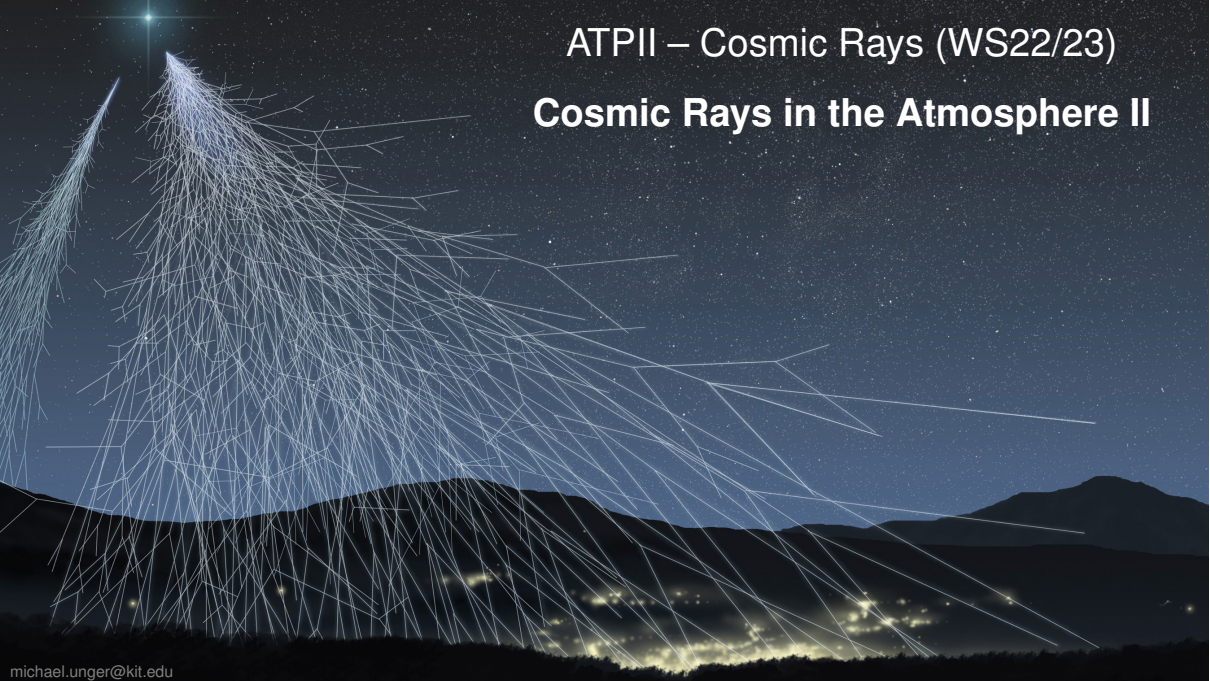
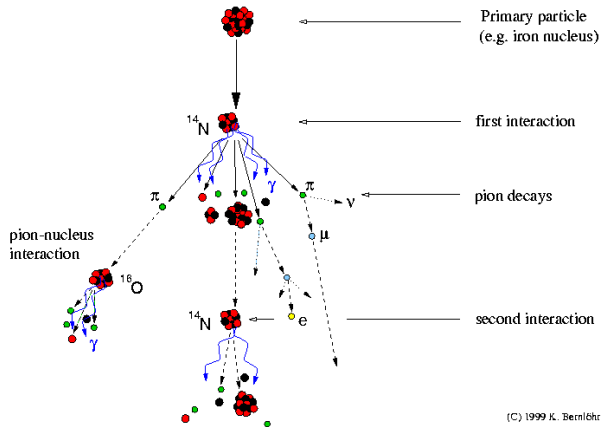
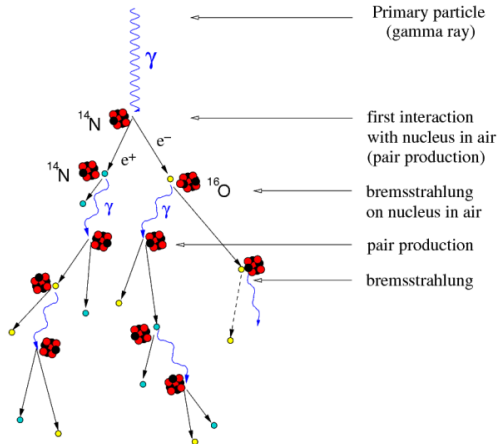


