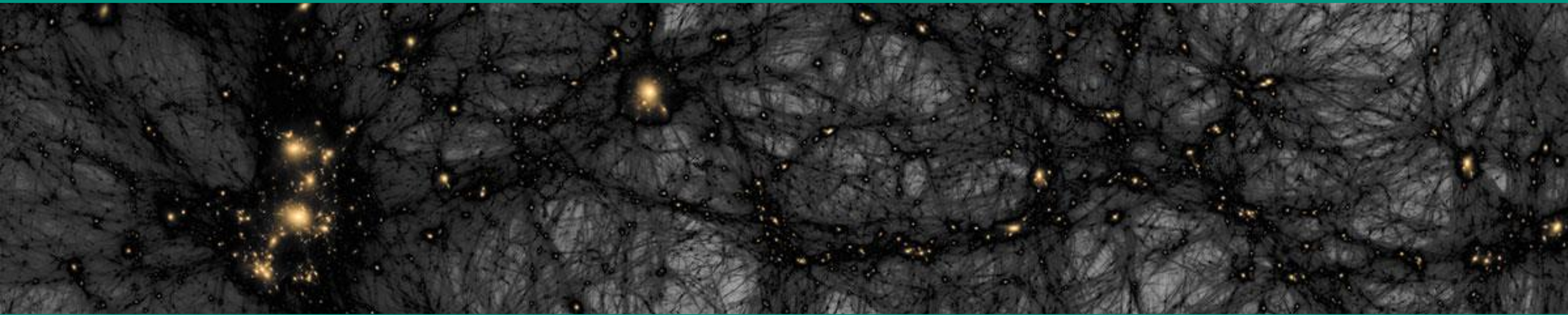


Astroparticle physics I – Dark Matter

WS22/23 Lecture 5

Nov. 10, 2022



Recap of Lecture 4

■ Extensive air showers: lateral & longitudinal distributions

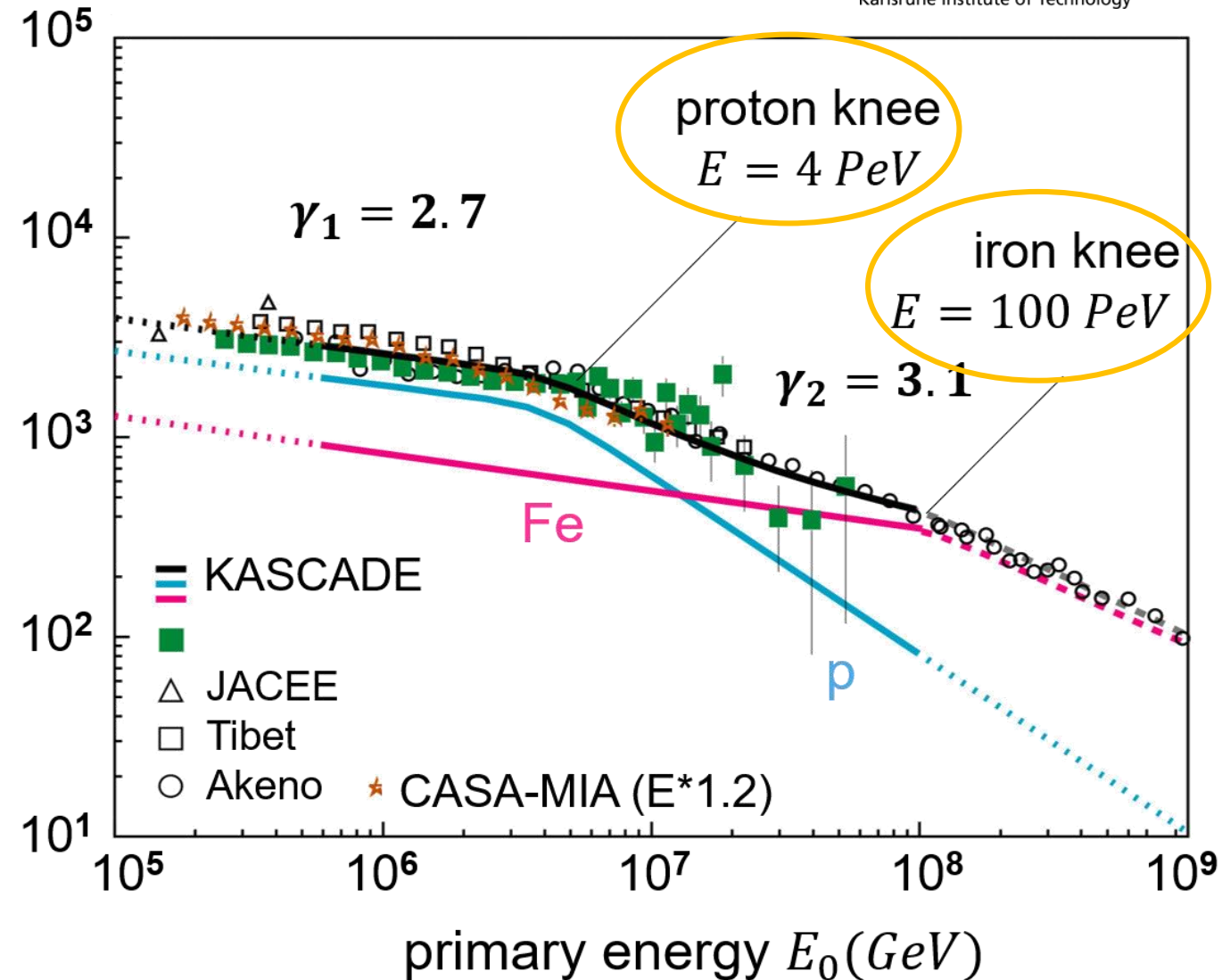
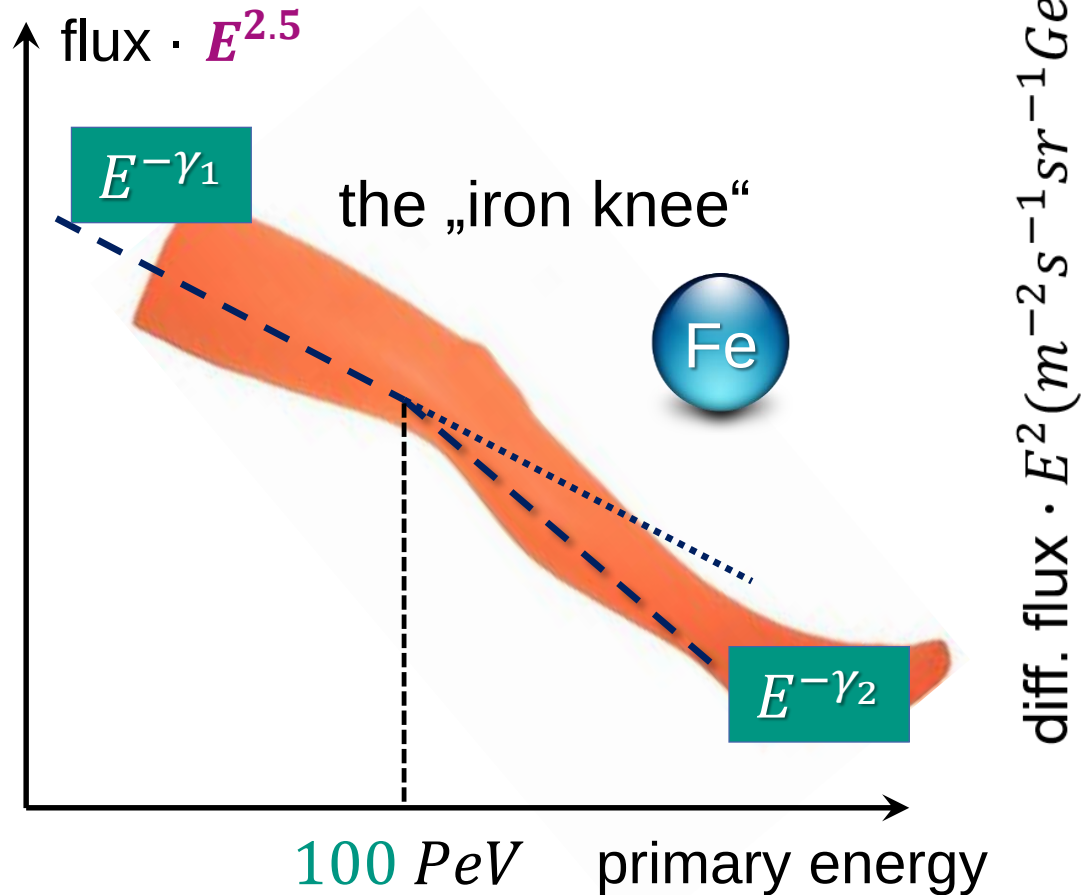
- primary CR **energy E_0** :
use integration over lateral distributions of $N_e + N_\mu$ (**‘foot print’**) as good indicator for E_0 (compare Greisen-fit to CORSIKA simulations)
- primary CR **mass M** :
use observed longitudinal distribution to determine shower maximum X_{max}
 - heavy CR (**^{56}Fe**): small $X_{max} \Rightarrow$ **small ratio N_e/N_μ**
 - light CR (**p**): large $X_{max} \Rightarrow$ **large ratio N_e/N_μ**

■ Pioneering air shower experiment KASCADE at KIT Campus North

- components: large scintillator array, muon tunnel, central hadron calorimeter
- extension to KASCADE-Grande: at the knee - change of mass composition

Power-Law Feature: element-specific 'knees'

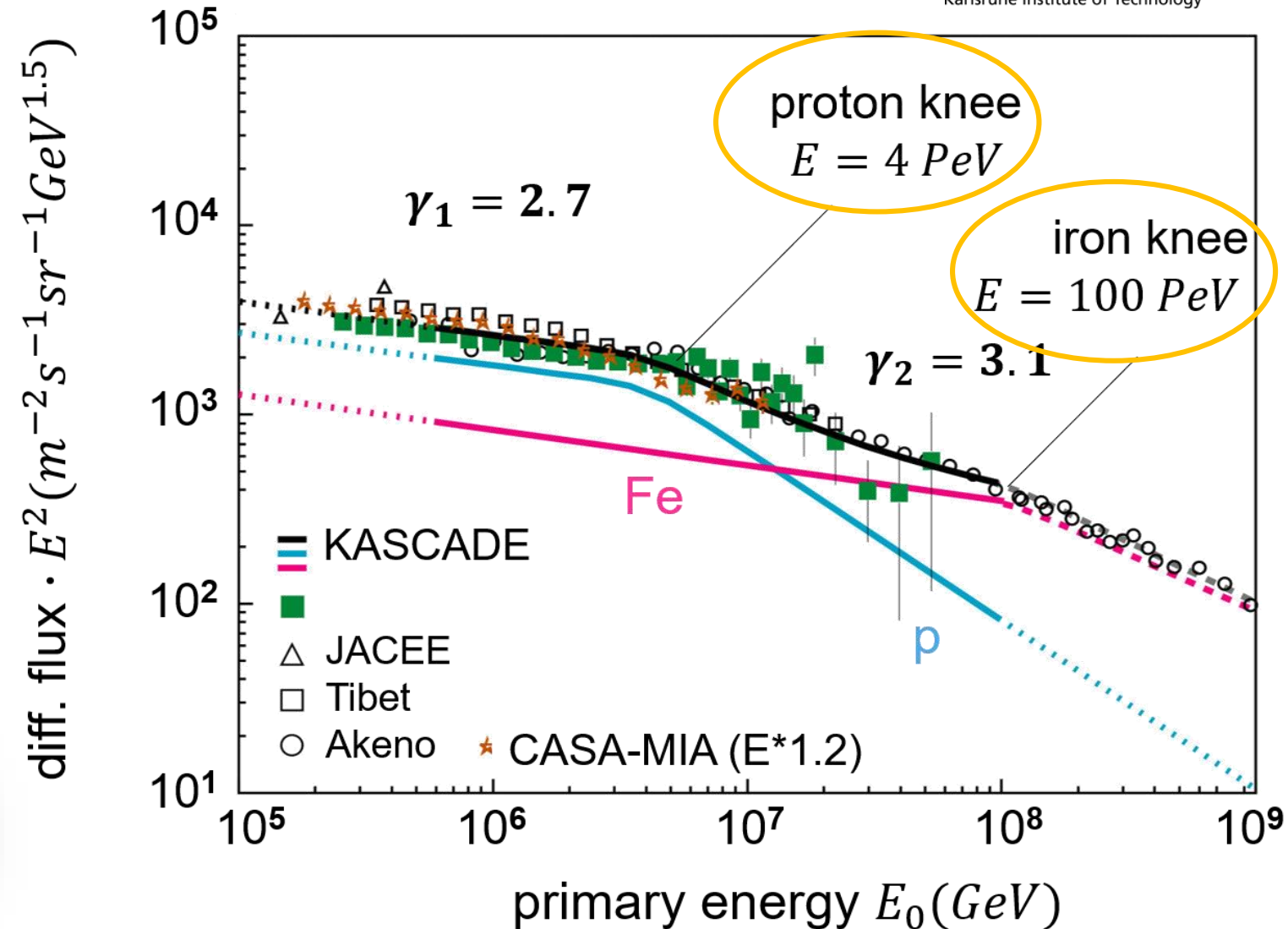
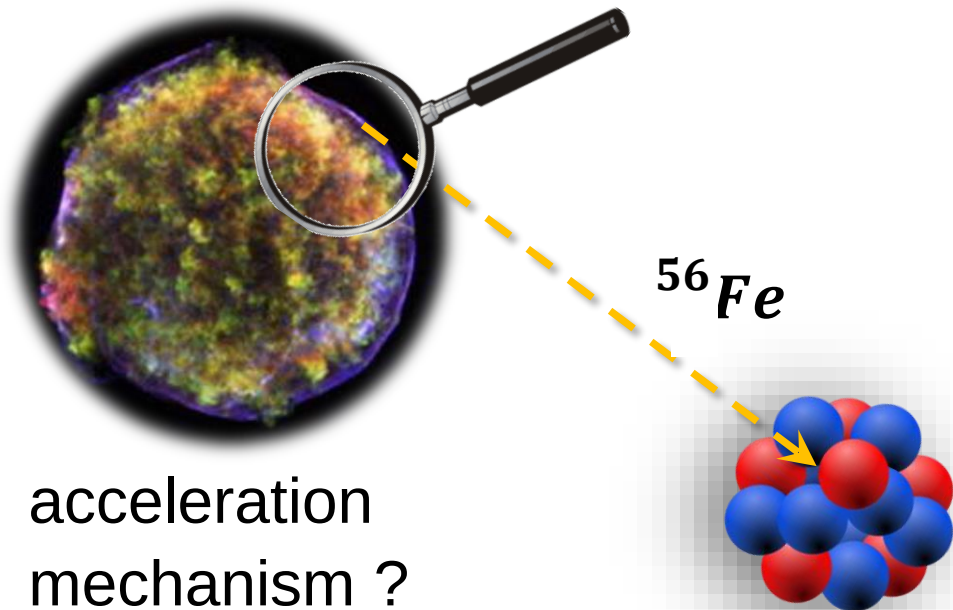
■ Iron knee at $E = 100 \text{ PeV}$



Power-Law Feature: p-knee & Fe-knee

■ Observations at 100 PeV

- KASCADE-Grande:
first ever observation
of **Fe-knee**



Knee due to particle losses during propagation?

■ Propagation of CRs with energy E & mass A

- GALPROP-MC-Code

transport of galactic CRs from the source to Earth (Diffusion):
diffusion of particles in 'leaky box'

- detailed modelling of effects due to CR propagation:

a) orientation of galactic B-fields

b) energy losses (light particles):

- inverse Compton effect
- emission of synchrotron radiation

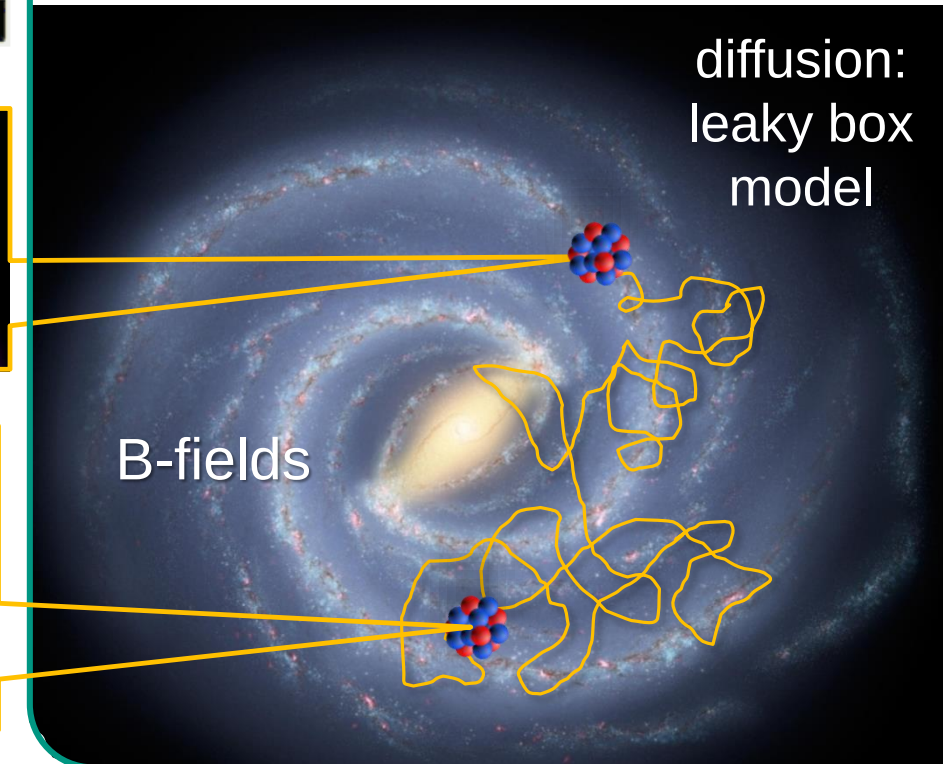


CR propagation in galaxis

- charged CRs (fully ionised nuclei): diffusion in $B_{galactic}$

diffusion:
leaky box
model

B-fields



Propagation of cosmic rays in our galaxy

■ Charged CRs: guided by galactic magnetic fields with $B \sim \text{few } \mu\text{G}$

- important parameter:

