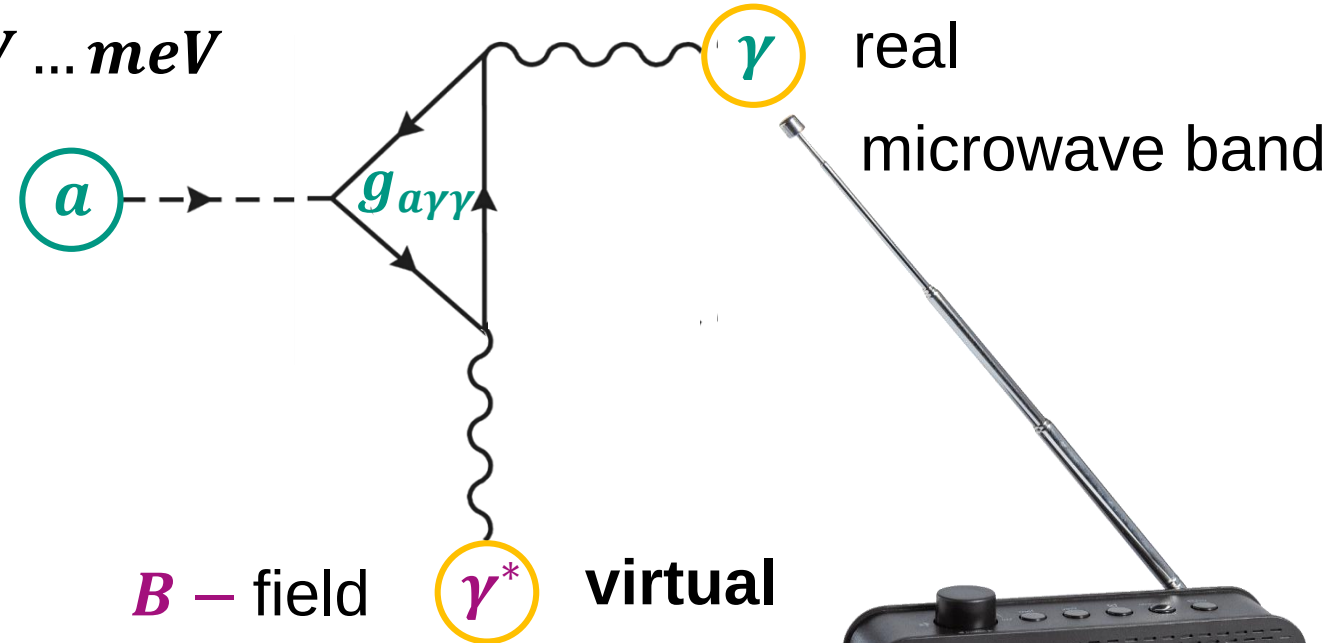


Axion experiments – haloscopes for galactic DM

■ Listening for low-mass axions from the galactic DM – halo

- sensitive mass range: $m_a = \mu eV \dots meV$

Microwave-cavities (haloscopes):
ADMX, HAYSTACK, ...



- on which frequency f
can I tune into the
galactic ' DM – radio'?

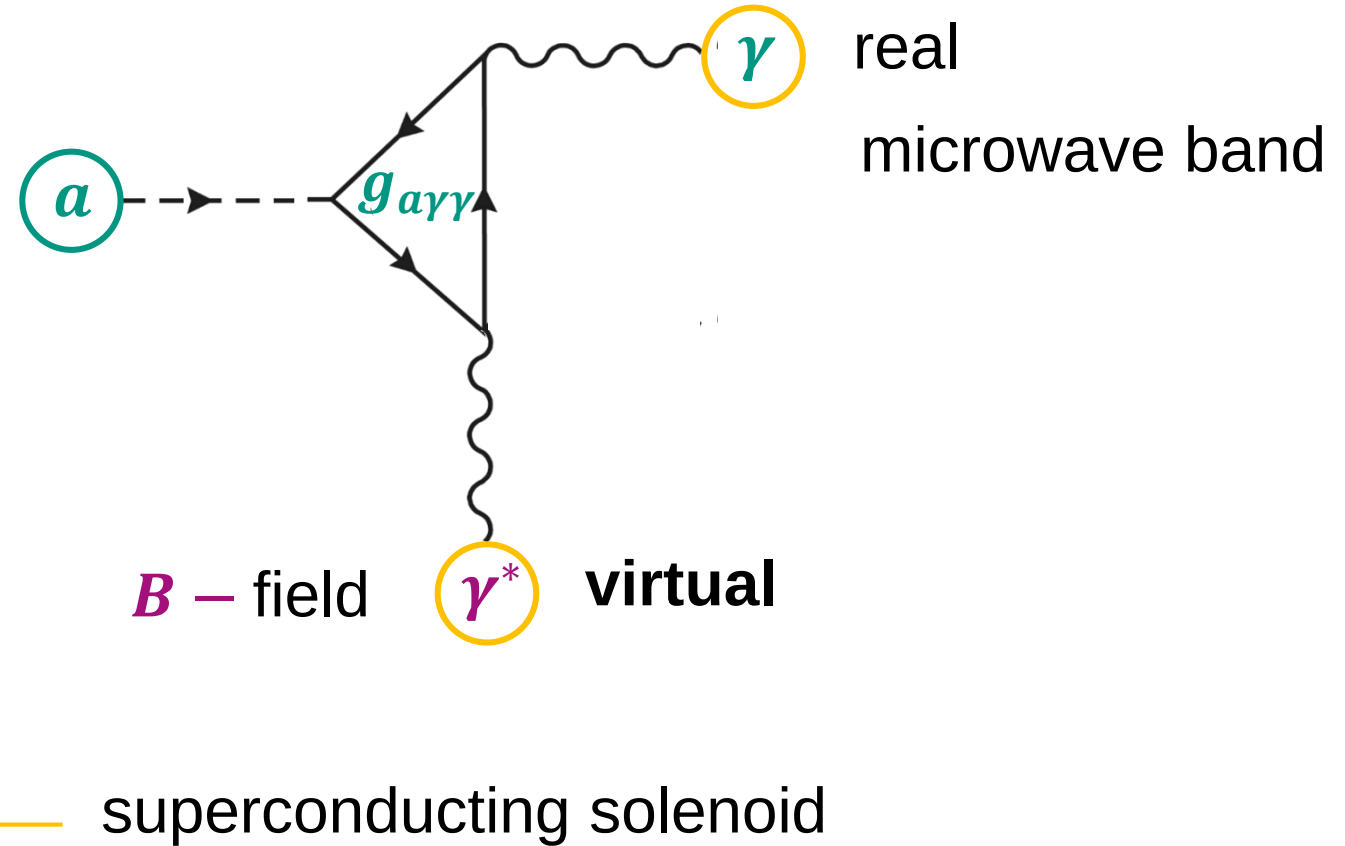
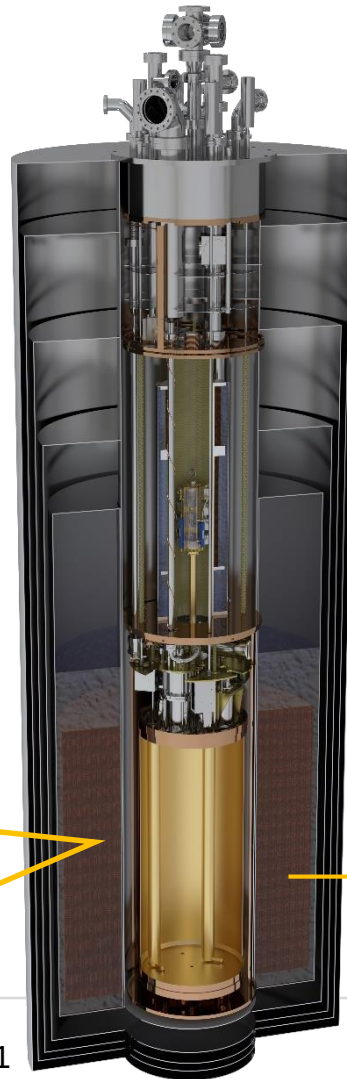
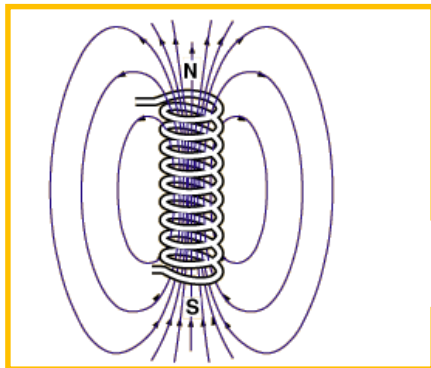


Axion experiments – haloscopes for galactic DM

■ Superconducting solenoid with $B = 6 \dots 9 \text{ T}$ & a large bore

- axion conversion rate (signal power) is a **coherent process**

$$P_{sig} \sim B^2$$

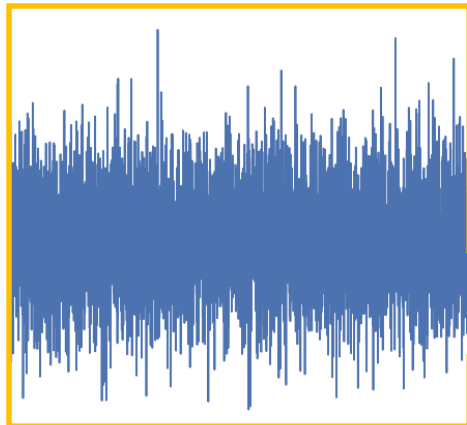


Axion experiments – haloscopes for DM – axions

- Superconducting solenoid with $B = 6 \dots 9 \text{ T}$ & a large bore at **cryogenic T**