ATP II – Cosmic Rays (WS22/23)

Cosmic Rays in the Atmosphere II



Particle Cascade in the Atmosphere / Air Shower



Recap: Atmosphere

- vertical depth X_v

$$X_v = \int_h^{\infty} S(h') dh'$$

$$[X_v] = \text{g/cm}^2 \Rightarrow \text{"grammage"}$$

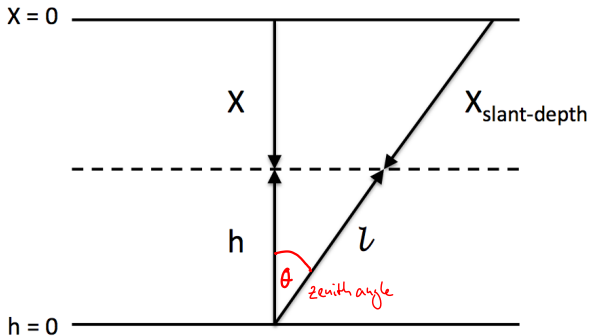
- isothermal atmosphere:

$$S(h) = S_0 e^{-h/h_0}$$

$$X_v = X_0 e^{-h/h_0}$$

- $X_0 \approx 1030 \text{ g/cm}^2$ at sea level

- scale height h_0 $\approx 8.4 \text{ km}$ at sea level, $\approx 6.4 \text{ km}$ high altitudes
above $h \approx 10 \text{ km}$



- slant depth:

$$X = \int_e^{\infty} S(h(e')) de'$$

- flat atmosphere approximation for $\theta \lesssim 65^\circ$

$$X = X_v / \cos \theta$$

Recap: Interactions and Decay

• decay length $d_j = \frac{E_j \cos \theta}{\epsilon_j}$; $\epsilon_j = m_j c^2 \frac{h_0}{c \tau_j} \leftarrow \text{lifetime}$ $[d] = \text{g/cm}^2$

Table 5.3 Decay constants for various particles

Particle	$c\tau$ (cm)	ϵ (GeV)
$\mu^\pm$	6.59×10^4	1.0
$\pi^\pm$	780	115
π^0	2.5×10^{-6}	3.5×10^{10}
$K^\pm$	371	850
K_S	2.68	1.2×10^5
K_L	1534	208
$D^\pm$	0.031	3.7×10^7
D^0	0.012	9.9×10^7

Table 5.4 Atmospheric interaction and attenuation lengths for $\gamma = 2.7$ in air, in units of g/cm^2 .

E_{lab} (GeV)	λ_N	Λ_N	λ_π	Λ_π	λ_K	Λ_K
100	88	120	116	155	134	160
1000	85	115	111	148	122	147
10000	79	106	101	135	110	133
100000	72	97	87	114	95	114

• interaction length: $j + \text{air} \rightarrow X$

mass density
 $\lambda_j = \ell_j S = \frac{S}{n_A \sigma_{j,\text{air}}} = \frac{\langle A \rangle m_p}{\sigma_{j,\text{air}}}$
 $[\lambda] = \text{g/cm}^2$ $[S] = \text{cm}$ number density σ cross section

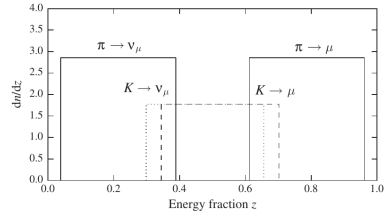
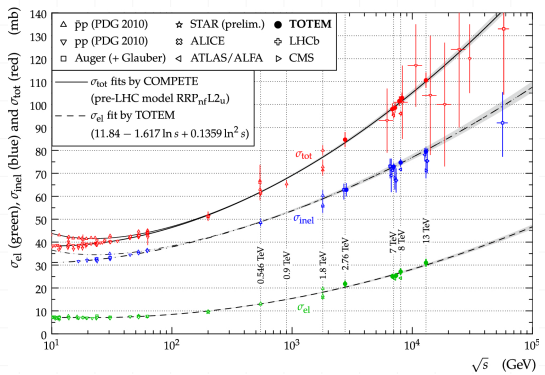