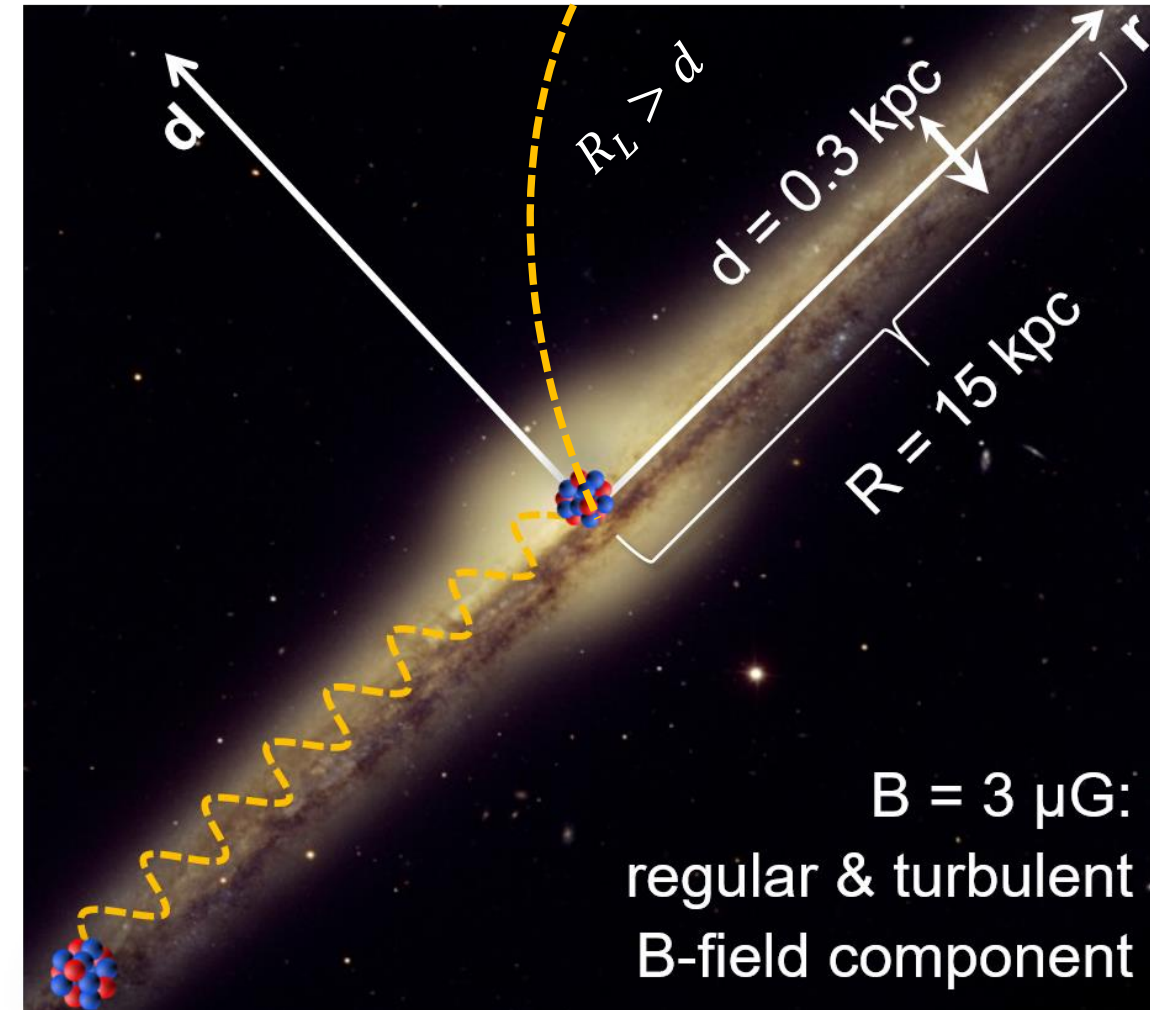
Larmor radius R_L

$$R_L \cong 1pc \cdot \left(\frac{E}{10^{15} eV} \right) \cdot \left(\frac{\mu G}{Z \cdot B} \right)$$

- high energies:

$$R_L > d = 0.3 \text{ kpc}$$

⇒ these CR nuclei will leave the galaxis



Propagation of cosmic rays in our galaxy

■ Charged CRs: guided by galactic magnetic fields with $B \sim \text{few } \mu\text{G}$

- important parameter:

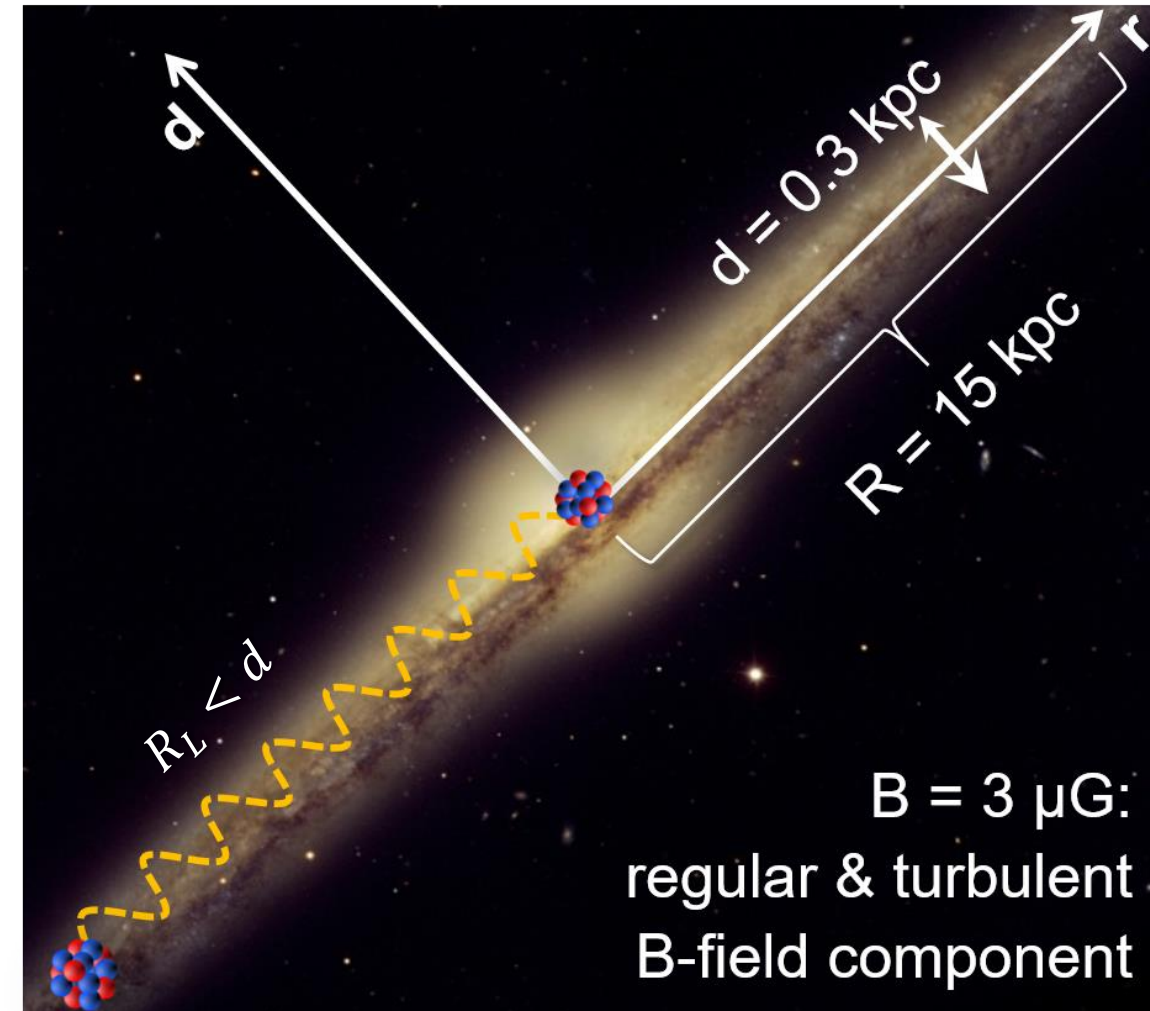
Larmor radius R_L

$$R_L \cong 1 \text{ pc} \cdot \left(\frac{E}{10^{15} \text{ eV}} \right) \cdot \left(\frac{\mu\text{G}}{Z \cdot B} \right)$$

- low energies:

$$R_L < d = 0.3 \text{ kpc}$$

⇒ these CR nuclei are trapped in galaxis
on time scales $\tau \sim 3 \dots 10 \cdot 10^6 \text{ y}$



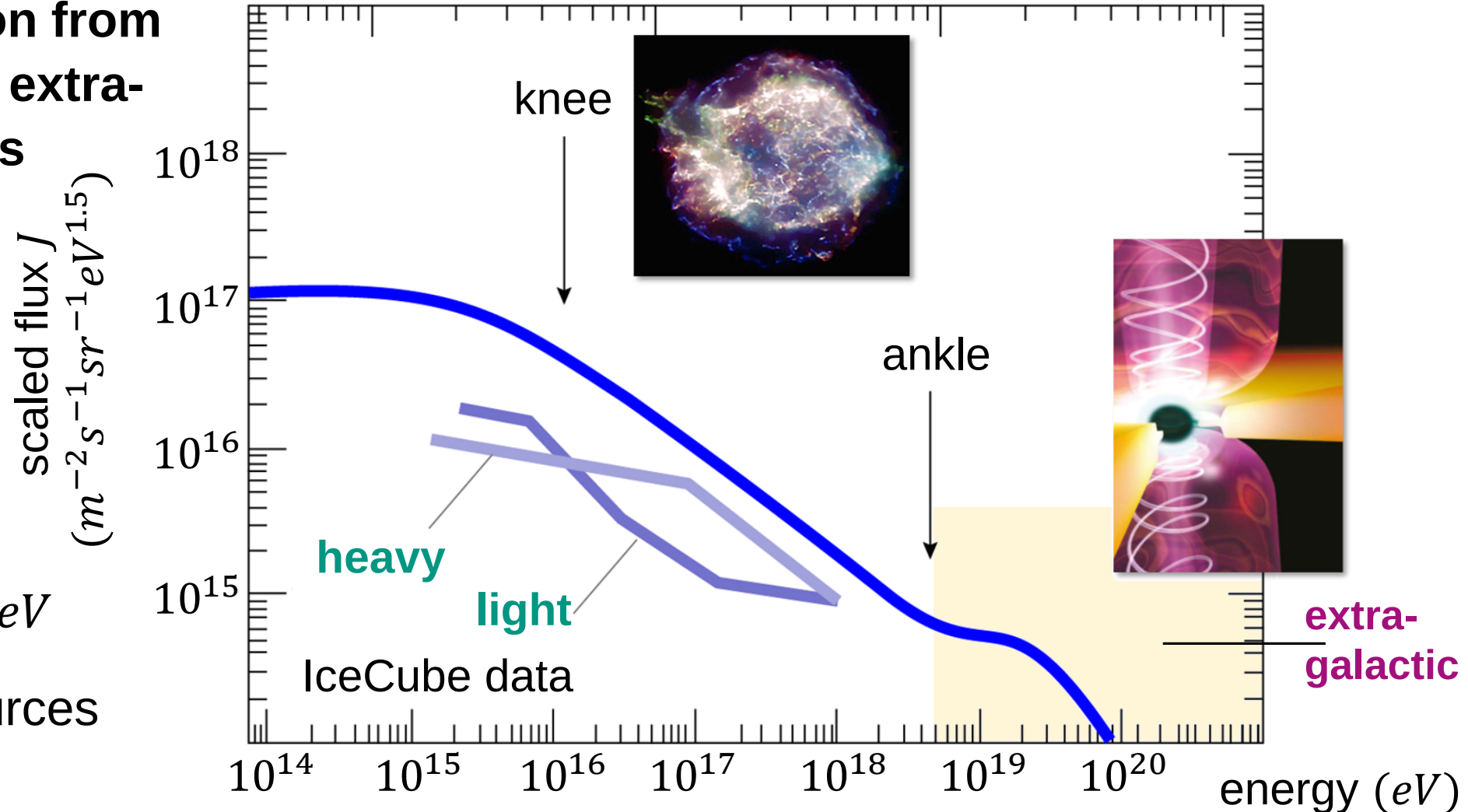
CR spectrum feature: the ankle

■ Ankle: transition from galactic CRs to extra-galactic sources

- IceCube*:
light + **heavy**
CRs (galaxis)

⇒ all galactic
accelerators
only up to $\sim 10^{18}$ eV

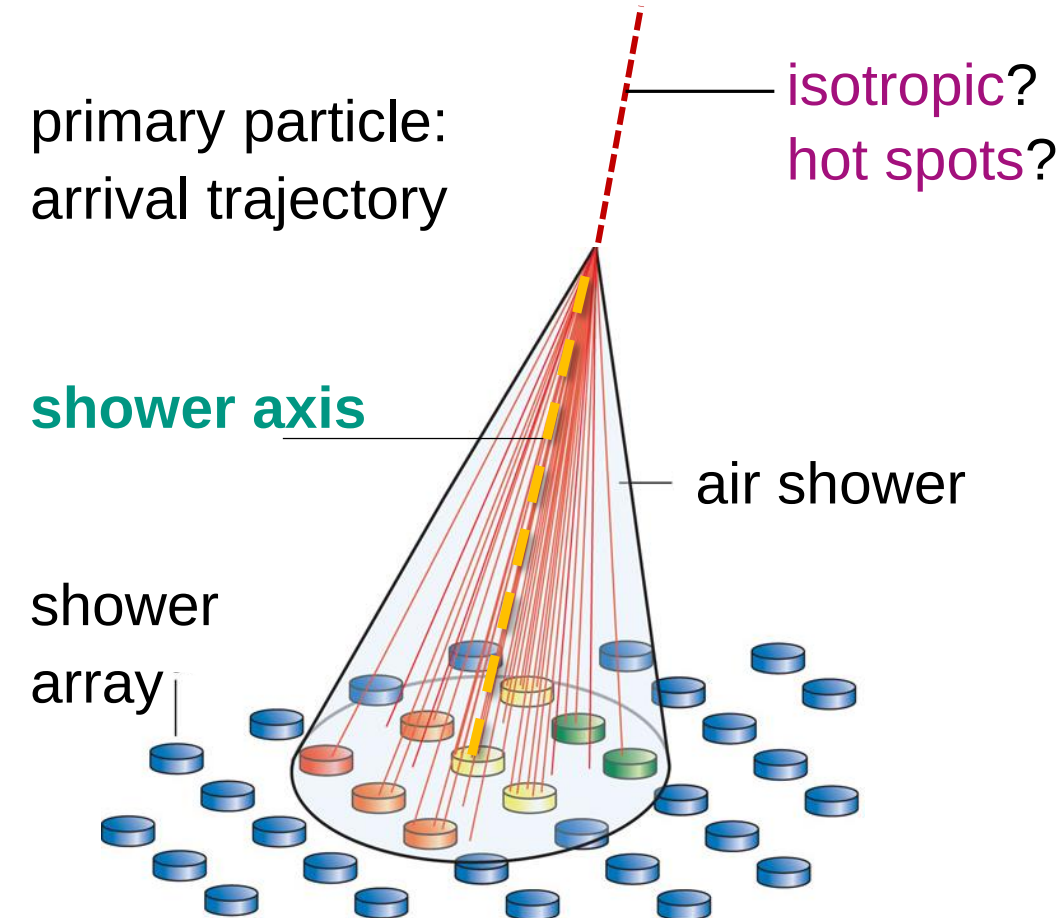
⇒ extra-galactic sources
must be different



What is the degree of isotropy of cosmic rays?

■ Reconstructing the arrival direction of CRs from the **shower axis**

- **galactic CRs:**
many sources, energies $\sim 10^{15} \text{ eV}$,
galactic B –fields result in deflection
 $\Rightarrow$ expect a high degree of **isotropy**
- **extra-galactic CRs:**
fewer sources at scales of UHECRs
 $\Rightarrow$ **nearby sources** could be identified
- the **axis of a large air shower**
allows to reconstruct the arrival
direction of the primary particle



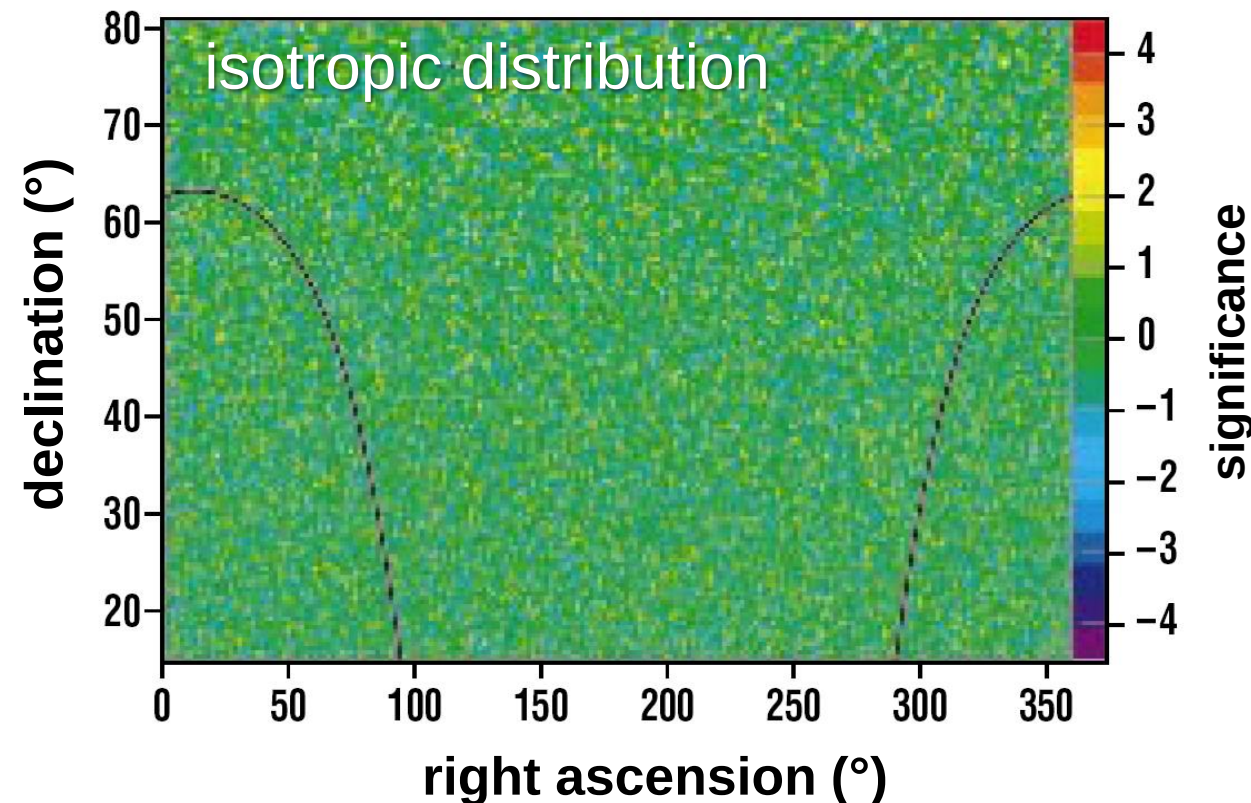
Results for galactic cosmic rays at $E \sim 10^{15} \text{ eV}$