- axion conversion rate (signal power) vs. **thermal noise**

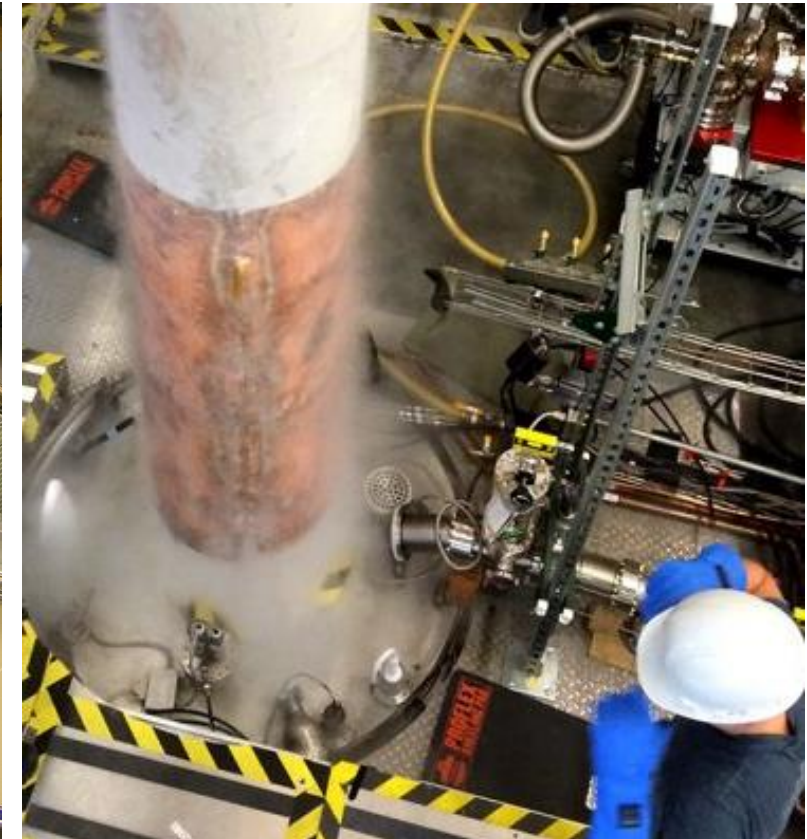
$$P_{noise} \sim T^4$$



4 K

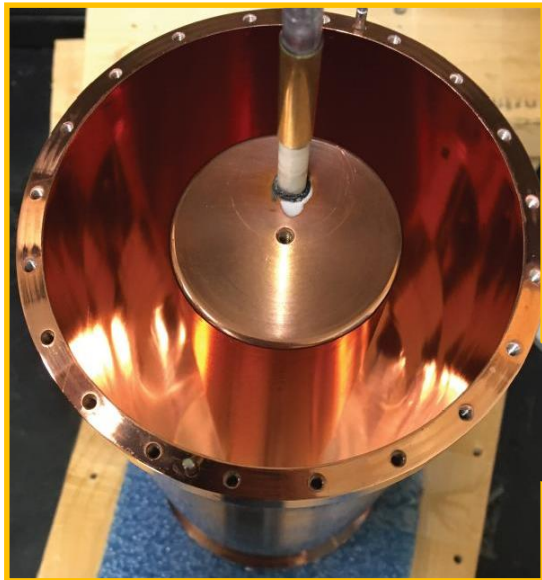
1 K

100 ...
250 mK

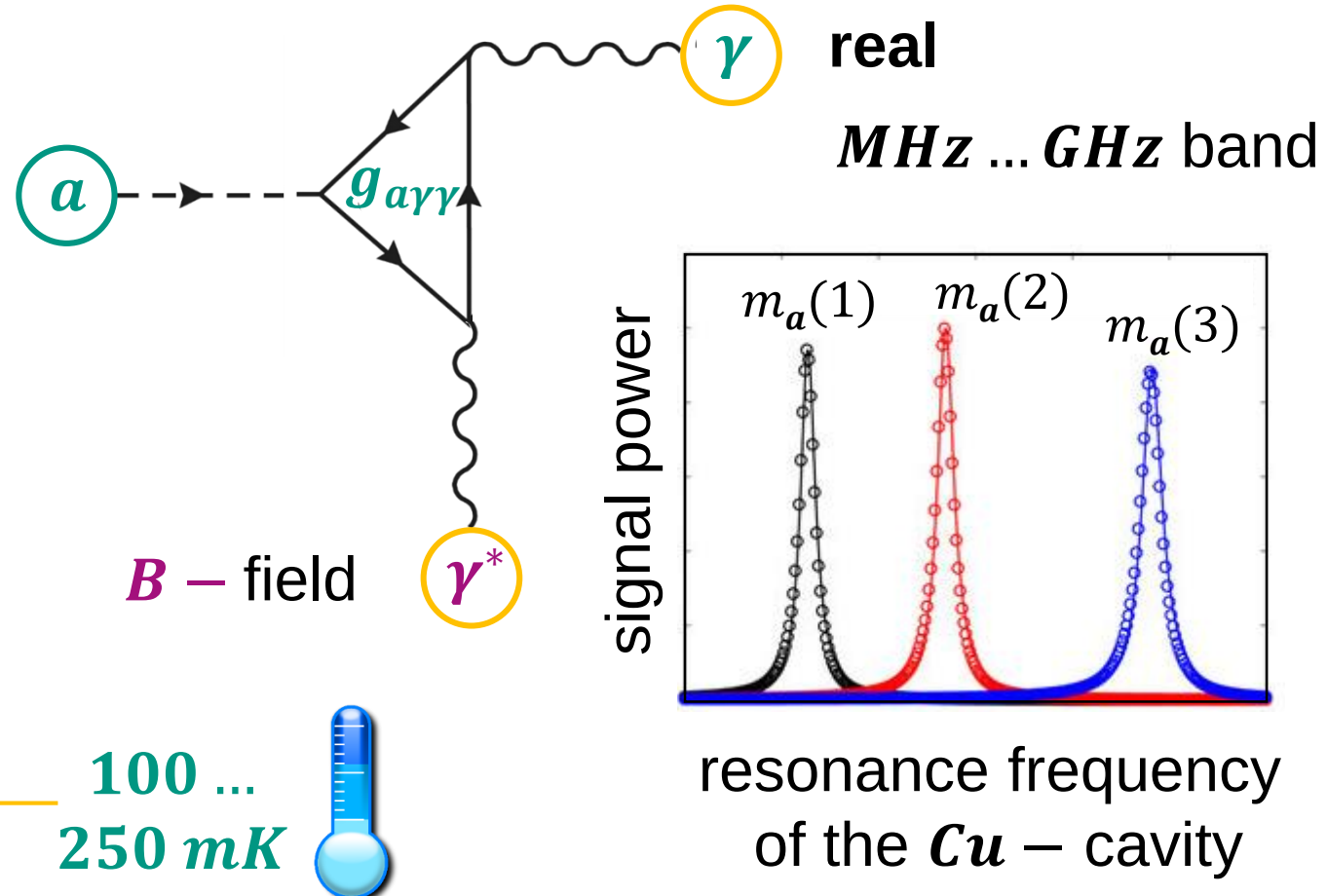
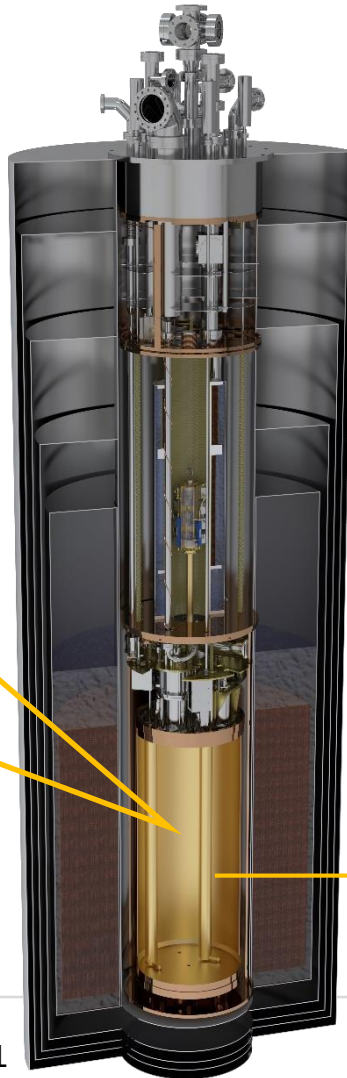


Axion experiments – haloscopes for DM – axions

- A Cu-based cylindrical resonance cavity: low thermal noise at cryogenic T



Cu – cavity & tuning rods for resonant amplification

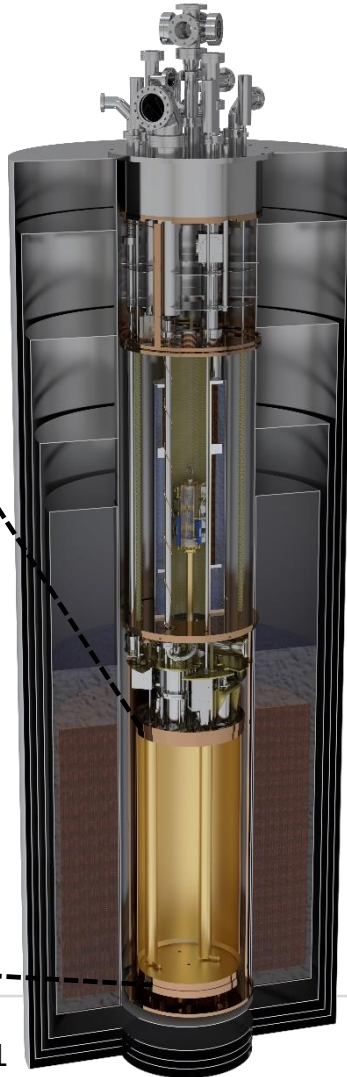


Haloscopes for DM – axions: where is f_{res} ?

■ A Cu-based cylindrical resonance cavity: modify frequencies over broad f_{res}

- theoretical axion masses:

$$m_a = 40 \dots 180 \mu eV$$



- experiment relies on **resonance** condition

axion energy $E_a = m_a \cdot c^2 + E_{kin}$



photon energy $E_\gamma = h \cdot f_{res}$

- examples

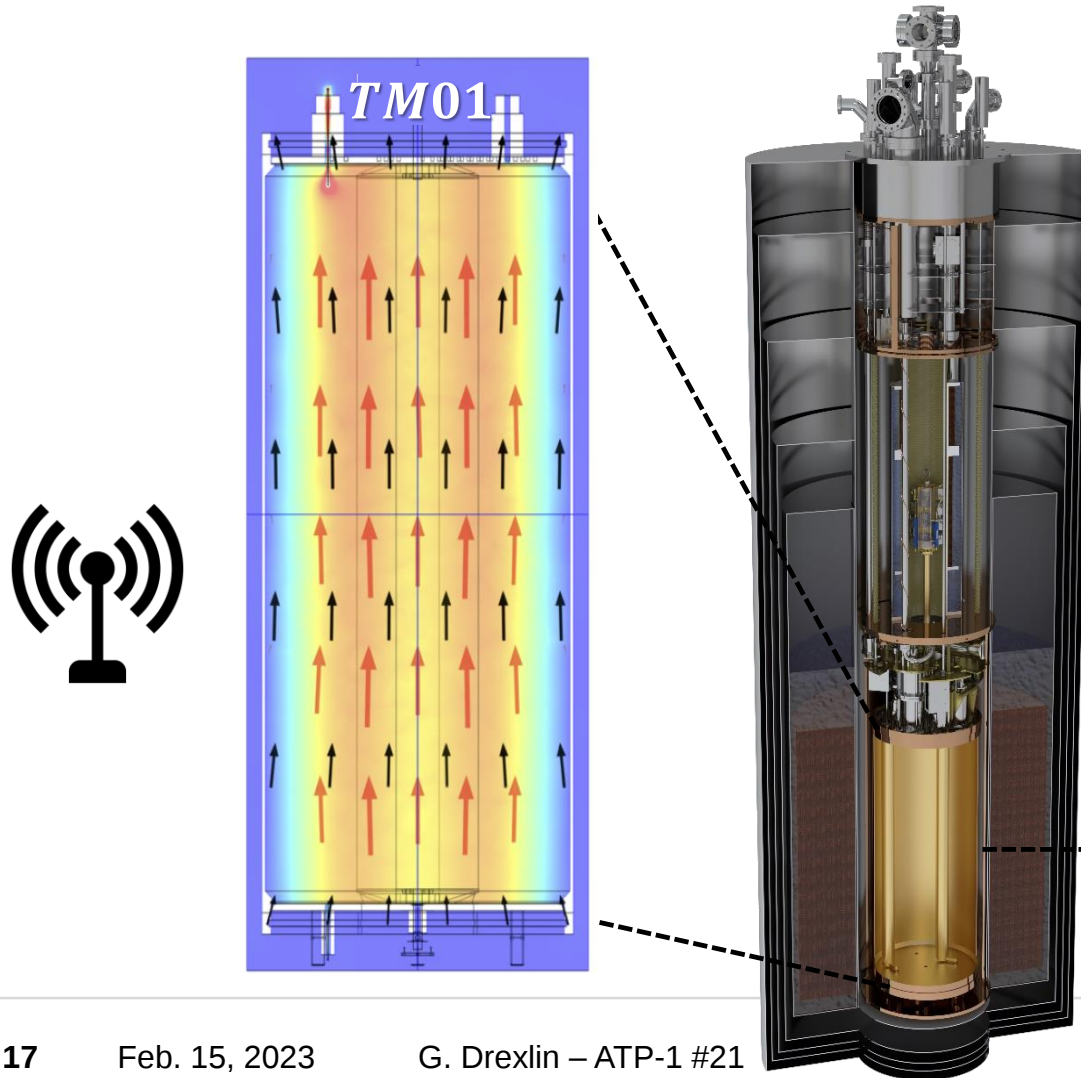
$$m_a = 10 \mu eV \Leftrightarrow f_{res} = 2.4 GHz$$

$$m_a = 100 \mu eV \Leftrightarrow f_{res} = 24 GHz$$

HF – mode

Haloscopes for DM – axions: use TM_{01} – mode

■ A Cu-based cylindrical resonance cavity with TM_{01} – mode



- excitation of the basic TM_{01} – mode
- TM_{01} – mode:
magnetic flux is transversal to propagation
- TM_{01} – mode:
 Cu – cavity acts as **wave-guide**