Figure 6.1 Decay distributions for π -decay and K -decay into $\mu \nu_\mu$ for 200 MeV/c parent mesons. z is the ratio of the total lab energy of the decay product to that of the parent.

Recap: Spectrum-weighted moments "Z-factors"

- inclusive cross section: $j + \text{air} \rightarrow a + X$
- inclusive energy distribution of particles of type a :

$$E_a \frac{dN(E_j, E_a)}{dE_a} \equiv F_{ja}(E_j, E_a) \approx F_{ja}\left(\frac{E_a}{E_j}\right)$$

'Feynman scaling'

• 'Z-factor': $Z_{pa} = \int_0^1 x^{\gamma-2} F_{ja}(x) dx \quad \left(x = \frac{E_a}{E_j}\right)$

- if energy spectrum j is a power law: $\phi_j \sim \frac{dN_j}{dE_j} = N_j^0 \cdot E^{-\gamma}$

$$\underline{\underline{\phi_a(E_a) = Z_{ja} \cdot \phi_j^0 E_a^{-\gamma}}}$$

Table 5.2 *Spectrum-weighted moments at 1 TeV*

Index $p\text{-air}$	$\gamma = 2$	$\gamma = 2.7$
$\pi^0(+\eta)$	0.206 ^I	0.0459
π^+	0.206	0.0489
π^-	0.156	0.0324
K^+	0.030	0.0071
K^-	0.018	0.0036
$K_L + K_S$	0.043	0.0092
$p + \bar{p}$	0.217	0.126
$n + \bar{n}$	0.114	0.052

Gaisser, Engel & Resconi, "Cosmic Rays & Particle Physics"

$$Z_{p\pi^+} > Z_{p\pi^-} \quad (\text{charge})$$

$$Z_{pp} \gg Z_{p\pi^\pm} \quad (\text{baryon number})$$

$$Z_{K^+} \gg Z_{K^-} \quad (\text{associate hyperon production})$$

$$\begin{aligned} p+N &\rightarrow K^+ + K^- + X \\ p+N &\rightarrow K^+ + \Lambda + X \end{aligned}$$

Cascade Equations

evolution of particle fluxes in the atmosphere:

$$\frac{dN_i(E_i, X)}{dX} = \underbrace{-\frac{N_i(E_i, X)}{\lambda_i}}_{\text{interaction}} - \underbrace{\frac{N_i(E_i, X)}{d_i}}_{\text{decay}} + \underbrace{\sum_{j=i}^N \int_E^{\infty} \frac{F_{ji}(E_i, E_j)}{E_i} \frac{N_j(E_j, X)}{\lambda_j} dE_j}_{\text{production}}$$

boundary conditions:

"inclusive particle flux": $N(E, 0) \equiv N_0(E) \approx 1.7 E^{-2.7} \frac{\text{nucleons}}{\text{cm}^2 \text{ sr s GeV/A}}$ nucleon flux at top of atmosphere

single air shower: $N(E, 0) = A \cdot S\left(E - \frac{E_0}{A}\right) \delta(t - t_0)$ (superposition approximation: $(E_0, A) \hat{=} A \times (E/A, 1)$)

Nucleon Flux $N_N(E, x) \equiv N(E, x)$

no decay, no production $j \neq i$ (e.g. no $\pi^+ + \text{air} \rightarrow p + X$; $Z_{\pi^+ N} \approx 0.02 \ll Z_{NN} \approx 0.18$)

$$\frac{dN(E, x)}{dx} = - \frac{N(E, x)}{\lambda_N} + \int_E^\infty \frac{N(E', x)}{\lambda_N(E')} F_{NN}(E, E') \frac{dE'}{E}$$