- The observed distribution of galactic CRs is indeed **isotropic** to first order

KASCADE results (>10 years):

- to first order ($\sim 10^{-3}$) one sees an **isotropic arrival distribution**
- conclusions: there are
 - ↪ **no nearby sources** of CRs
 - ↪ **no neutral primaries** (gammas) which initiate showers and which would point back to sources of CRs

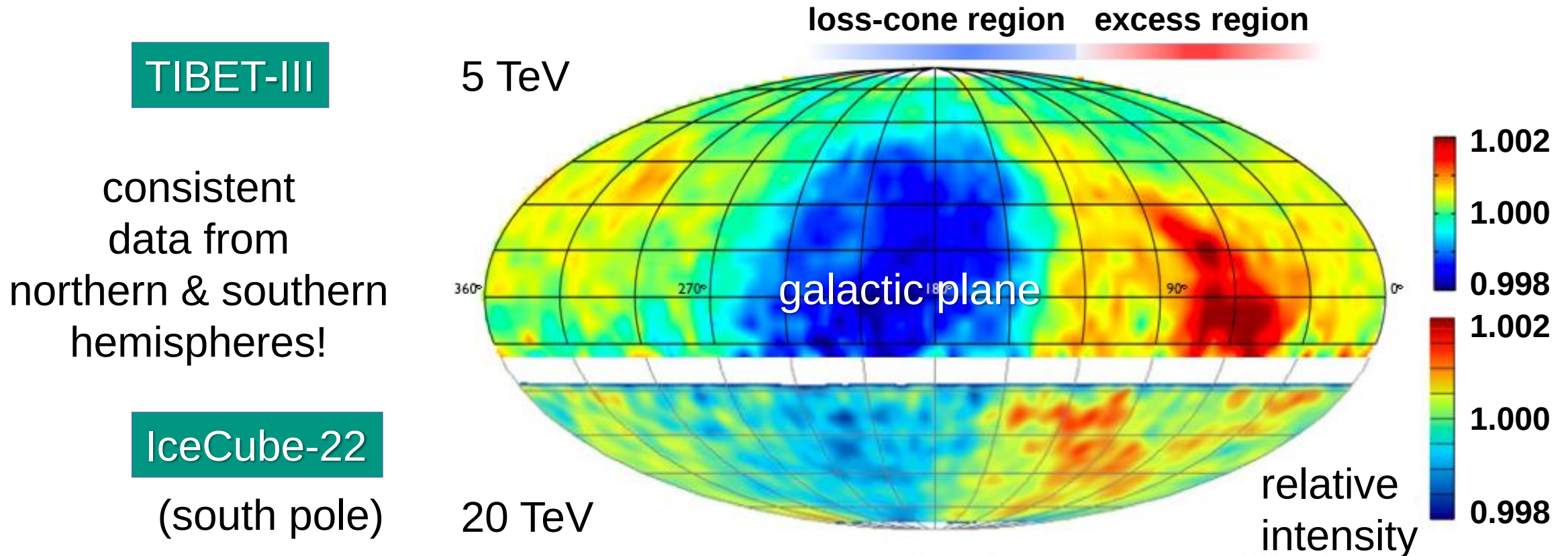
KASCADE: arrival distribution of CR



Looking very precisely: a very small anisotropy!

■ new results of arrival distributions reveal a tiny anisotropy of galactic CRs

- anisotropic arrival distribution with small dipole amplitude $A \sim 6 \cdot 10^{-4}$



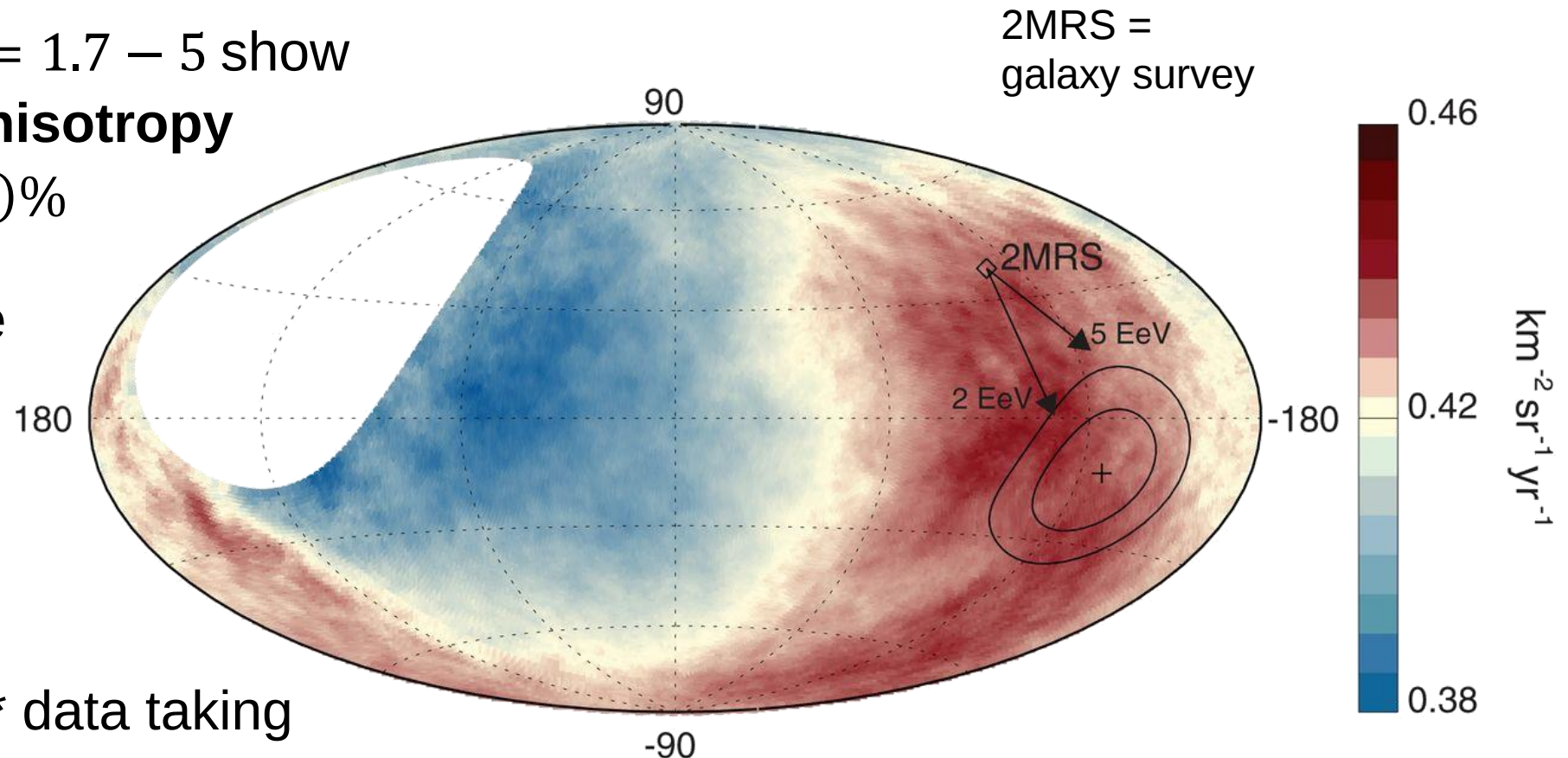
Anisotropy at ultra-high energies: southern view

■ Pierre Auger Observatory* reveals spatial anisotropy at energies $E > 8 \text{ EeV}$

- at the highest EeV-scales: UHECRs with nuclear charge $Z = 1.7 - 5$ show a **small dipole anisotropy** with $A = (6.5 \pm 1)\%$

- UHECRs must be nearby, hot spot despite extended **intergalactic B – fields**

~ 13 years of PAO* data taking

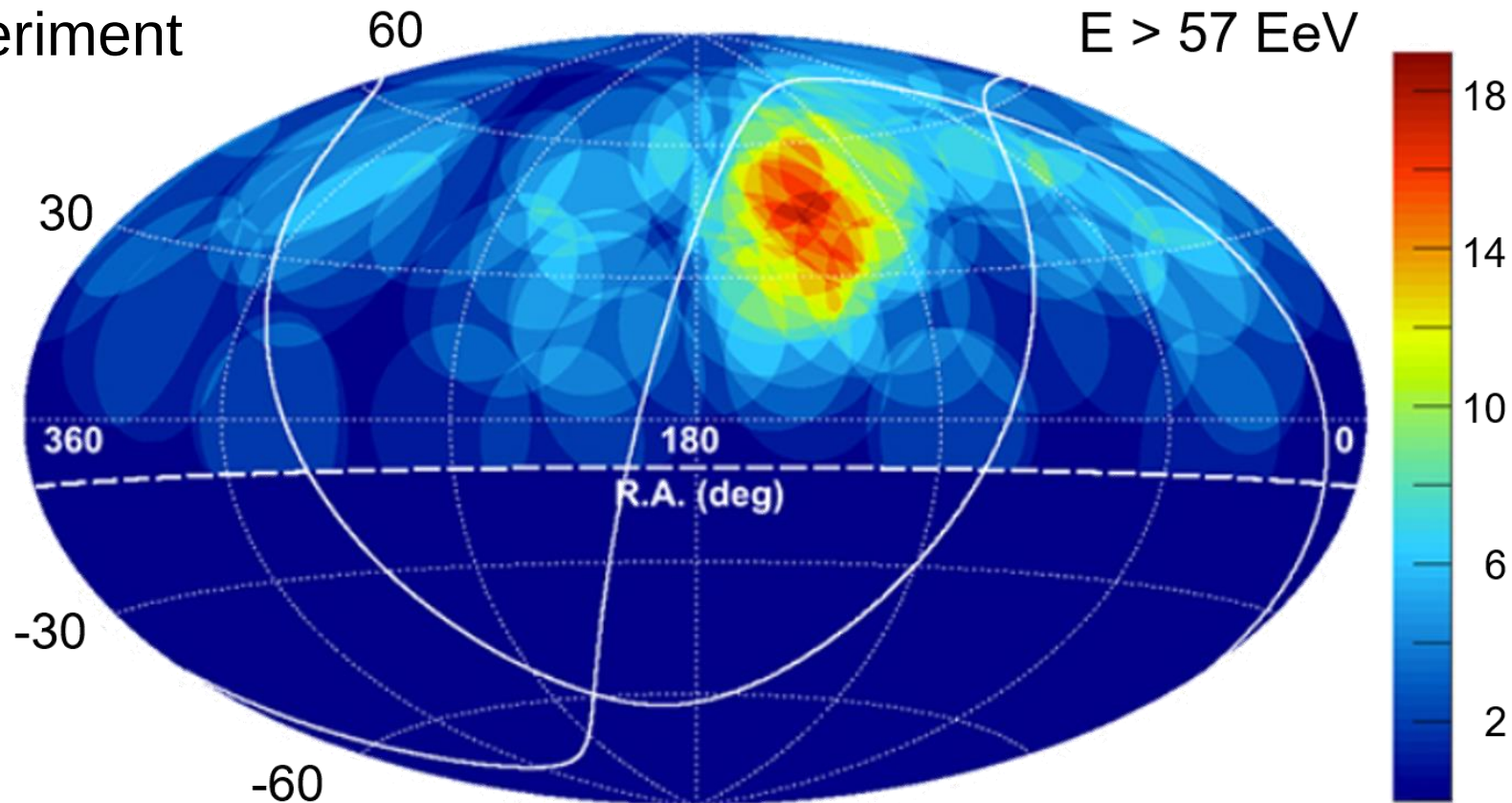


Anisotropy at ultra-high energies: northern view

- The Telescope Array reveals a **localised 'hot spot'** above $E = 5.7 \cdot 10^{19} \text{ eV}$

Telescope Array (TA) experiment
in Utah (USA) observes
UHECRs in northern
hemisphere

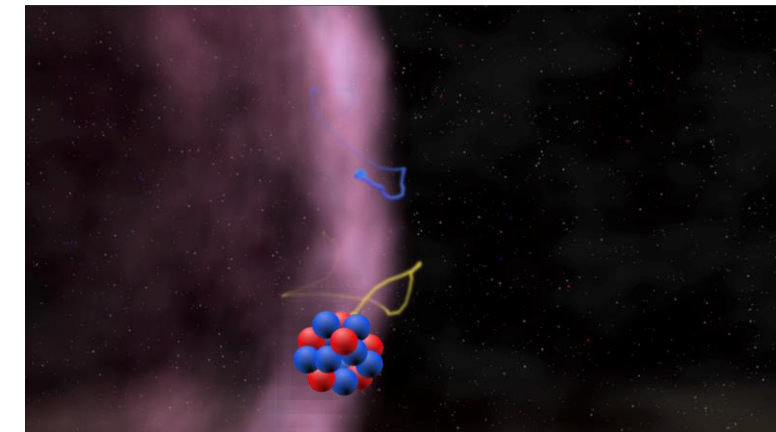
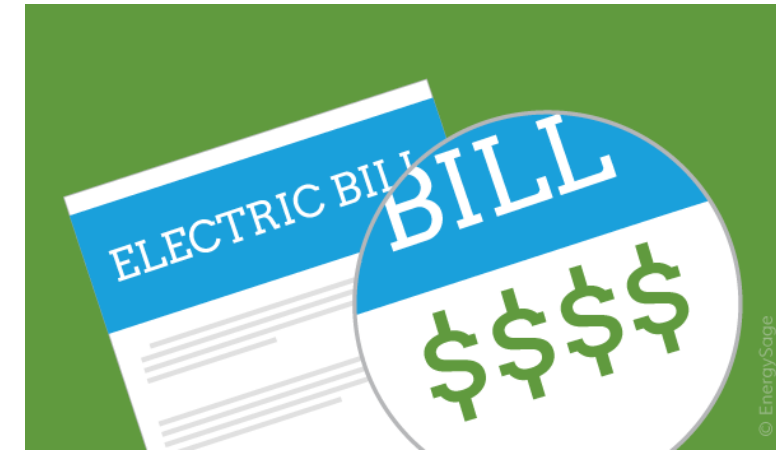
- data taking since 2008:
27 events from only
6% of the area being
surveyed: a **hot spot of
UHECR**



Cosmic accelerators: principles

■ galactic/extra-galactic accelerators: who powers them & how do they do it?

- where does the energy for CRs come from?
- what is the efficiency of cosmic accelerators?
- can cosmic accelerators work over long time scales?
- how do cosmic accelerators work up to $10^{15} \dots 10^{20} \text{ eV}$?
- what limits the energy of galactic accelerators?
- why do extra-galactic accelerators go beyond?

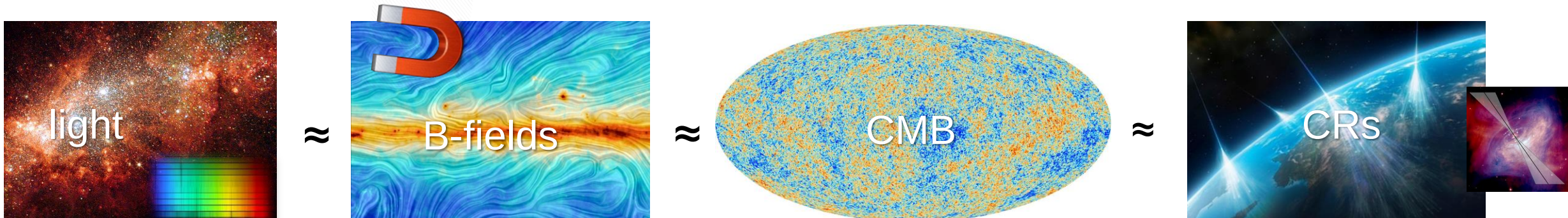


Local (galactic) energy densities: an overview

- CR energy density: **galactic sources** must be able to provide this value

local energy densities in the Milky Way

electromagnetic radiation (star light)	$\sim 0.6 \text{ eV/cm}^3$
galactic magnetic field $\sim B^2 / 2\mu_0$ (3 μ G)	$\sim 0.25 \text{ eV/cm}^3$
cosmic microwave background (CMB)	$\sim 0.26 \text{ eV/cm}^3$
cosmic rays	$\sim 1 \text{ eV/cm}^3$
local matter density (WIMPs)	$\sim 0.3 \text{ GeV/cm}^3$



comparison of W_{SN} & W_{CR} in galaxis

■ Do galactic supernovae provide enough power W to account for CRs?

