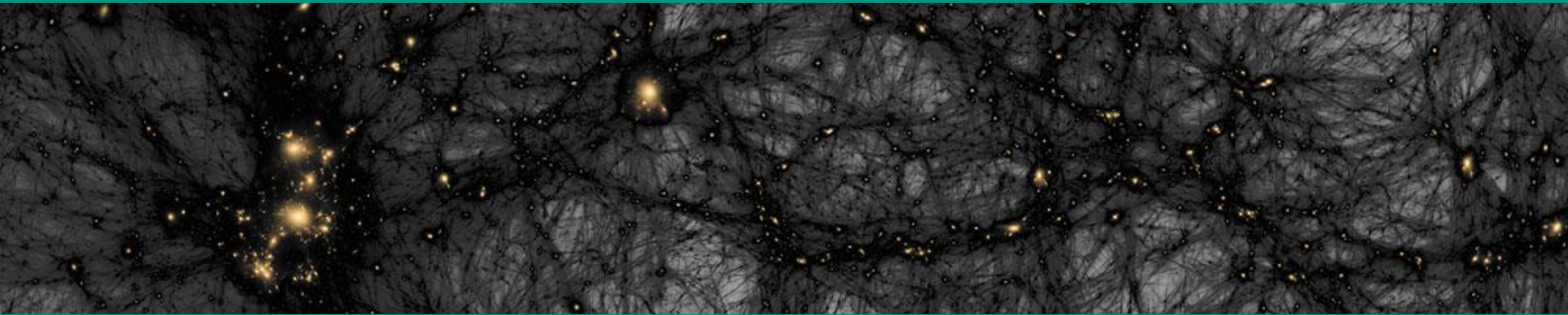


Astroparticle physics *I* – Dark Matter

Winter term 23/24

Lecture 4

Nov. 8, 2023



Recap of Lecture 3

- multi-messenger observations: all-sky maps in *CRs*, *gammas*, *ν 's* & *GW*
 - **Mollweide projection** with *GC* at the centre
 - sources of messengers: *SNR*, *GRB*, *AGN*, pulsars, binary compact objects,...
- studies of charged cosmic rays (*CR*)
 - *V. Hess* (**balloon**), *P. Auger* (**shower**), *J. Linsley* (**arrays**, up to 10^{20} eV)
 - energy spectrum of charged *CRs*: characteristic **power law**
 - direct *CR* detection via balloons & satellites (up to $\sim 10^{14} \text{ eV}$):
element composition allows to determine **galactic storage time** of *CRs*
 - indirect *CR* detection (above $\sim 10^{13} \text{ eV}$) via air shower arrays: **3 components**

Air showers & jet physics at accelerators

■ synergy

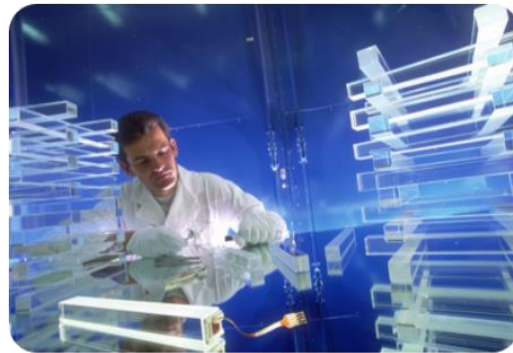
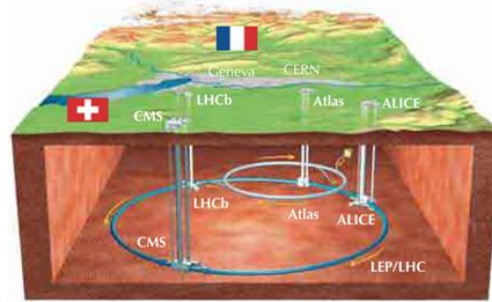
- in the **calorimetry** of hadrons & electrons

ATP & TP!

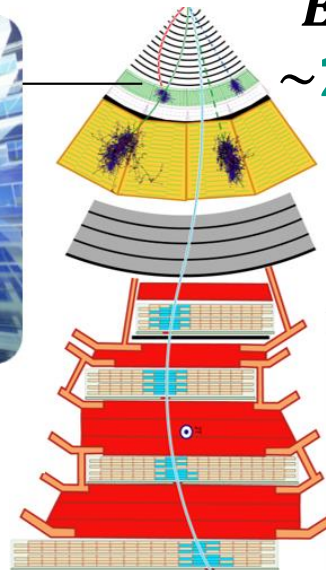


shower physics at *LHC*

accelerators:
protons with up
to $\sqrt{s} = 7 \text{ TeV}$
($\sqrt{s}$ is known)



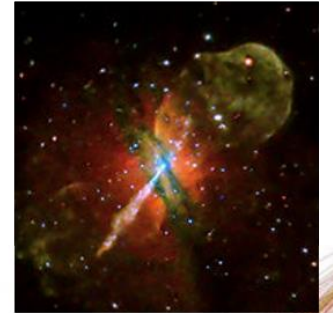
61.200 PbWO_4
 $X_0 = 0.89 \text{ cm}$



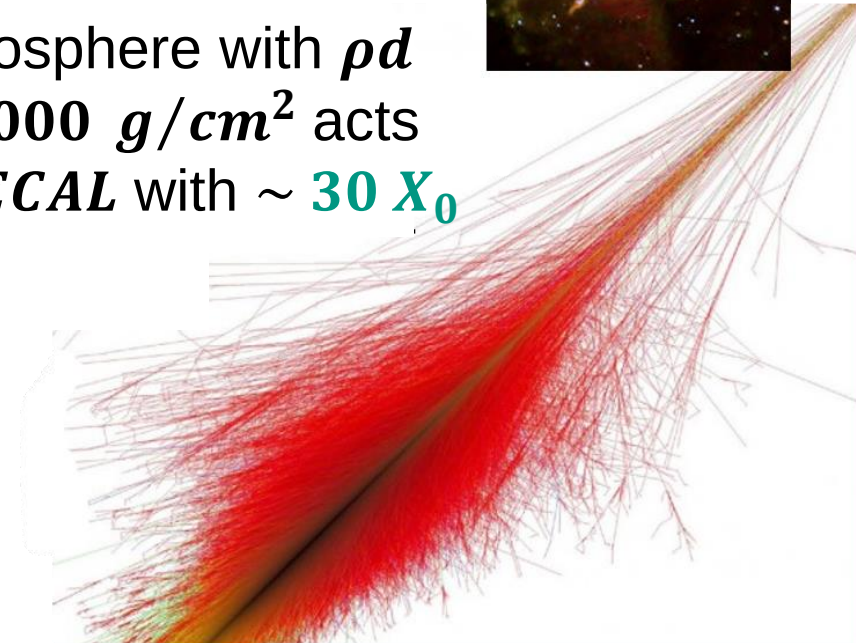
ECAL
 $\sim 20 X_0$

showers from *AGN* accelerators

cosmic accelerators:
protons (nuclei) with
up to $E = 10^{20} \text{ eV}$



atmosphere with ρd
 $= 1000 \text{ g/cm}^2$ acts
as *ECAL* with $\sim 30 X_0$

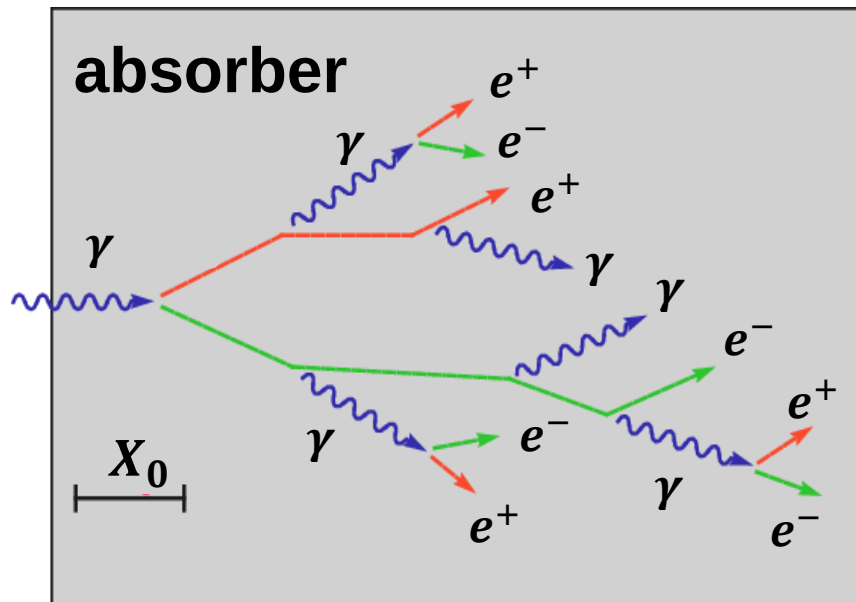


Electromagnetic component of showers

■ majority ($\sim 98\%$) of shower particles: from cascade $e^- e^+ \gamma$

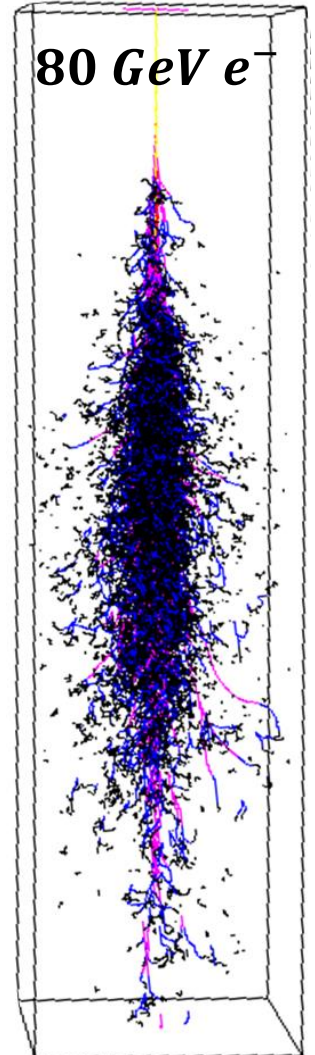
- cascade processes: **pair production** $\Leftrightarrow$ **bremsstrahlung**

- *Heitler* model for electromagnetic cascades:



Q: aanda.org

- initial increase of # of particles (successive particle ‘generations’)
- mean energy per particle decreases
- huge number of secondary, low-energy particles: $e^- e^+ \gamma$
- shower ‘dies out’ once $E < E_c$ with: $E_c = \text{critical energy}$



Electromagnetic component of showers

■ majority ($\sim 98\%$) of shower particles: from cascade $e^- e^+ \gamma$

- two processes: **pair production** $\Leftrightarrow$ **bremsstrahlung**

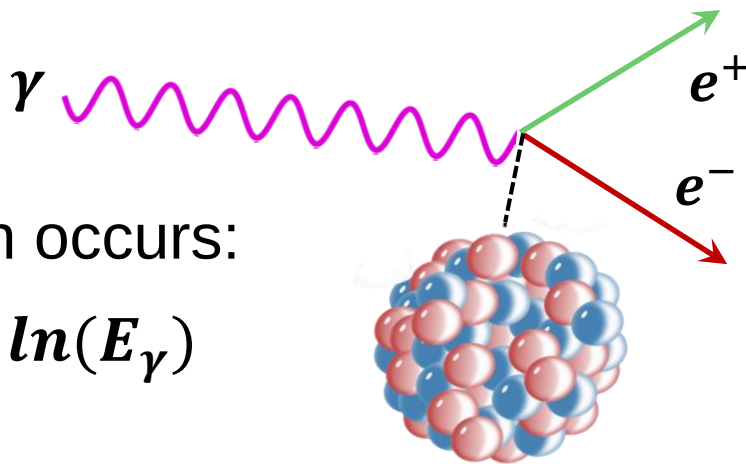
- **threshold** energy:

$$E_{\gamma, \text{thres}} \cong 2 \cdot m_e = 1.02 \text{ MeV}$$

- saturation occurs:

$$\sigma_{\gamma}(E_{\gamma}) \sim \ln(E_{\gamma})$$

kinematics of nucleus: $E_{\text{rec}}, \vec{p}$

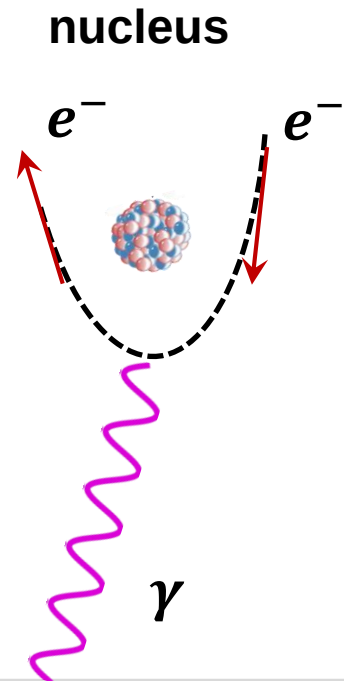


- radiative energy losses at high energies:

$$\left(\frac{dE}{dX} \right)_{\text{brems}} = \frac{1}{X_0} \cdot E$$

- important only for very **light particles** (e^-, e^+):

$$\sigma_{\text{Brems}} \sim 1/m^2$$



Electromagnetic component of showers

■ development of electromagnetic showers in the atmosphere

- key parameter: **radiation length X_0**

$$E(x) = E_0 \cdot e^{-(x/X_0)}$$

- for highly **relativistic electrons**:

i.e. after the passing of an atmospheric thickness $X = X_0$,
the **energy E** of an e^- has dropped from E_0 to a fraction of **$1/e$**

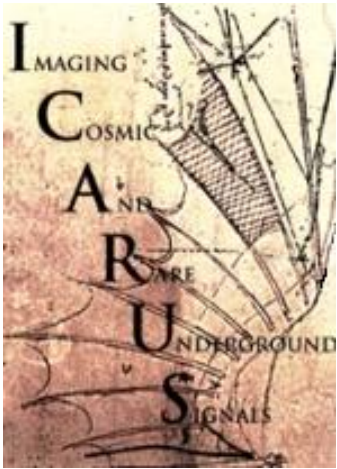
- for **high–energy photons**:

i.e. the **mean free path λ_{pair}** of high–energy photons
for **pair production** in the atmosphere is given by

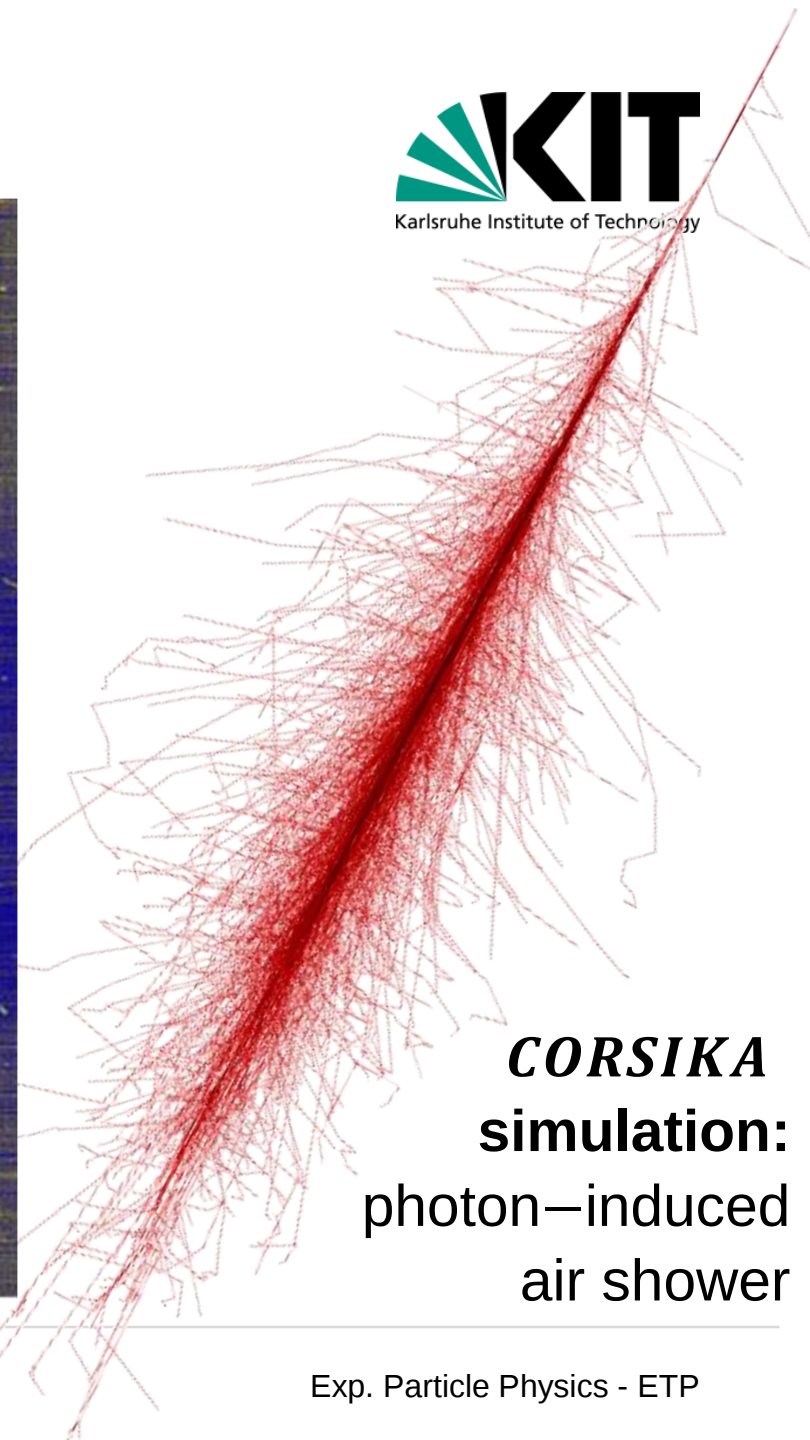
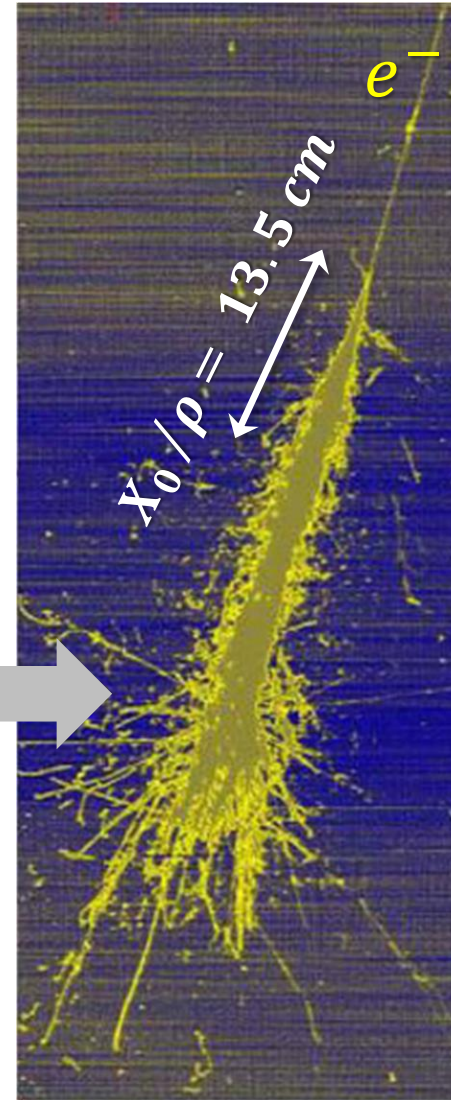
$$\lambda_{pair} = \frac{9}{7} \cdot X_0$$

Electromagnetic component of showers

- electromagnetic showers:
experimental signature & modelling
experimental shower signatures &
modelling via simulation–codes*



ICARUS experiment
neutrino–induced
particle shower in a
liquid argon calorimeter:
interactions of a primary
high–energy electron