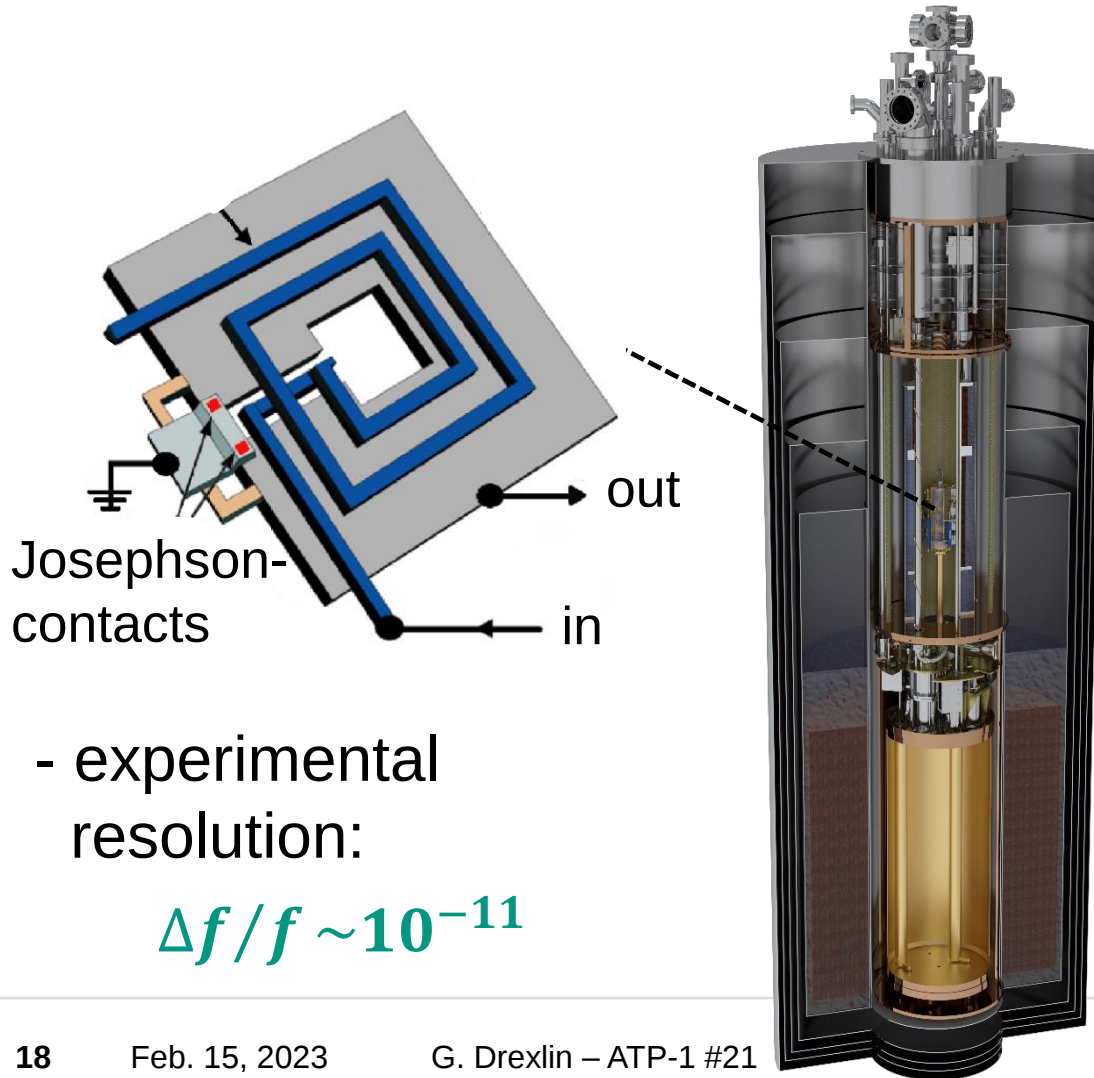
1981



Arthur Schawlow:
„Never measure anything
but **frequency**“

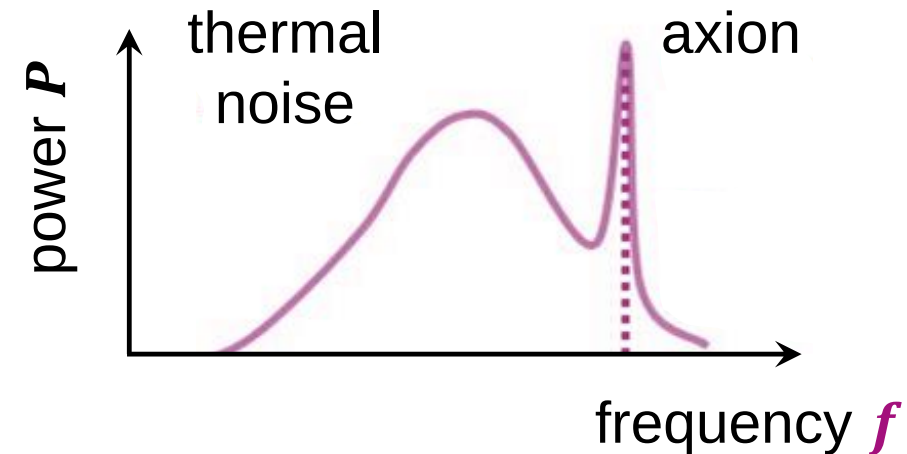
Haloscopes for DM – axions: signal & noise

■ Read-out of haloscopes via a *SQUID* – unit*



- expected signal of a haloscope:

$$m_a \cdot c^2 + E_{kin} = h \cdot f$$

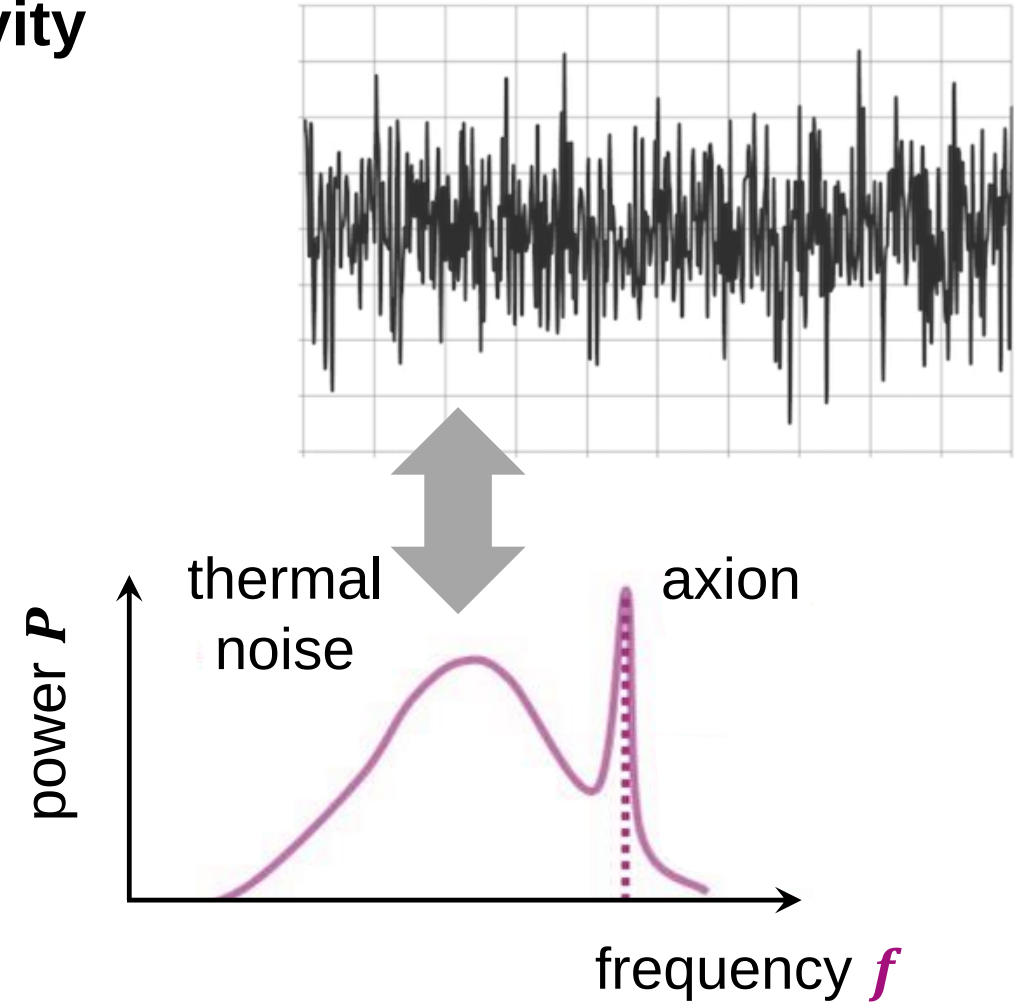
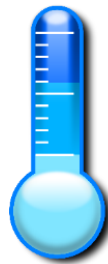


Haloscopes for DM – axions: signal & noise

- haloscopes: cooling of microwave cavity is essential to limit the overall noise

- thermal microwave photons are the dominant source of noise: $\Rightarrow$ **cooling**

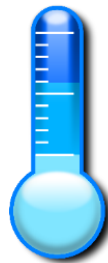
$$P_{\text{noise}} \sim T^4$$



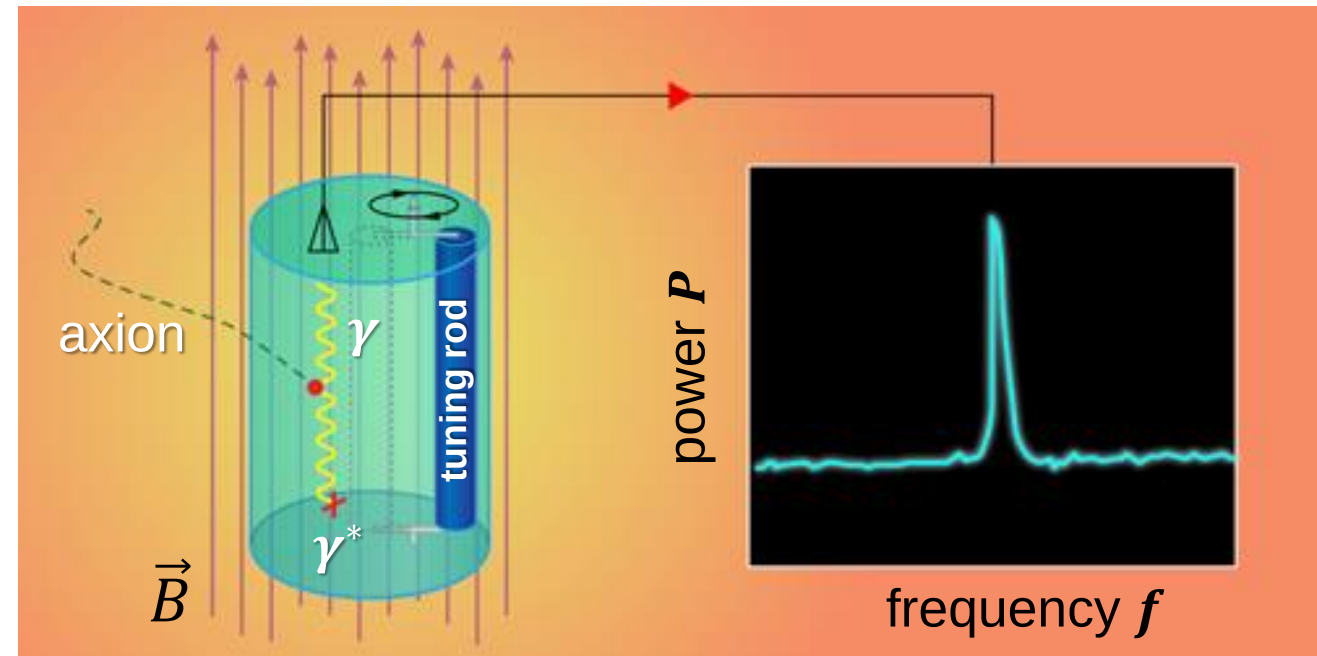
Haloscopes for DM – axions: signal & noise

- haloscopes: detection via advanced sensors (*SQUID*) & *FFT* & strong B & cooling (T)

what are the most important parameters?



signal power $P_{sig} \sim 10^{-22} \text{ W}$



Haloscopes for DM – axions: figure–of–merit

■ haloscopes: tiny (!!) expected axion **signal power** $P_{sig} \sim 10^{-22} \text{ W}$

- *figure-of-merit*:

strong B – field

$$B \approx 7.16 \text{ T}$$

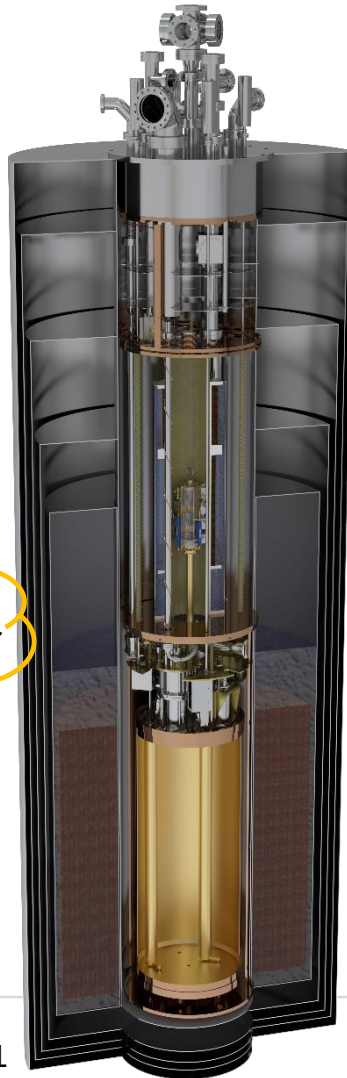
large volume V

$$V \approx 0.15 \text{ m}^3$$

high quality factor Q

$$Q \approx 10^5$$

Oops, this
is challenging



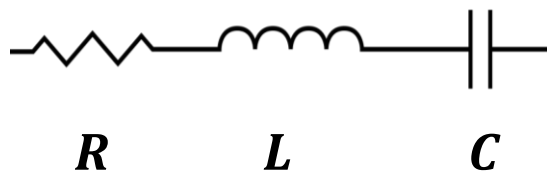
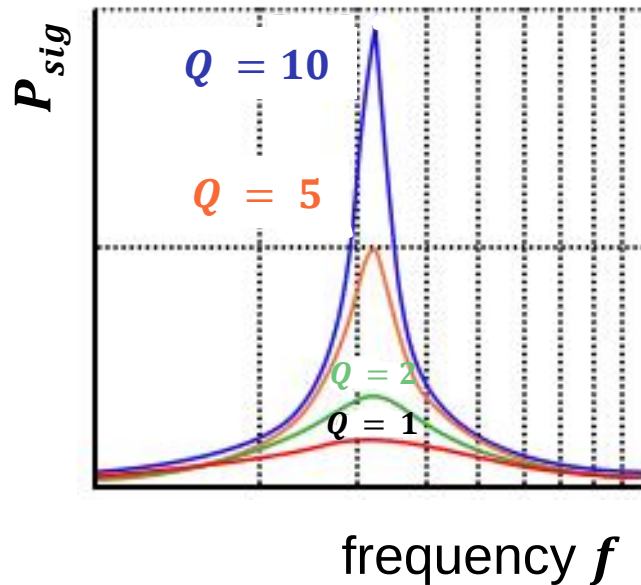
local DM – density

experimental parameters

$$P_{sig} \sim \underbrace{g_{a\gamma\gamma}^2 \cdot \frac{\rho_{a,local}}{m_a}}_{\text{axion parameters}} \cdot \underbrace{B^2 \cdot V \cdot Q}_{\text{experimental parameters}}$$