- Total energy/year W_{CR} going into galactic cosmic rays

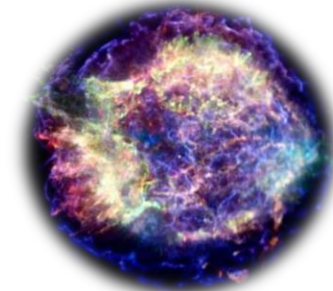
$$W_{CR} = \rho_{CR} \cdot \pi \cdot R^2 \cdot d \cdot \tau^{-1} = 2 \cdot 10^{41} \text{ J/year}$$



CR

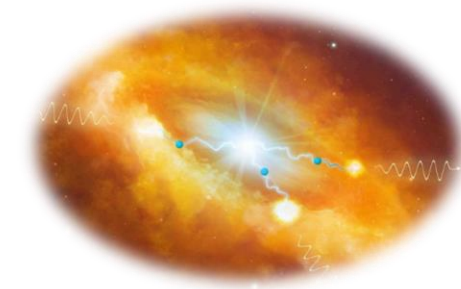
- Total energy/year W_{SN} produced by supernovae in galaxis

$$W_{SN} \sim 3 \text{ SN-explosions/century} = 5 \cdot 10^{42} \text{ J/year}$$



SNR

- **Supernovae** = ideal candidate sites* for galactic CR
required efficiency for $W_{CR} \rightarrow W_{SN} : < 5\%$

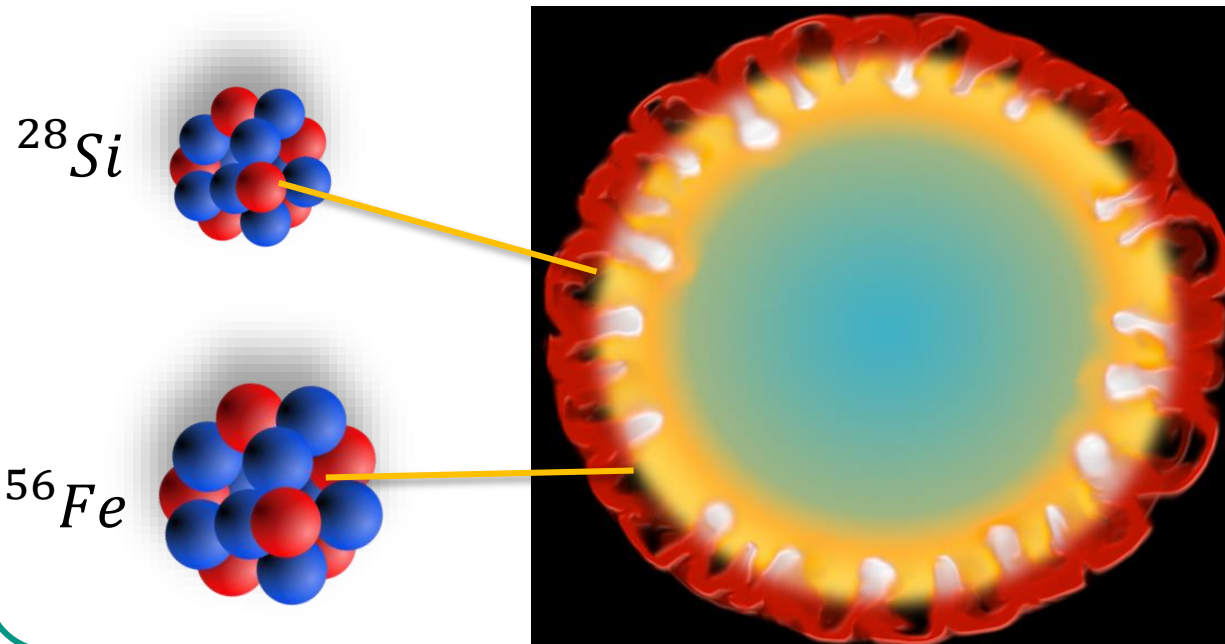


Sources of Cosmic Rays in our galaxy

■ Most likely Cosmic Ray sources in our galaxy at $E \sim 10^{15} \text{ eV}$

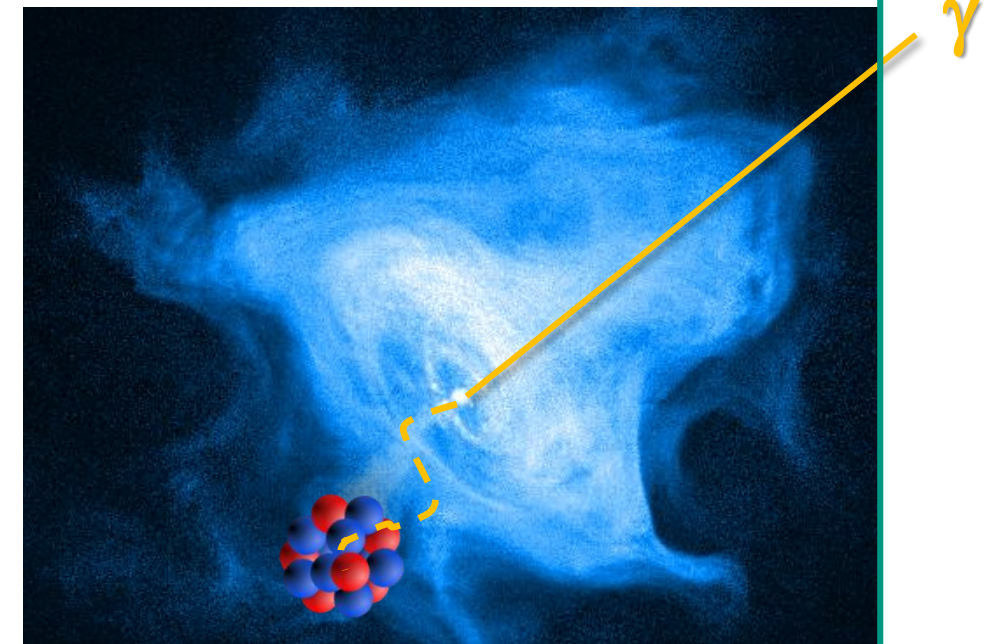
Supernova shock fronts

- shock fronts in SNR: energy distribution ?
maximum energy E_0 ? what nuclei A_Z ?



Pulsars, pulsar wind nebulae

- acceleration mechanism:
hadronic / leptonic schemes ?



Accleration mechanism for CRs

■ Fundamental physics principles to accelerate particles

- dynamical

scattering (reflection) of particles in magnetic clouds (Fermi-acceleration)

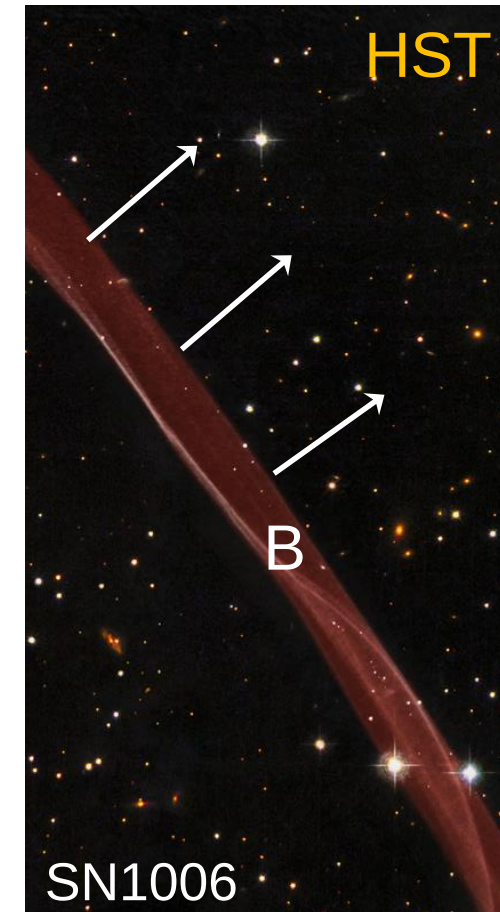
- hydrodynamic

accleration in plasma sheets

- electromagnetic

time-variable E, B – fields

$$\frac{d}{dt}(\gamma \cdot m \cdot \vec{v}) = e \cdot (\vec{E} + \vec{v} \times \vec{B}) \quad \text{with} \quad \vec{v} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

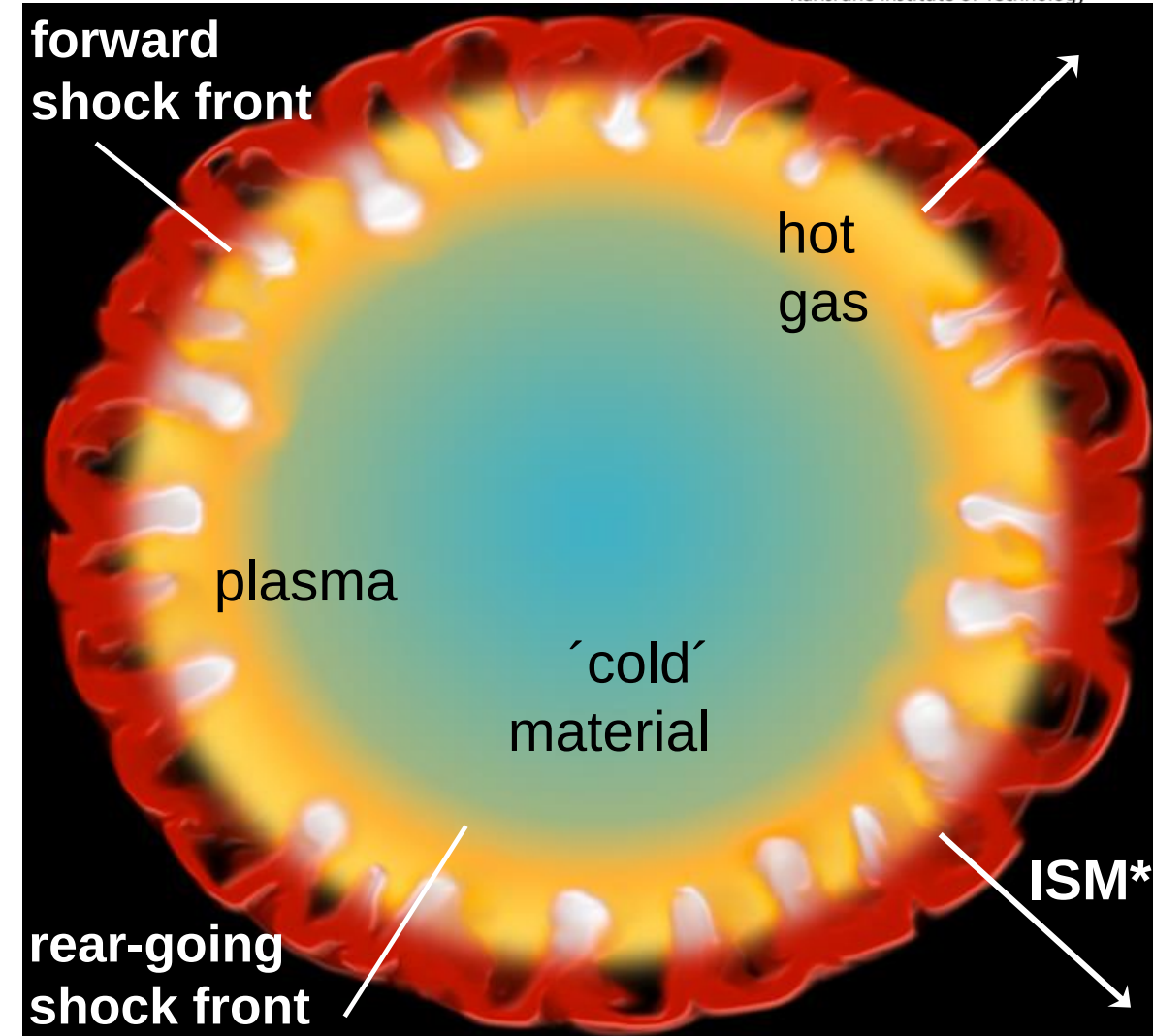
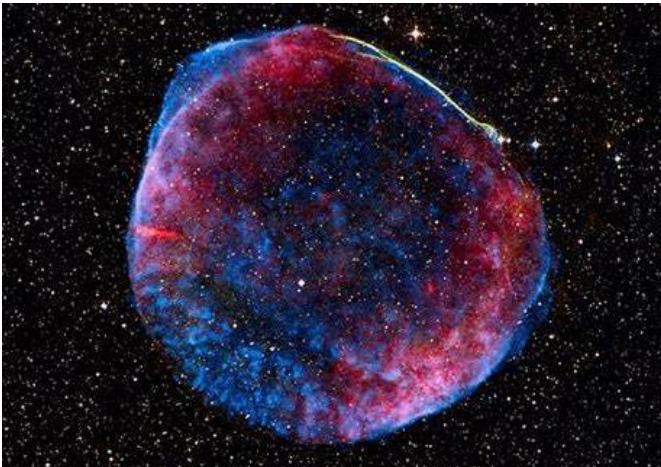


Fermi acceleration in shock fronts of SN remnants

■ Simplified 1D-scenario

- SN-shock front:

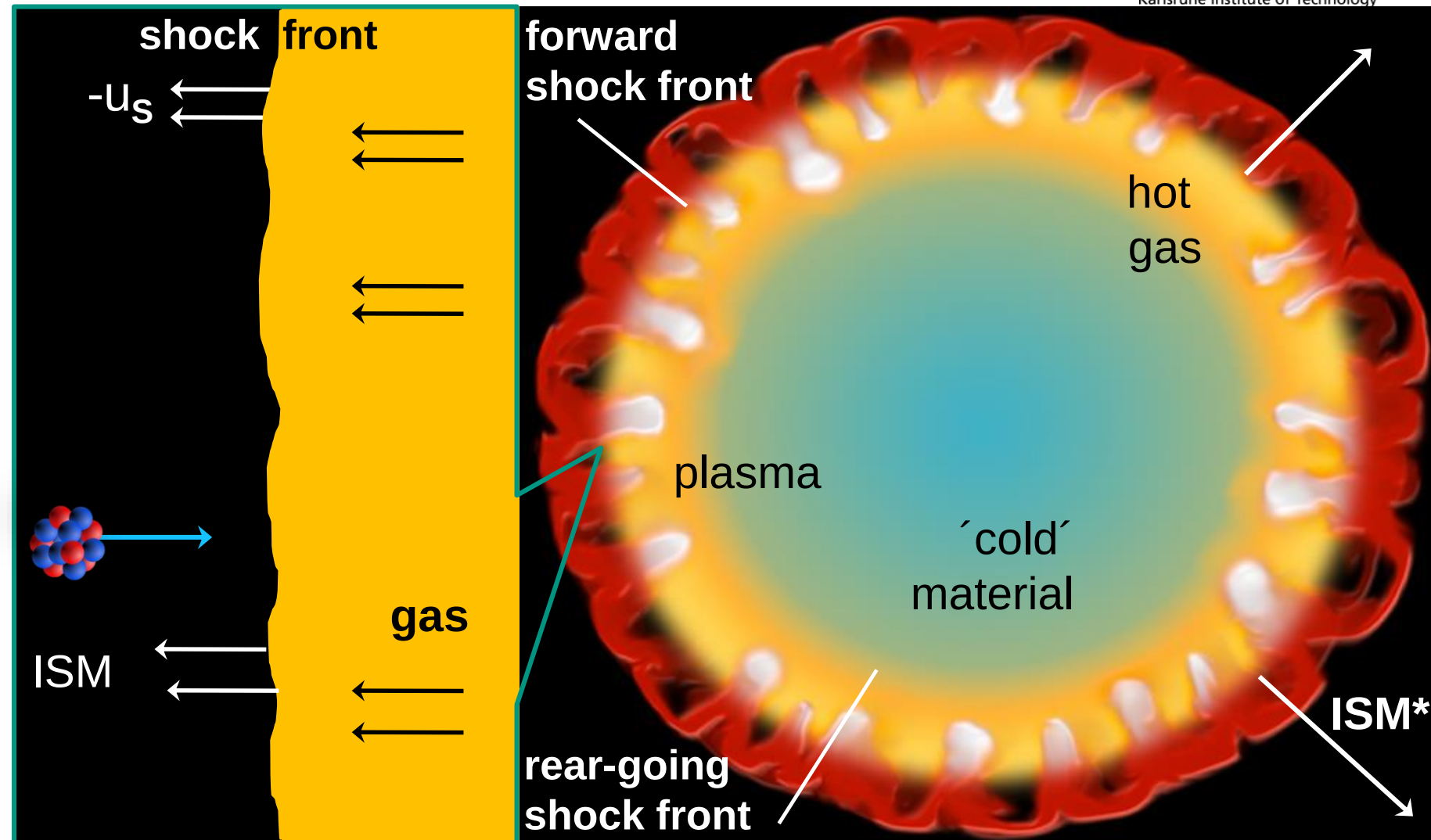
spreads out into the very thin ISM*
over extended time scales of
 $t = 10^4 \dots 10^5 y$.



Fermi acceleration in shock fronts of SN remnants

■ Simplified 1D-scenario

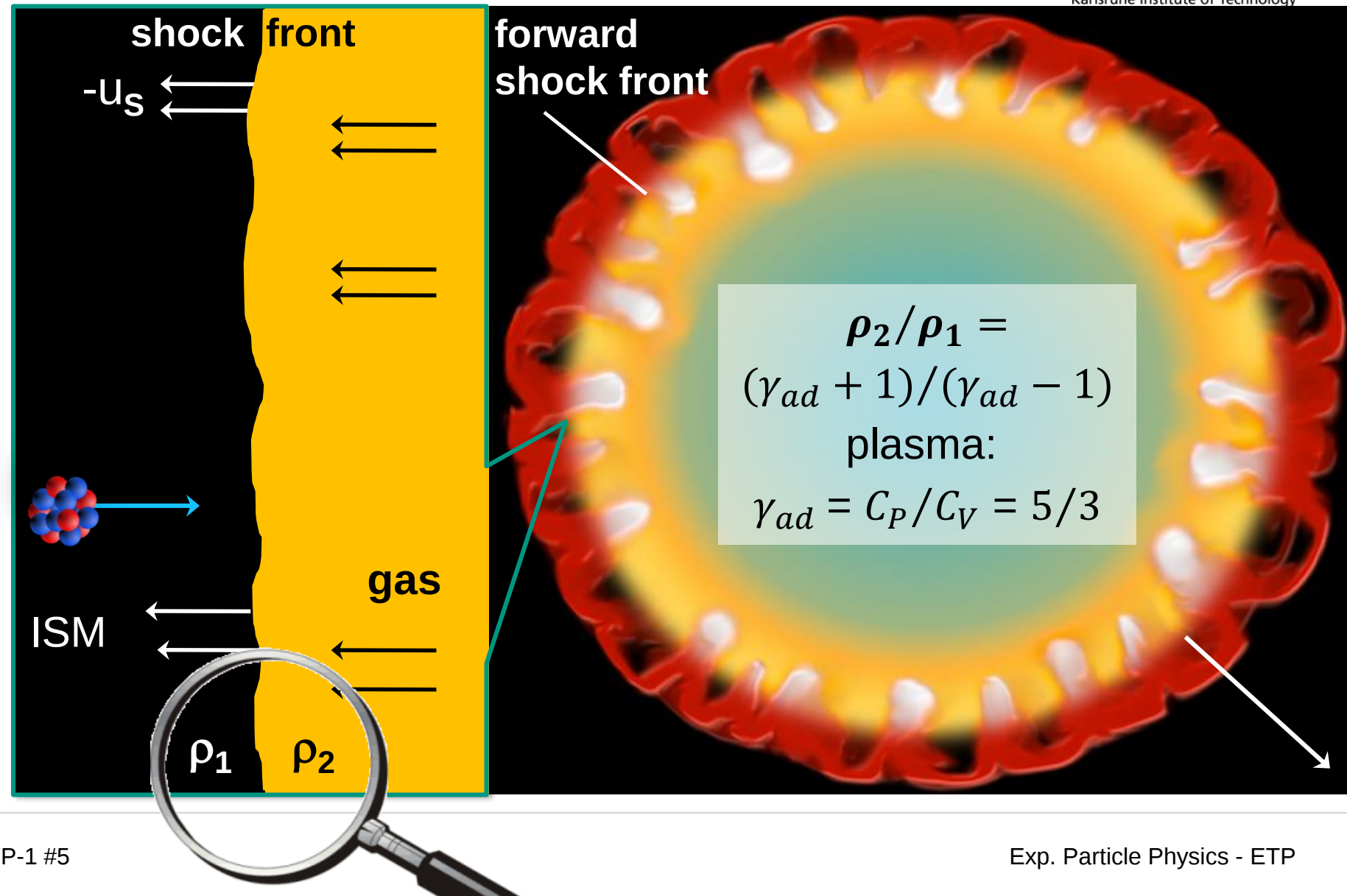
- SN-shockfront propagates with $u_s \cong 10^4 \text{ km/s}$ or $\beta_s = 0.03$
- β_s decreases over time