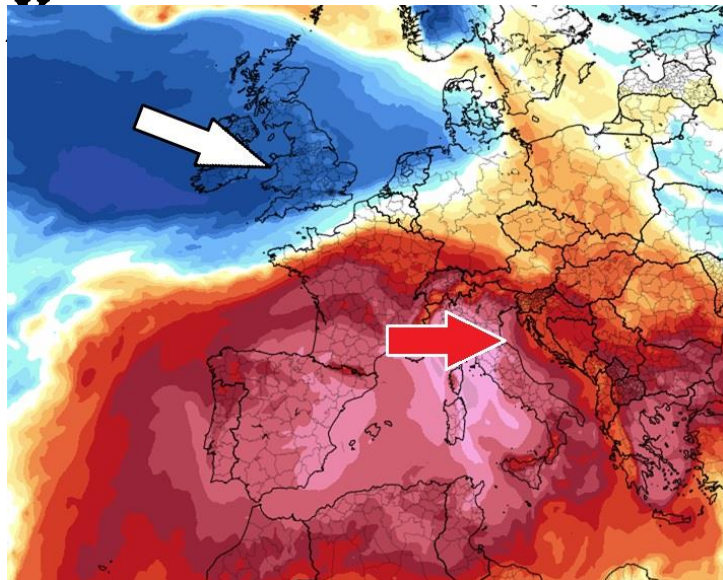


CORSIKA
simulation:
photon–induced
air shower

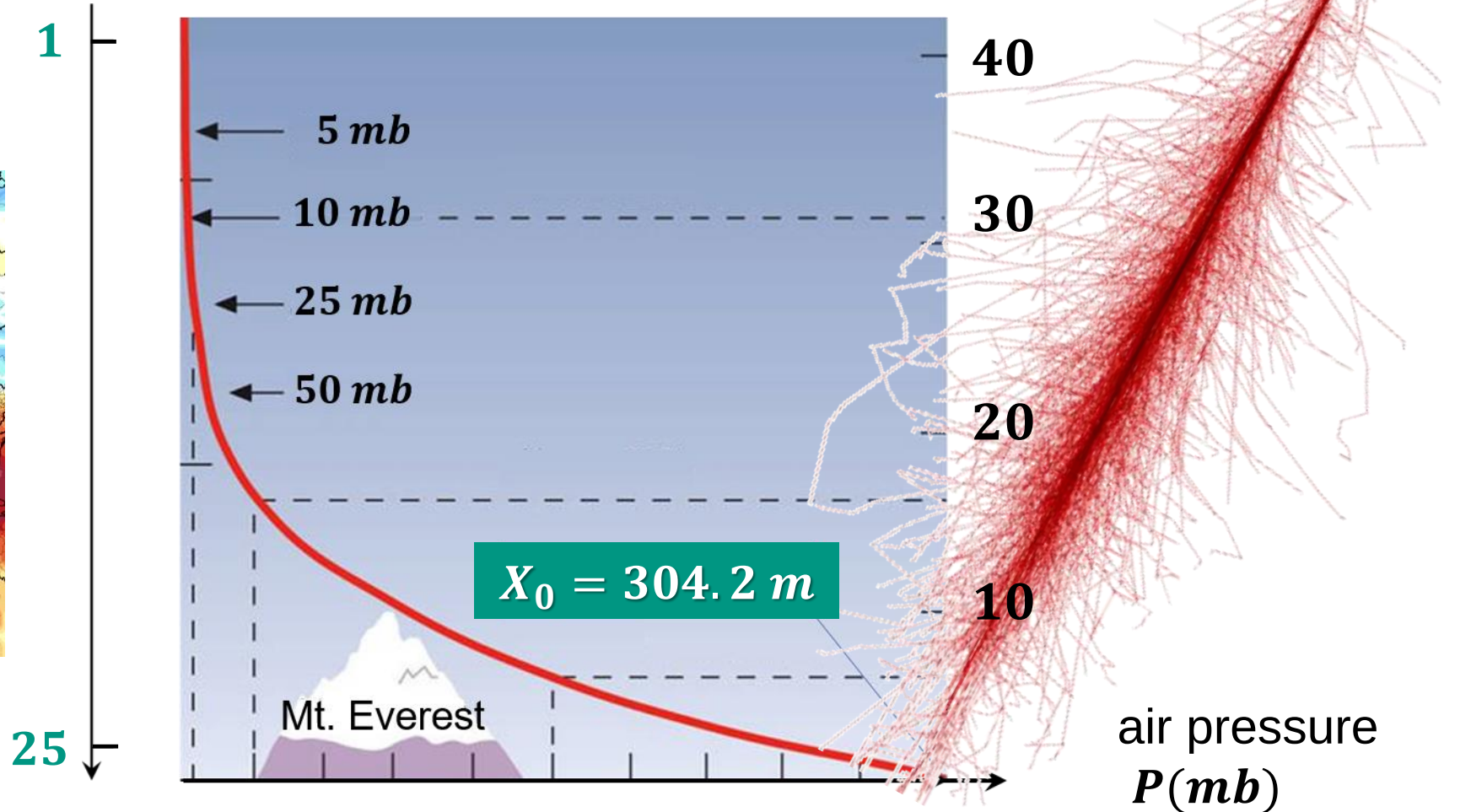
Electromagnetic component of showers

■ atmosphere as electromagnetic calorimeter: $\sim 25 X_0$

- changing atmospheric weather: changes of

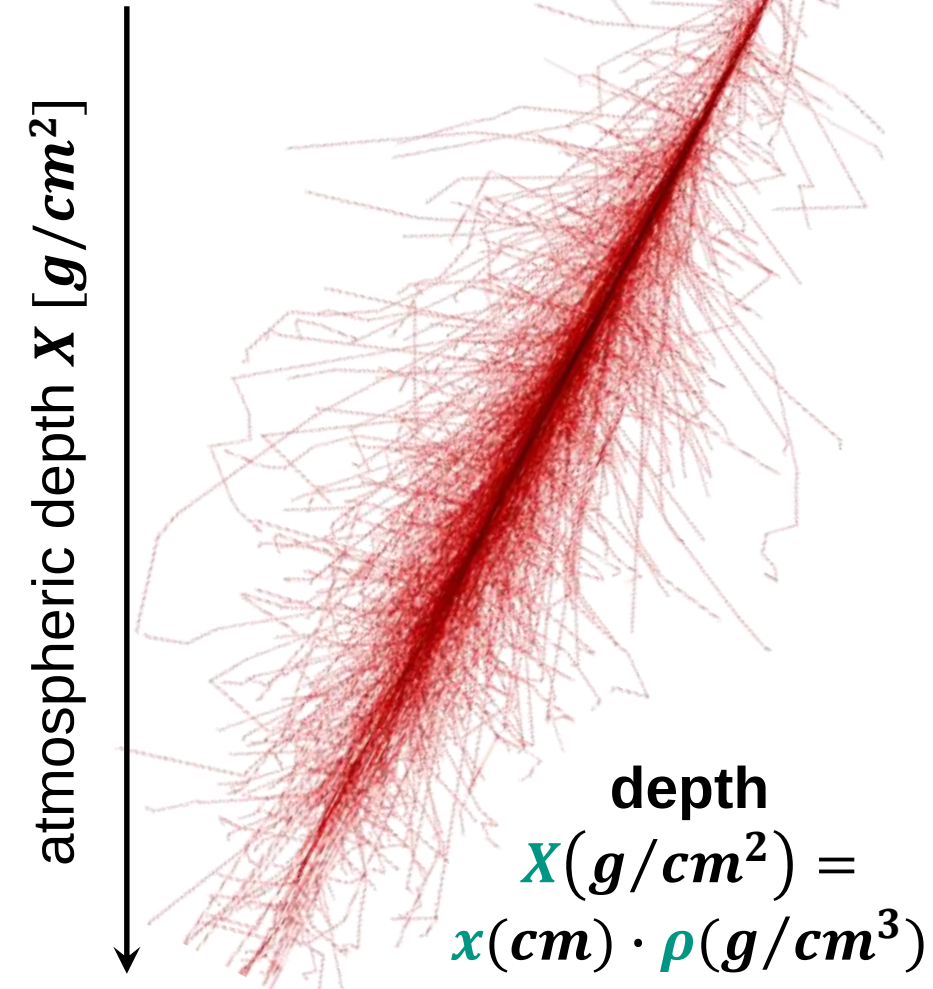
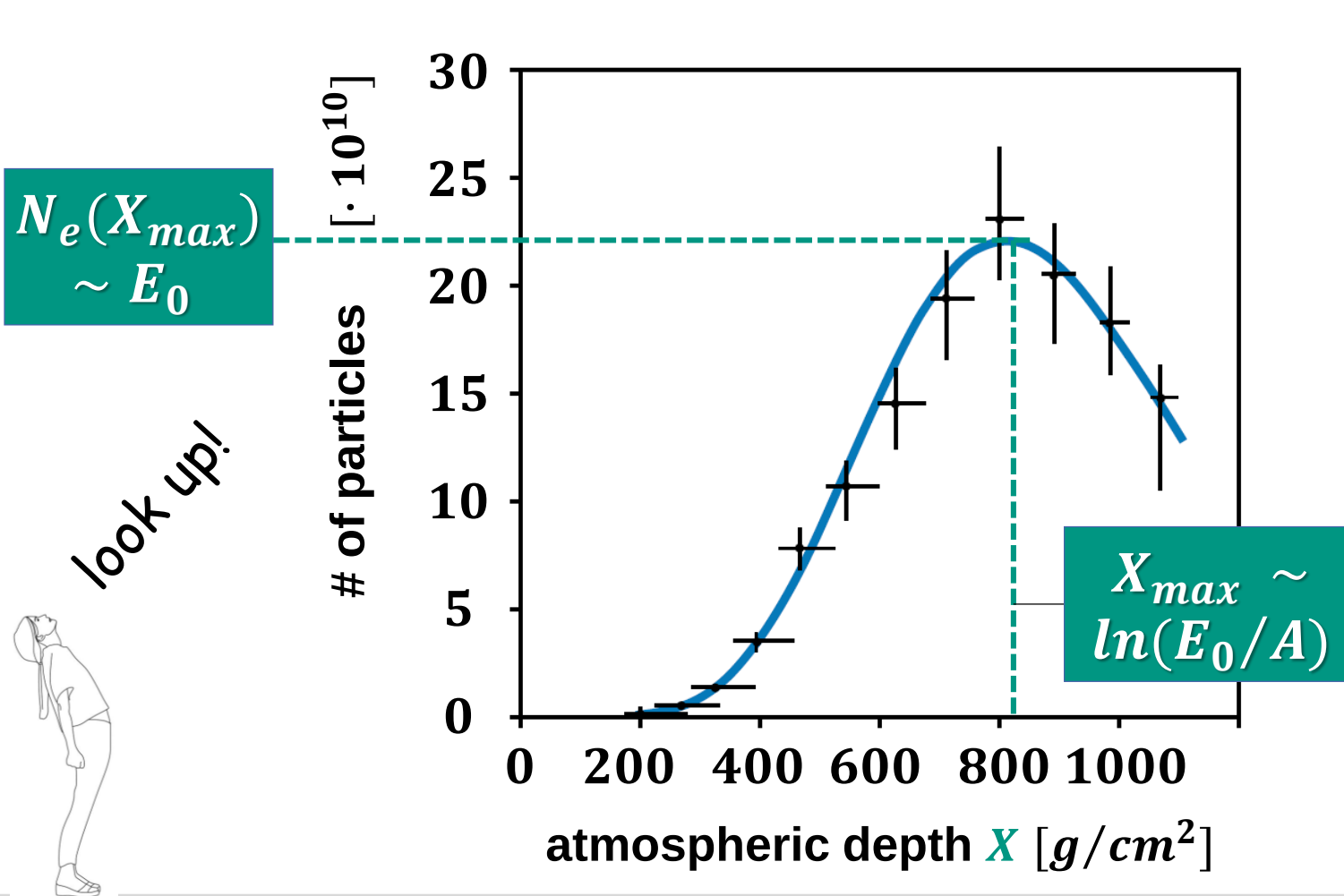


atmosphere (# of X_0)



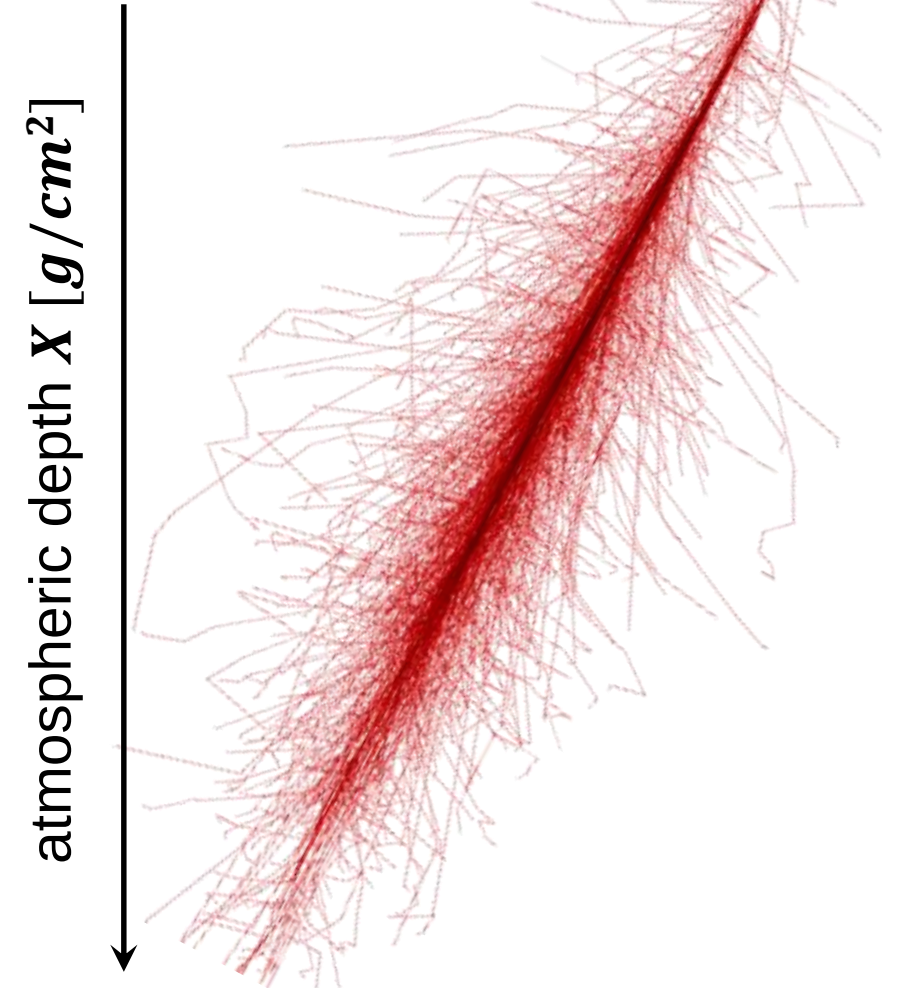
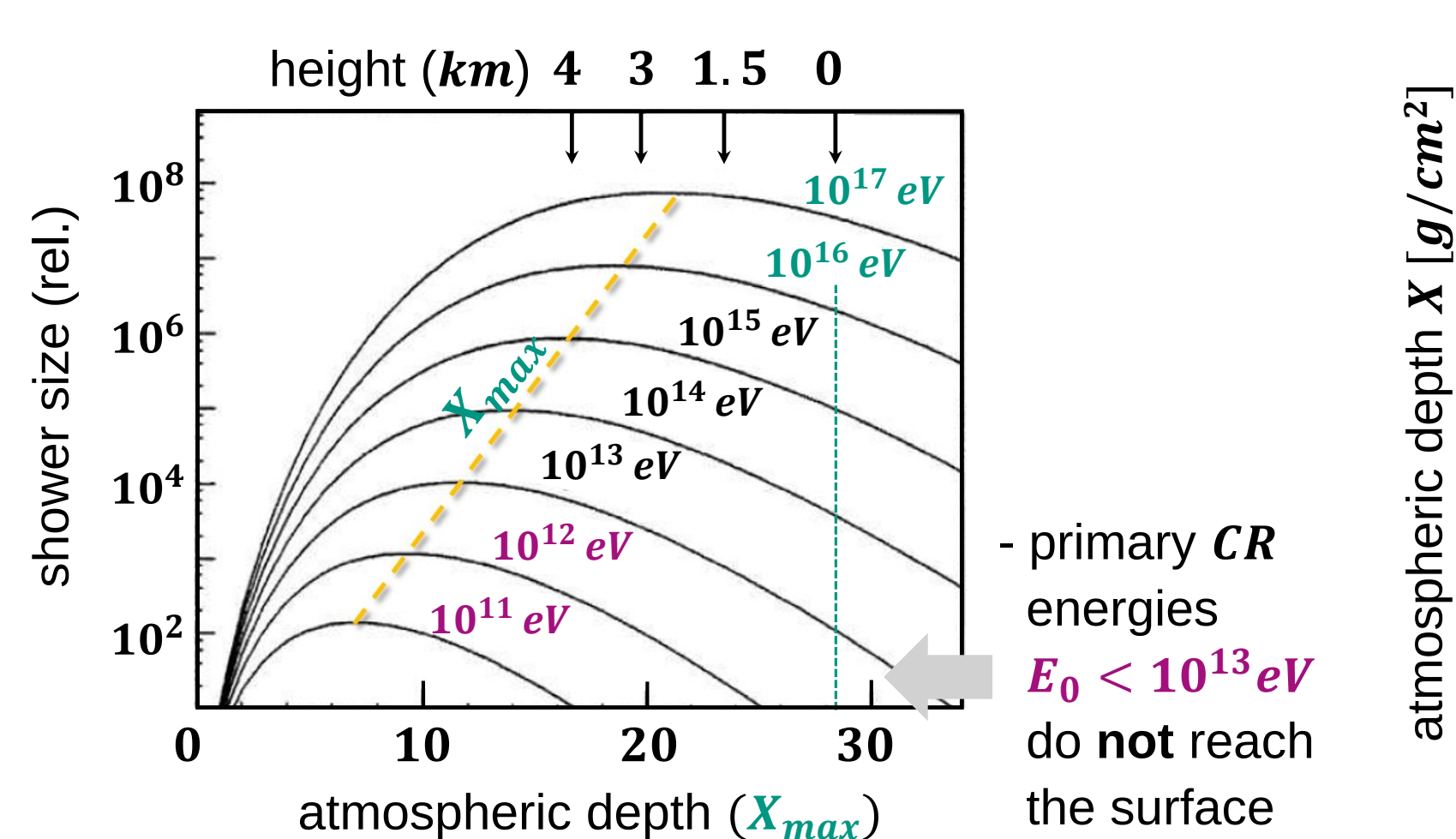
Electromagnetic showers in the atmosphere

■ Key parameter: **shower maximum** X_{max} for mass A of CRs



Electromagnetic showers in the atmosphere

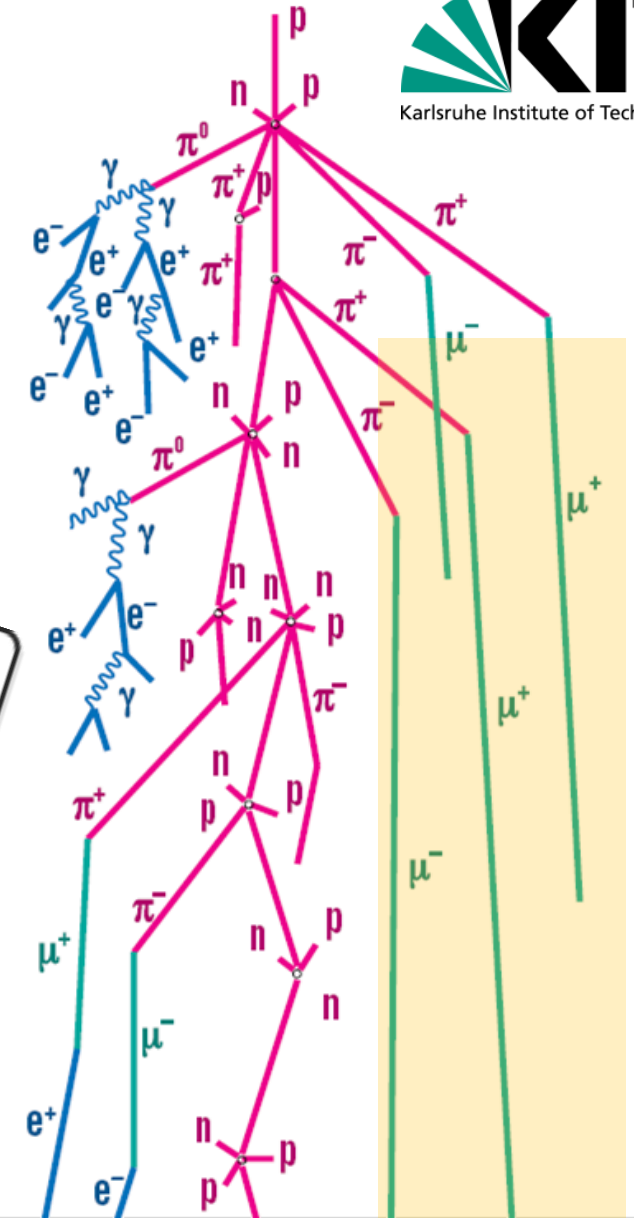
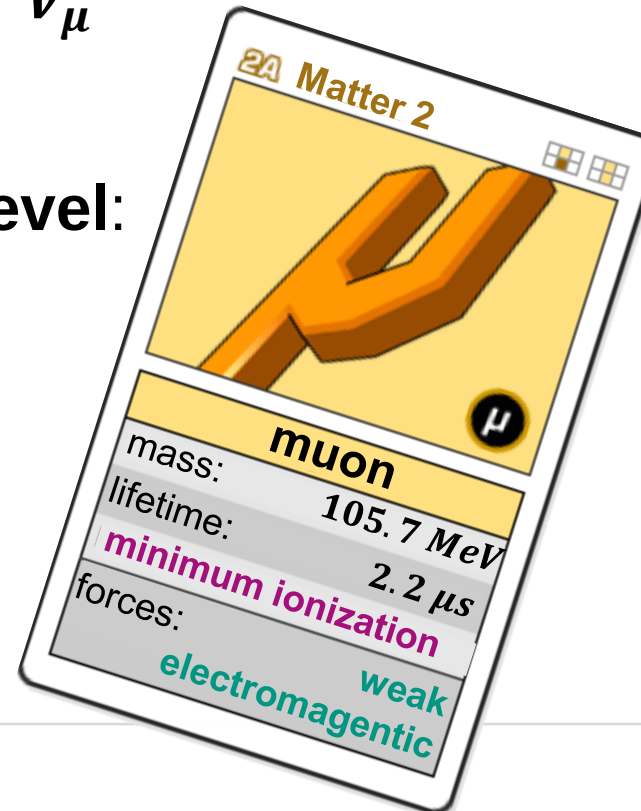
- shower maximum X_{max} varies as function of primary energy E_0



Muonic component: *mips*

■ muons are *m*inimum *i*onizing *p*articles (*m.i.p.s*)

- ‘cosmic’ muons (slang: ‘**cosmics**’)
- from **pion decays** $\pi^\pm \rightarrow \mu^\pm + (\bar{\nu}_\mu)$ in upper atmosphere
- flux density of muons at **sea level**:
 $d\Phi/dA \sim 100 \mu's m^{-2}s^{-1}$
- charge ratio: $\mu^+/\mu^- = 1.2$
- overall fraction of secondary particles $\varepsilon_\mu \sim 1.7\%$



application of air showers: **muography**

- muography: muon rate depends on air pressure

physicsworld

Cosmic muons probe the interiors of tropical

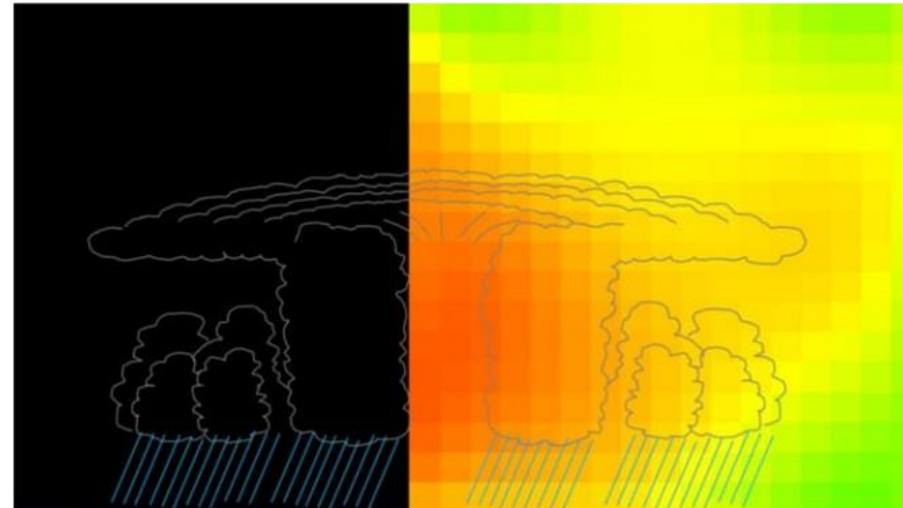
cyclones

18 Oct 2022

PARTICLE AND NUCLEAR

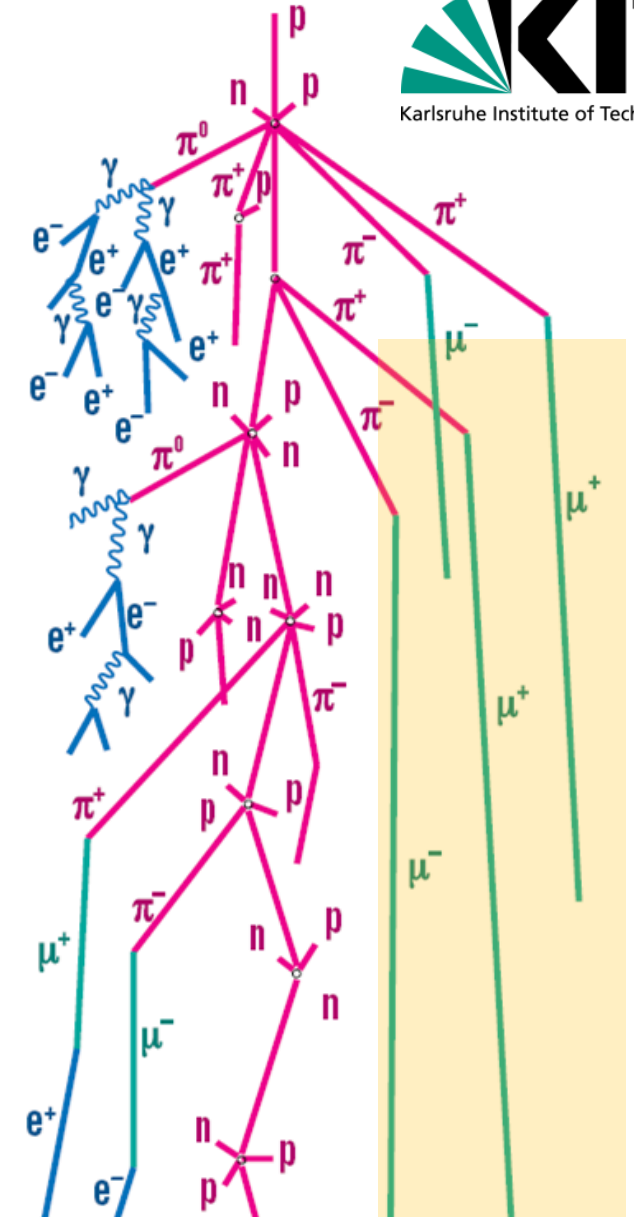
RESEARCH UPDATE

- advantage of using muons from showers: **long range of μ** in the atmosphere



Peering inside: the image on the right shows the interior of a cyclone. The redder areas are regions of lower pressure, and the greener areas are higher pressure. (Courtesy: Hiroyuki KM Tanaka)

Cosmic muons have been used to image structures deep within tropical storms, according to an international team of researchers. Led by [Hiroyuki Tanaka](#) at the University of Tokyo, the team used a network of muon detectors to identify differences in air density within several typhoons.