Haloscopes for DM – axions: figure–of–merit

- haloscopes: a **high Q – value** implies a very **narrow bandwidth Δf**



local DM –
density

experimental
parameter: Q

$$P_{sig} \sim g_{a\gamma\gamma}^2 \cdot \frac{\rho_{a,local}}{m_a} \cdot B^2 \cdot V \cdot Q$$

axion
parameters

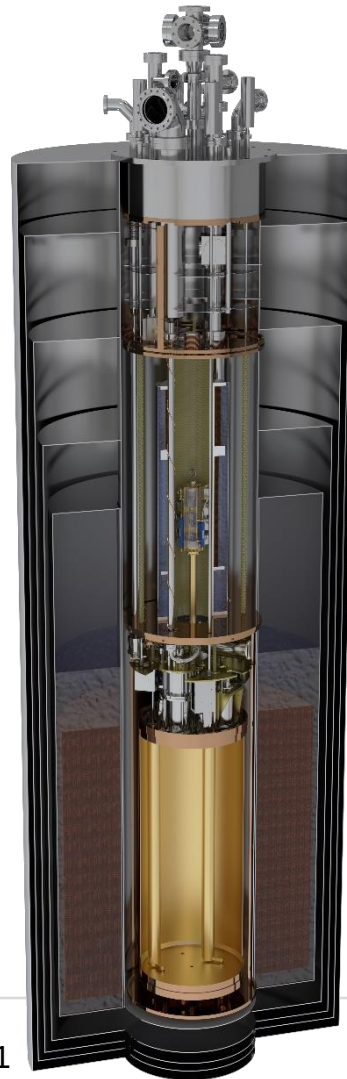
$$\frac{1}{Q} = \frac{\Delta f}{f_{res}}$$

Axion *D*ark *M*atter *E*Xperiment - *ADMX*

■ A haloscope now based at *UW* Seattle: measurements since 1995

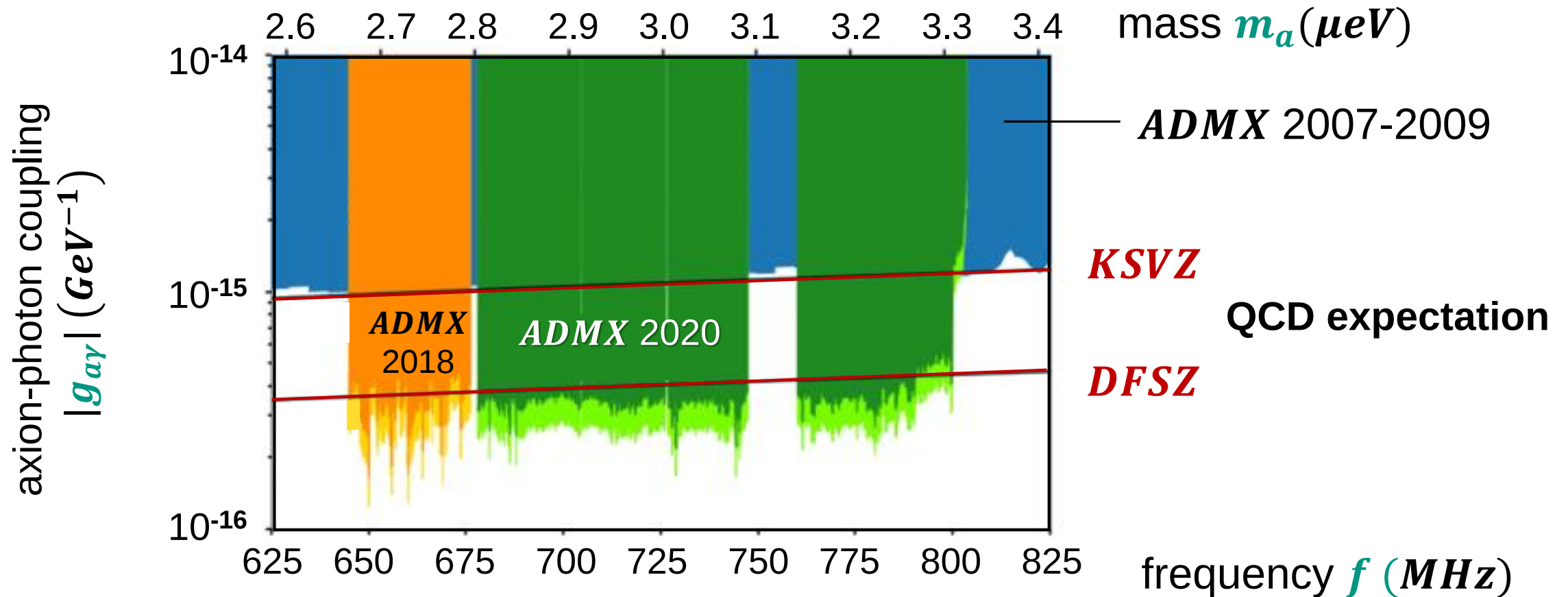
- 1994-2004:
measurements
with conventional
radio amplifiers
- 2007-....:
measurements
with *SQUIDS*
exclusion limits
for axion masses

$$m_a = (1.9 \dots 3.65) \mu\text{eV}$$



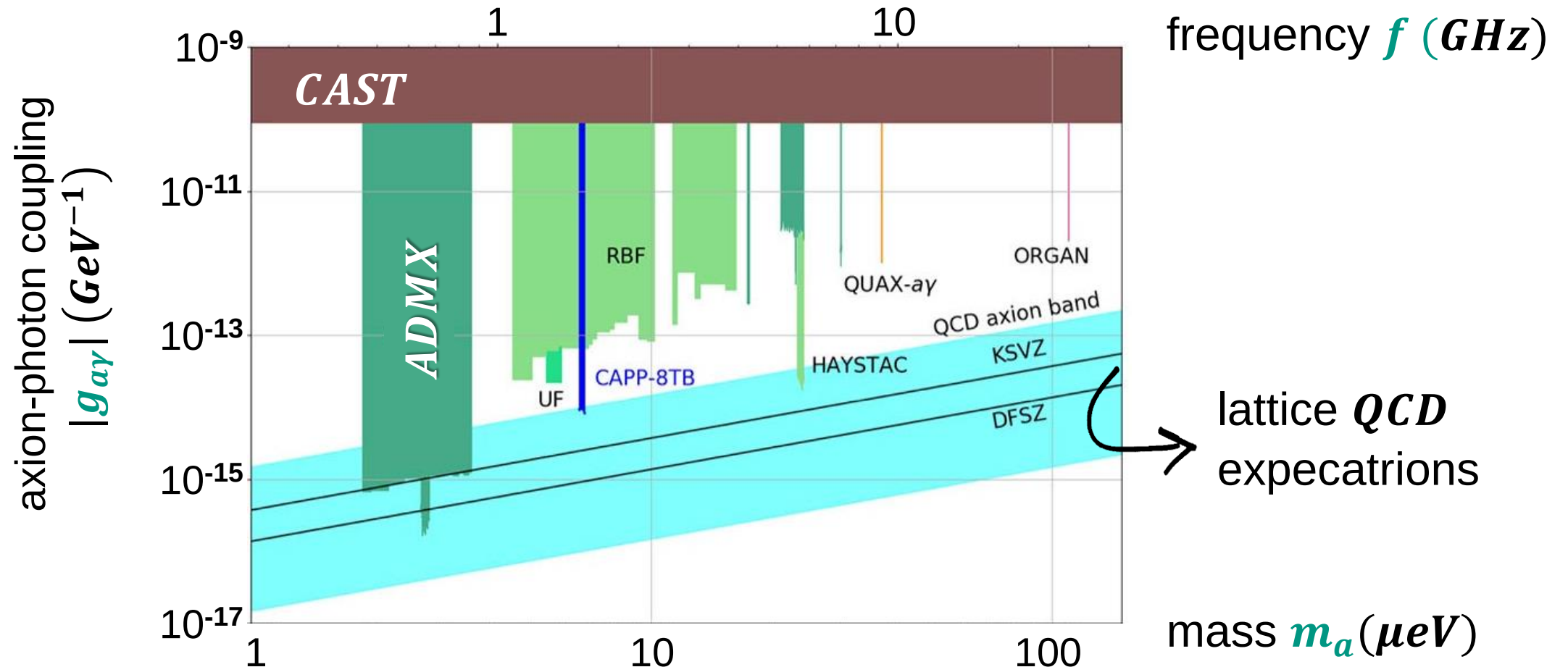
ADMX : exclusion limits for axions

■ Overview of different axion campaigns over the years



ADMX & HAYSTACK: future campaigns

- Overview of different axion campaigns: a lot more parameter space to cover



ADMX & HAYSTACK: axion alert

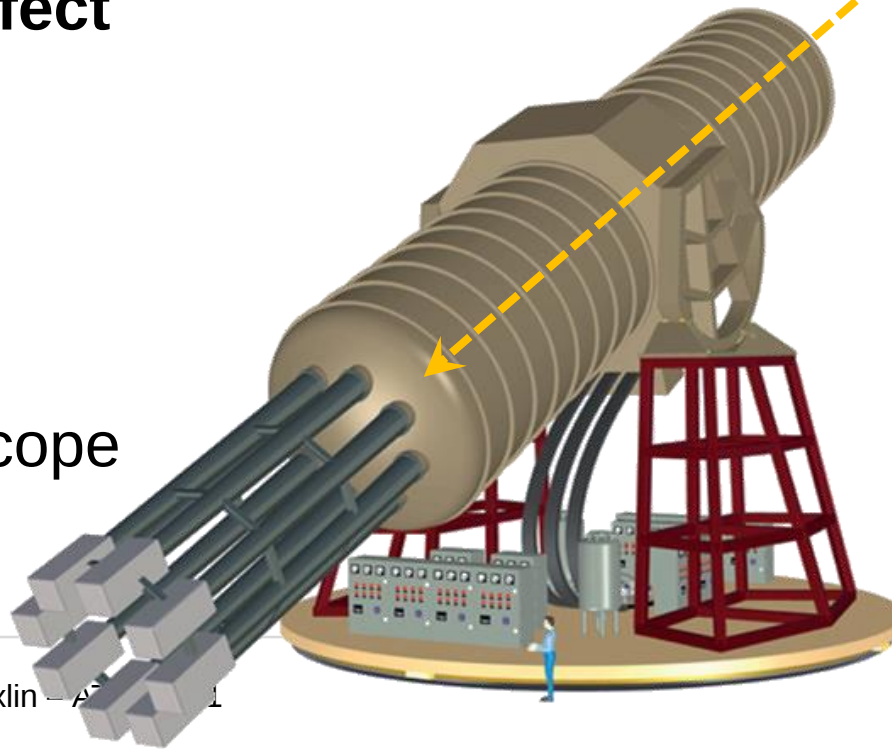
- Overview of different axion campaigns: a lot more parameter space to cover



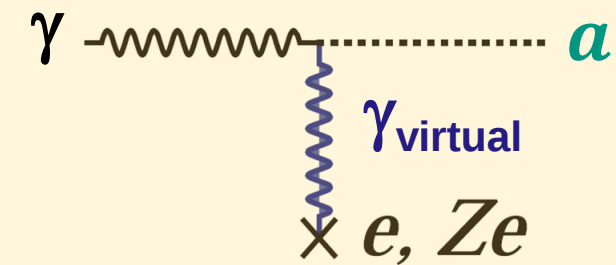
Axion experiments: helioscopes

- Detecting **solar axions** from the magnetized plasma on $E_a \approx \text{keV}$ – scale
- **high mass axions** (eV) could be produced in the solar interior via the **incoherent Primakoff effect**