Fermi acceleration in shock fronts of SN remnants

■ important: there is a density gradient at the transition to ISM

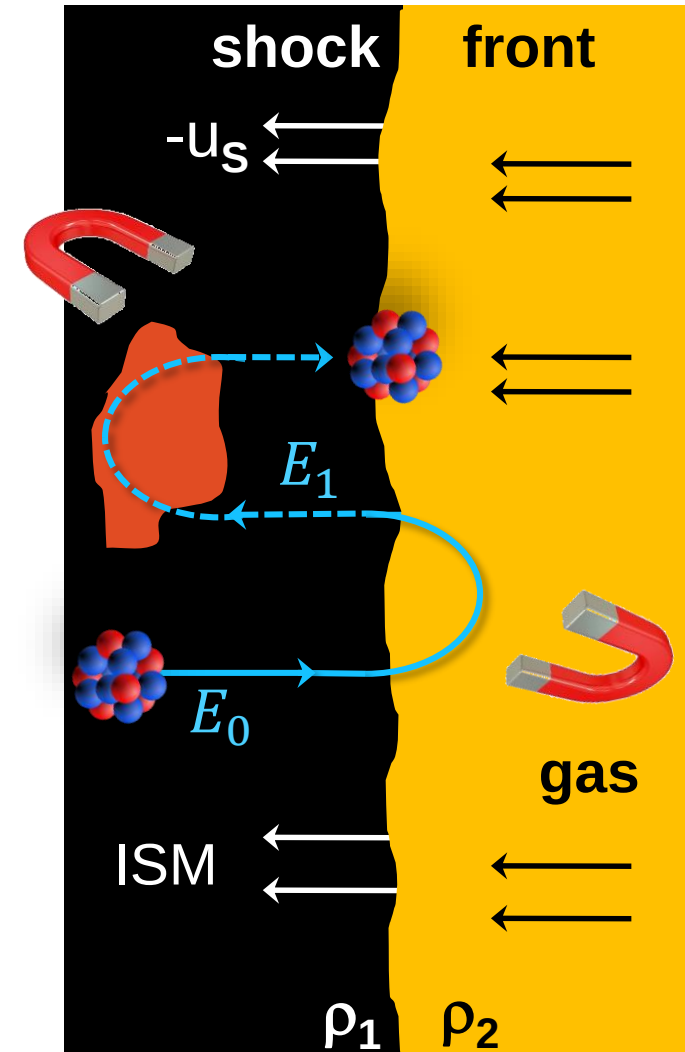
- ISM: density ρ_1 , shock: density ρ_2
- plasma physics: ratio ρ_2/ρ_1 depends on the adiabatic coefficient γ_{ad}



Fermi acceleration in shock fronts of SN remnants

■ The process of Fermi acceleration of CRs comprises several key steps:

- charged nucleus from ISM with **primary energy E_0** passes the SN-shock front: **energy gain ΔE**
- B-fields within the gas of the shock front back-reflect the nucleus: **adiabatic process** without energy loss!
- nucleus again passes pressure gradient (in opposite direction): again **energy gain ΔE**
- further adiabatic backscattering of nucleus due to electro-magnetic fields in the ISM

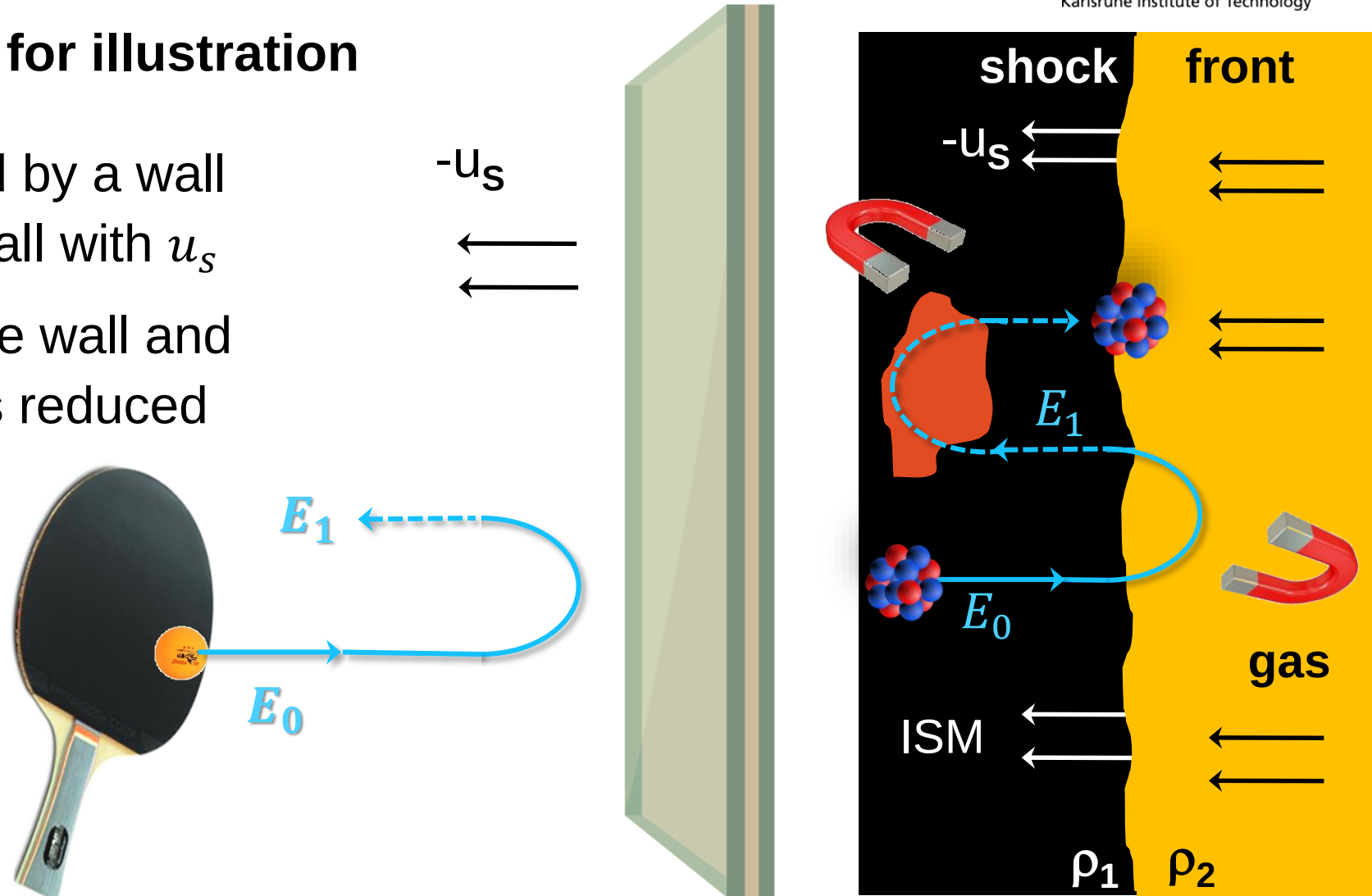


Fermi acceleration in shock fronts: an analogon

■ a mechanic analogon for illustration

- ping pong ball reflected by a wall which moves against ball with u_s
- ⇒ distance between the wall and table tennis racket is reduced
- ⇒ ball is accelerated, and (if this can be repeated) it will be ever faster

$n = 1$



Fermi acceleration in shock fronts: a net gain

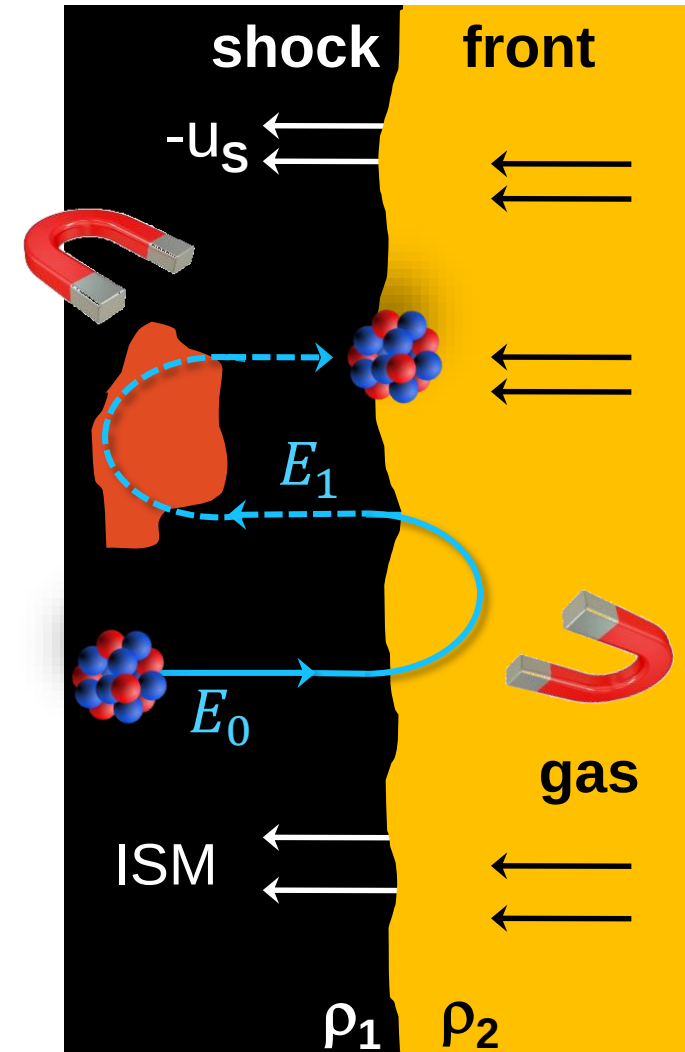
■ Summing up: the CR nucleus gains energy

- by Fermi acceleration, which is ...
- ... based on the **multiple passing** of a SN-shockfront
- ... a collision-less process, i.e. **no energy losses** due to inelastic scattering in the gas of the shock front

■ Energy gain ΔE

for a single acceleration cycle (& independent of the direction of the nucleus) one obtains a net energy gain

$$\frac{\Delta E}{E} = \frac{u_s}{c} = \beta_s$$



Fermi acceleration in shock fronts: many cycles

■ Summing up: single cycle vs. multiple cycles

- net energy gain ΔE per single acceleration cycle

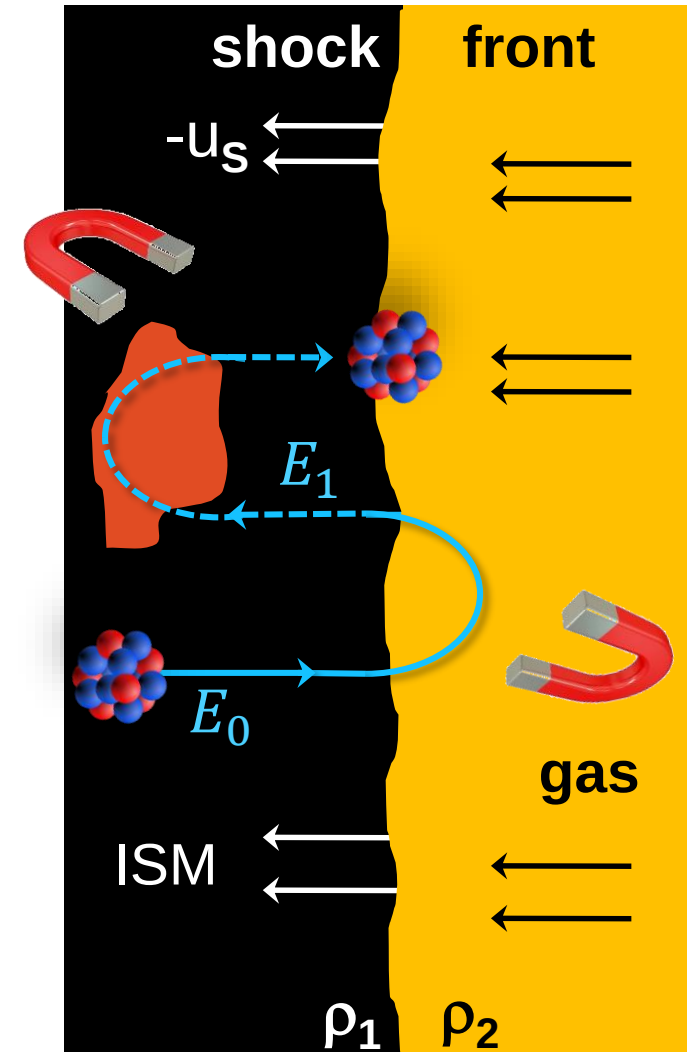
$$\Delta E = \alpha \cdot E$$

- CR energy E after n acceleration cycles (for starting energy E_0)

$$E = E_0 \cdot (1 + \alpha)^n$$

- number n of acceleration cycles to reach max. energy E :

$$n = \frac{\ln(E / E_0)}{\ln(1 + \alpha)}$$



Origin of the CR power law distribution

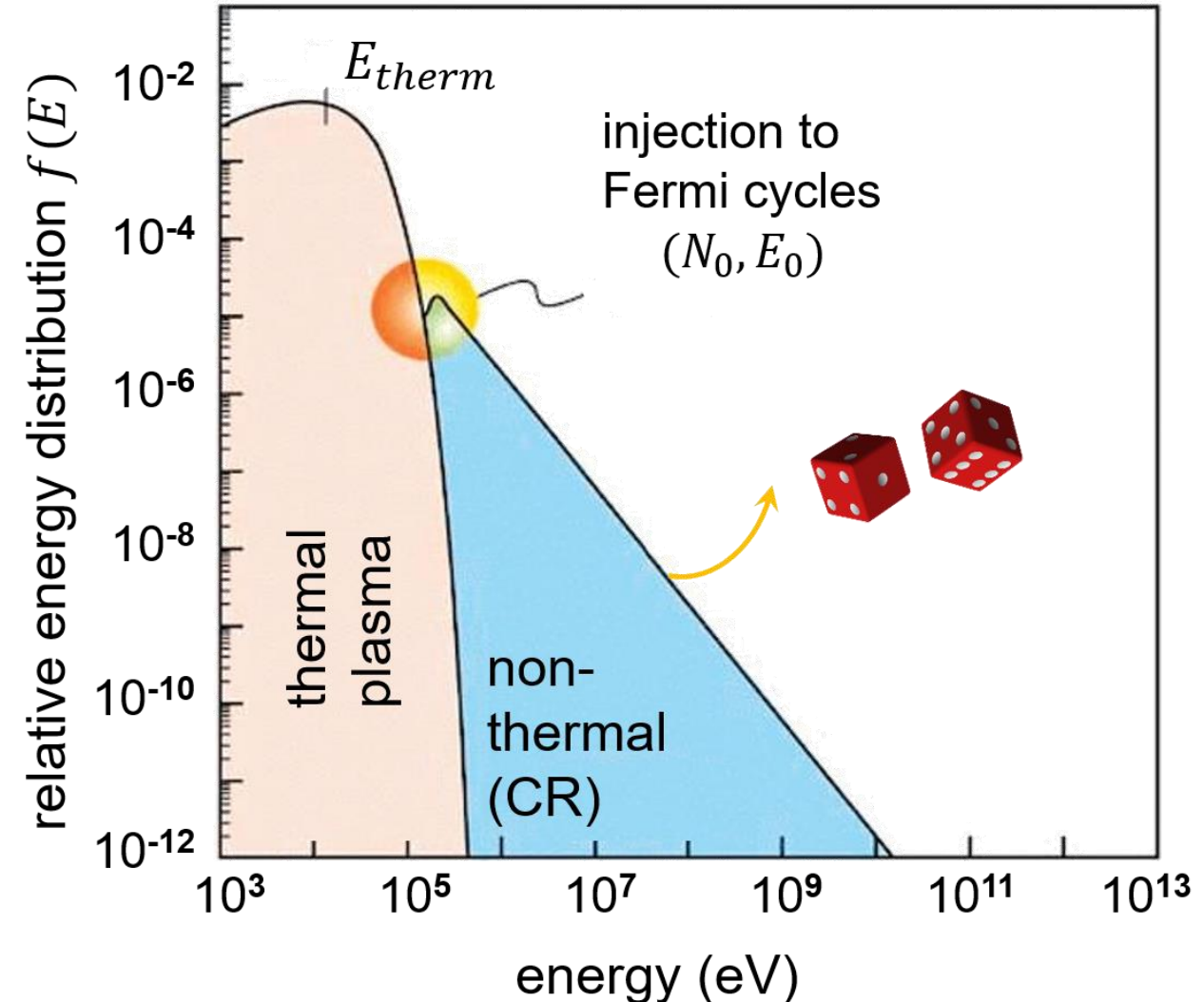
■ power law distribution due to Fermi acceleration

- taking into account an energy-dependent probability P for particle losses during the acceleration cycle, we obtain a power-law distribution