Muonic component: underground location!

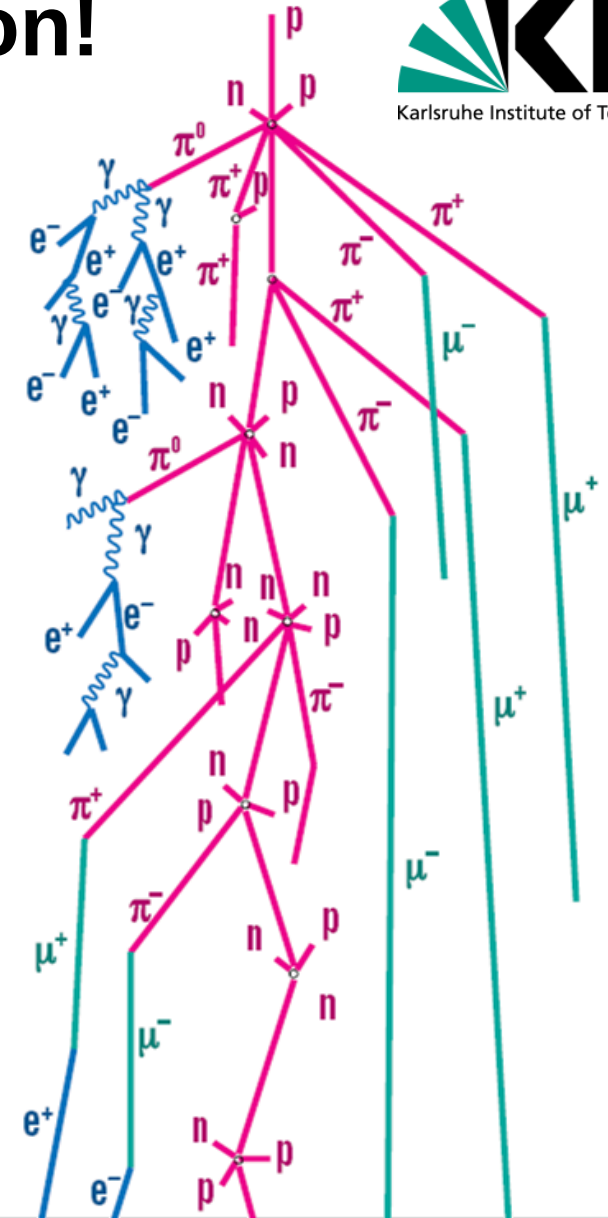
■ muons are penetrating particles

- relativistic lifetime: $\tau = \gamma \cdot 2.2 \mu s$
- energy loss: $dE/dx = 2 \text{ MeV/cm}$ (H_2O)
- very penetrating shower component, despite massive shielding $\Rightarrow$ underground laboratory

deep-underground lab: small μ –rate!



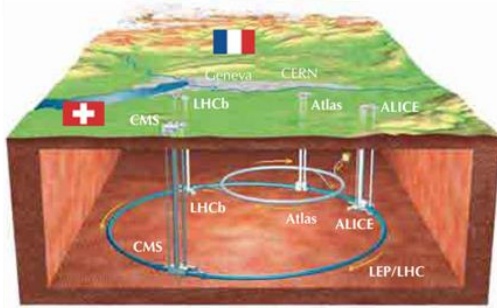
$\gamma \cdot 2.2 \mu s$



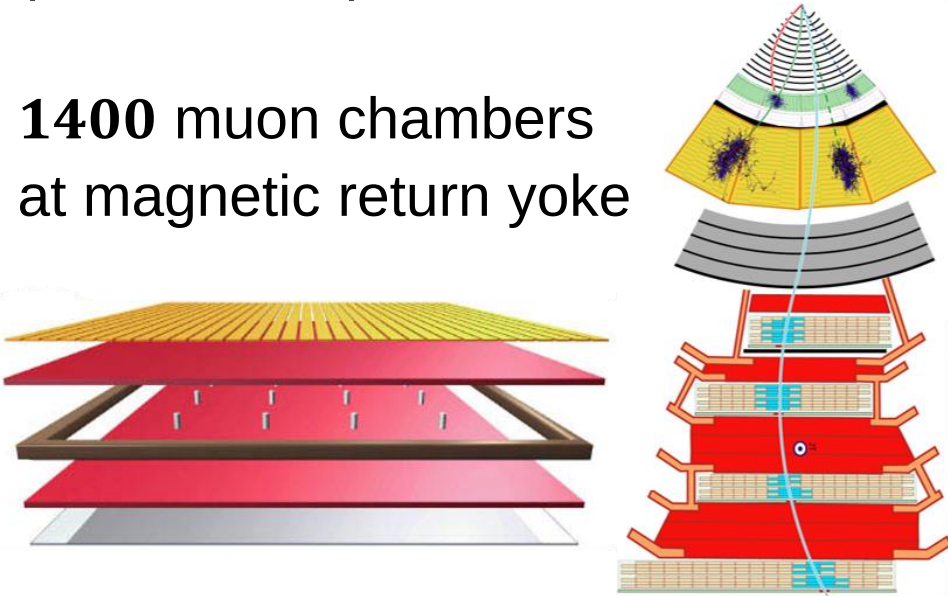
Muons at the *LHC* and in astroparticle physics

muons from an accelerator (LHC)

accelerators:
protons with up
to $\sqrt{s} = 7 \text{ TeV}$
($\sqrt{s}$ is known)



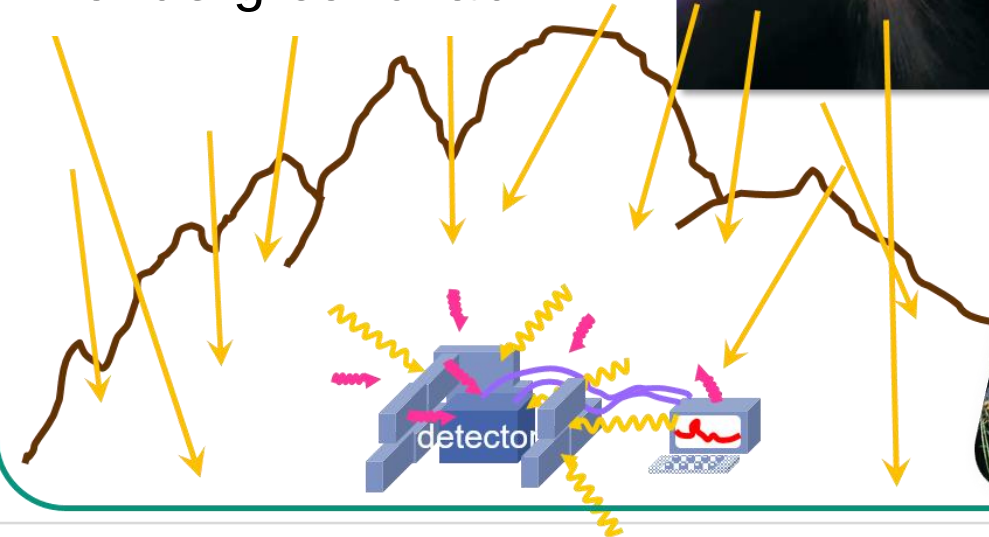
1400 muon chambers
at magnetic return yoke



muons from cosmic accelerators

cosmic accelerators:
protons (nuclei) with
up to $E = 10^{20} \text{ eV}$

shielding against muons
in underground lab



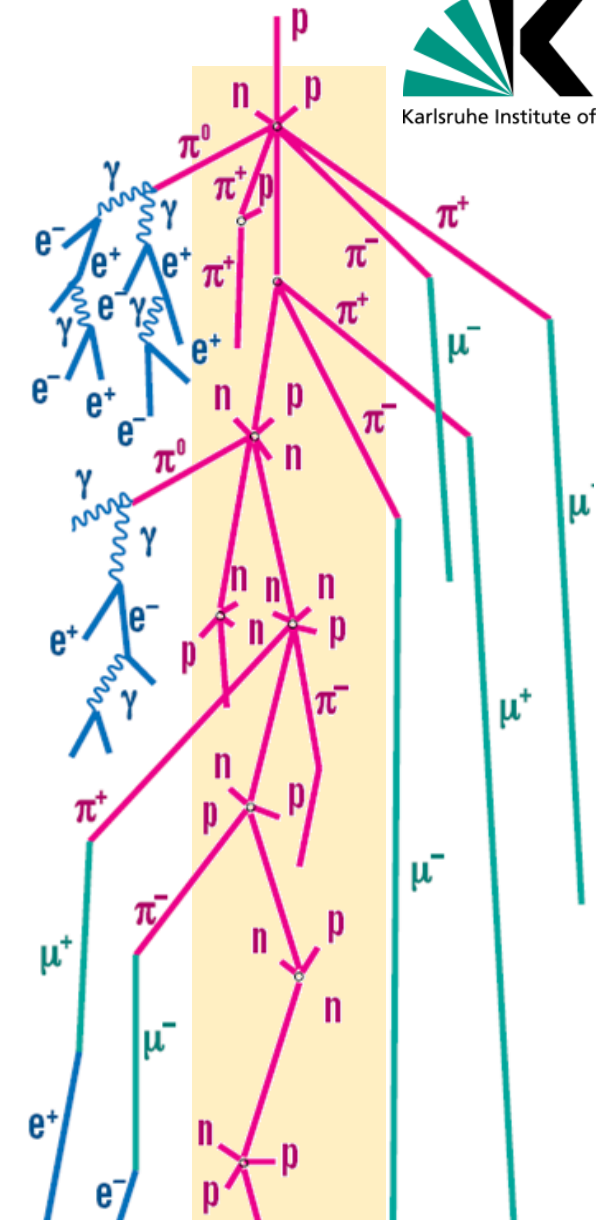
Laboratori
Nazionali
del
Gran Sasso



Hadronic component of air showers

■ Hadronic interactions in the atmosphere

- **all** hadrons ($p, n, \pi^+, \pi^-, \pi^0, K^+, K^-, K^0, \dots$): processes of the **strong** interaction, decays:
→ **hadronic** (& electromagnetic) **shower**
- **charged** hadrons ($p, \pi^+, \pi^-, K^+, K^-, \dots$): also
→ **electromagnetic interaction** (i.e. ionisation)
- hadrons are localized in **shower core** ($\sim 0.3\%$ fraction)
- hadrons are characterized by their **large dE/dx**



Hadronic component of air showers

■ Hadronic interactions in the atmosphere

- **hadronic interaction length Λ**
measured in $[g/cm^2 \text{ or } cm]$

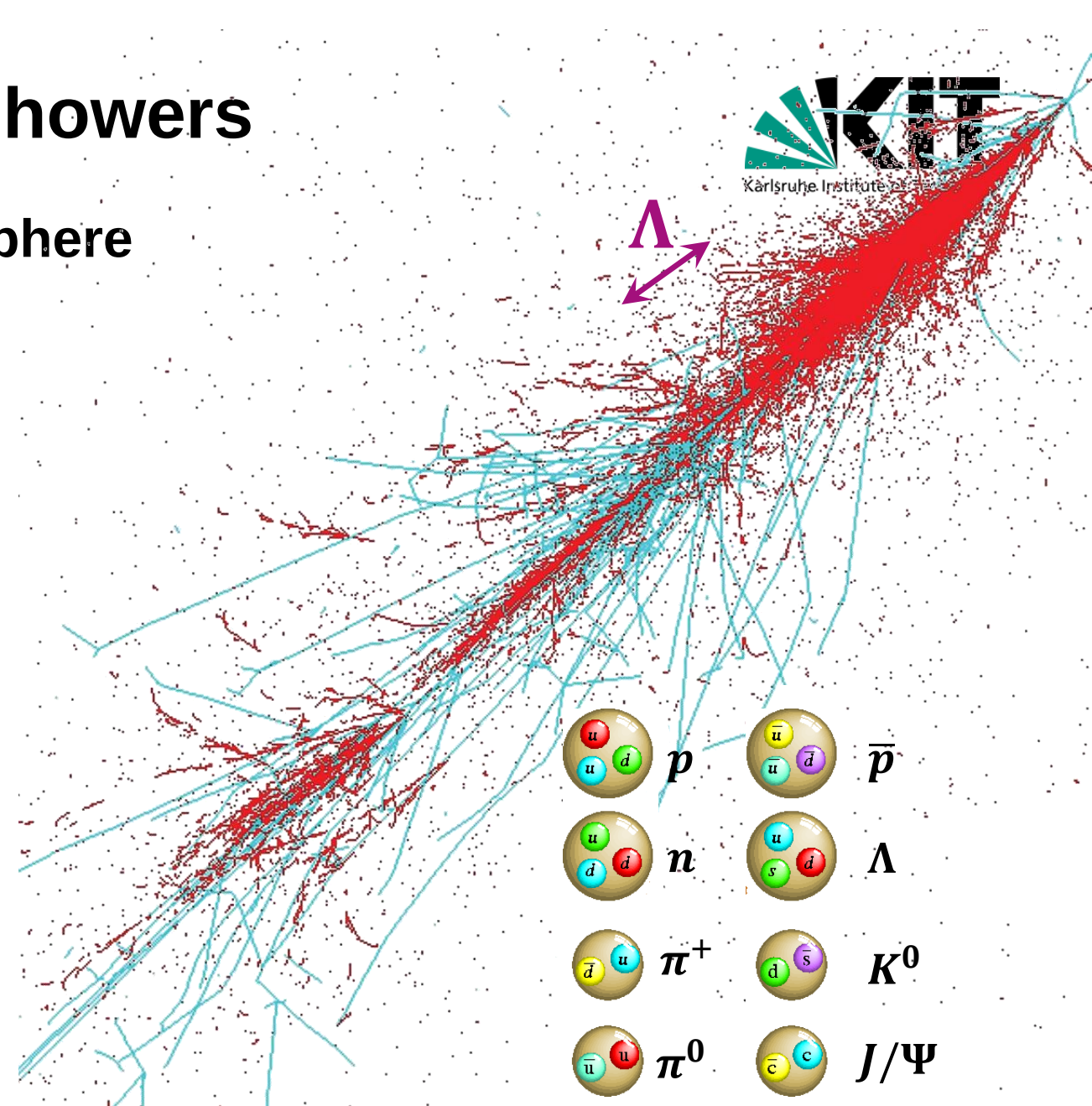
= mean length for inelastic scattering

$$\Lambda = \frac{1}{\sigma \cdot n}$$

σ : cross section with nucleus

n : number of scatter centres

- atmospheric thickness corresponds
~ **11 hadronic interaction lengths Λ**

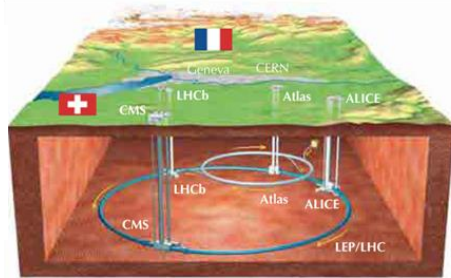


Hadronic component of air showers

■ Hadronic interactions: *HCAL* characteristics at the *LHC* & in the atmosphere

HCAL for terrestrial accelerators

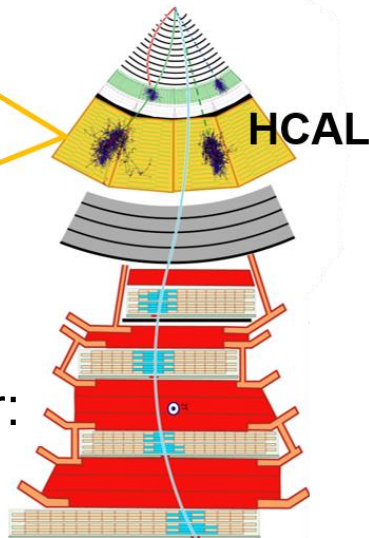
accelerators:
protons with up
to $\sqrt{s} = 7 \text{ TeV}$



iron (*Fe*) as absorber:

$$\Lambda = 16.8 \text{ cm}$$

$$X_0 = 1.76 \text{ cm}$$



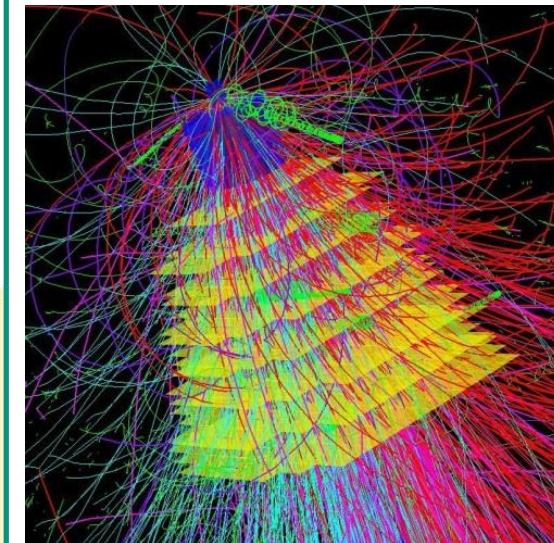
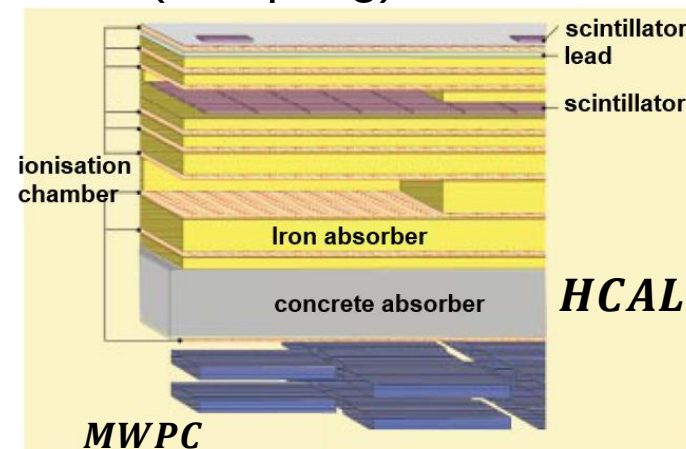
HCAL for cosmic accelerators

cosmic accelerators:
protons (nuclei) with
up to $E = 10^{20} \text{ eV}$



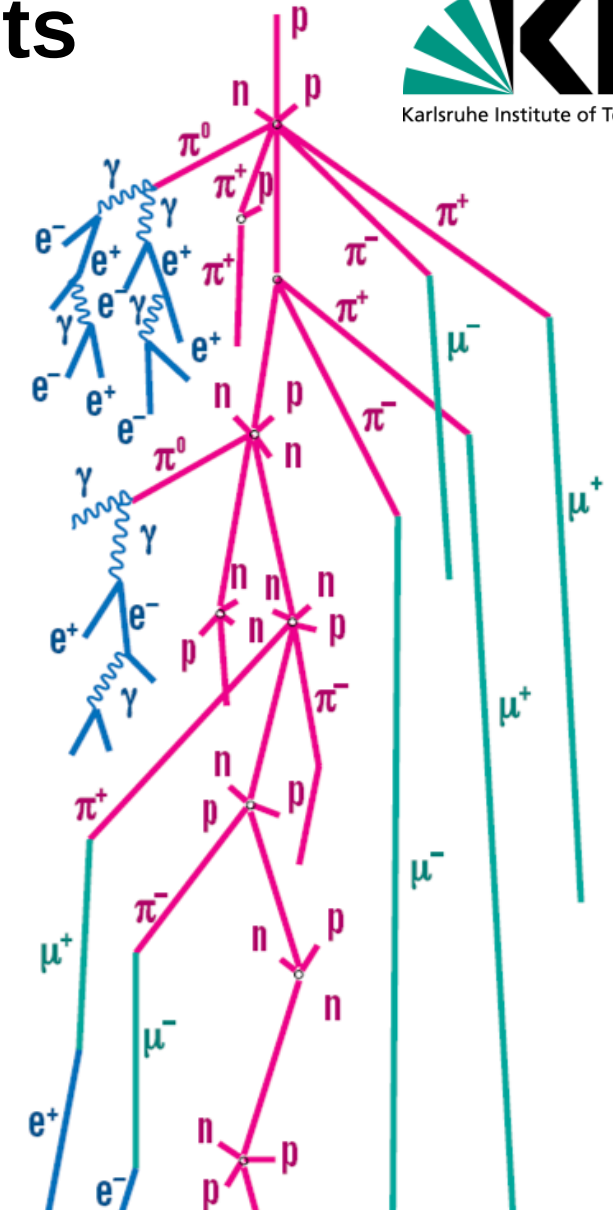
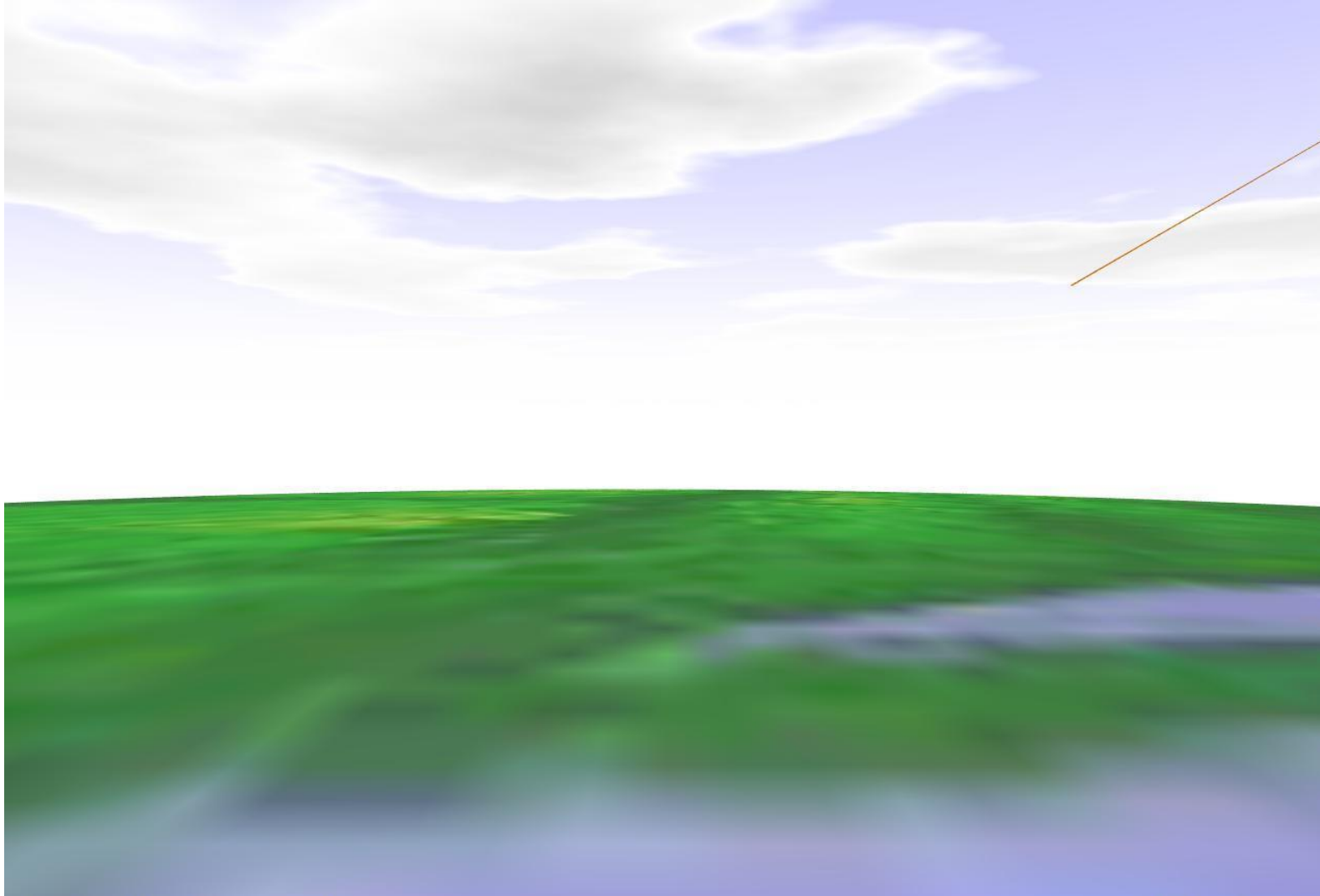
Fe – hadron– (sampling) calorimeter