pointing a helioscope towards to sun

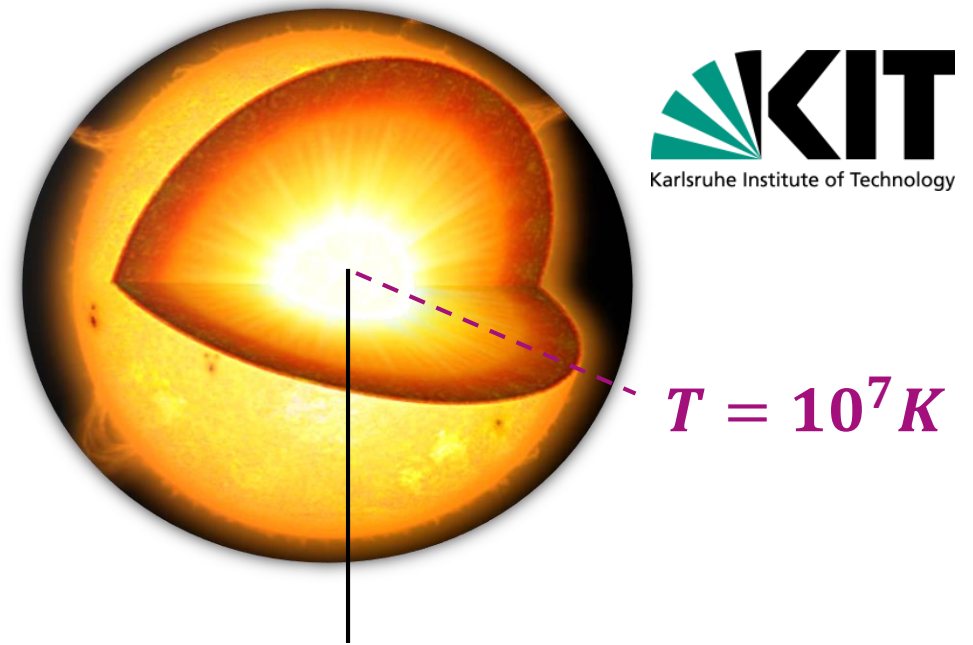
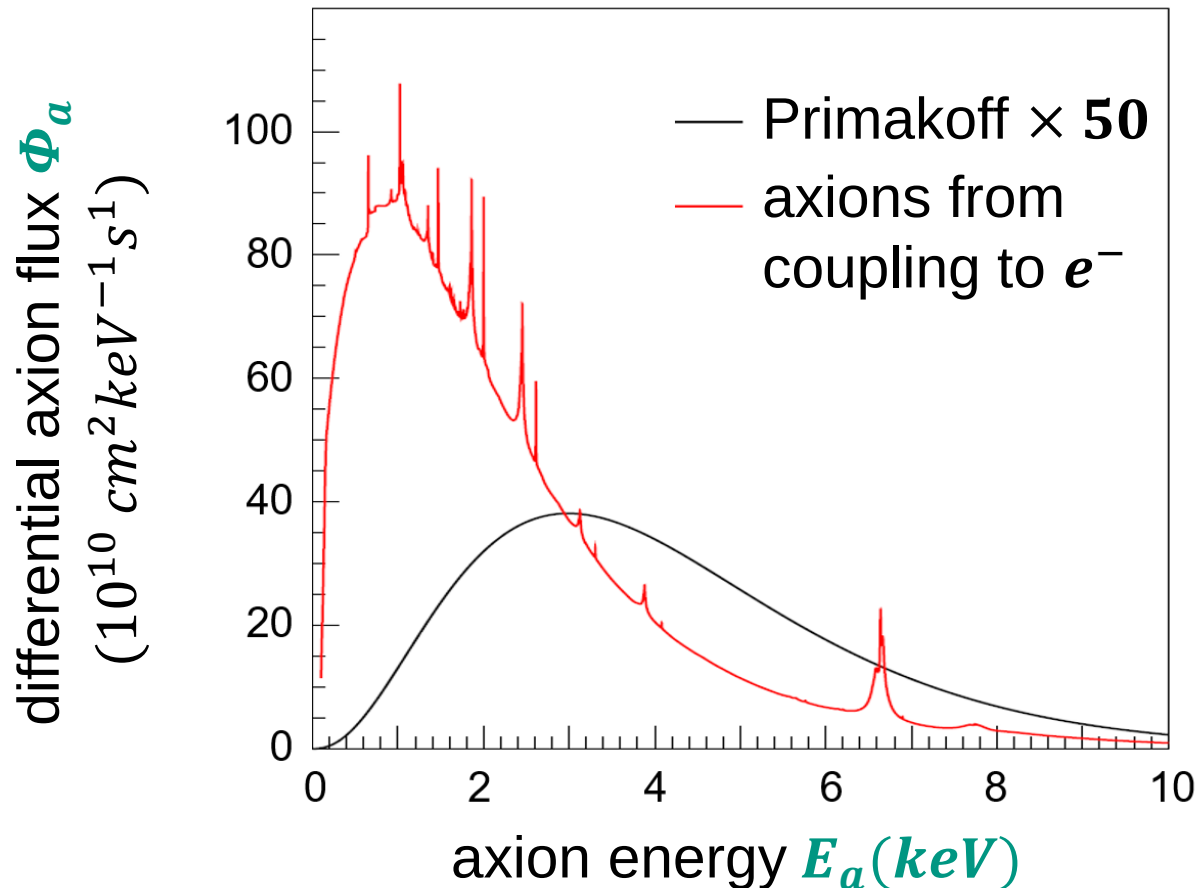


solar plasma with electrons & ions:



Axion experiments: helioscopes

- Detecting **solar axions** from the magnetized plasma on $E_a \approx \text{keV}$ – scale



- dominant (*incoherent*) process:
axions generated in **electric field**
of magnetized plasma (ions, e^-)

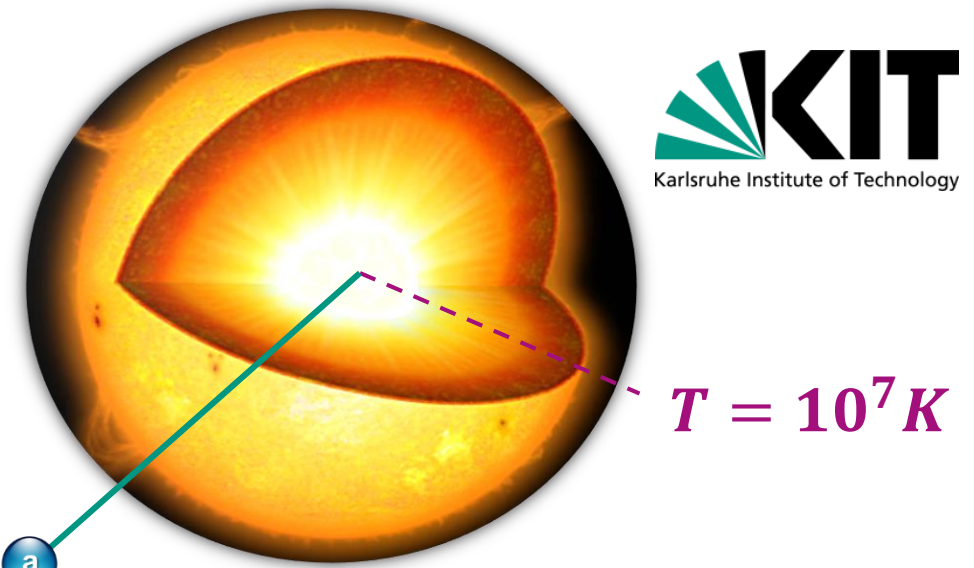
flux

$$\Phi_a = \left(\frac{g_{a\gamma\gamma}}{10^{-10} \text{ GeV}^{-1}} \right)^2$$

Axion experiments: helioscopes

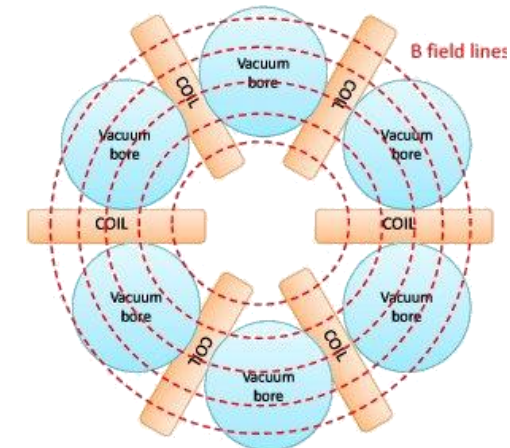
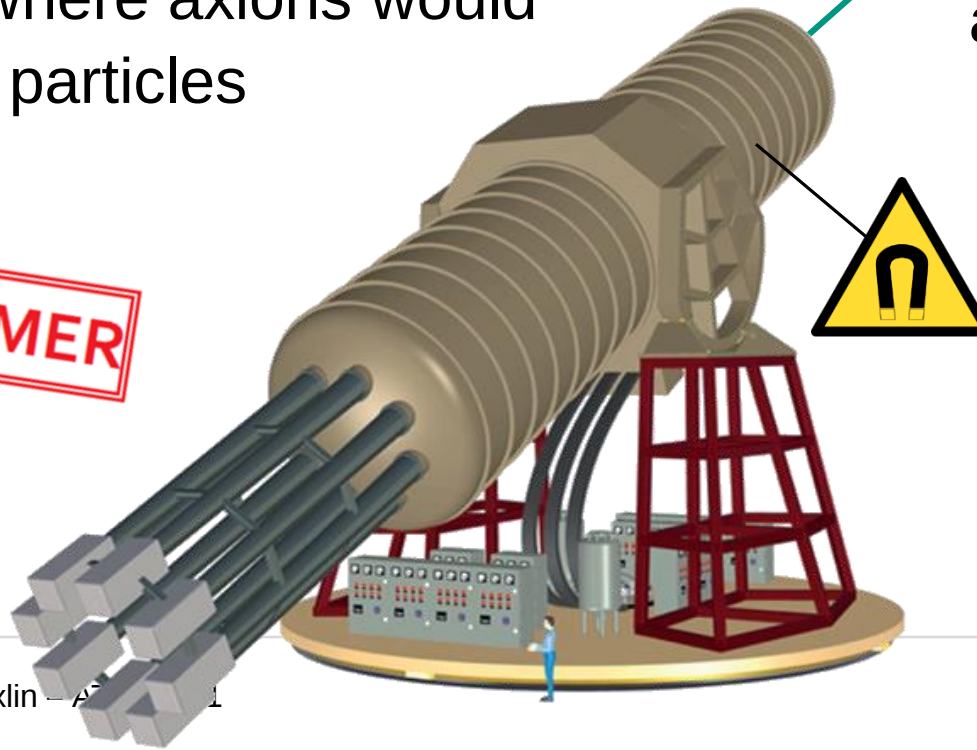
■ Detecting **solar axions** from the magnetized plasma on $E_a \approx \text{keV}$ – scale

- helioscopes are only sensitive to axion parameters (mass m_a & coupling regime $g_{a\gamma\gamma}$) where axions would not be **CDM** – particles



axion- conversion into X – rays
in a strong dipole B – field

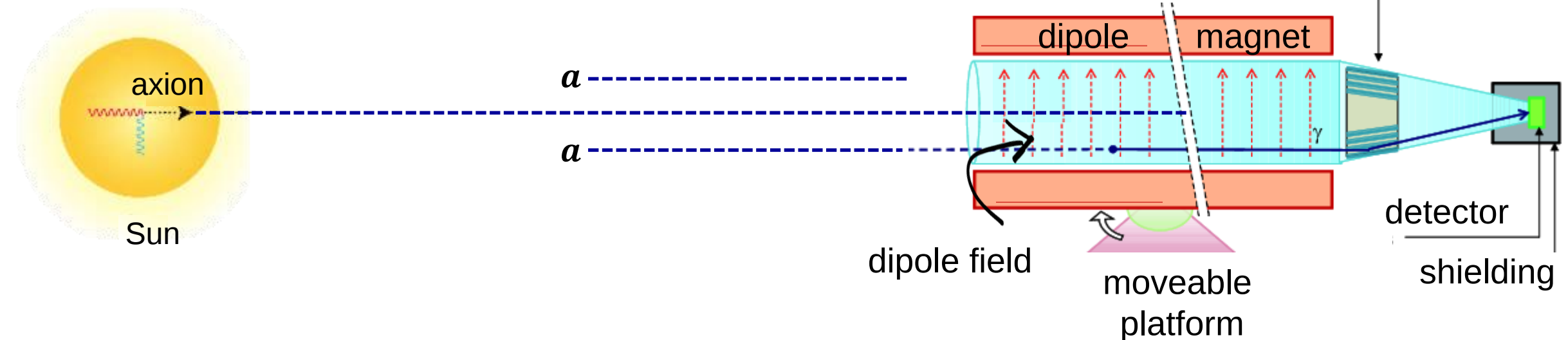
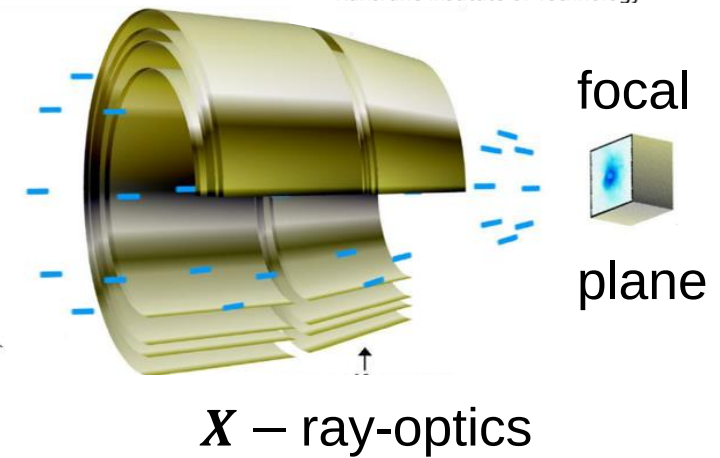
DISCLAIMER