Ansatz: $N(E, x) = G(E) g(x)$ factorization and change of variables $x = E/E'$

$$\begin{aligned} (E' = E/x, \frac{dE'}{dx} = -\frac{E}{x^2}) \\ \rightarrow dE' = -\frac{E}{x^2} dx \\ x(E' = \infty) = 0 \\ x(E' = E) = 1 \end{aligned}$$

$$\Rightarrow G g' = - \frac{G g}{\lambda_N} + g \int_0^1 \frac{G(E/x) F_{NN}(x, E)}{\lambda_N(E/x)} \frac{dx}{x^2}$$

$$\Rightarrow \frac{g'}{g} = - \frac{1}{\lambda_N(E)} + \frac{1}{G(E)} \int_0^1 \frac{G(E/x) F_{NN}(x, E)}{\lambda_N(E/x)} \frac{dx}{x^2}$$

$$\approx - \frac{1}{\lambda_N} \left(1 - \int_0^1 x^{E-2} F_{NN}(x) dx \right) = - \frac{1 - Z_{NN}}{\lambda_N}$$

$\lambda_N \approx \text{const}$, $F(x, E) \approx F(x)$ (Feynman scaling)

+ boundary condition: $N(E, 0) = E^{-\gamma} g(0)$

$$Z_{NN} = Z_{pp} + Z_{pn}$$

$$(Z_{pp} = Z_{un}, Z_{pn} = Z_{np})$$

Solution:

$$N(E, x) = N_0 e^{-\frac{x}{\Lambda}} E^{-\gamma}$$

attenuation length

$$\Lambda = \frac{\lambda_N}{1 - Z_{NN}}$$

"regeneration"

Meson Flux

$$N_\pi(E, X) \equiv \Pi(E, X)$$

similar for $K^\pm, D^\pm$

$$\frac{d\Pi}{dX} = -\left(\frac{1}{\lambda_\pi} + \frac{1}{d_\pi}\right)\Pi + \int_0^1 \frac{\Pi(E/x) F_{\pi\pi}(x, E)}{\lambda_\pi(E/x)} \frac{dx}{x^2} + \int_0^1 \frac{N(E/x) F_{N\pi}(x, E)}{\lambda_N(E/x)} \frac{dx}{x^2}$$

(note $x = \frac{E}{E'}$, X : slant depth)

• boundary condition: $\Pi(0, E) = 0 \Leftrightarrow$ no primary mesons

• extreme cases:

a) $E \gg \varepsilon_\pi \Rightarrow$ no decay $\Pi(E, X) = N(E, 0) \frac{Z_{N\pi}}{1 - Z_{NN}} \frac{\Lambda_\pi}{\Lambda_\pi - \Lambda_N} (e^{-X/\Lambda_\pi} - e^{-X/\Lambda_N}) \sim E^{-\gamma} \quad \Lambda_\pi > \Lambda_N$
 $\Rightarrow$ same shape as primary spectrum

b) $E \ll \varepsilon_\pi \Rightarrow$ decay dominates $\Pi(E, X) = N(E, 0) \frac{Z_{N\pi}}{\lambda_N} e^{-X/\lambda_N} \frac{X E \cos \theta}{\varepsilon_\pi} \sim E^{-\gamma+1}$