K
A
S
C
A
D
E

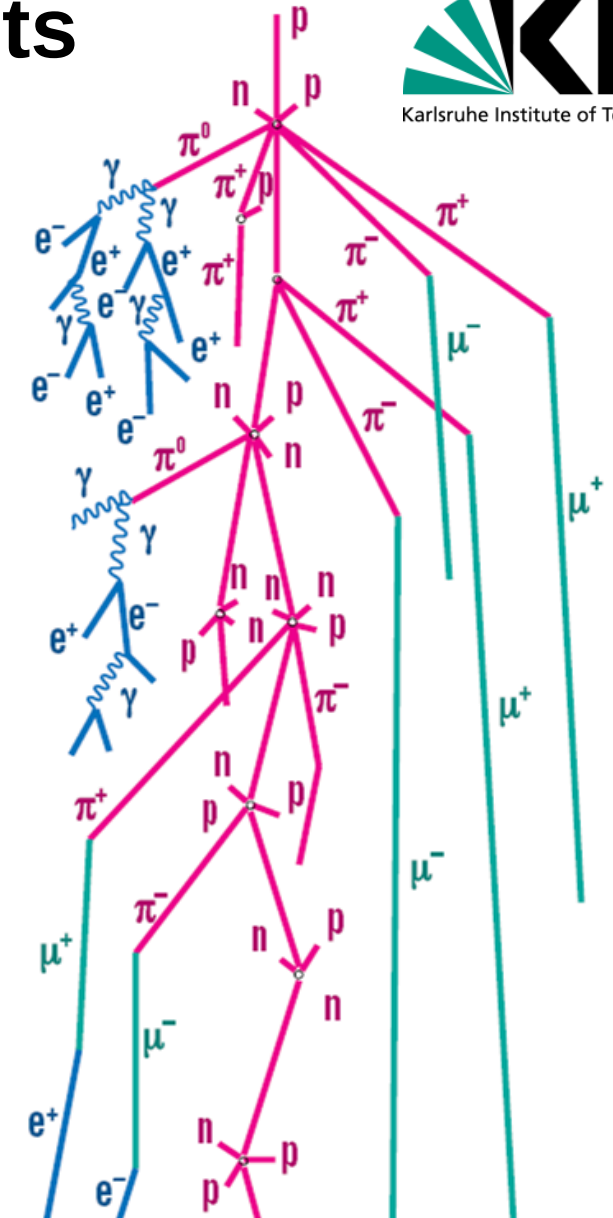
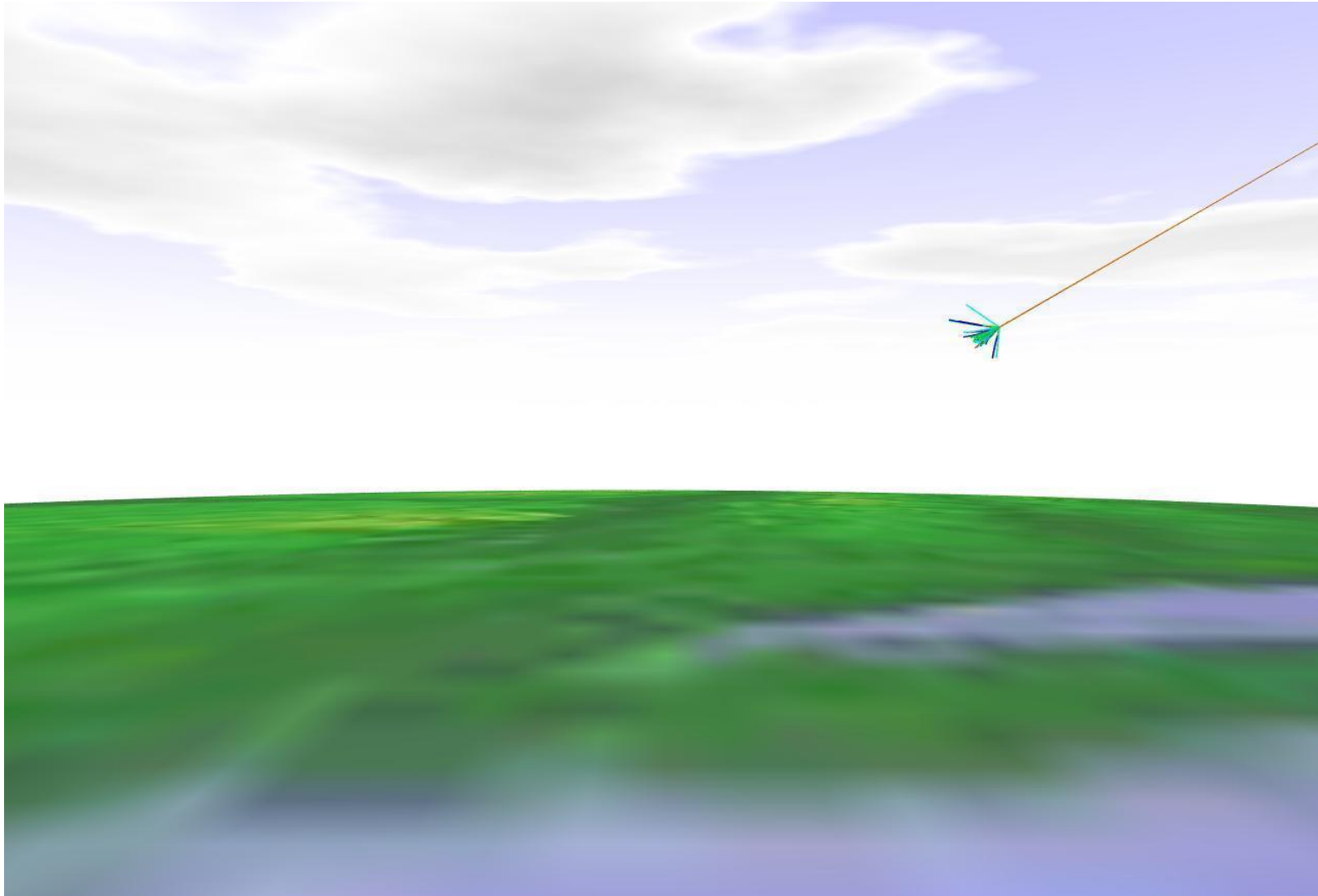


hadronic shower
(*MC*) in an *HCAL*

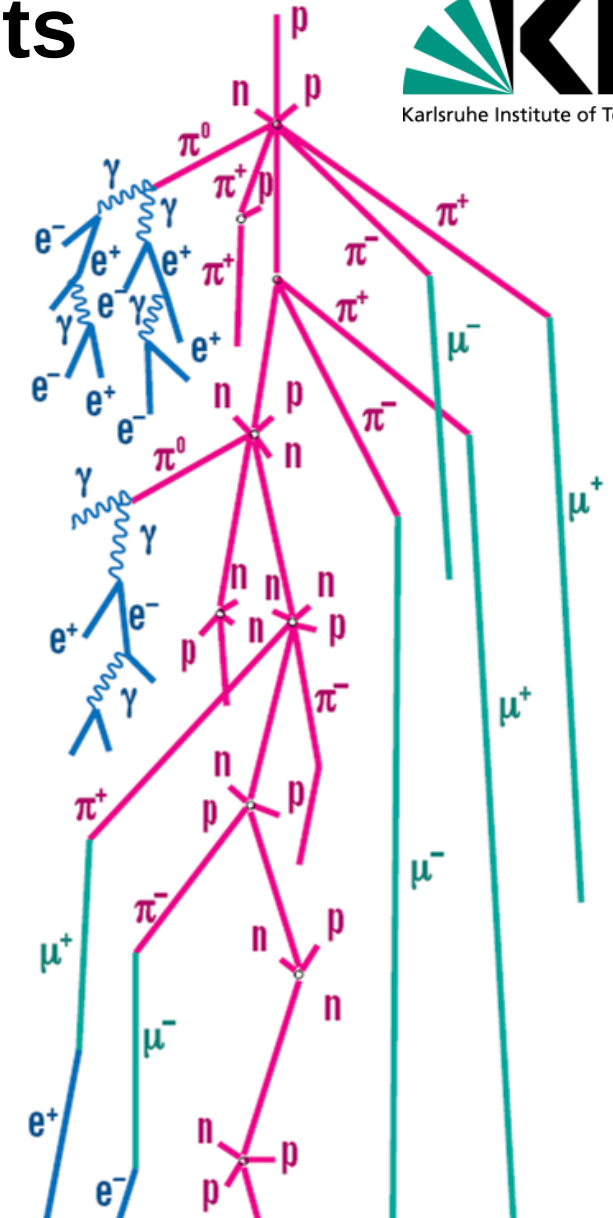
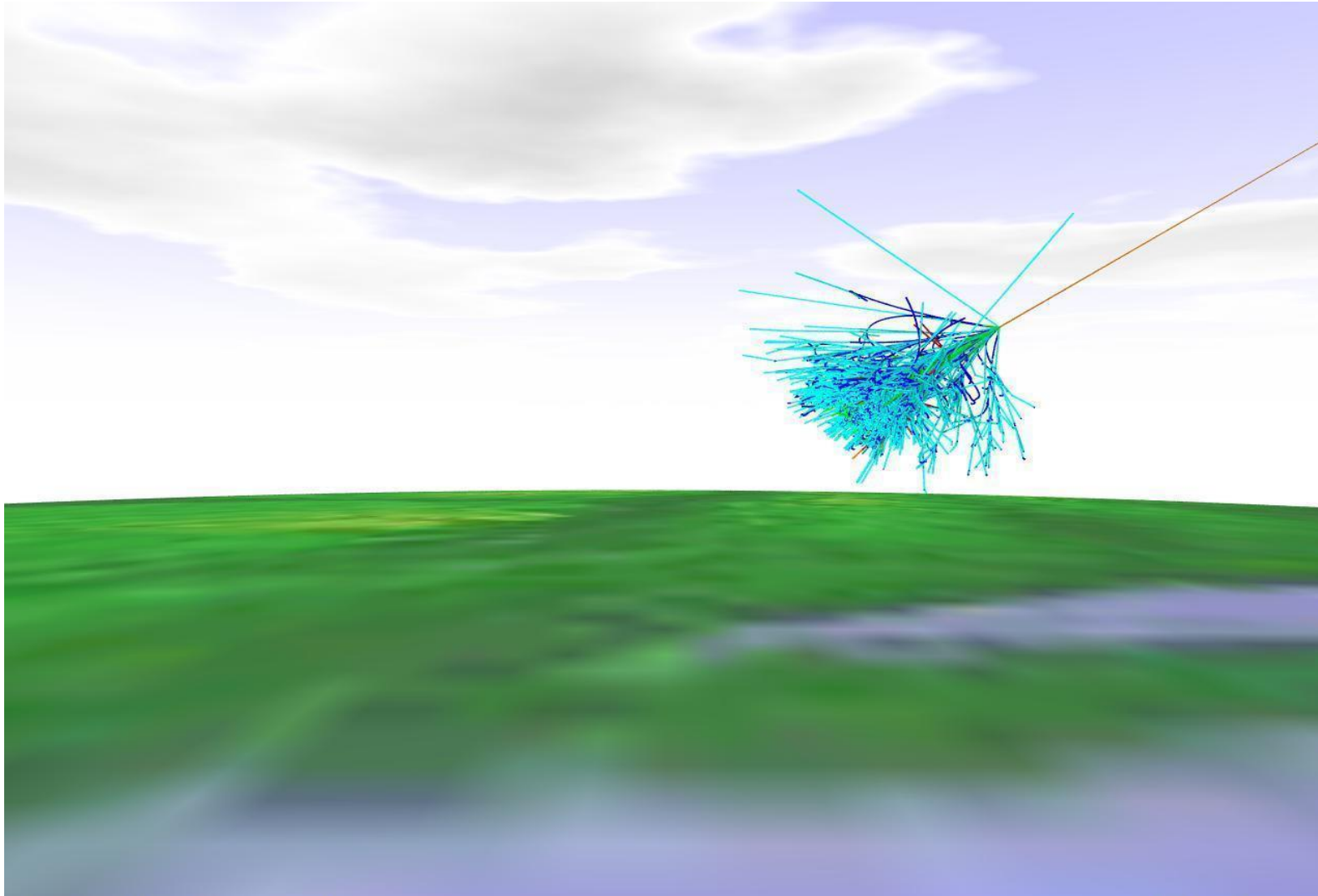
Simulation of air showers: all components



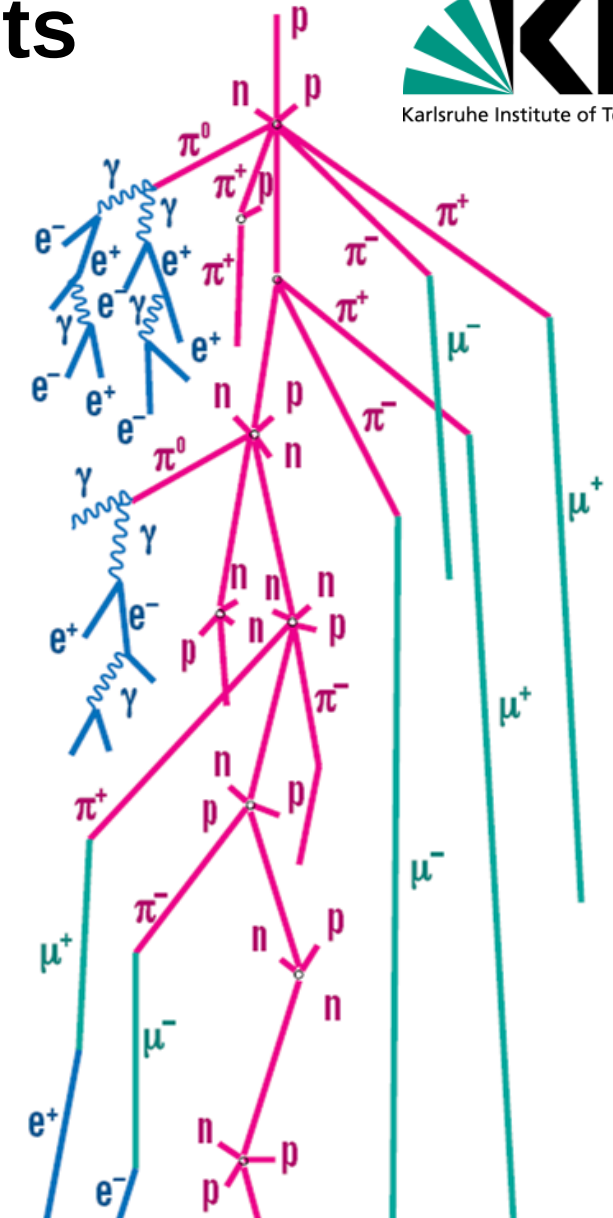
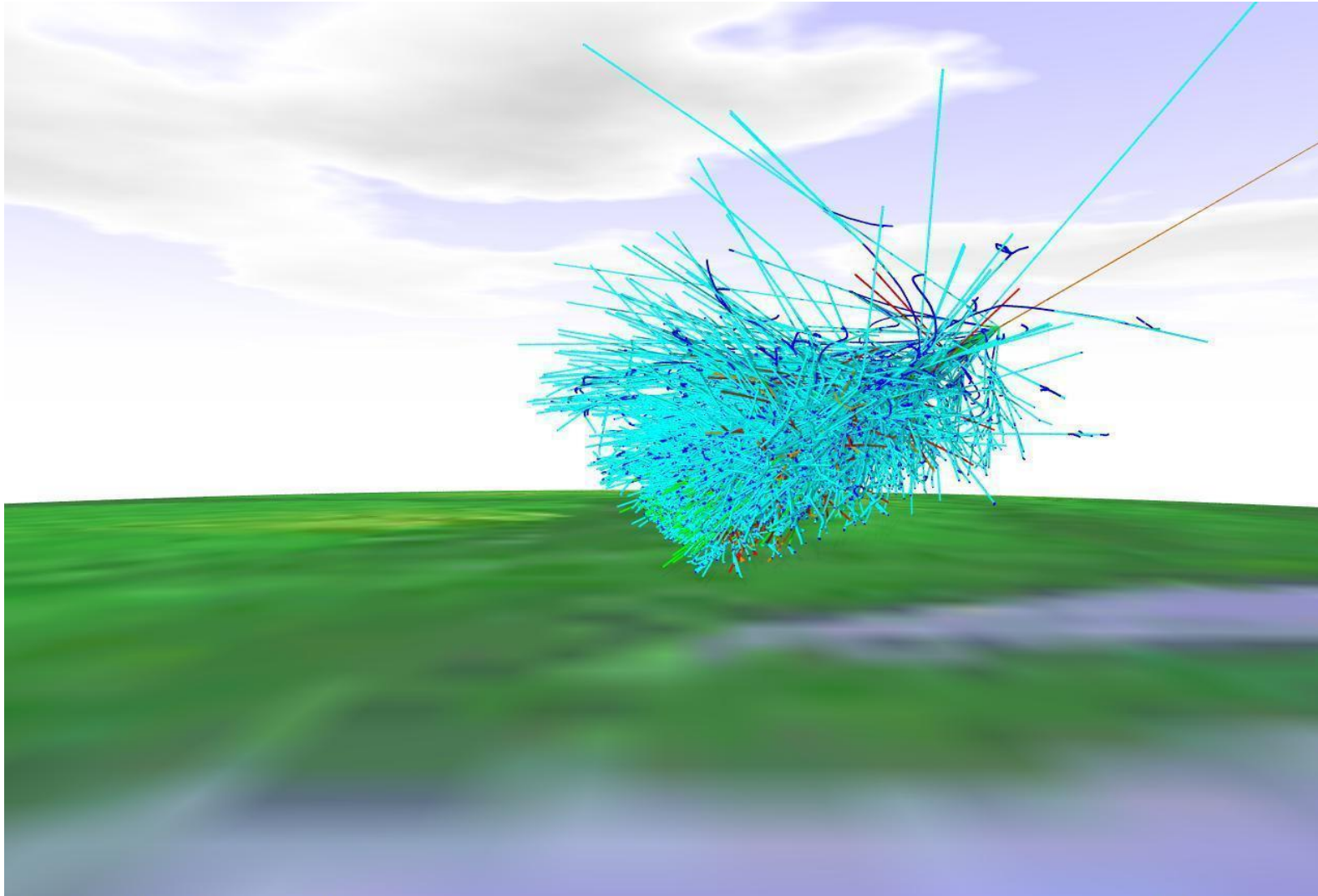
Simulation of air showers: all components



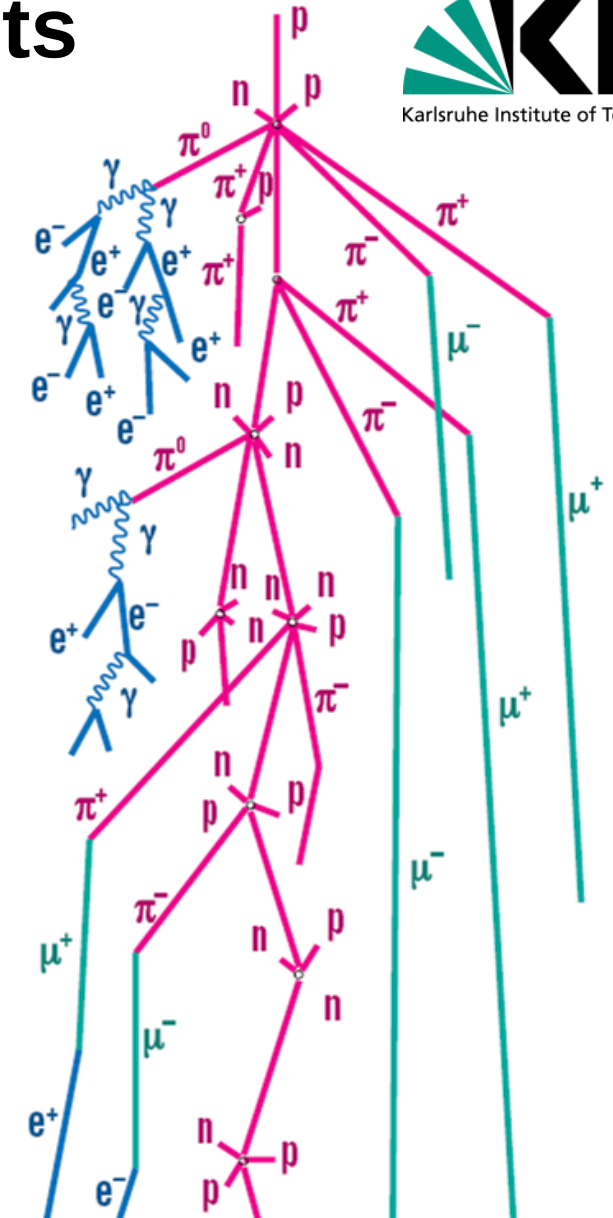
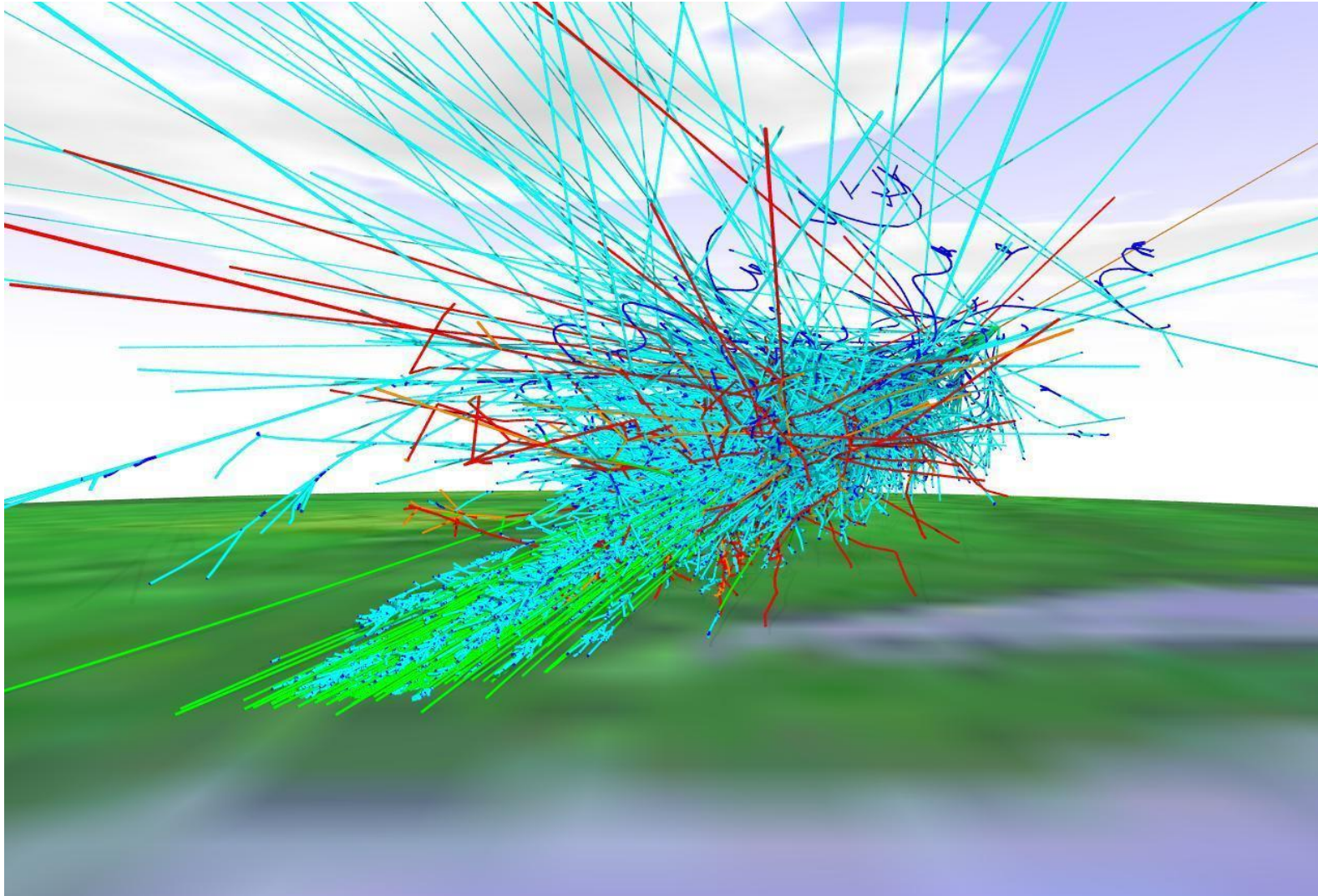
Simulation of air showers: all components