Schematic set-up of a helioscope

■ Reconverting solar axions into X – ray photons

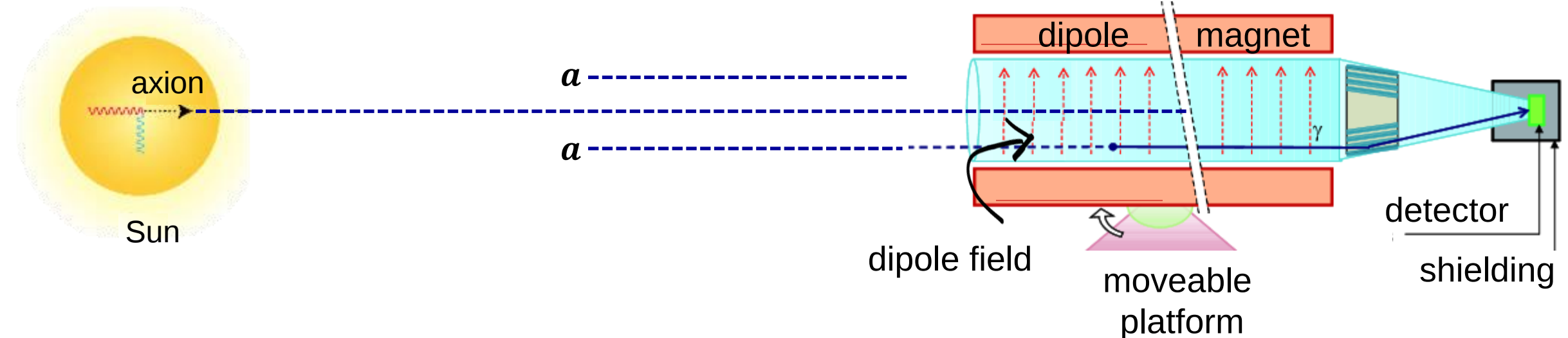
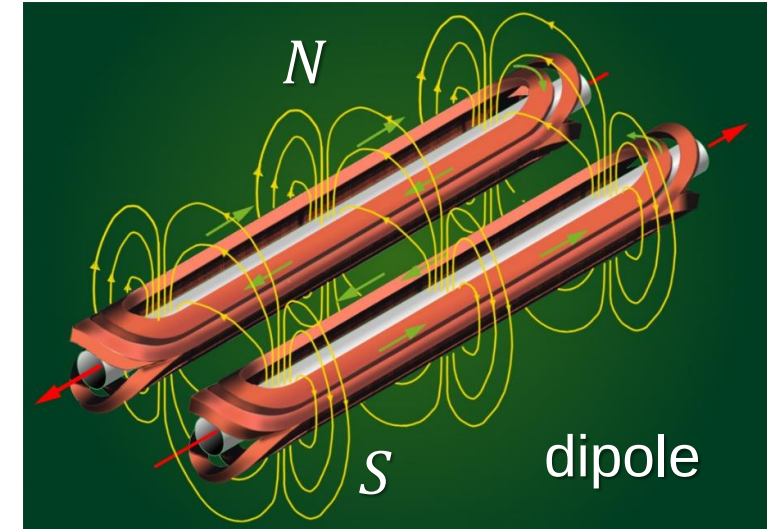
- **inverted Primakoff-effect** is used to convert the high-mass axions back into energetic X – rays (keV)
- X – rays (keV) have to be focused via ‘mirrors’ (Wolter telescope) onto the focal plane detector



Schematic set-up of a helioscope

■ We can make use of unused *LHC* dipoles

- we again need a **magnetic dipole field**
 - ⇒ only γ 's with polarisation parallel to the $\mathbf{B}$ – field mix with axions
 - ⇒ photon propagation transversal to $\mathbf{B}$ – field

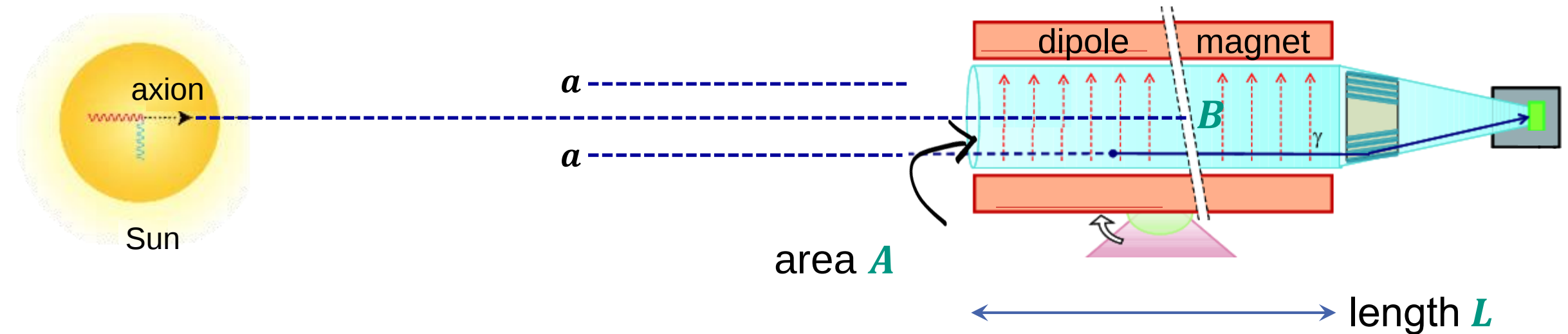


Schematic set-up of a helioscope

- Conversion probability $P_{a\gamma}$ to transform an axion into an X – ray photon

$$P_{a\gamma} = 2.6 \times 10^{-17} \cdot \left(\frac{B}{10T} \right)^2 \cdot \left(\frac{L}{10m} \right)^2 \cdot (g_{a\gamma\gamma} \times 10^{10} \text{ GeV})^2 \cdot F$$

- experimental „figure-of-merit“ $\sim B^2 \cdot L^2 \cdot A$ (magnetic field B , length L , cross section A)

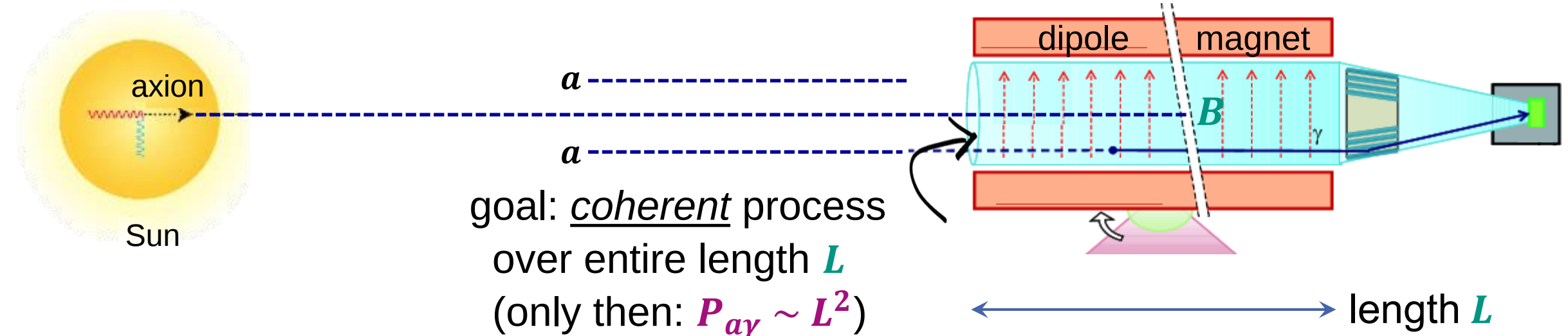


Schematic set-up of a helioscope

- Conversion probability $P_{a\gamma}$: the importance of **coherence** – factor F

$$P_{a\gamma} = 2.6 \times 10^{-17} \cdot \left(\frac{B}{10T} \right)^2 \cdot \left(\frac{L}{10m} \right)^2 \cdot (g_{a\gamma\gamma} \times 10^{10} \text{ GeV})^2 \cdot F$$

- conversion is **coherent** over entire length L if the **phase** of the massive axion and the massive photon remain the same!



Maintaining coherence in a helioscope

■ Enhancing the coherence – factor F via the injection of buffer gas

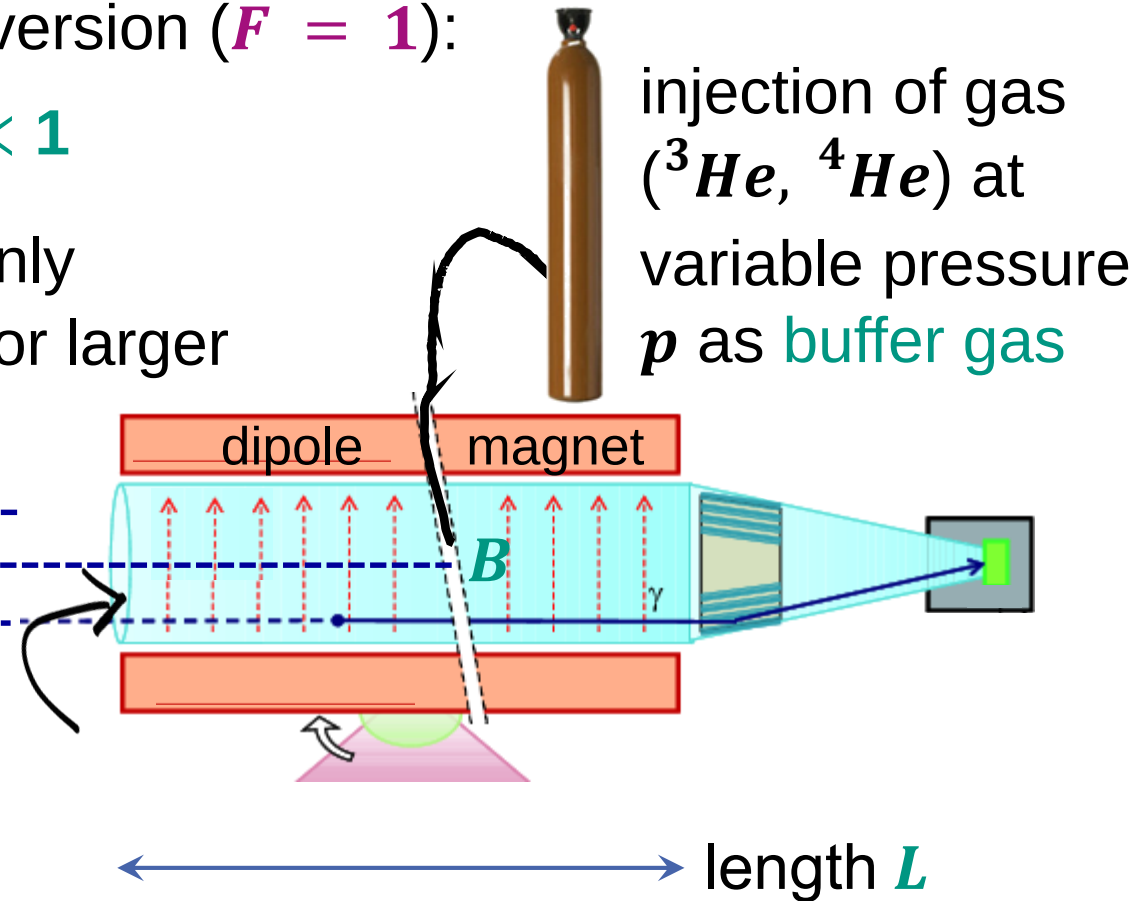
- **coherence** requirement for full axion conversion ($F = 1$):

momentum transfer $q \cdot \text{length } L \ll 1$

- for axions with $E_a \approx \text{keV}$ – **scale** this is only fulfilled for small masses $m_a < 10^{-2} \text{ eV}$, for larger masses we need

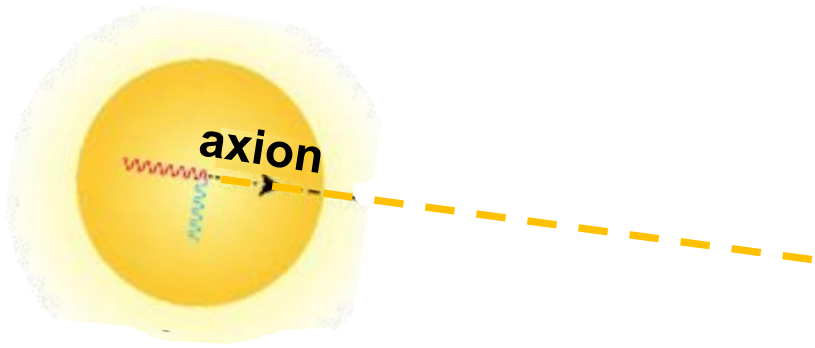
to add a buffer gas to adjust the **effective mass** of X –ray photons!

goal: coherent process over entire length L
(then: $P_{a\gamma} \sim L^2$)



Helioscopes: *CAST** experiment at *CERN*

- Making use of a spare $L = 9.3\text{ m}$ long dipole magnet ($B = 9.5\text{ T}$) of the *LHC*