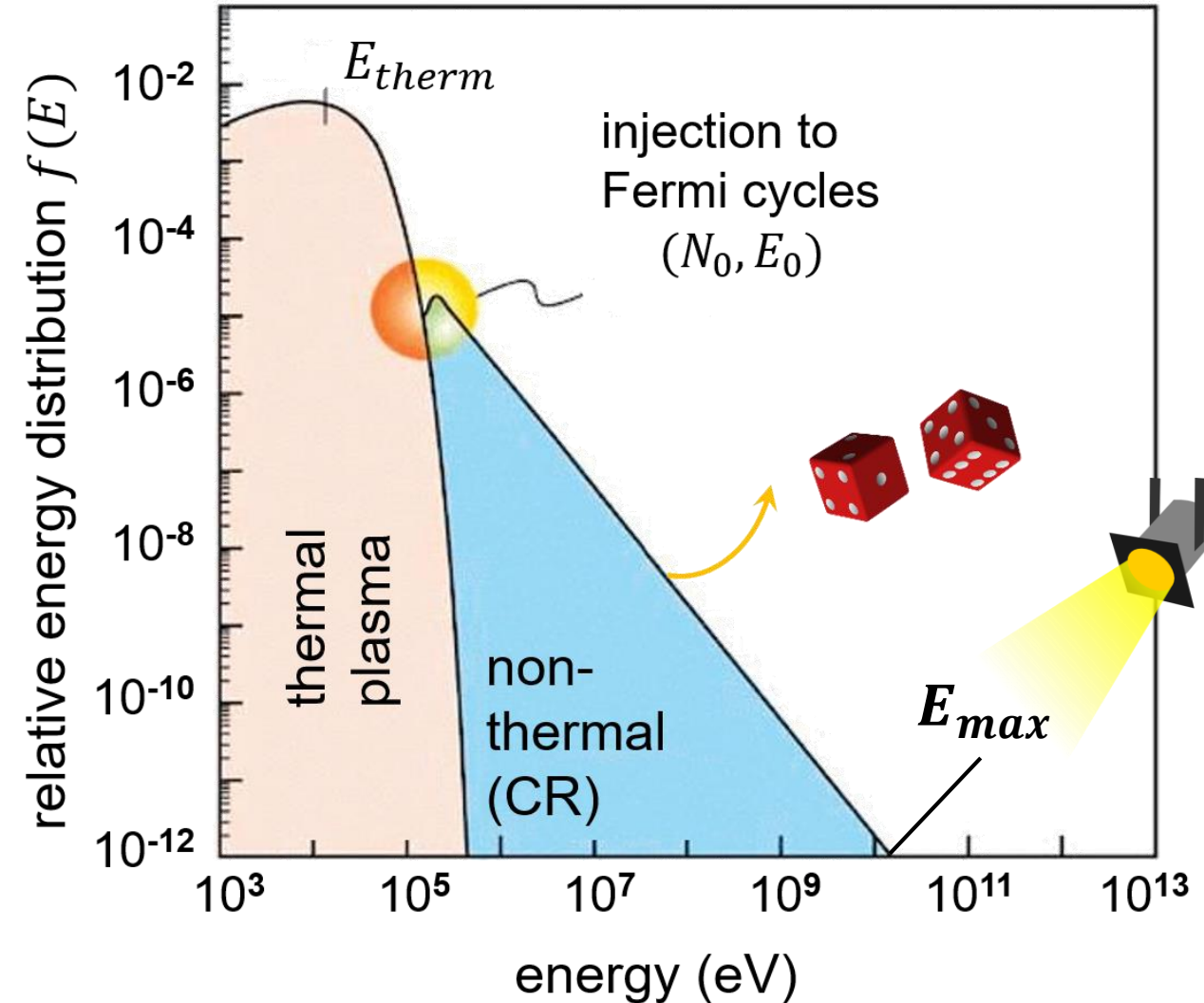
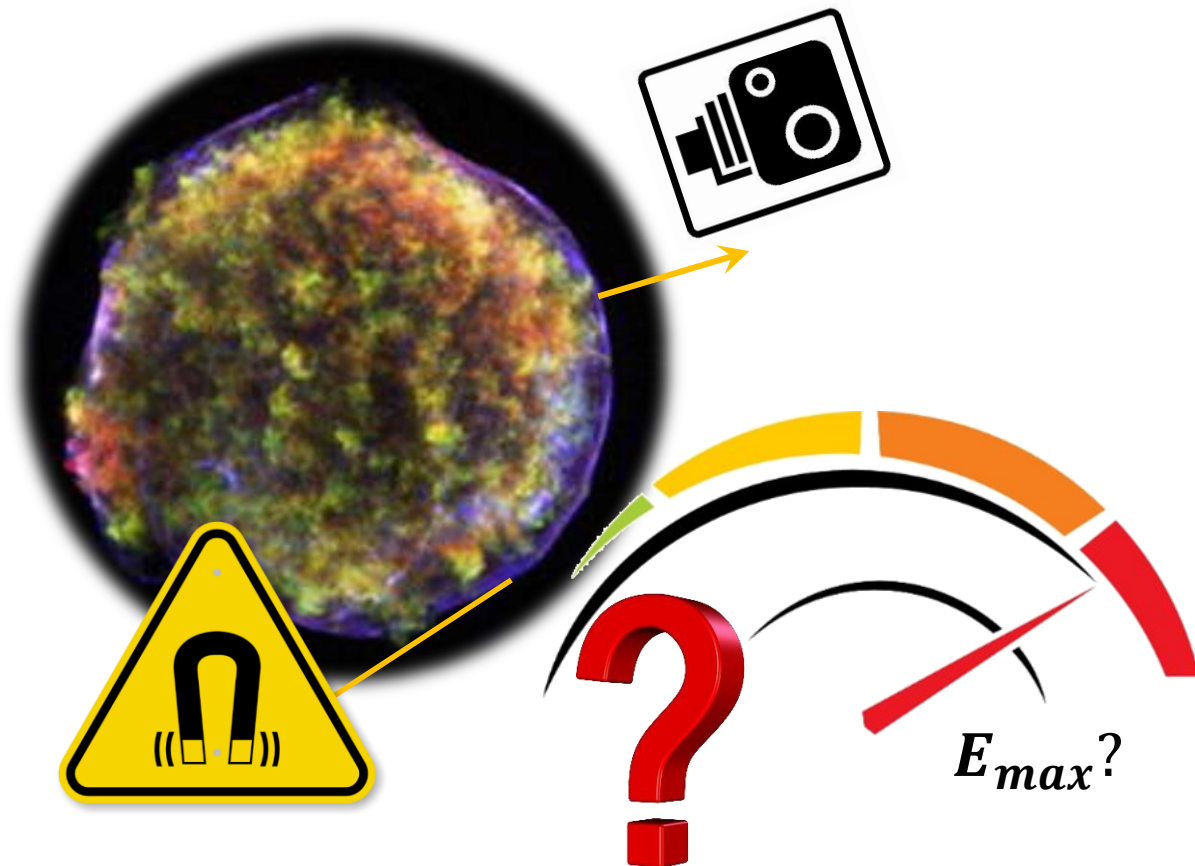
$$\frac{dN(E)}{dE} \sim E^{-2.7}$$

power law spectrum



Maximum energy of CR due to SN shocks?

- What are the parameters defining the maximum CR energy E_{max} ?

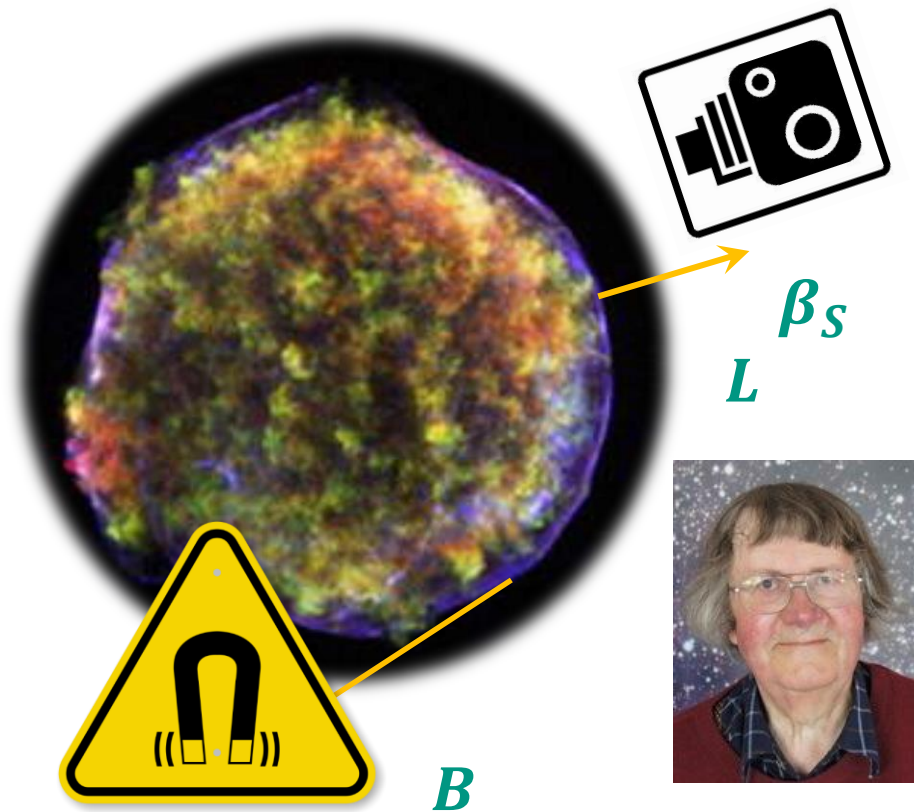


Maximum energy of CR via famous Hillas formula

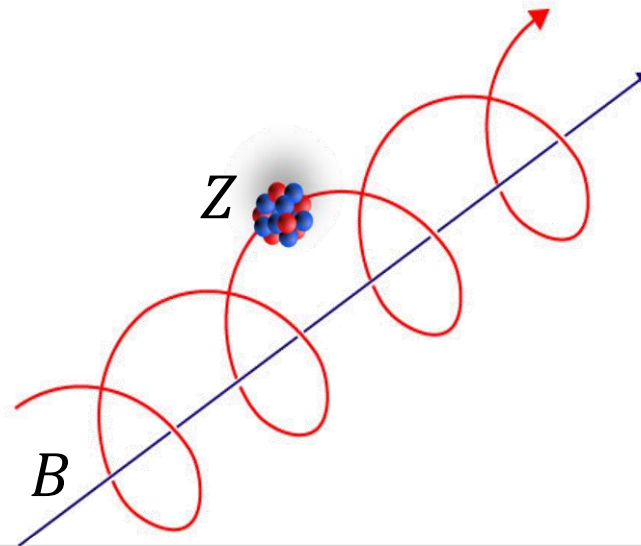
■ A.M. Hillas: cosmic accelerators use a B-field of size L to guide particles

- maximum energy E_{max} of a particle of **charge Z** in a SN shock front (see dimensions!)

$$E_{max} \sim \beta_s \cdot Z \cdot B \cdot L$$



A. M. Hillas



particle:

- nuclear charge Z

source:

- field strength B

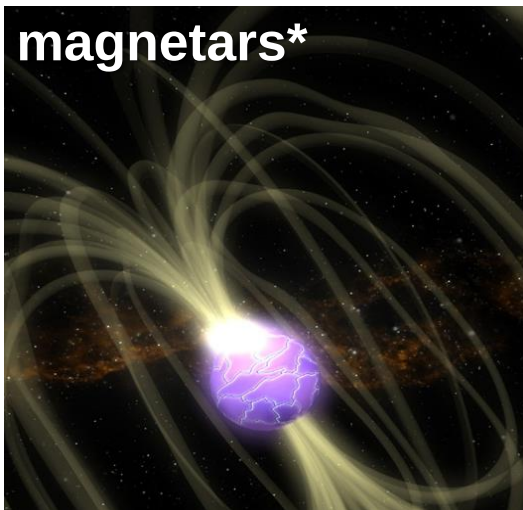
- size L

- shock velocity β_s

Hillas plot for UHECRs up to $E_{max} = 10^{20} \text{ eV}$

■ candidate sites

small L / large B



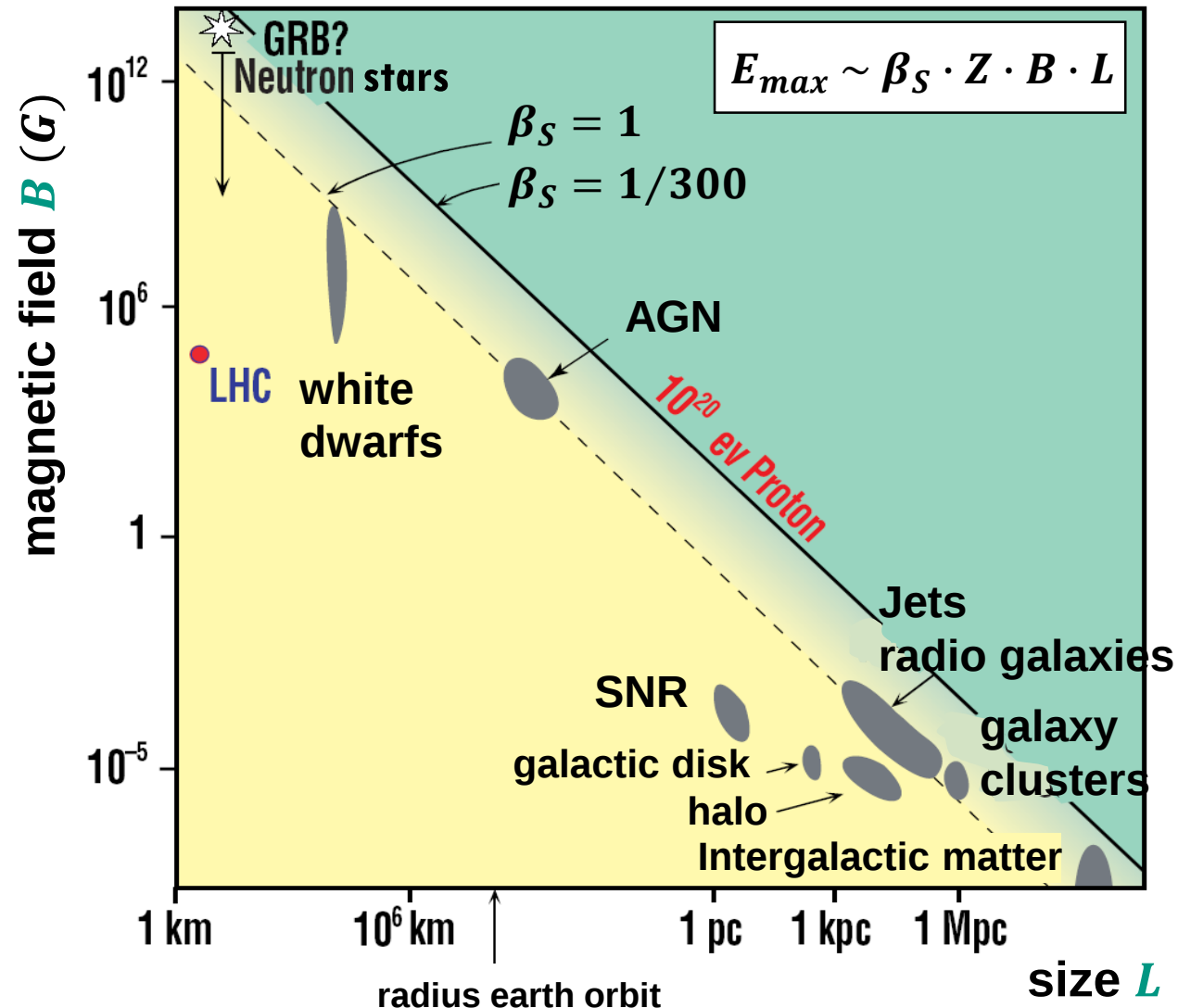
magnetars*

B-field

... 10^9 T

dimension

10 km ...



large L / small B



B-field

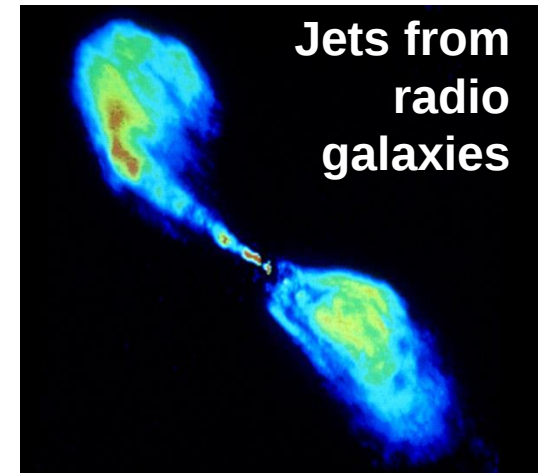
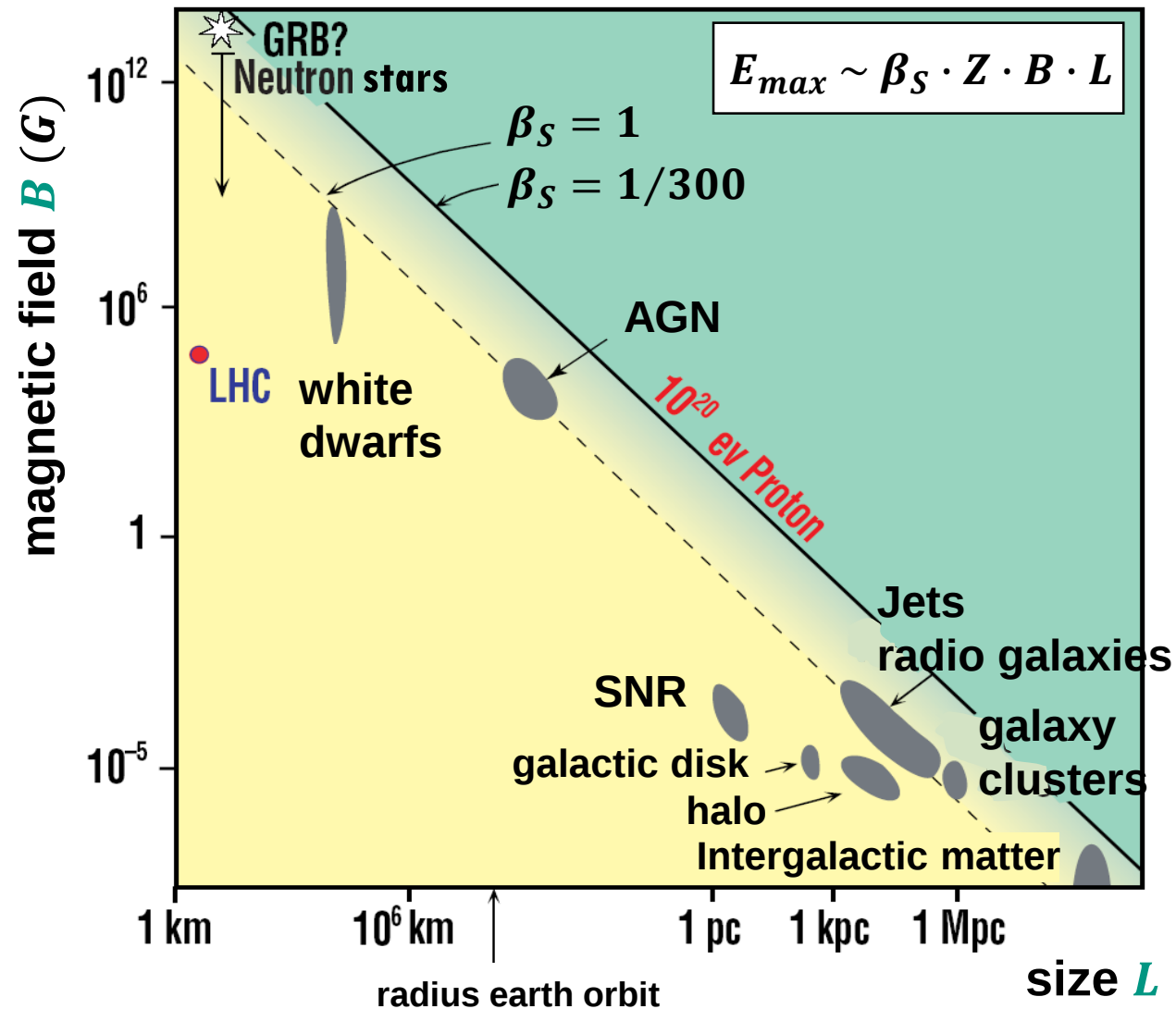
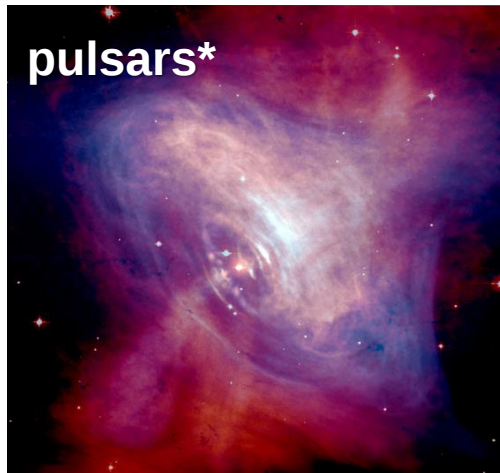
... 10^{-9} T