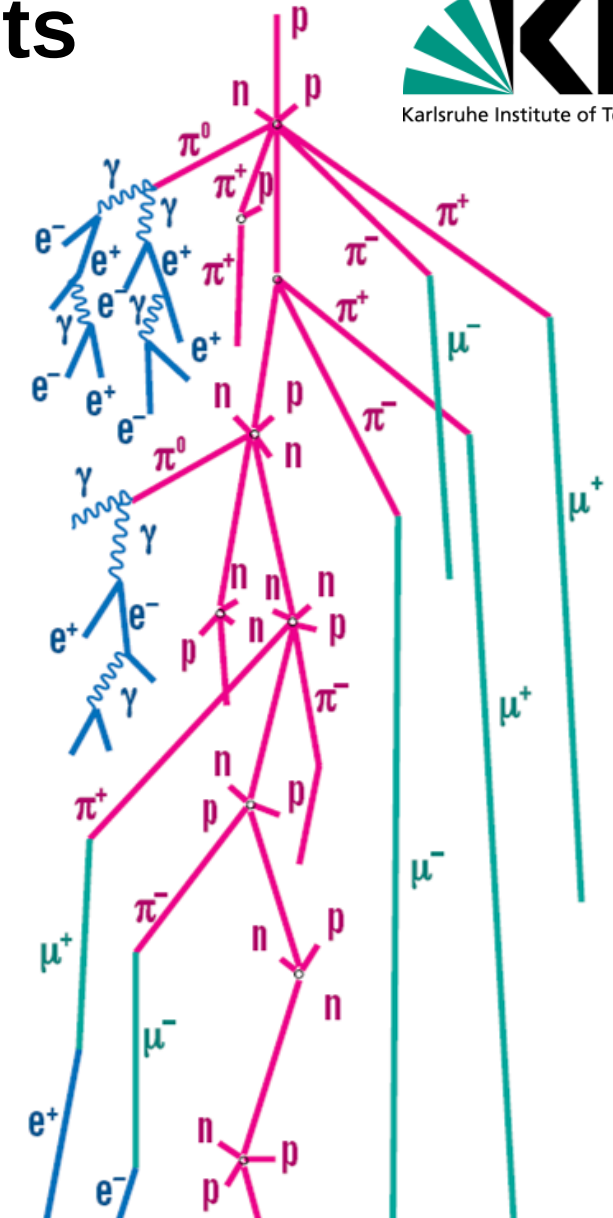
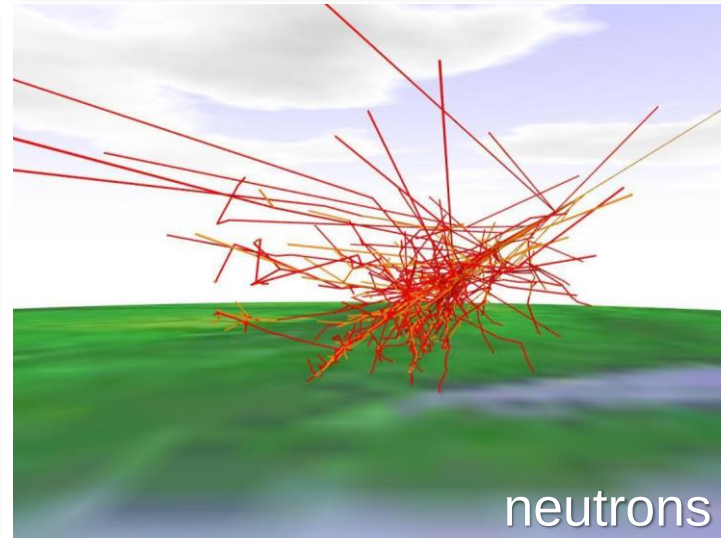
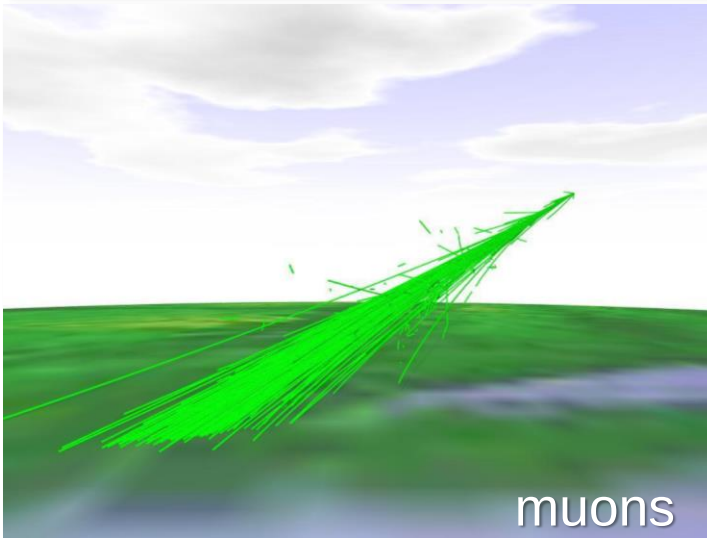
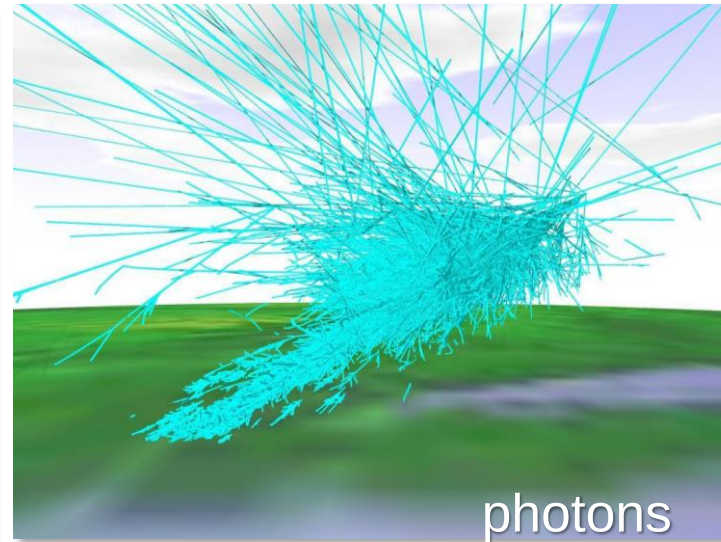
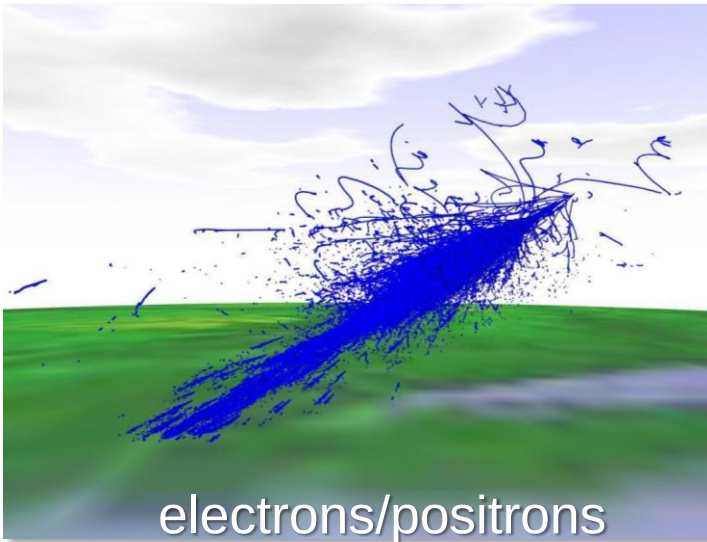
Simulation of air showers: all components



Simulation of air showers: all components

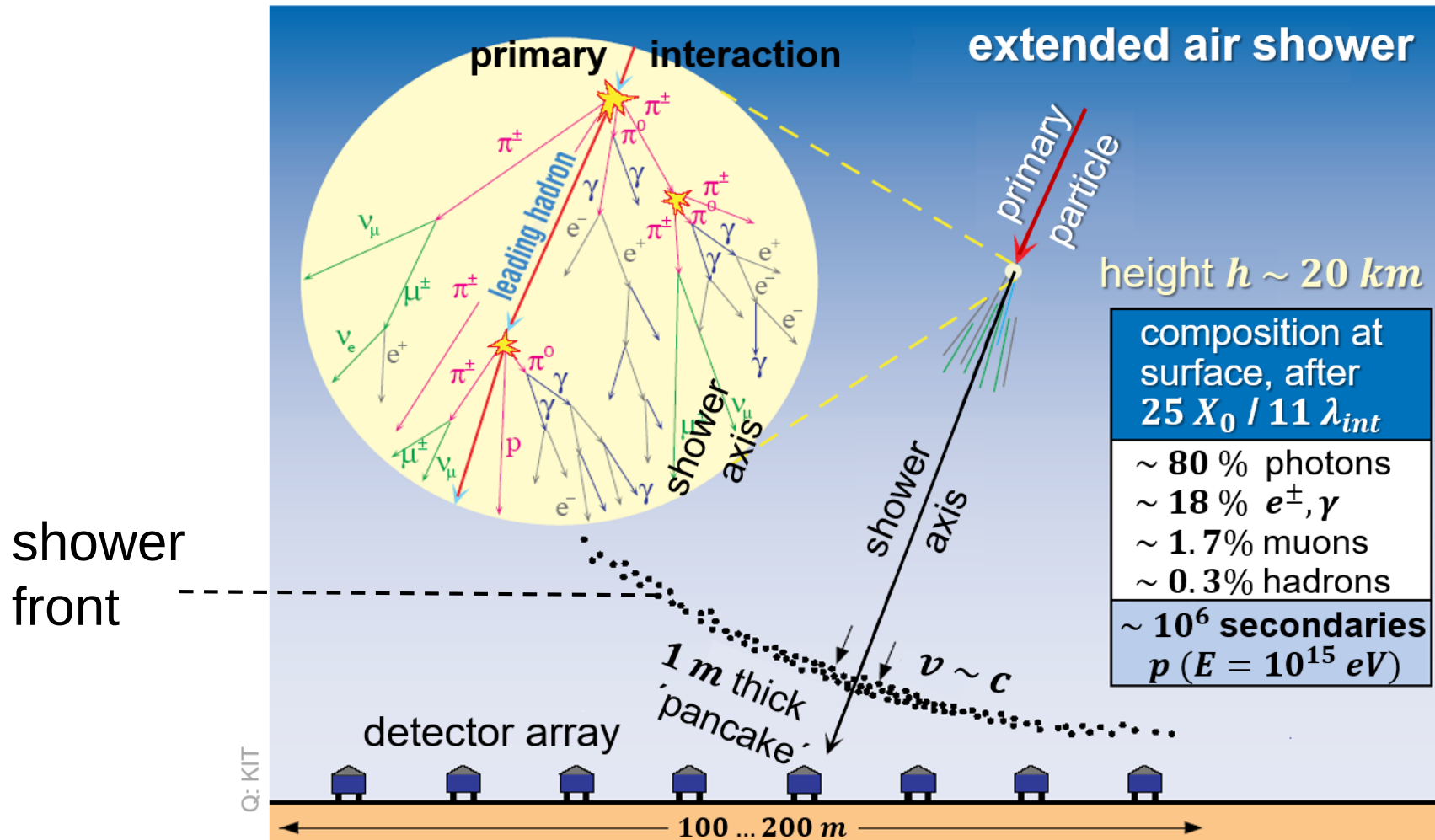


Simulation of air showers: all components

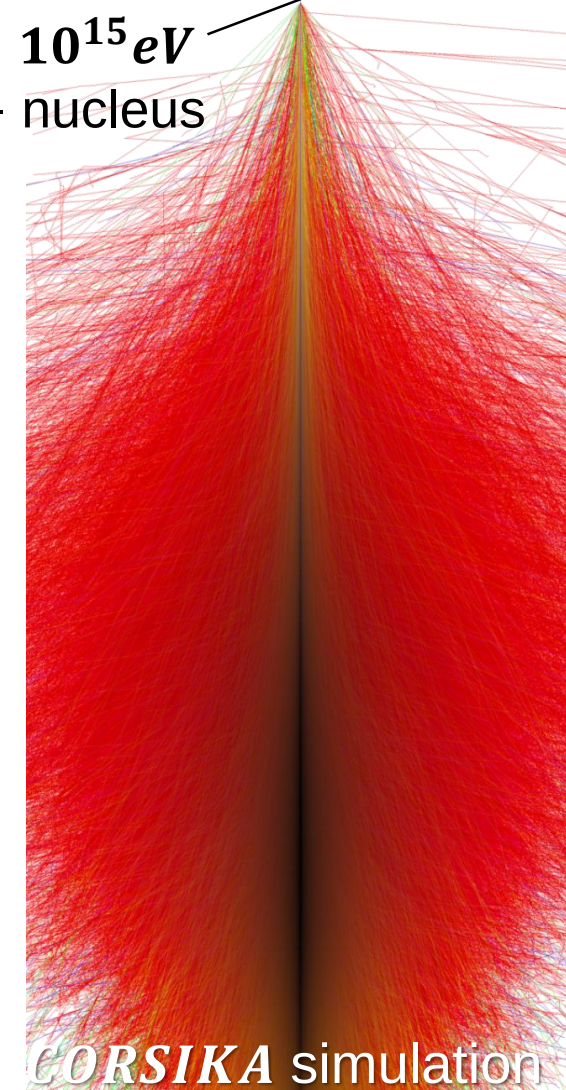


Features of an extended air shower

- primary *CR*: Earth's atmosphere acts as **calorimeter**



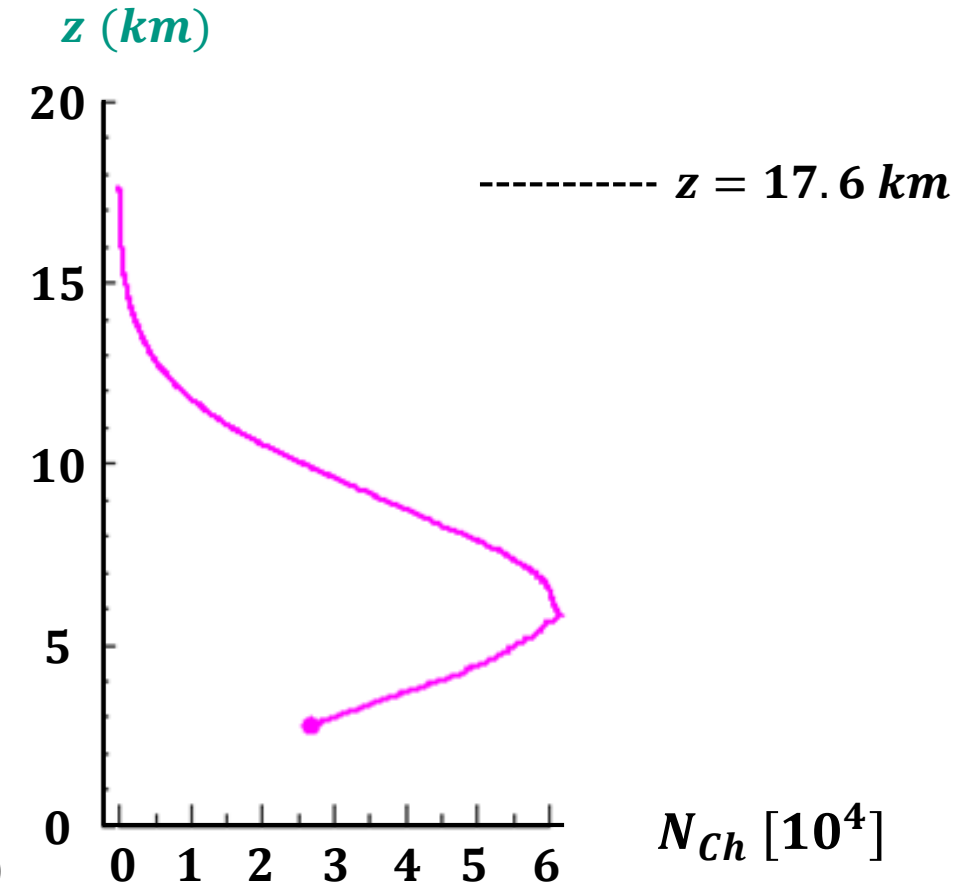
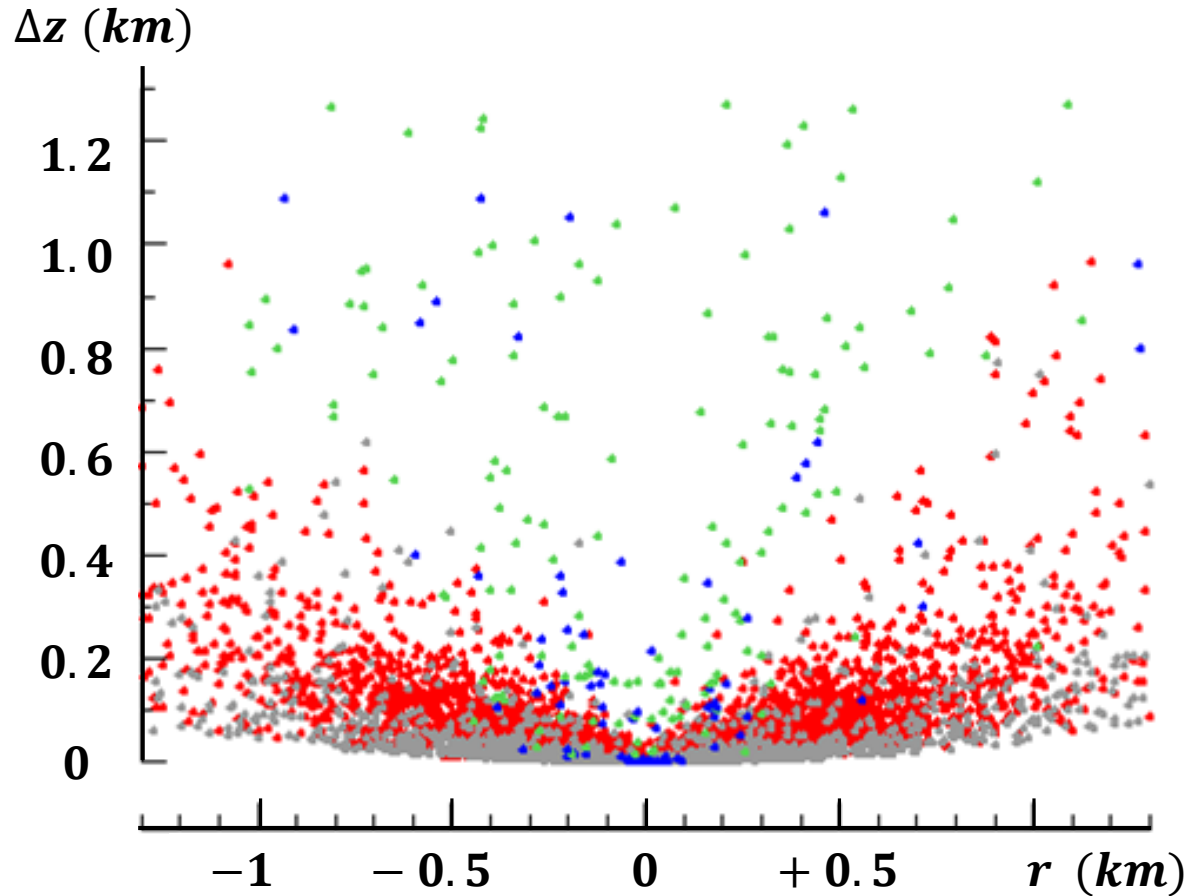
$E_0 = 10^{15} \text{ eV}$
 $Fe - \text{nucleus}$



Extended air showers – time development

- Longitudinal air shower profile = shower particles as function of height z

hadrons
neutrons
muons
electrons



Q: KIT

Shower parameters: extracting primary energy E_0

■ **Lateral distribution** of air showers = the 'footprint' at the surface, which allows to determine energy E_0