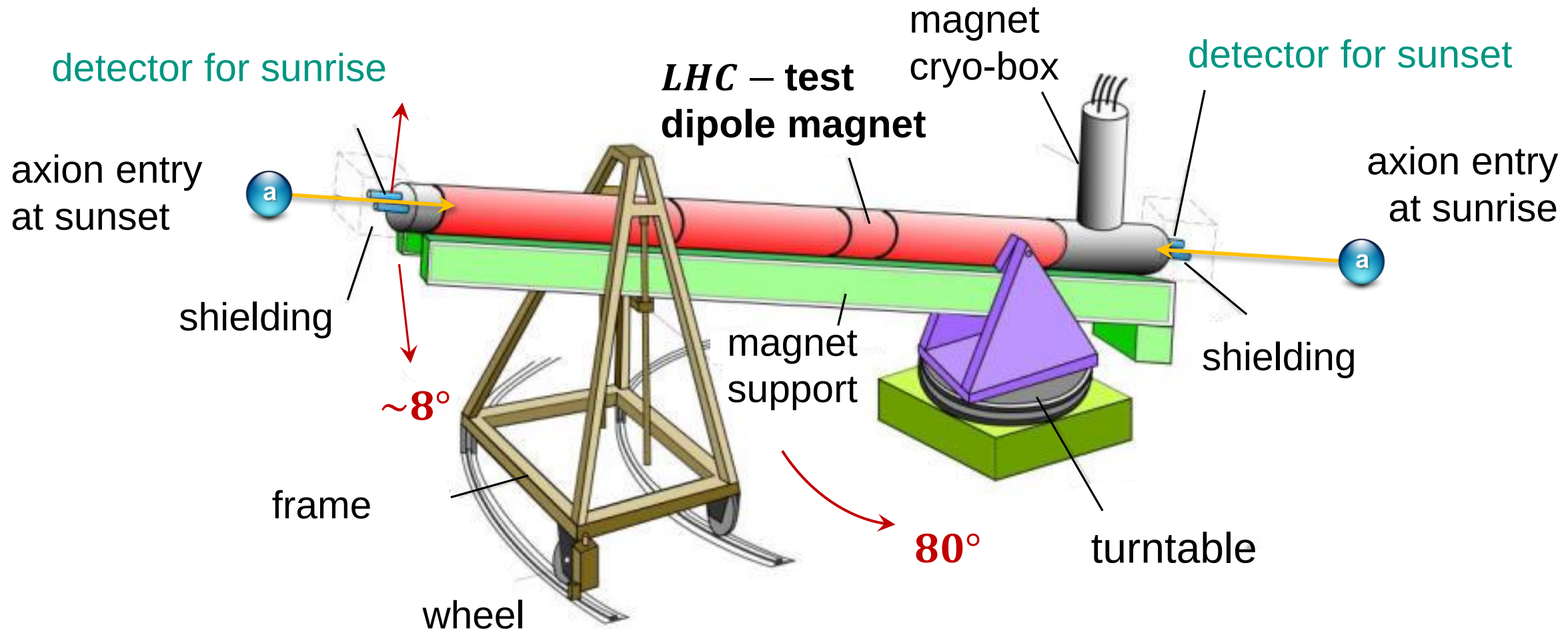
- magnet axis follows path of the Sun over several hours per day
- both magnet entries (side of the setting / rising sun) are fully instrumented with modern X – ray detector systems



Q: CERN Courier

Helioscopes: *CAST* experiment at *CERN*

- A moveable dipole magnet that can also be tilted to follow the Sun



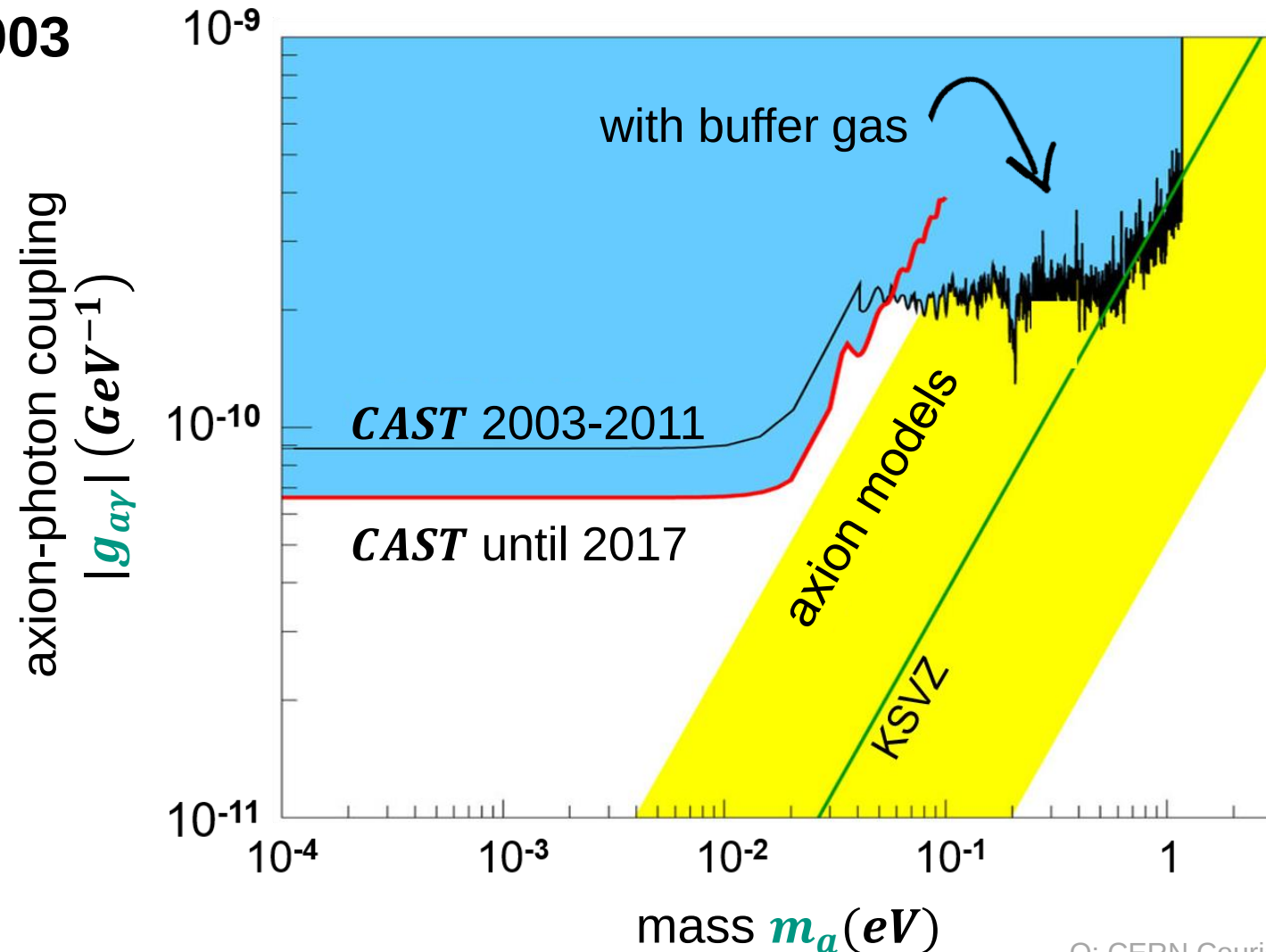
Helioscopes: *CAST* results – use of buffer gas

■ Axion measurements since 2003

2005-2006:
measurements with ^4He

2008-2011:
measurements with ^3He
first push forward into the
region of *QCD* – axions

since 2012:
improved measurements
with ^4He & in vacuum

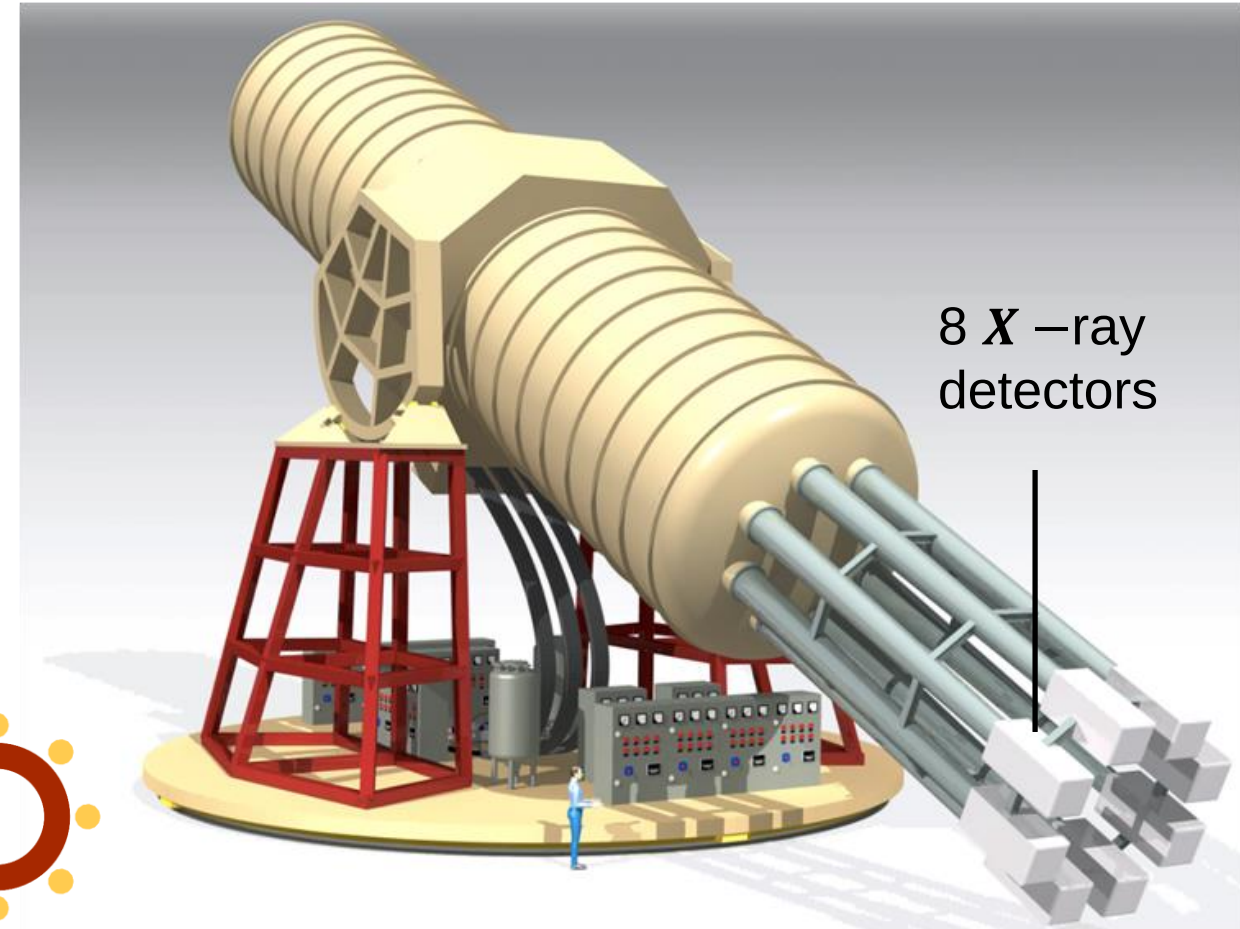
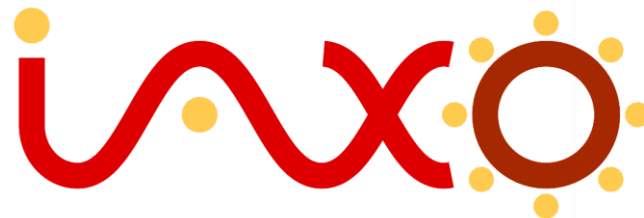


Q: CERN Courier

Helioscopes: *IAXO** – an outlook

■ Future search for solar axions with a large **toroidal magnet system**

- setup: $L = 20\text{ m}$ long **toroidal magnet** with $\varnothing = 5.2\text{ m}$
- toroidal magnet system formed by 8 coils ($\varnothing = 0.6\text{ m}$) with $B = 2.5\text{ T}$
- **factor 20** improved sensitivity relative to *CAST*