dimension

1 Mpc ...

Sources of UHECRs up to $E_{max} = 10^{20} \text{ eV}$

■ candidate sites



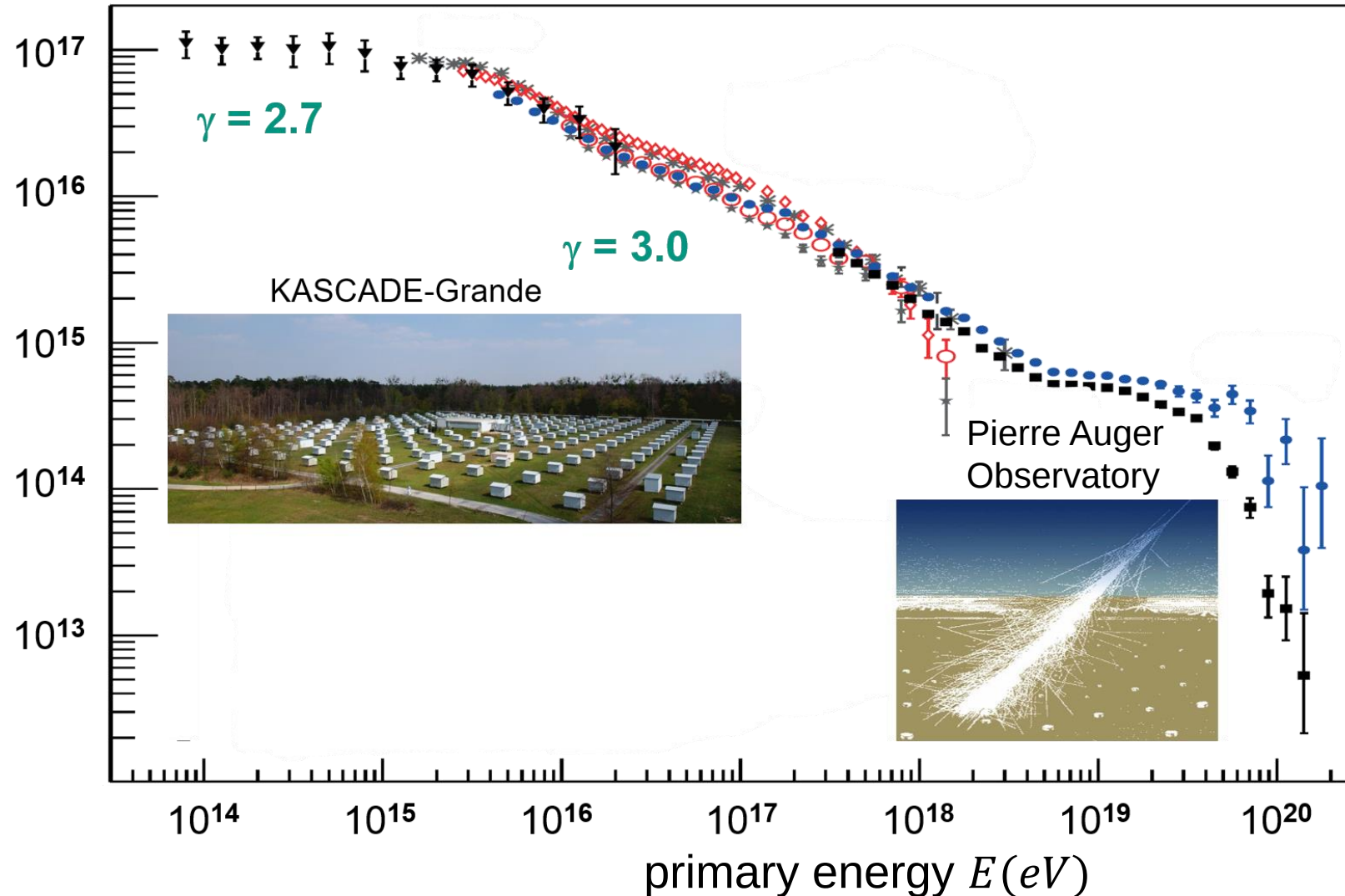
Energy spectrum of UHECR

■ Overview of results

- measurements of air showers arrays with KIT participation

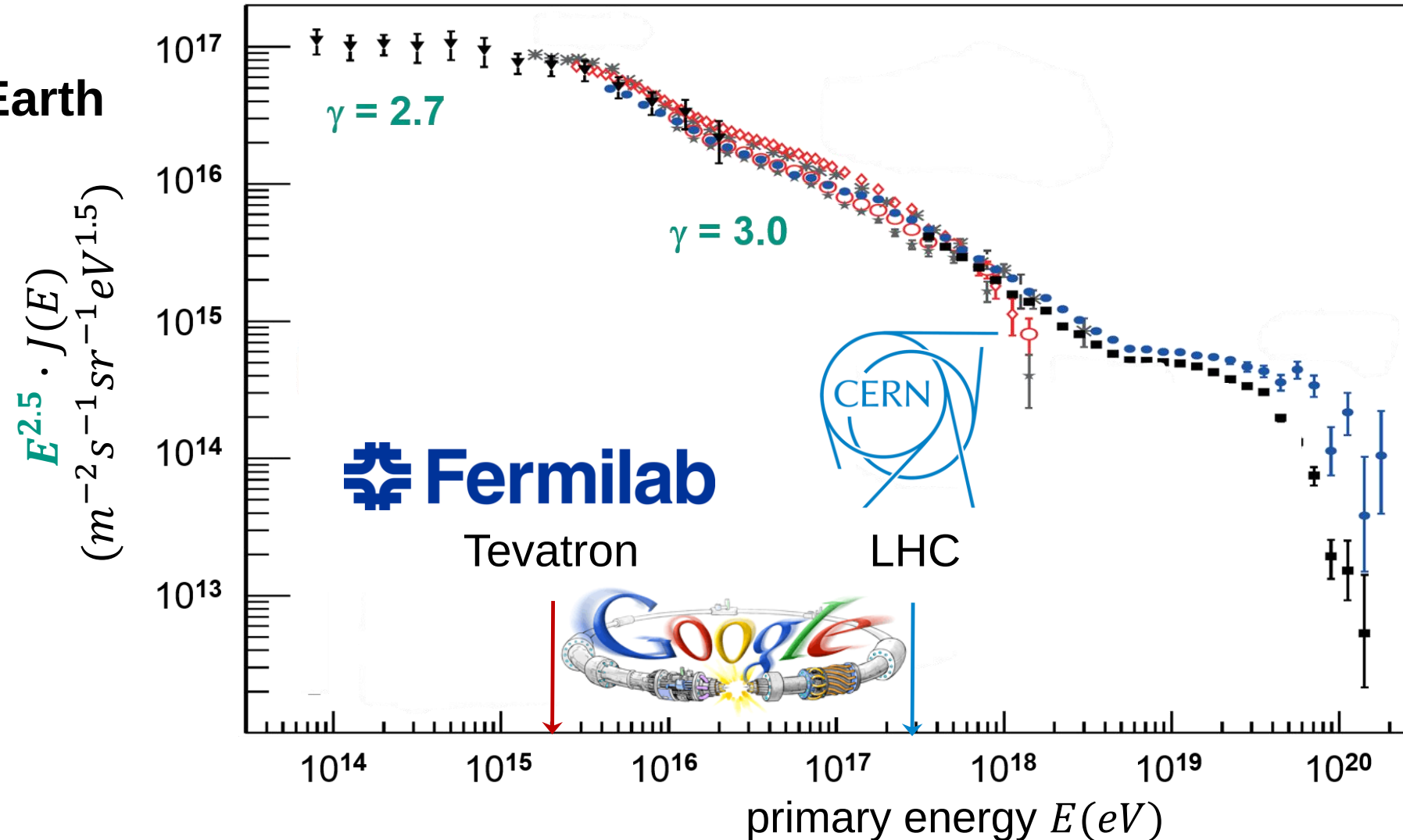
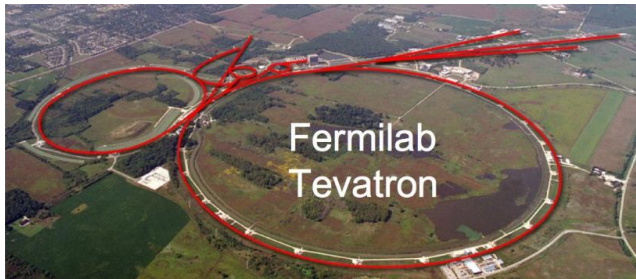


$$E^{2.5} \cdot J(E) \quad (m^{-2} s^{-1} sr^{-1} eV^{1.5})$$



Energy spectrum of UHECR & accelerators

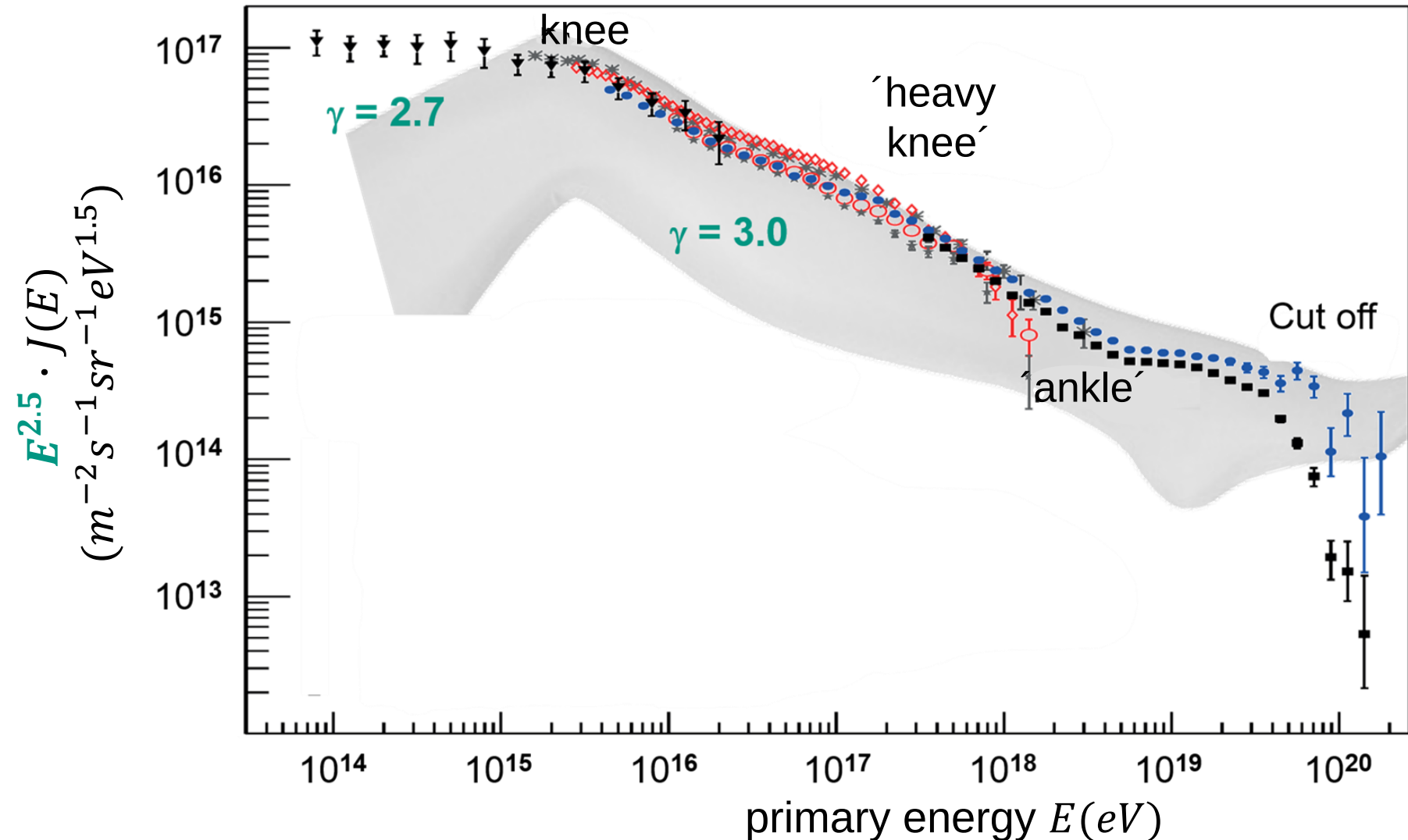
■ Comparison to accelerators on Earth



Energy spectrum of high-energy CRs & features

■ Energy spectrum influenced by:

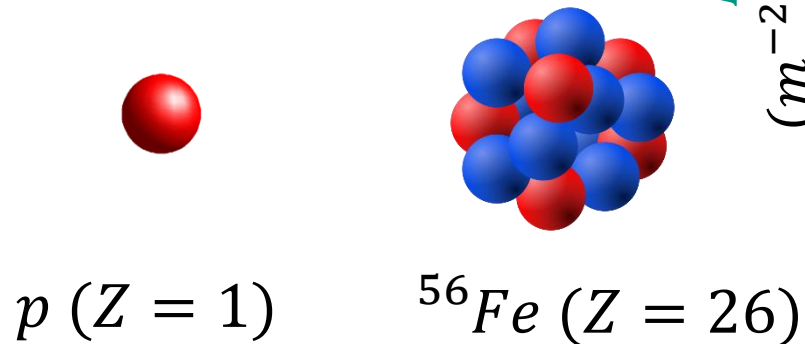
- change of mass composition (**light** → **heavy**): Hillas relation
- change of sources & CR acceleration mechanisms (**galactic** → **extra-galactic**)



Energy spectrum of high-energy CRs & features

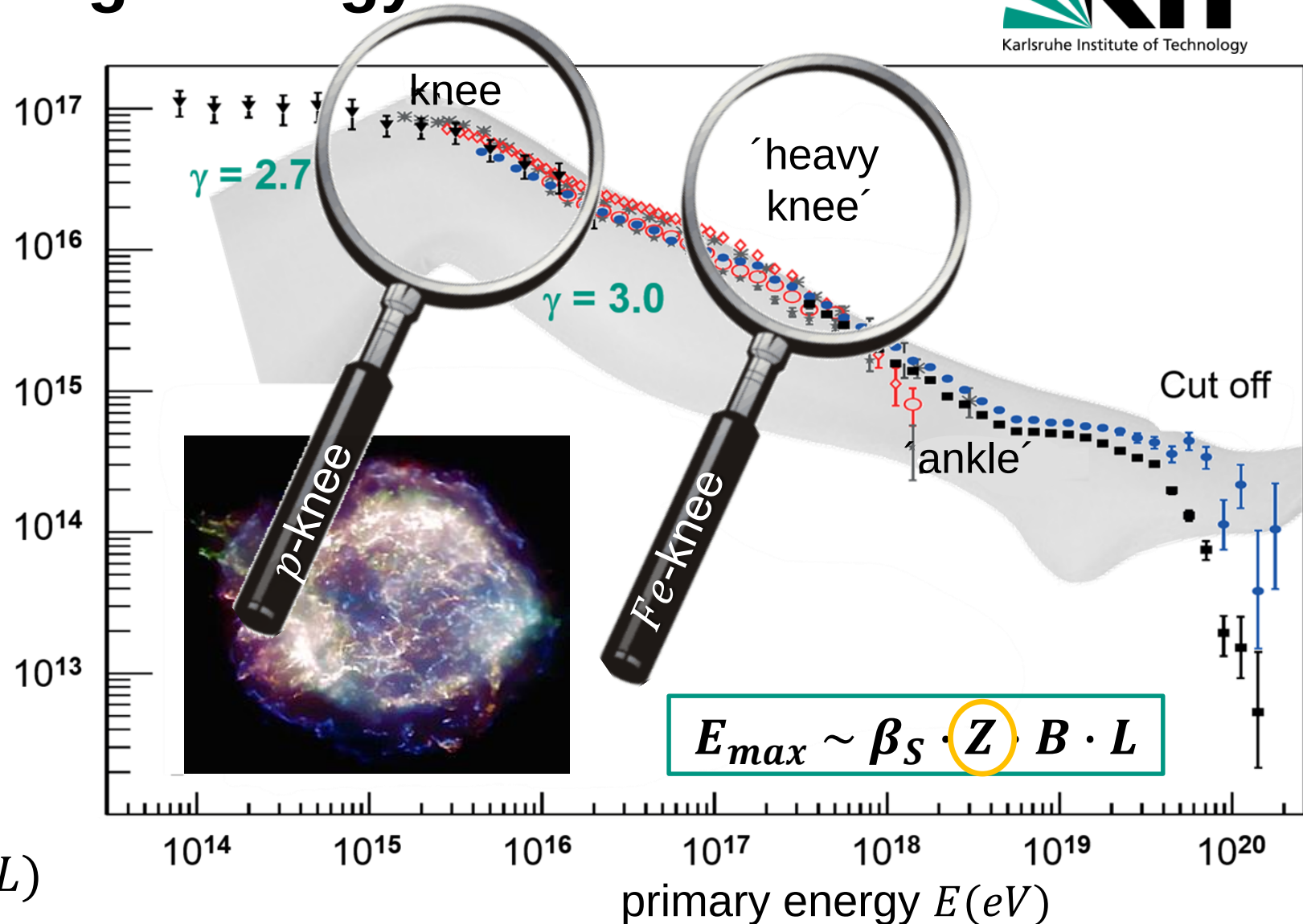
■ Energy spectrum of galactic CRs:

- change of mass composition (**light** → **heavy**):
- Hillas relation



sources: SNR in galaxis (B, L)

$$E^{2.5} \cdot J(E) \quad (m^{-2} s^{-1} sr^{-1} eV^{1.5})$$



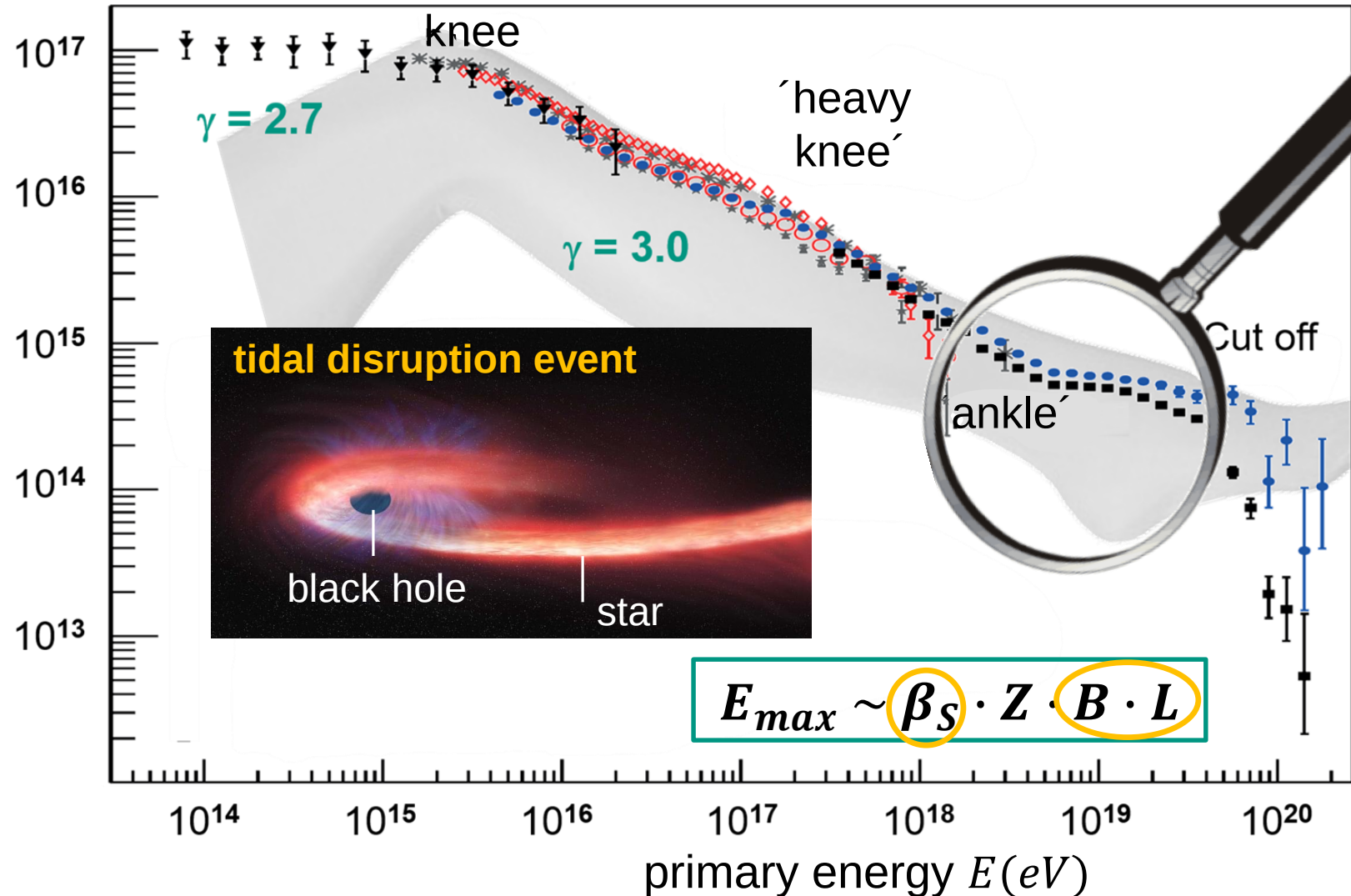
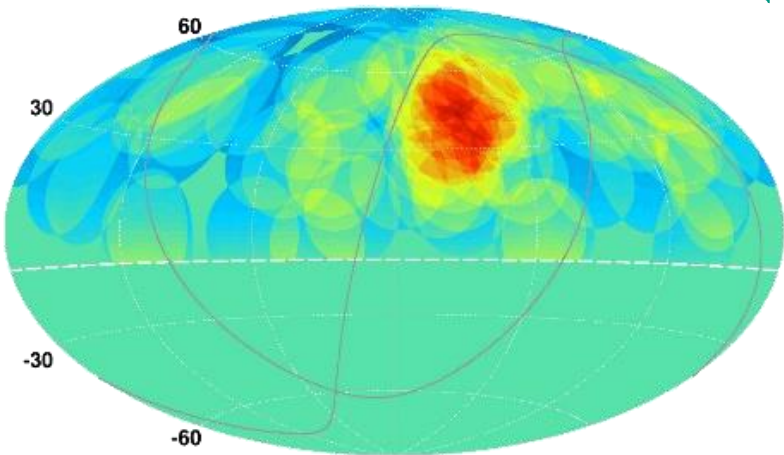
Energy spectrum of high-energy CRs & features

■ Energy spectrum of **extra-galactic CRs**:

- change of acceleration sites

tidal disruptions?

$$E^{2.5} \cdot J(E) \quad (m^{-2} s^{-1} sr^{-1} eV^{1.5})$$



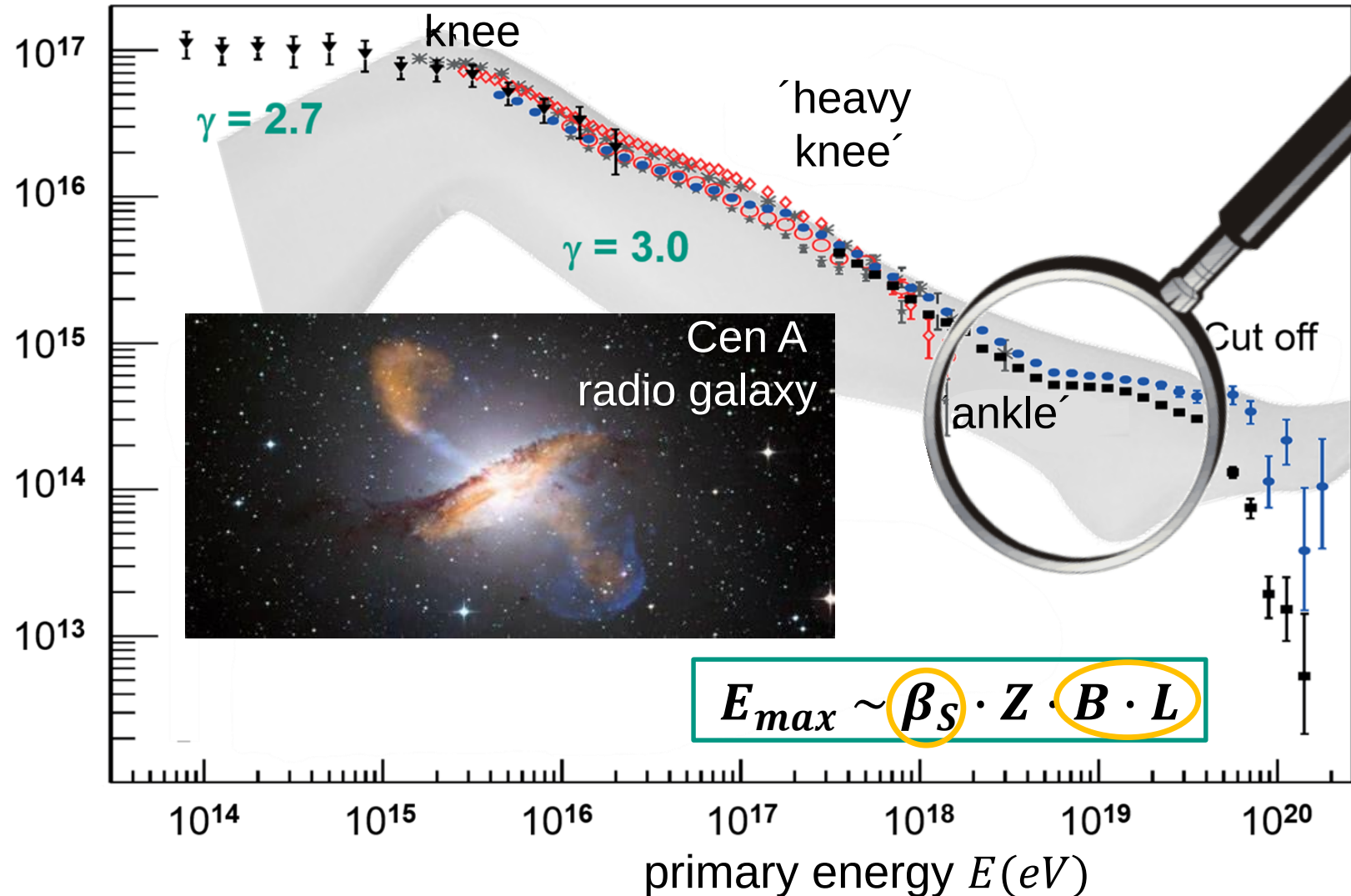
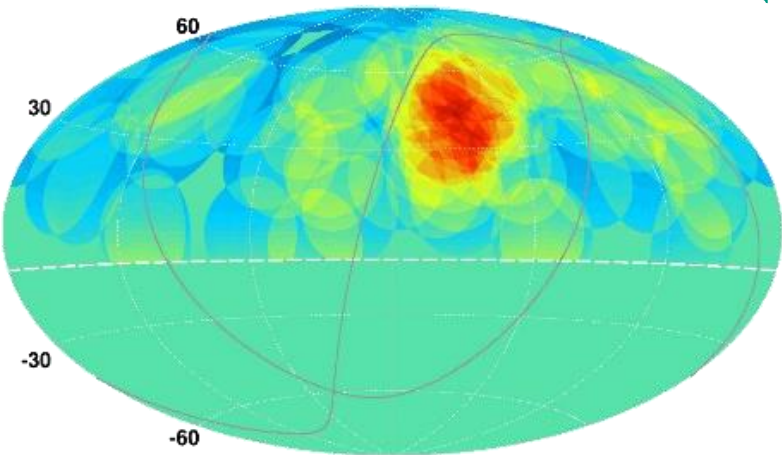
Energy spectrum of high-energy CRs & features

■ Energy spectrum of **extra-galactic CRs**:

- change of acceleration sites

radio galaxies?

$$E^{2.5} \cdot J(E) \quad (m^{-2} s^{-1} sr^{-1} eV^{1.5})$$

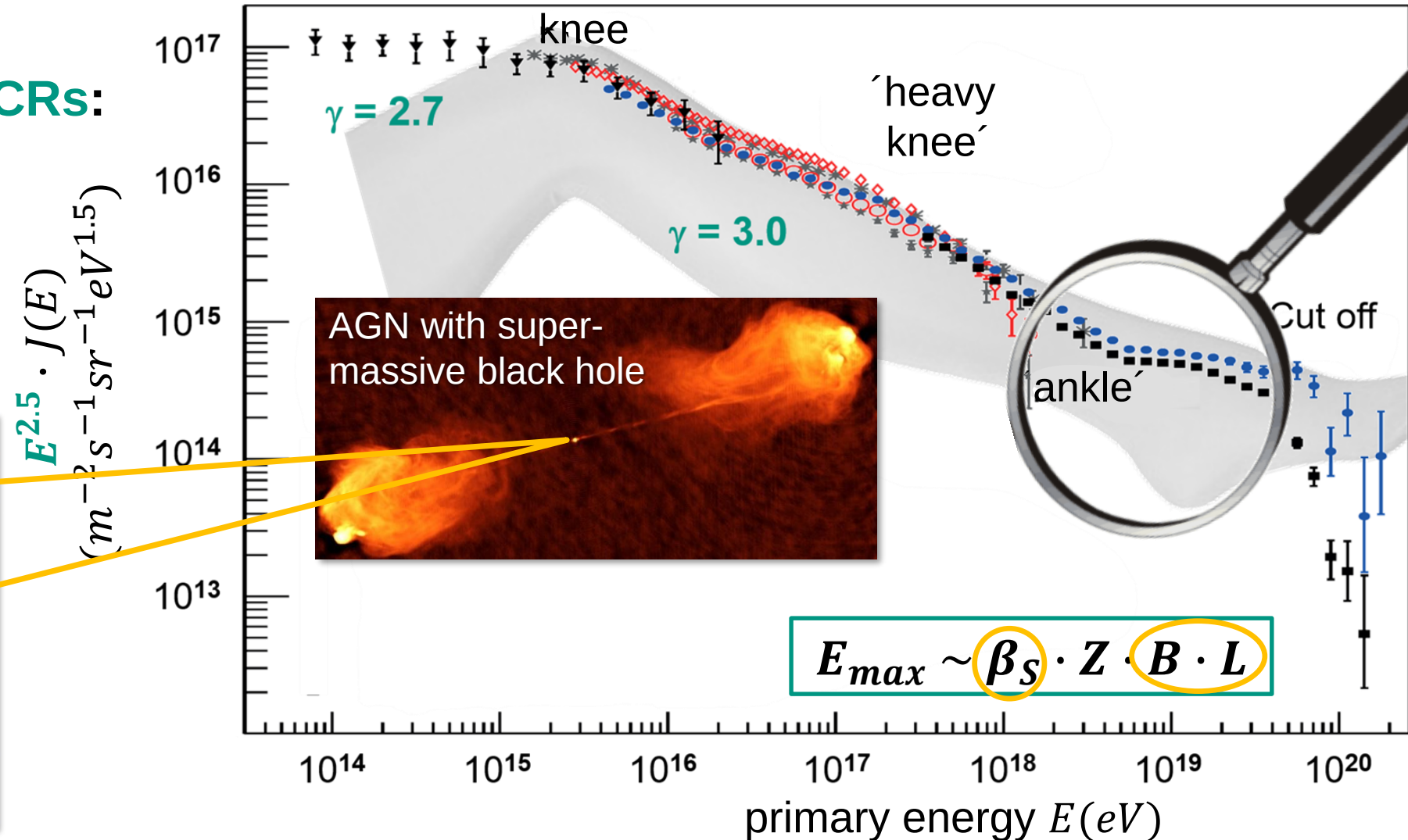
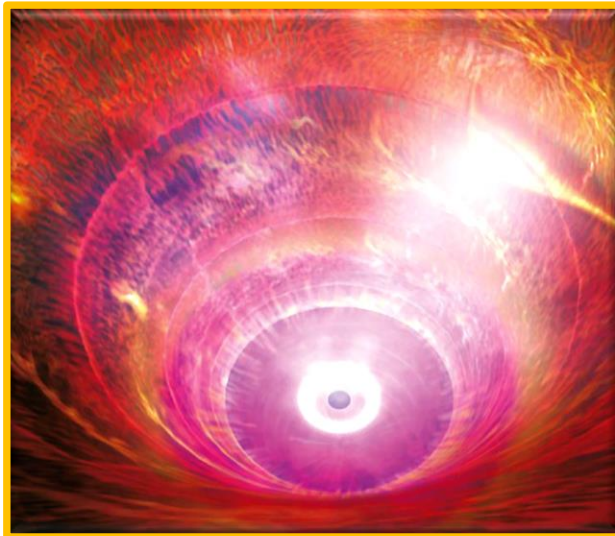


Energy spectrum of high-energy CRs & features

■ Energy spectrum of **extra-galactic CRs**:

- change of acceleration sites

AGNs?



Energy spectrum of UHECRs & cut-off

■ Energy spectrum of **extra-galactic CRs**:

- what causes the cut-off at $\sim 10^{20} \text{ eV}$?

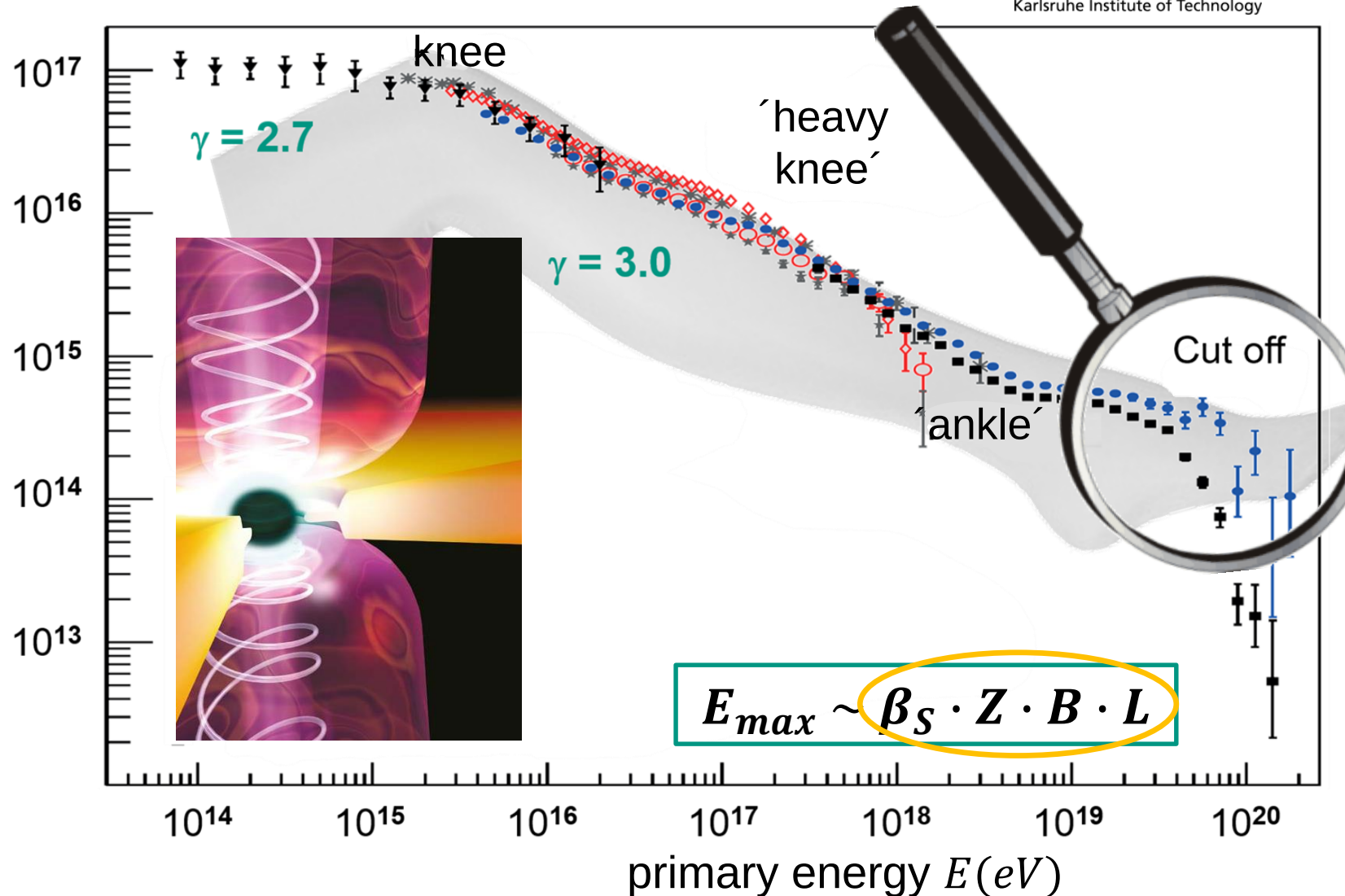
a) the accelerators 'run out of steam'

$$E^{2.5} \cdot J(E) \quad (m^{-2} s^{-1} sr^{-1} eV^{1.5})$$



„there simply is an upper end...“

A. M. Hillas



Energy spectrum of UHECRs & cut-off

■ Energy spectrum of **extra-galactic CRs**:

- what causes the cut-off at $\sim 10^{20} \text{ eV}$?

b) cosmic background fields (CMB) limit range

$$E^{2.5} \cdot J(E) \quad (m^{-2} s^{-1} sr^{-1} eV^{1.5})$$



Greisen, Zatsepin & Kuzmin (GZK)