- extraction of E_0 **via**

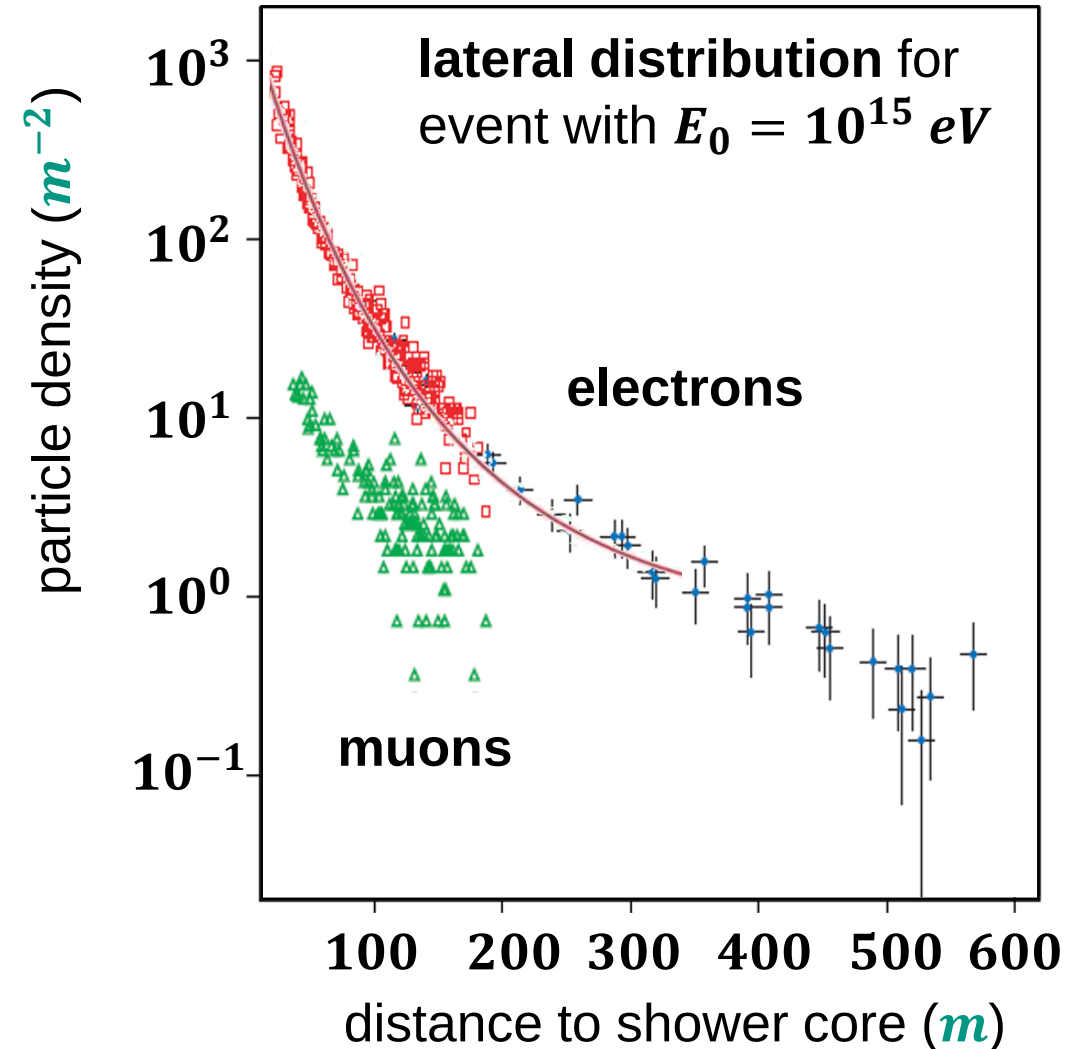
N_e = number of **electrons**
 N_μ = number of **muons** } in a shower

via integration of **lateral distribution** $\rho(r)$

- analytical parameterisation of $\rho(r)$ via a so-called **Greisen** – fit*

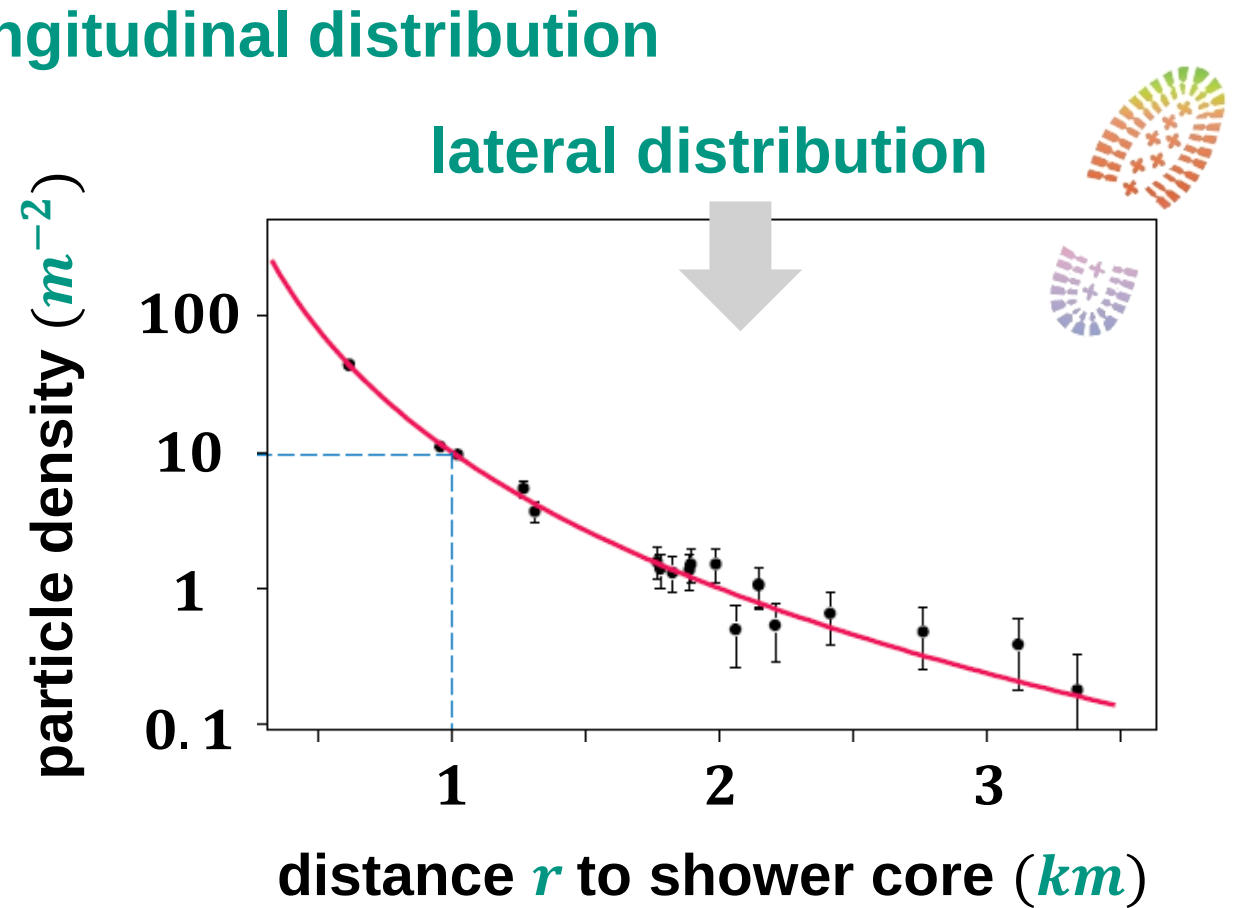
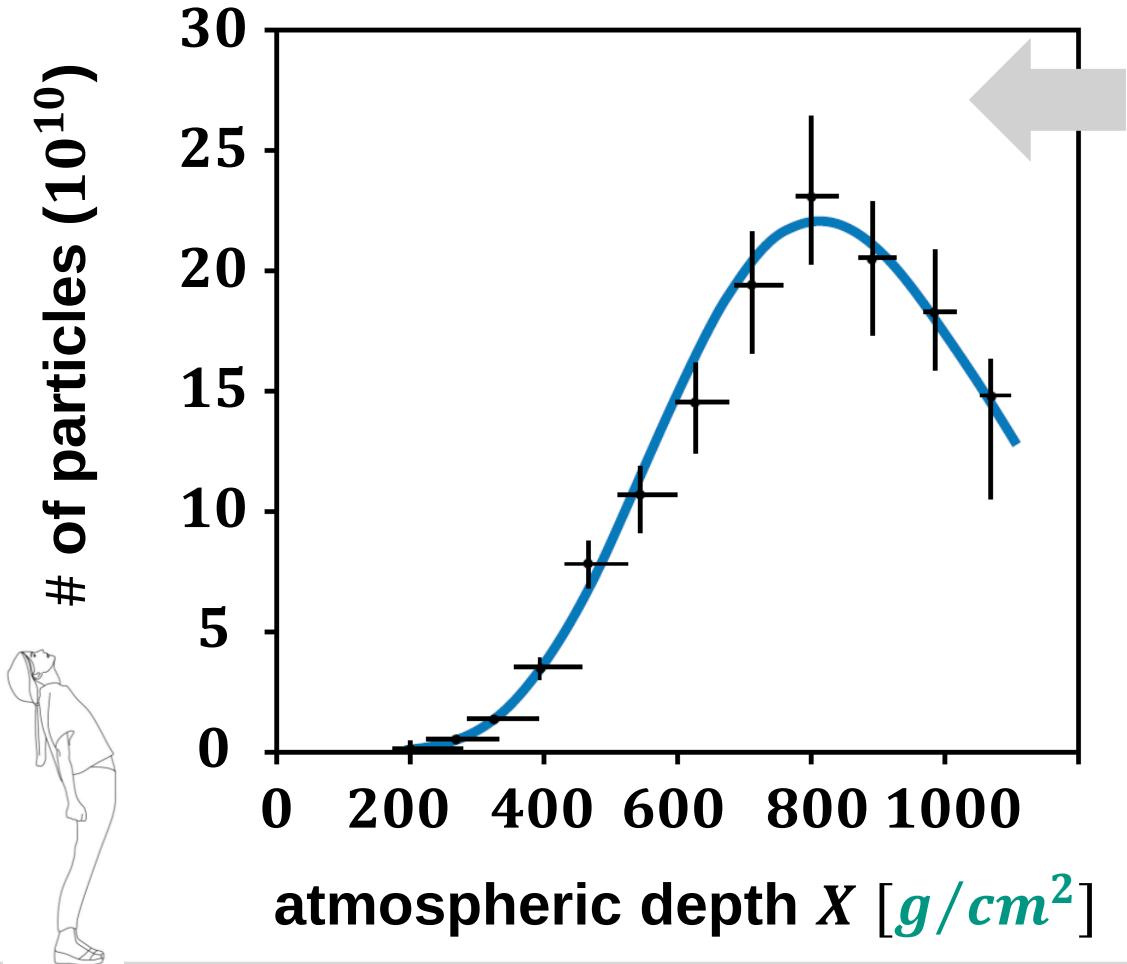


Kenneth Greisen



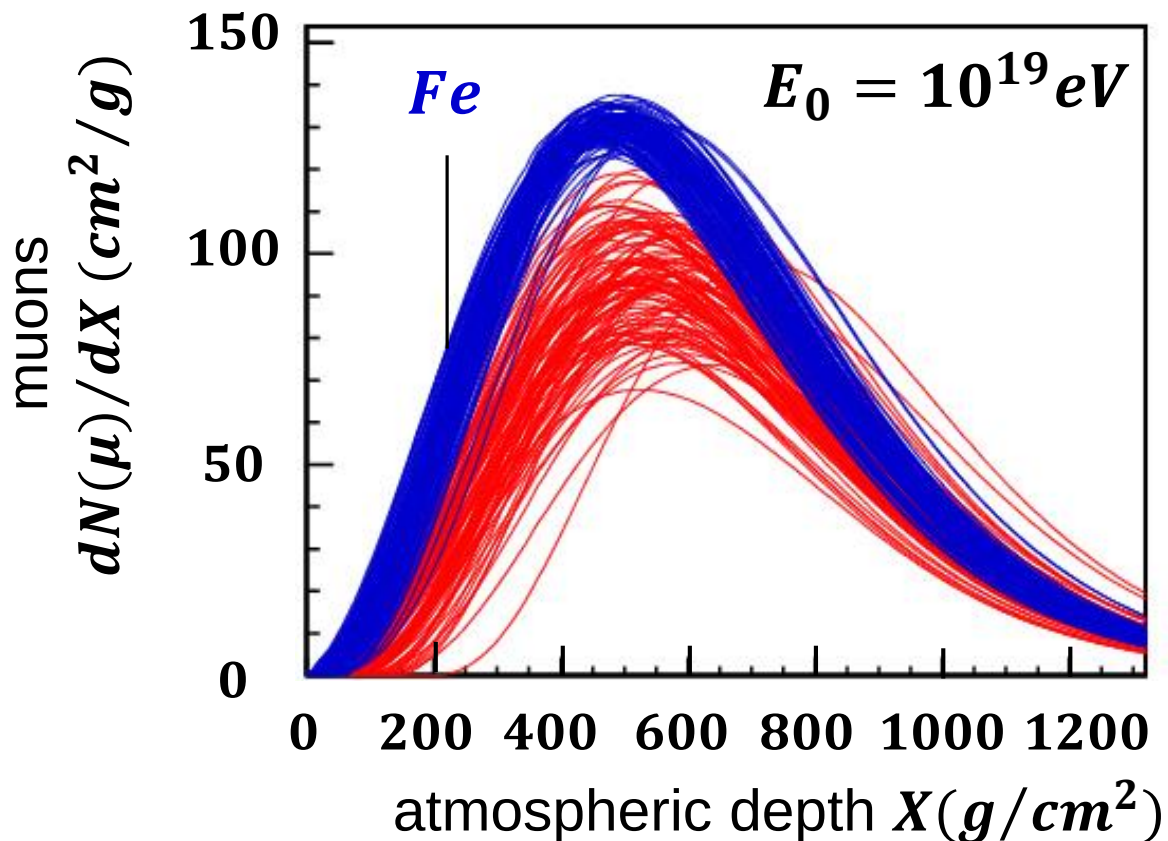
Combining lateral $\rho(r)$ with longitudinal profile

- Reconstructing two key shower parameters: primary energy E_0 & mass A

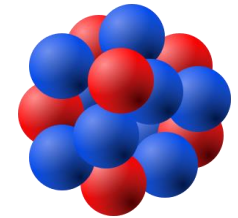


Air showers: determining the primary mass A

- Method 1: longitudinal distributions – example nucleus Fe vs. proton p
 - longitudinal profiles allow to discriminate **heavy nuclei** from light ones & protons



heavy nuclei (${}^{56}_{26}Fe$):



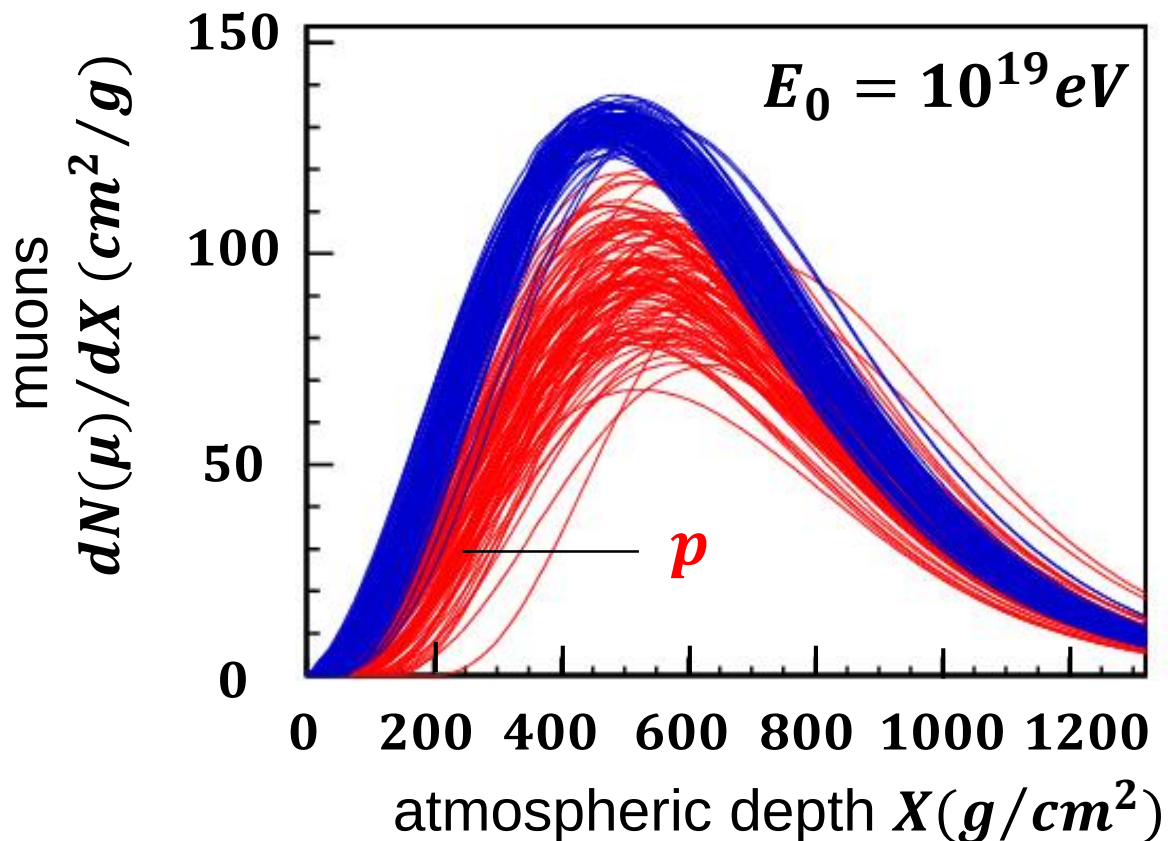
- **extended size** of Coulomb fields
($Z = 26$)

⇒ nuclei already interact in
topmost atmospheric layers

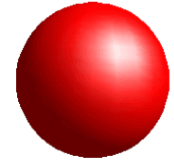
⇒ **smaller X_{max}**

Air showers: determining the primary mass A

- Method 1: longitudinal distributions – example nucleus Fe vs. proton p
 - longitudinal profiles allow to discriminate heavy nuclei from light ones & protons



lightest nuclei (incl. p):



- smaller size of Coulomb fields
($Z = 1, \dots$)

⇒ light nuclei penetrate to
deeper atmospheric layers

⇒ larger X_{max} & large fluctuations

Air showers: determining the primary mass A

■ Method 2: lateral distributions – heavy nuclei Fe vs. light ones (proton p)

- surface measurements allow to determine the ratio N_e/N_μ :

electron number N_e –
muon number N_μ

electrons (e):

- #: in general much **more e than μ**
- e have a **short range** in atmosphere
⇒ reach surface only from **larger X**

muons (μ):

- #: in general much **less μ than e**
- μ have a **large range** in atmosphere
⇒ reach surface also from **smaller X**

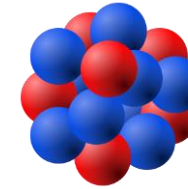
Air showers: determining the primary mass A

■ Method 2: lateral distributions – heavy nuclei Fe vs. light ones (proton p)

- surface measurements allow to determine the ratio N_e/N_μ :



X_{max} : $N_e - N_\mu$ ratio



lightest nuclei (p):

- small nuclear charge ($Z = 1, \dots$)

⇒ larger X_{max} & large fluctuations

⇒ large ratio of N_e/N_μ

heavy nuclei (${}^{56}_{26}Fe$):

- large nuclear charge ($Z = 26$)

⇒ smaller X_{max}

⇒ small ratio of N_e/N_μ

CR – detector technologies

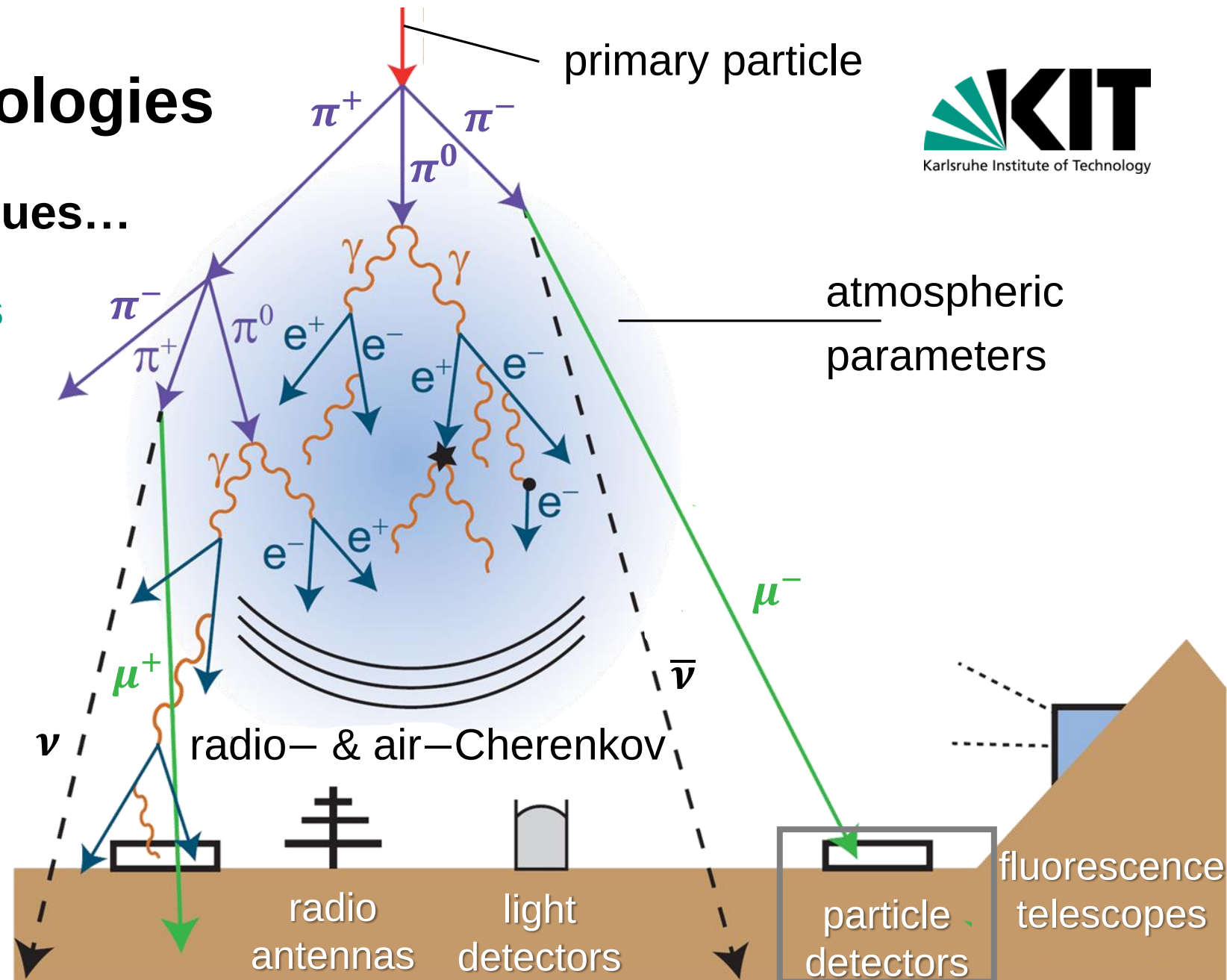
■ A wide range of techniques...

longitudinal distributions

- telescopes for
 - fluorescence light
 - Cherenkov light
- radio antennas

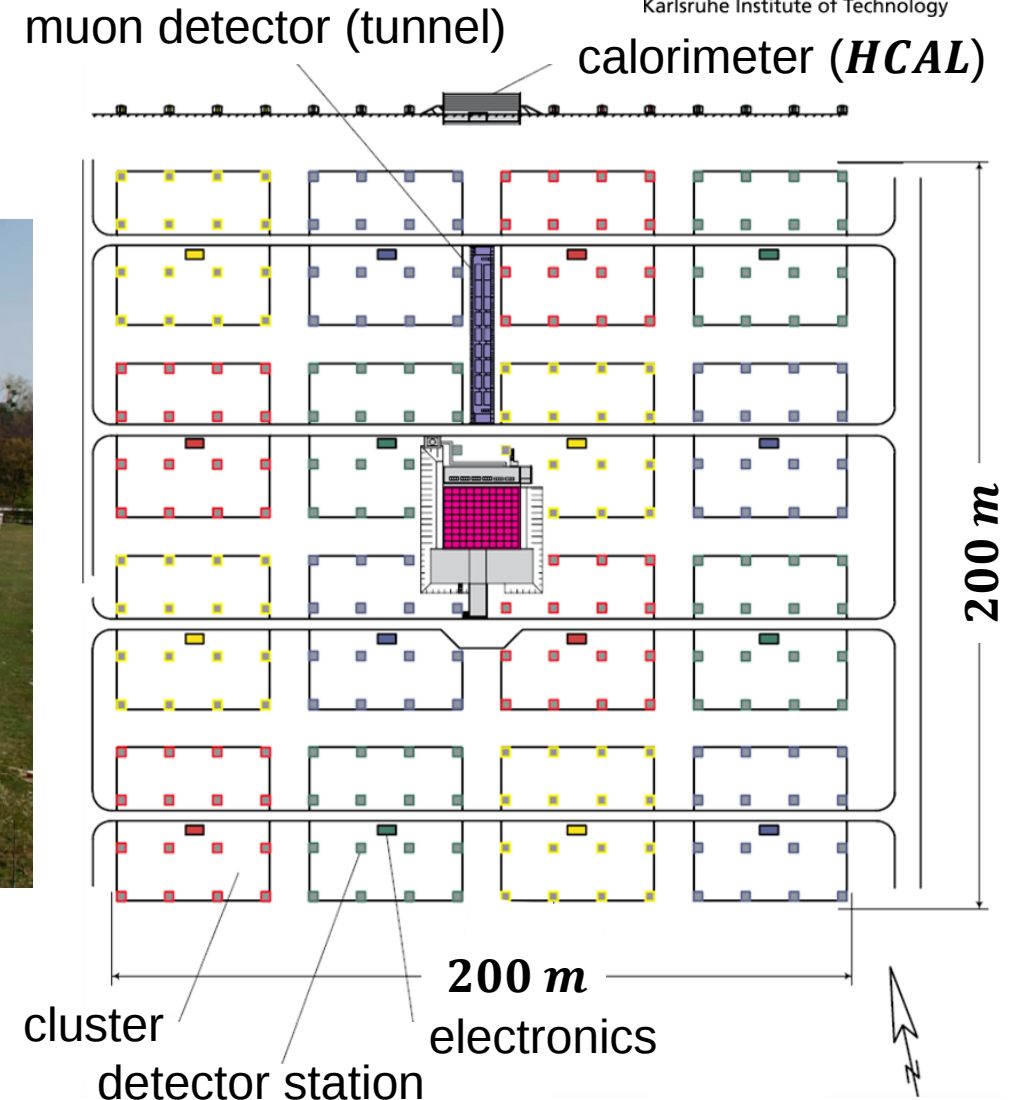
lateral distributions

- scintillators
- *MWPC**
- water Cherenkov



KASCADE experiment at KIT

■ **KASCADE** – **K**arlsruhe **S**hower **C**ore
and **A**rray **D**etector (1996 ... 2009)



total area: $A \sim 200 \times 200 \text{ m}^2$
coverage: $\varepsilon \sim 2 \%$ of total area