Helioscopes: *IAXO* – an outlook

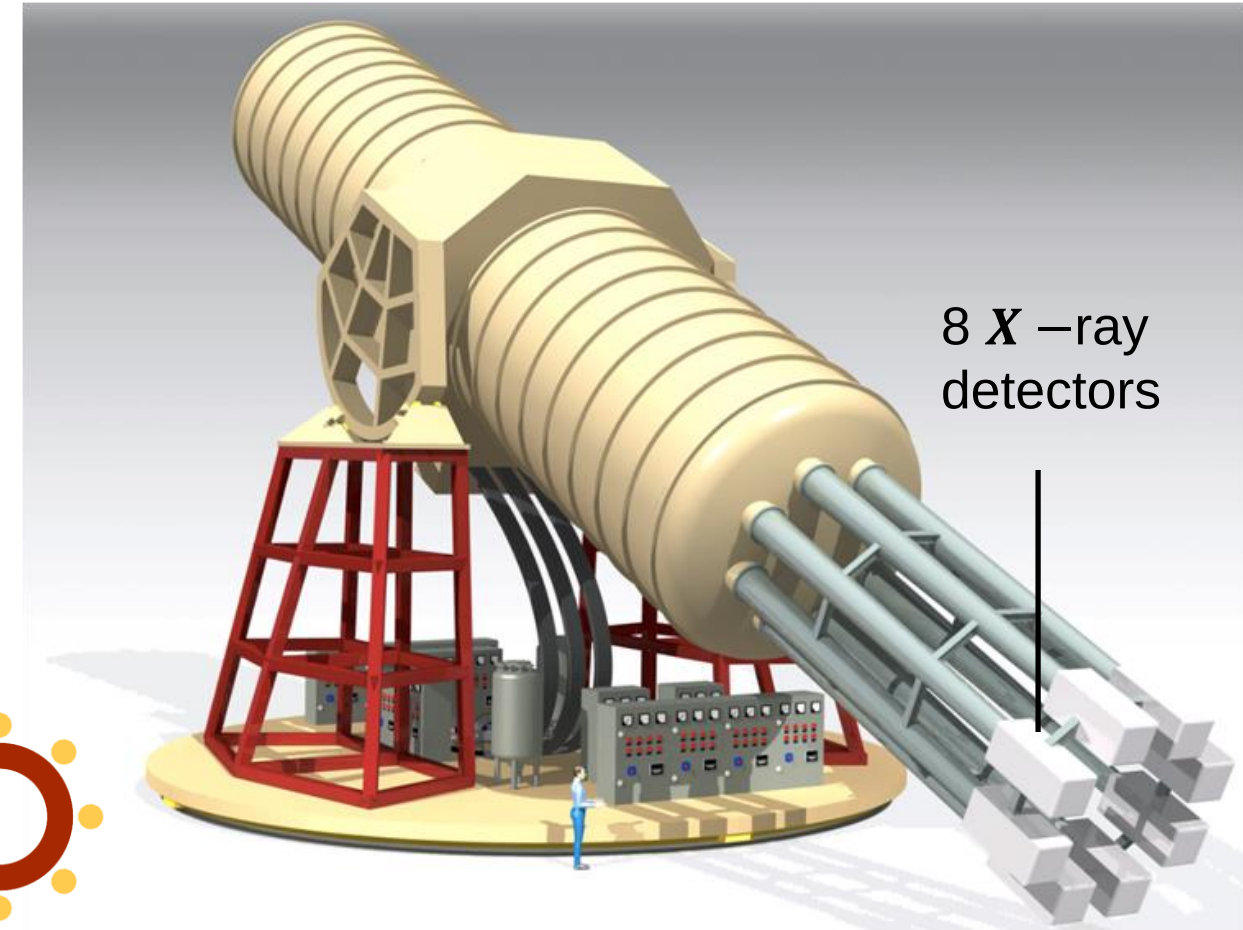
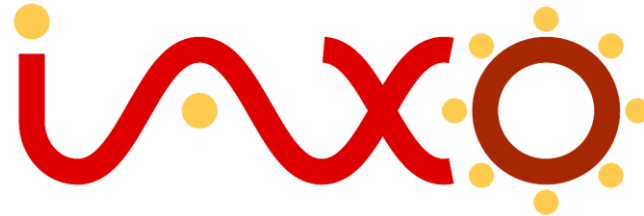
■ Future search for solar axions with a large **toroidal magnet system**

- figure-of-merit: $\sim B^2 \cdot L^2 \cdot A$

- *CAST*: $21 T^2 m^4$

- *IAXO*: $\sim 6000 T^2 m^4$

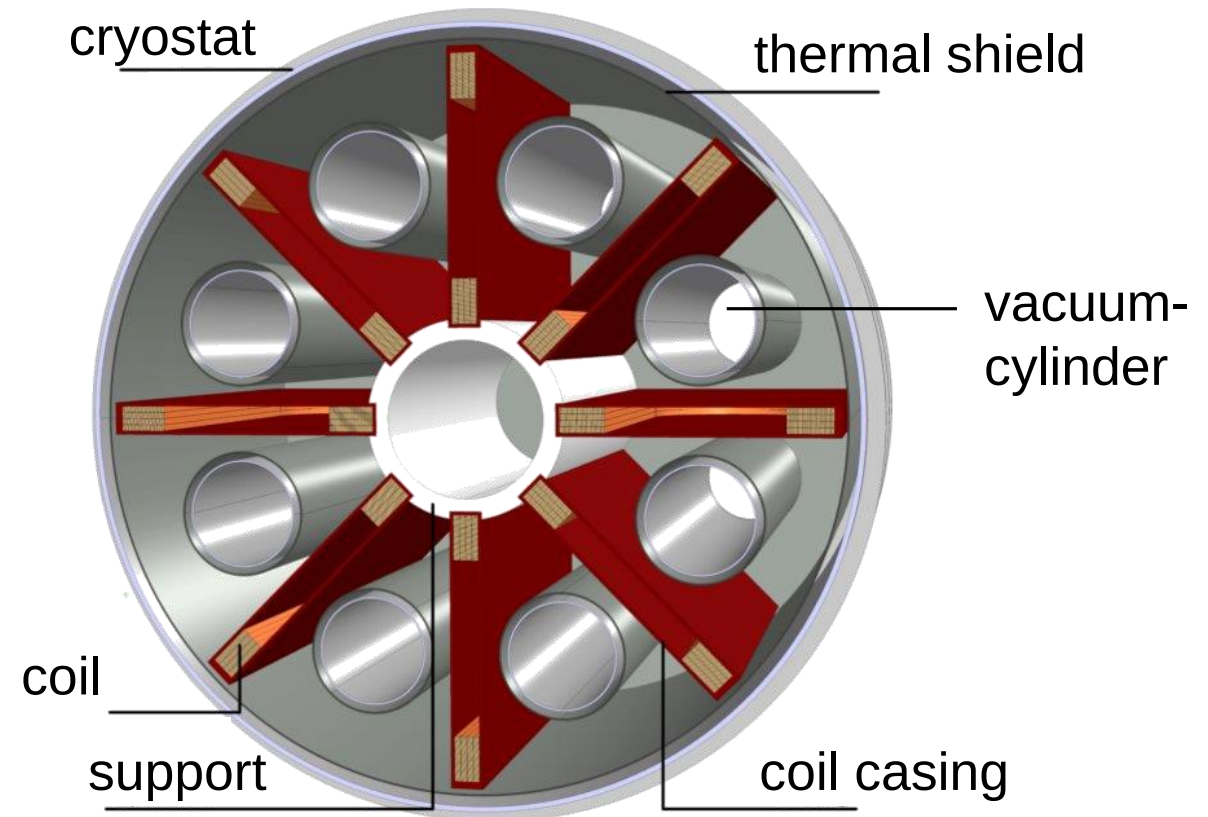
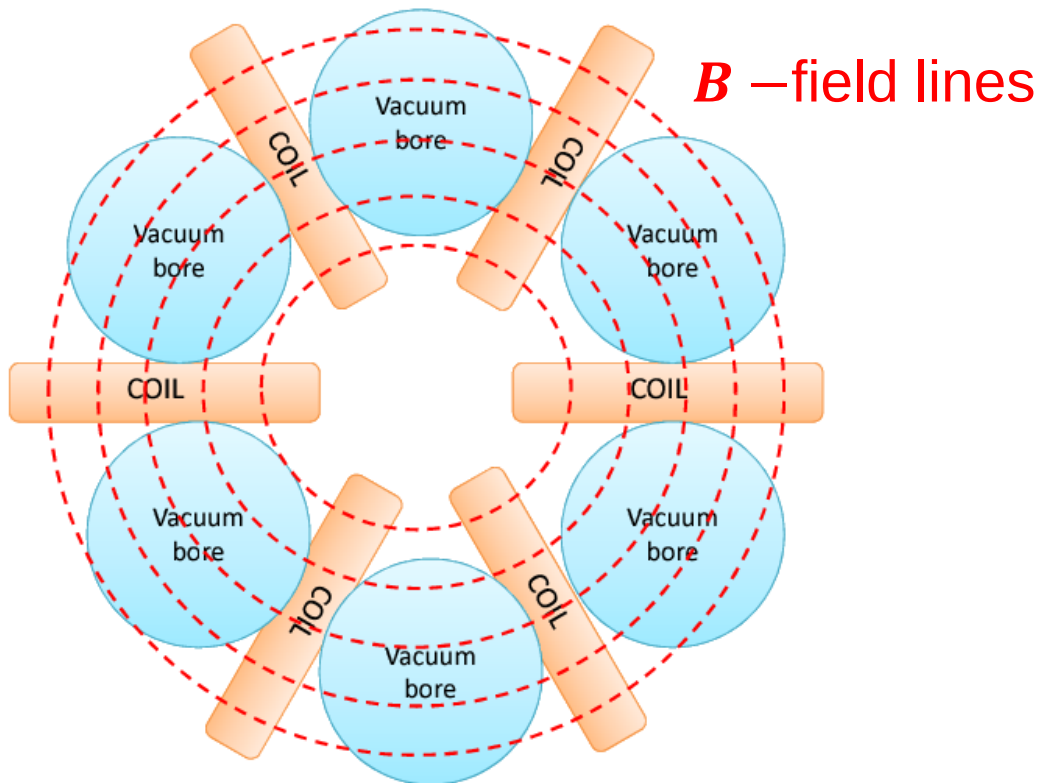
- **factor 20** improved sensitivity
relative to *CAST*



Helioscopes: *IAXO* – an outlook

■ Future search for solar axions with a large **toroidal magnet system**

- 8 (6) vacuum cylinders interspersed with 8 (6) coils in a toroidal set-up



Helioscopes: *IAXO* – an outlook

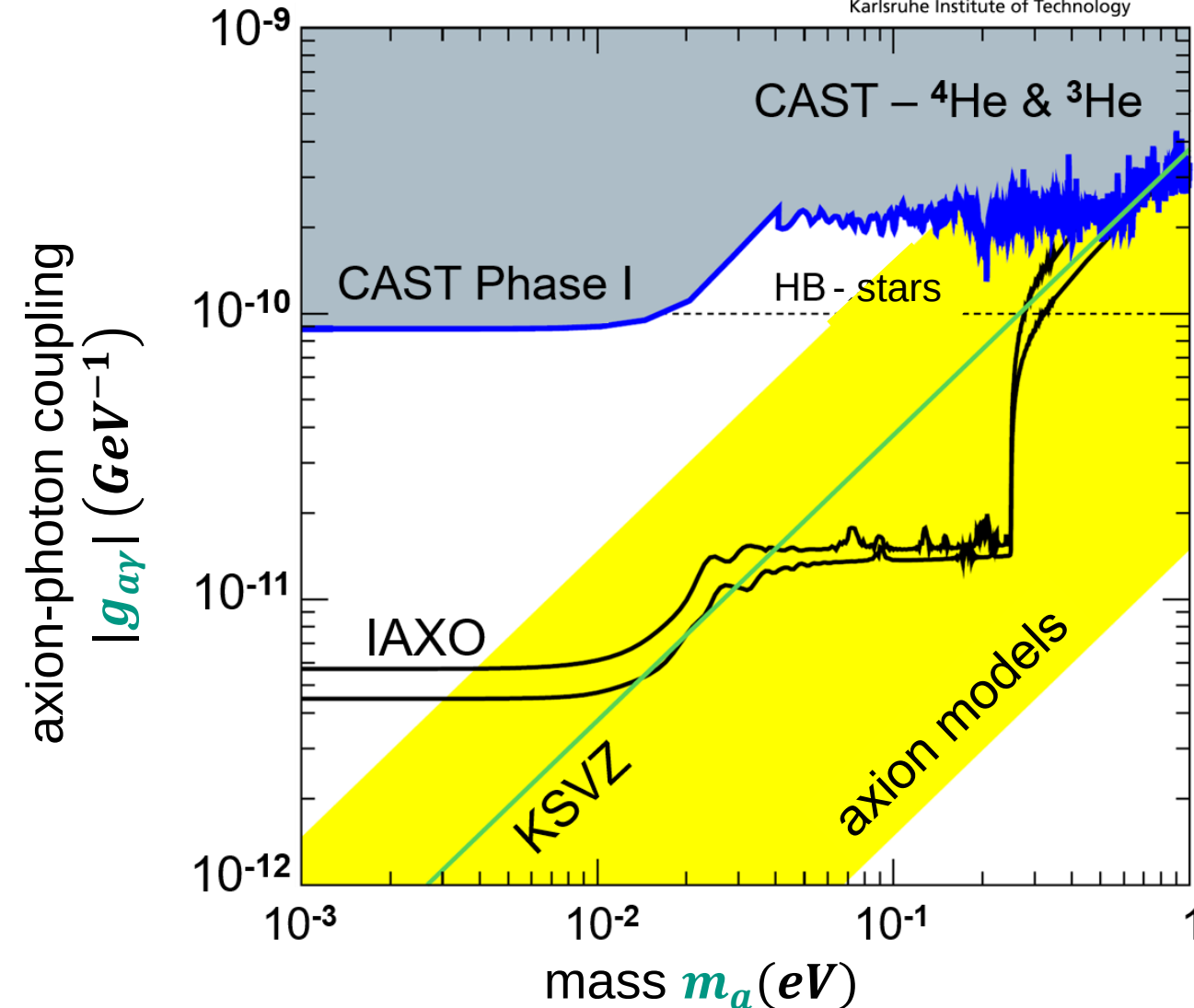
■ Future search for solar axions

CAST

- results after ~20 year long data taking (vacuum / buffer gas)

IAXO

- expected for 3 years data taking (vacuum / buffer gas)
- test of realistic *QCD* – axion parameter regions



Helioscopes: first step - BabyIAXO

■ Search for solar axions with a smaller toroidal magnet at *DESY*

- figure-of-merit: $\sim B^2 \cdot L^2 \cdot A$
 - *CAST*: $21 \text{ T}^2 \text{ m}^4$
 - *BabyIAXO*: $\sim 230 \text{ T}^2 \text{ m}^4$
- follows path of Sun $t = 24 \text{ h}$ a day
- magnet design: uses expertise from the *ATLAS* experiment (*CERN*)
- Wolter-Telescope (X-Ray): from *ESA* (for the *XMM – Newton* telescope)