KASCADE experiment at KIT

■ KASCADE – three main components

Array:

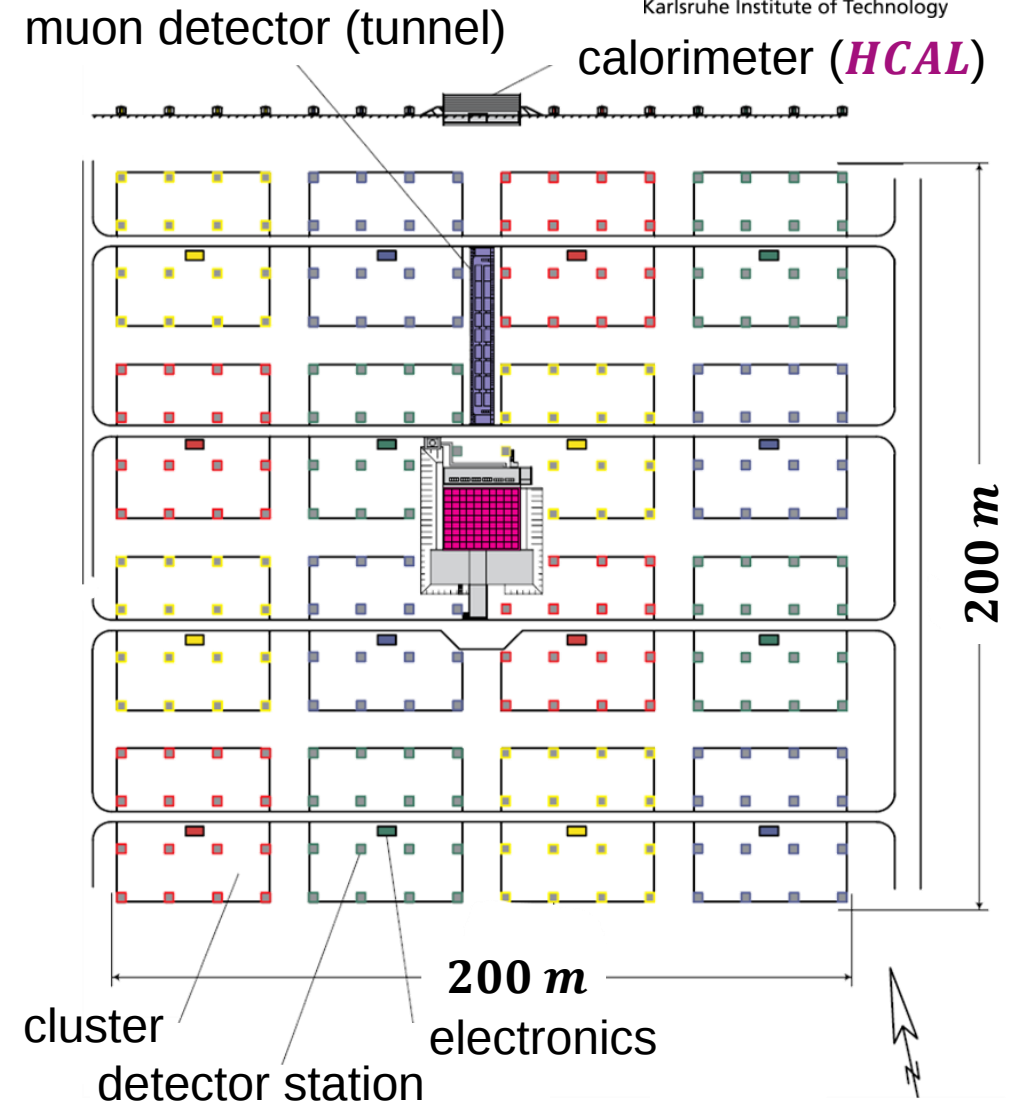
16 × 16 detector stations
detectors for muons / electrons
trigger electronics for entire shower

central calorimeter (HCAL):

*TMS** ionisation chambers
⇒ detection of hadronic shower core

muon tunnel:

sampling of muon distribution



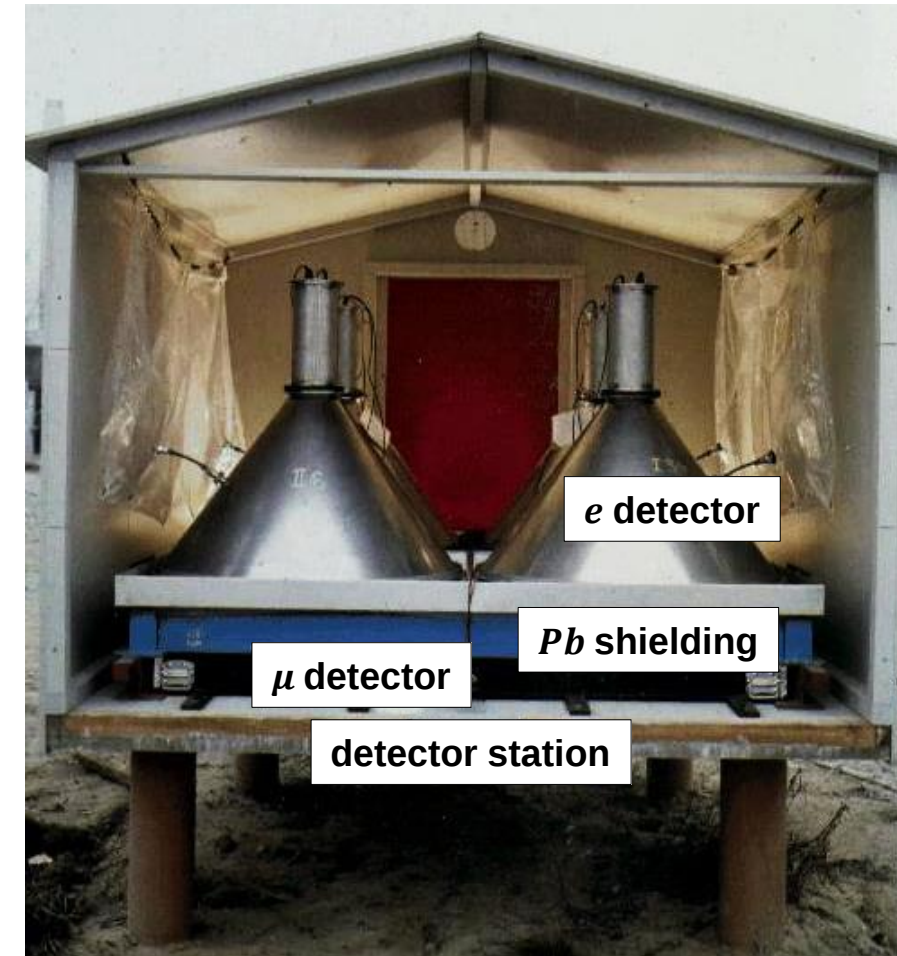
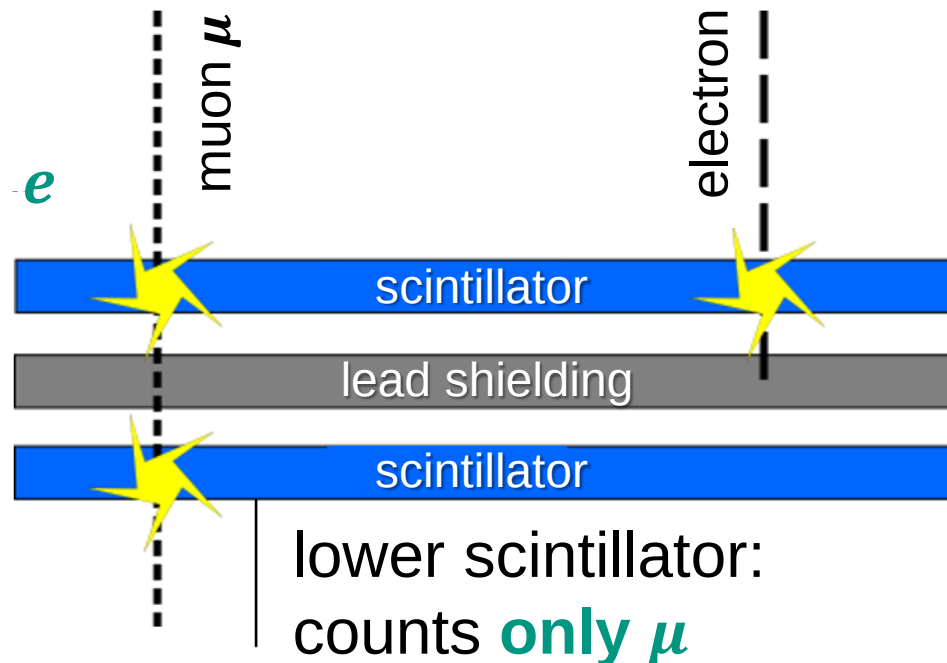
KASCADE experiment

- The detector array to measure electron & muon distributions **separately**

Muon number N_μ and electron number N_e :

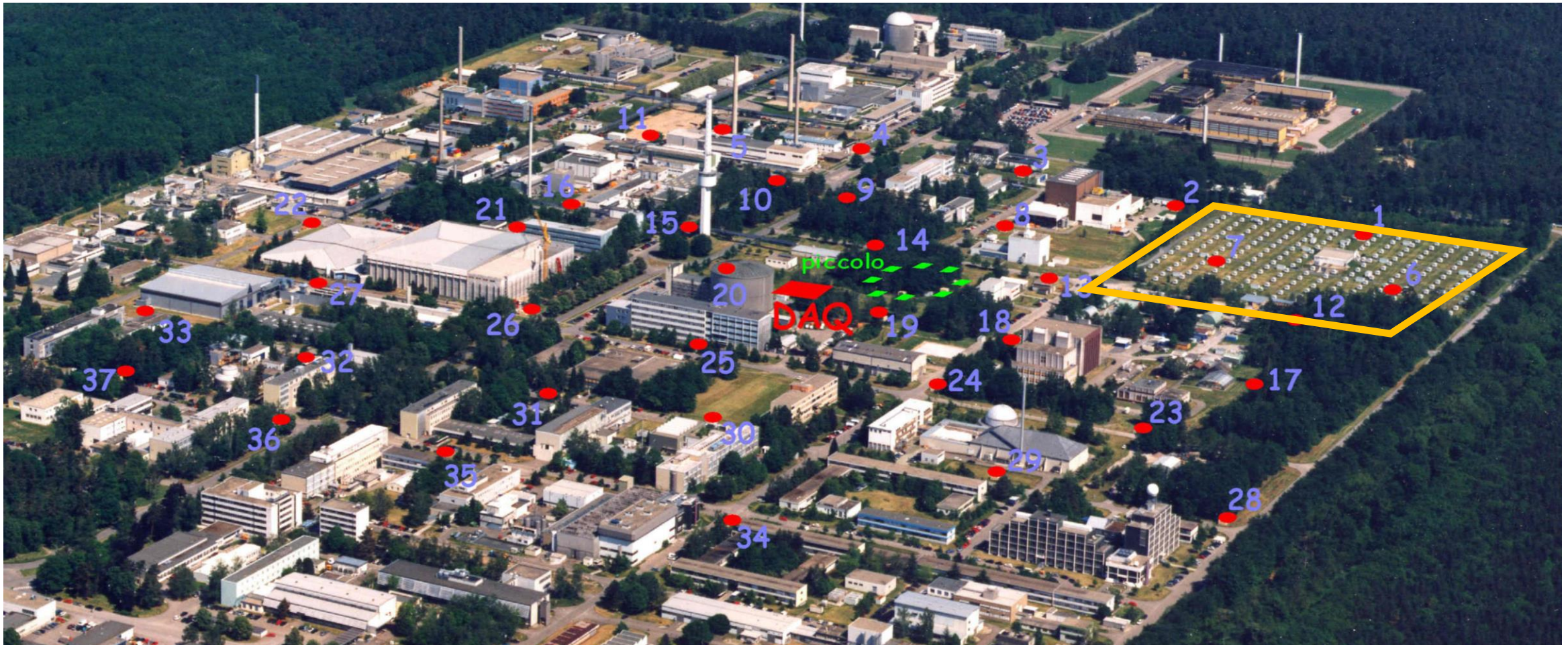
- separate measurement via 2 scintillator layers with **massive Pb – shielding** in between

upper scintillator:
counts both μ and e



KASCADE – Grande experiment at *KIT*

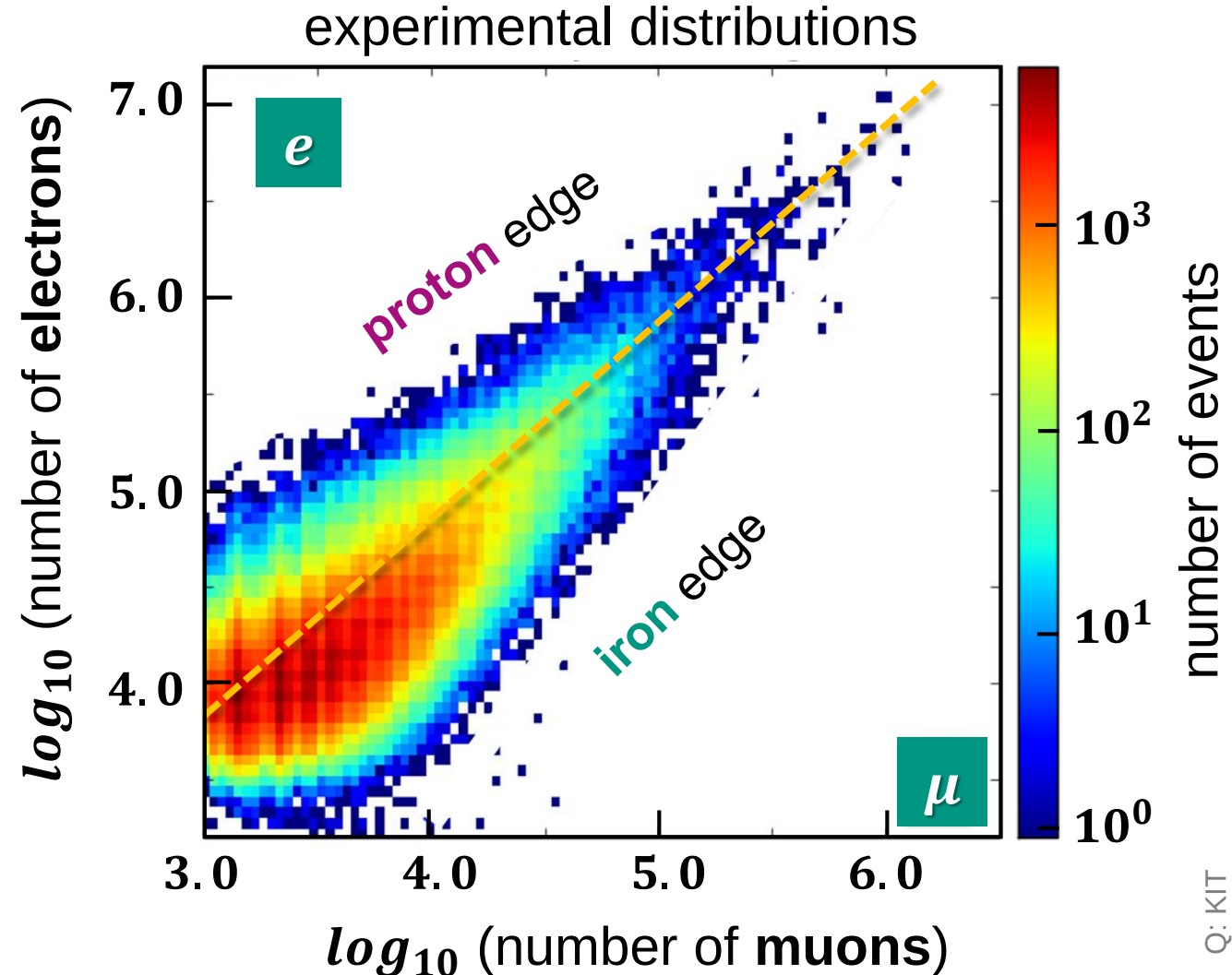
- Extending the detector array to measure up to **higher CR energies**



KASCADE – measuring mass A of primary nuclei

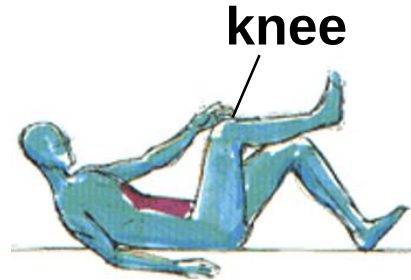
■ measuring N_e and N_μ : using
sum $N_e + N_\mu$ & ratio N_e/N_μ

- good correlation of N_e and N_μ
↳ sum $N_e + N_\mu$ as indicator
for primary energy E_0
- light nuclei ($p, \alpha, {}^{12}\text{C}$):
reactions start deep in atmosphere
↳ ratio N_e/N_μ large: small A
- heavy nuclei (${}^{56}\text{Fe}$):
reactions start at top of atmosphere
↳ ratio N_e/N_μ small: large A



KASCADE – experimental observations

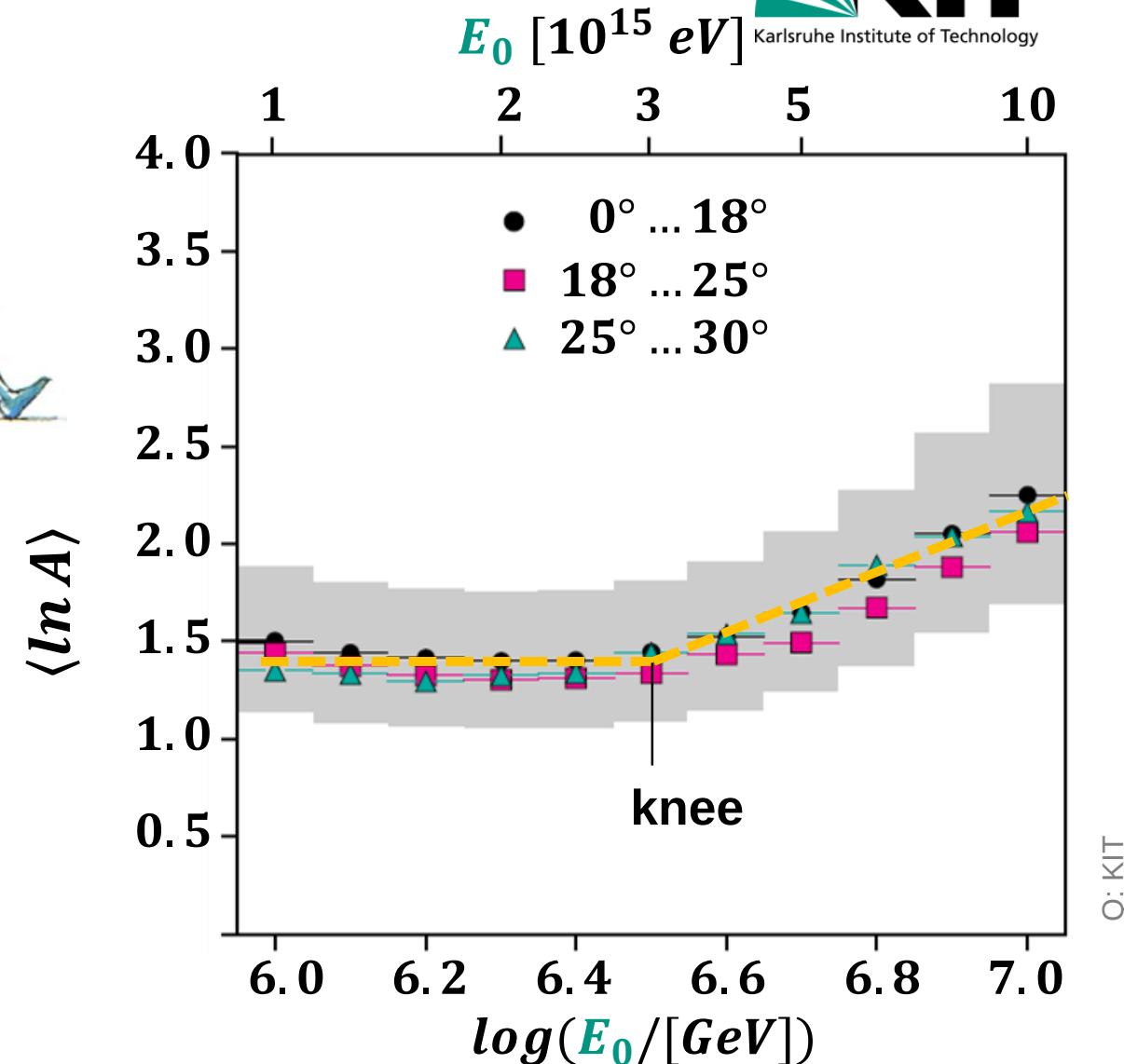
■ Chemical composition of CR changes as function of energy E_0



- **knee**: a spectral feature at energies $E_0 \sim 3 - 4 \cdot 10^{15} \text{ eV}$
- above the **knee**:

charged CR get 'heavier'

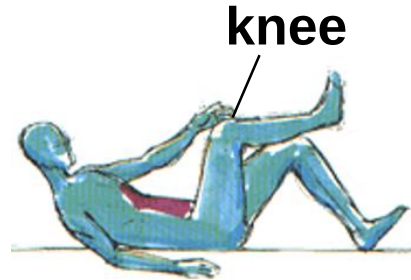
mean mass number $\langle A \rangle$ of primary increases with increasing energy E_0



KASCADE – experimental observations

■ Chemical composition of *CR* changes as function of energy E_0

- 2 resulting questions as to the **root cause of the knee**:

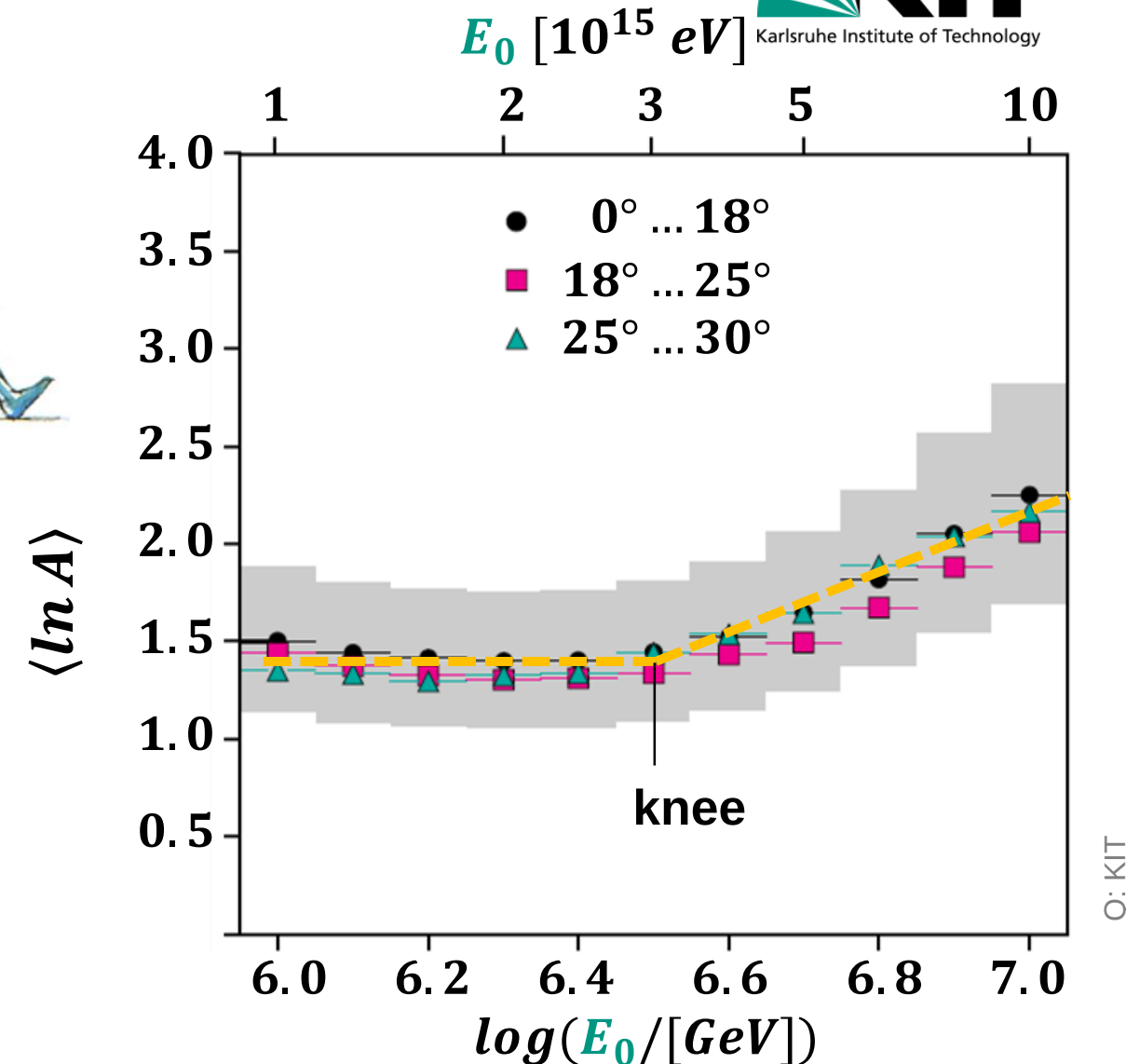


a) **cosmic accelerators**:

do they reach the end of their 'acceleration power' for light nuclei?

b) **propagation effects**:

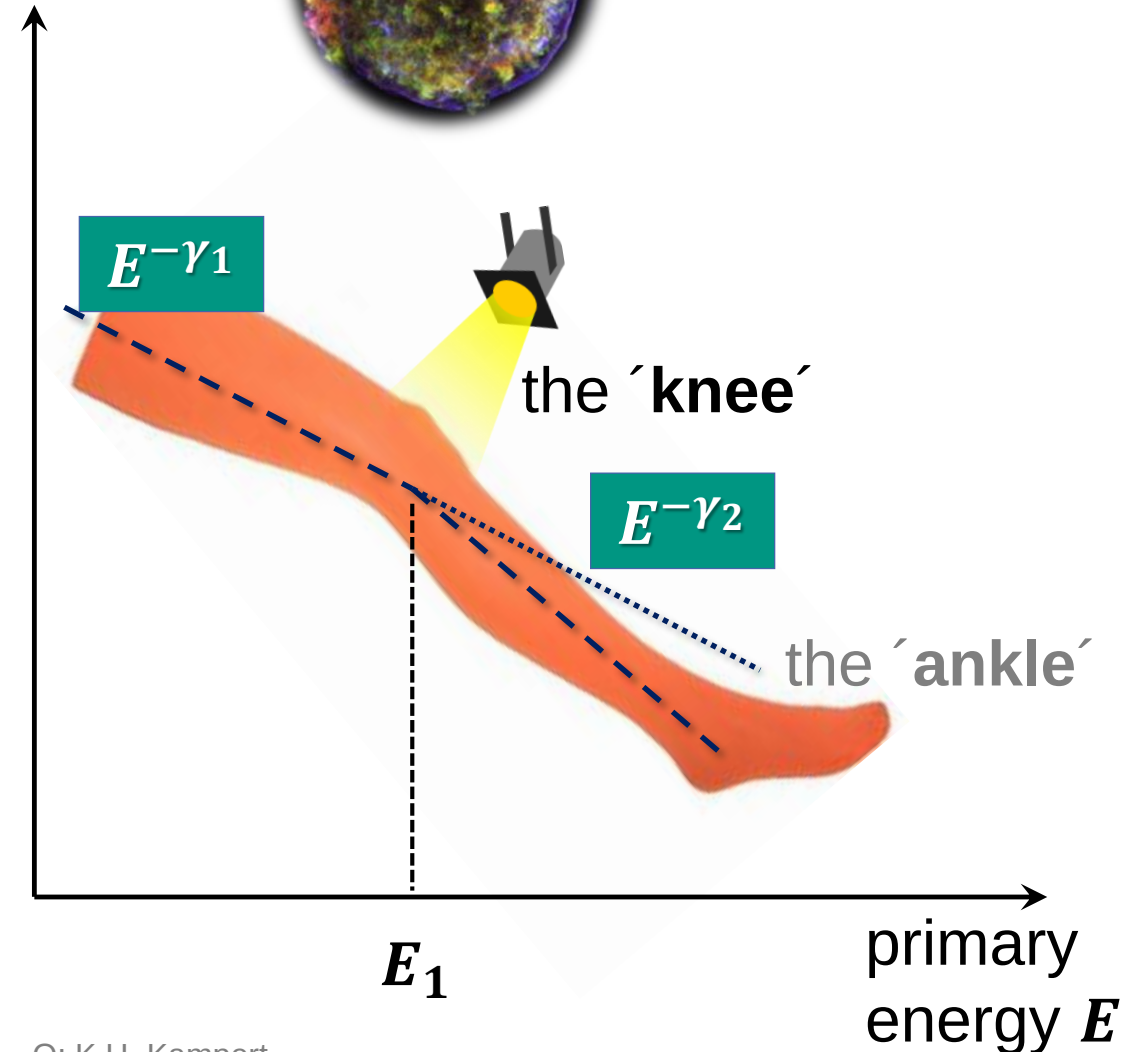
does the *CR* propagation in magnetic fields (of our galaxy) cause a **loss of light nuclei** from source to Earth?



CR energy spectra – the knee

- Distinct feature: change of spectral index & mass composition of CRs
 - at characteristic energy E_1 :
 - change of energy spectral index γ
 - change of mass composition:
light (p) $\rightarrow$ **heavy nuclei** (^{56}Fe)
 - galactic CR accelerators = SN shocks
 $\rightarrow$ SNR reach their maximum energy E_0
 - the 'ankle': change of acceleration sites -
galactic $\leftrightarrow$ **extragalactic**

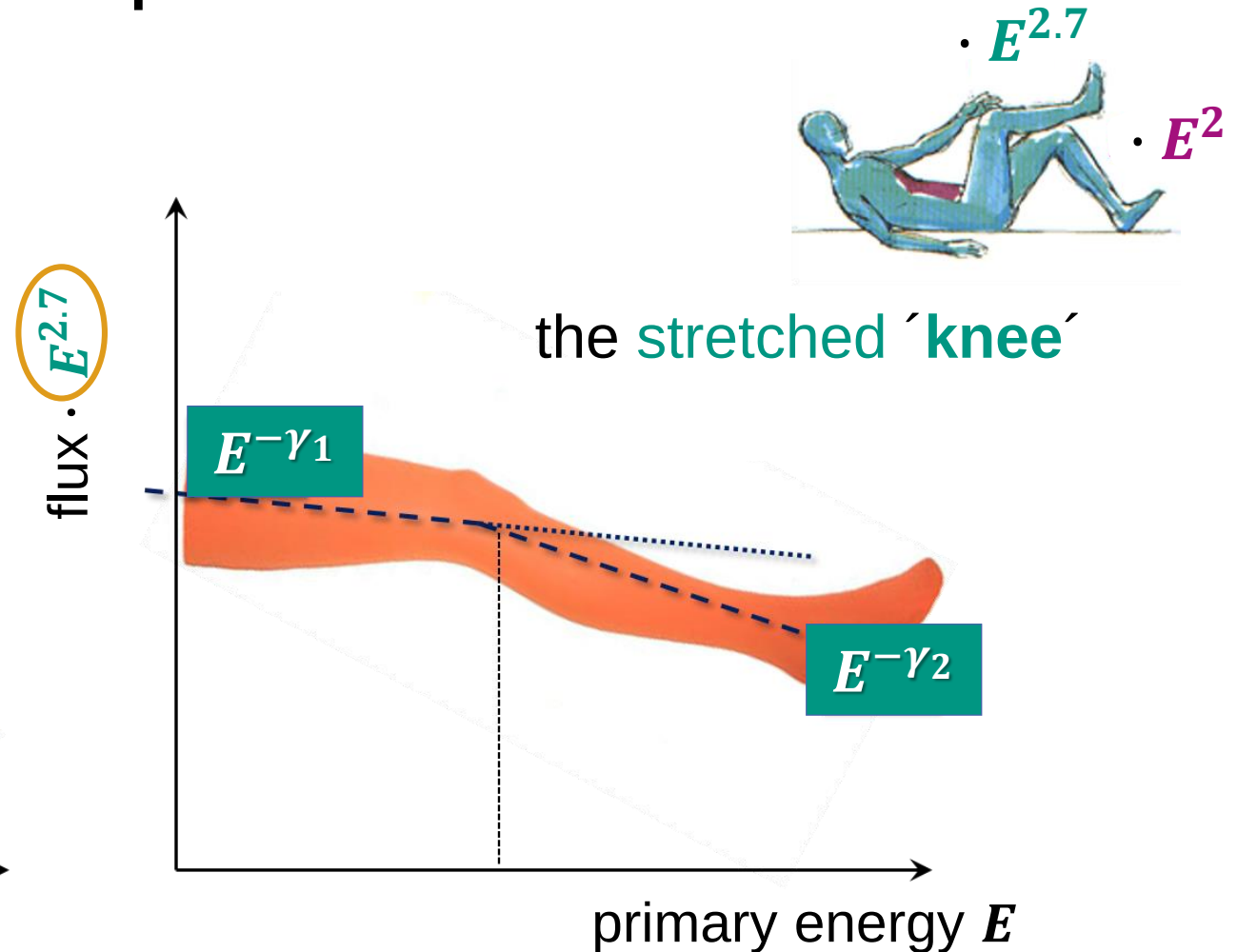
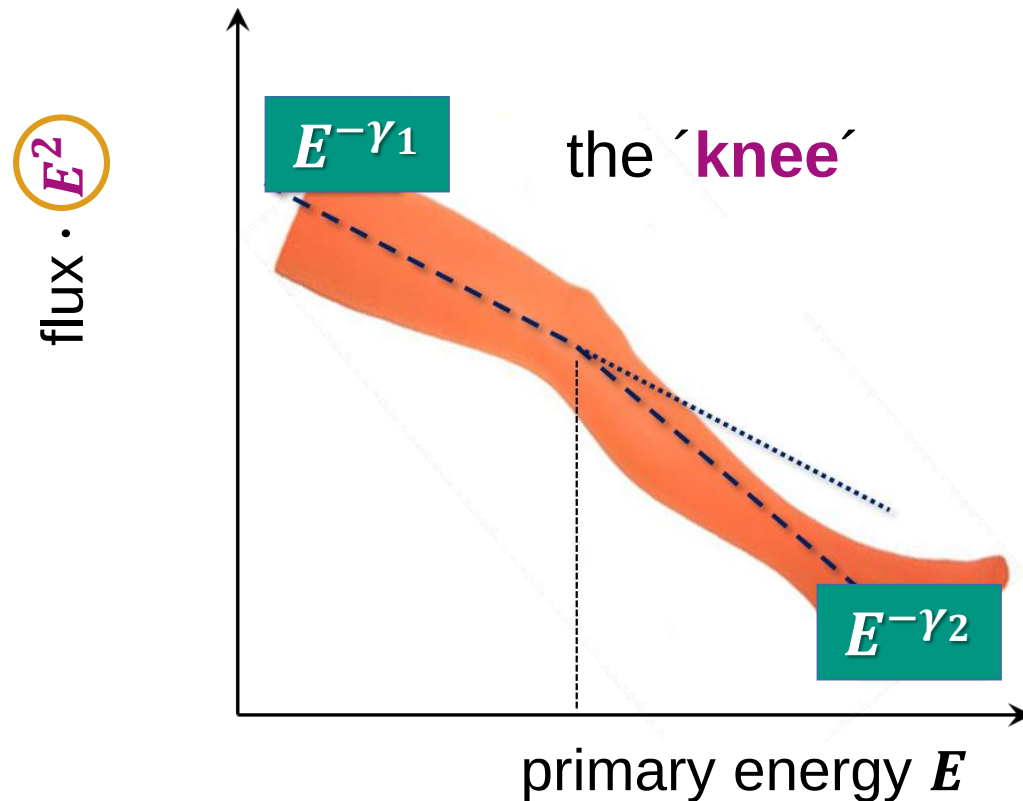
flux $\cdot E^2$



Q: K.H. Kampert

CR – energy spectra multiplied by energy factors

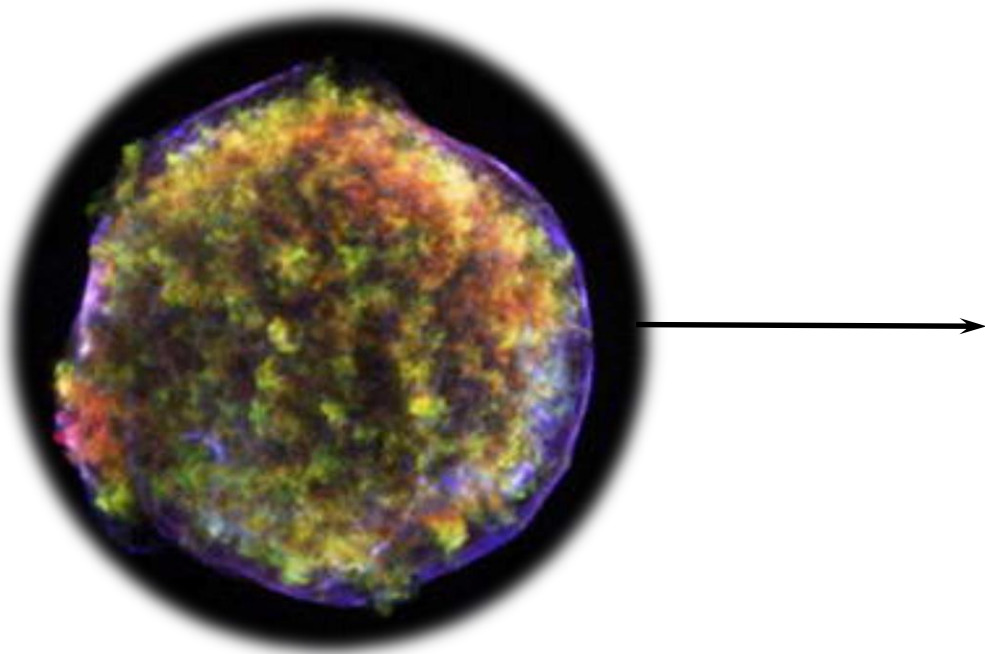
- CR community routinely multiplies spectra with an energy factor ($E^2 \dots E^3$)



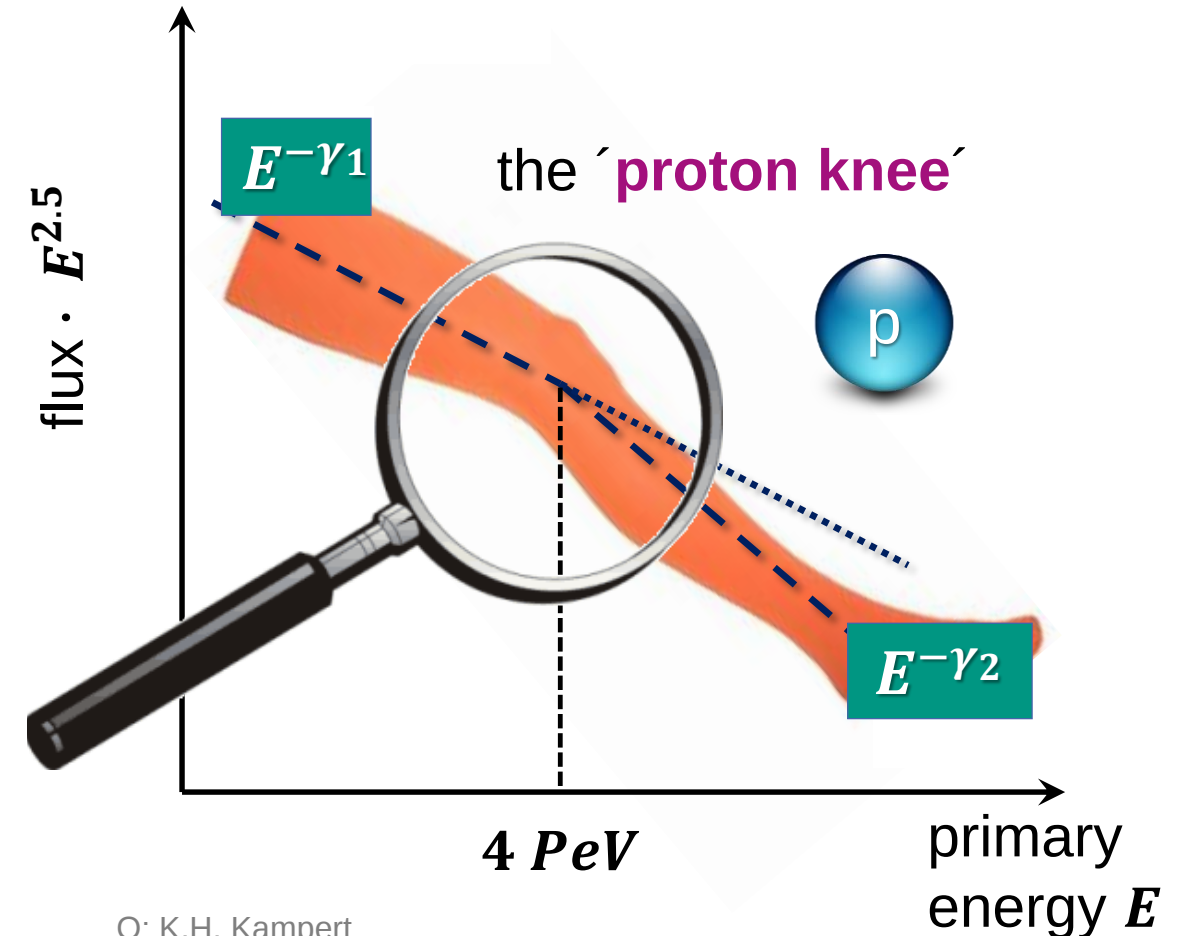
CR energy spectra – the p, Fe knee

■ observation of several ‘knees’: from the p – knee up to the Fe – knee

- important: different nuclear charges



proton acceleration in SN – shock front

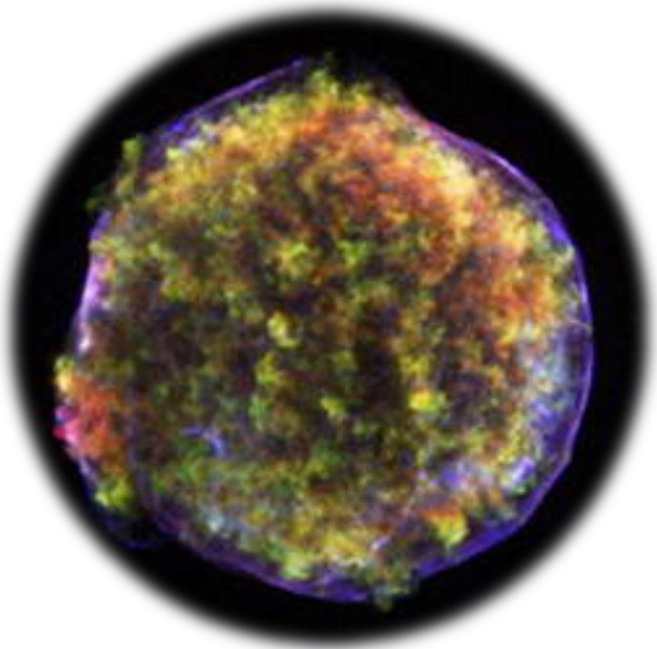


Q: K.H. Kampert

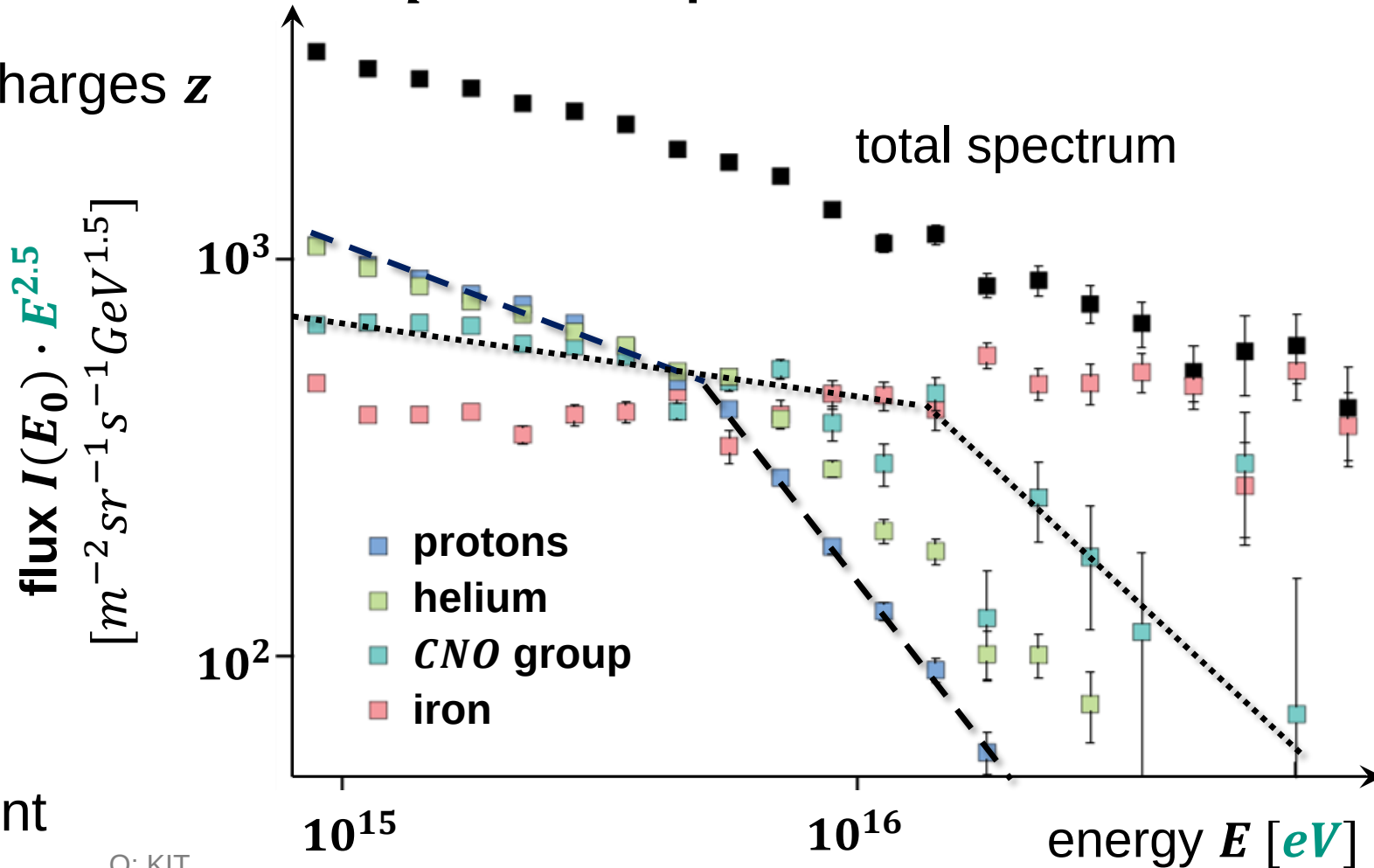
CR energy spectra – the p, Fe knee

■ observation of several ‘knees’: from the p – knee up to the Fe – knee

- important: **different** nuclear charges z



acceleration in SN – shock front



Q: KIT