$\Rightarrow$ harder at low energies $E < \varepsilon_\pi$

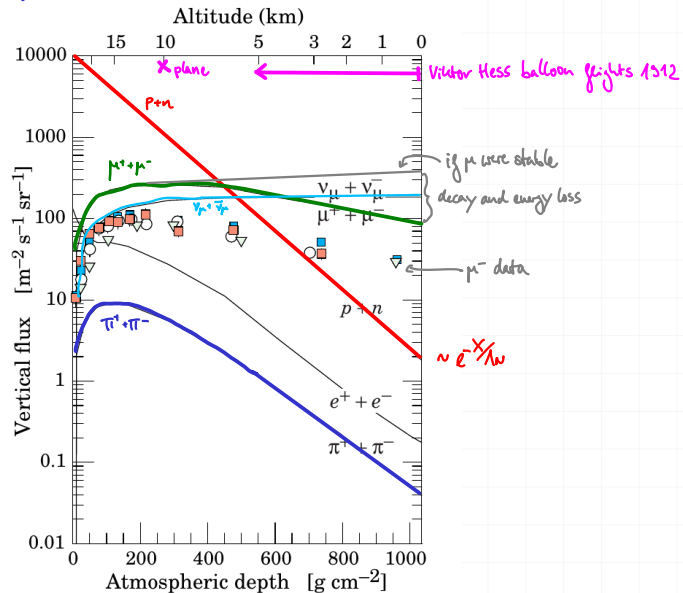
• decaying mesons $\Rightarrow$ source term for $i = \mu$ and ν_μ

$$P_{i\pi} = \int_{E_{\min}}^{E_{\max}} \frac{dn_{i\pi}(E, E')}{dE} \frac{\Pi(E', X)}{d_\pi(E')} dE'$$

($E_{\min/\max}$, $\frac{dn}{dE}$ from kinematics)



Inclusive Particle Fluxes in the Atmosphere



Muon Spectrum at Sea Level

Muons from $\pi^\pm$ and $K^\pm$

$\Rightarrow$ important prediction of atmospheric ν $\rightarrow$ atmospheric ν -oscillation
 $\rightarrow$ background for astrophysical ν s
 $\Rightarrow$ see summer

$$\frac{dN_\mu}{dE_\mu} \simeq S_\mu(E_\mu) \frac{N_0(E_\mu)}{1 - Z_{NN}} \left\{ A_{\pi\mu} \frac{1}{1 + B_{\pi\mu} \cos \theta E_\mu / \epsilon_\pi} + 0.635 A_{K\mu} \frac{1}{1 + B_{K\mu} \cos \theta E_\mu / \epsilon_K} \right\} \quad (6.36)$$

(Similar for ν , see Eq.6.29 in CRPP)

nucleon spectrum (top of atmosphere)

branching ratio $K \rightarrow \mu + \nu$

zenith

$\rightarrow 115 \text{ GeV}$

$\rightarrow 850 \text{ GeV}$

$A_{\pi\mu} \sim Z_{N\pi}$ (meson production)

$B_{\pi\mu}$: kinematics etc.

higher energy fraction to ν in $K \rightarrow \mu + \nu$

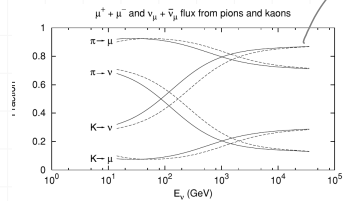
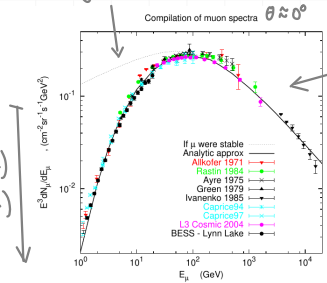


Fig. 1. Fraction of muons and muon neutrinos from pion decay and from kaon decay vs. neutrino energy. Solid lines for vertical, dashed lines for zenith angle 60° .

$S_\mu(E_\mu)$
 (decay + dE_μ/dX)

follows $\sim$ primary spectrum



softening due to $\lambda_i < d_i$

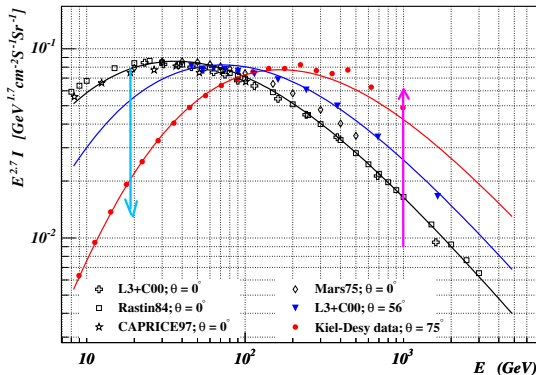
Figure 8: Summary of measurements of the vertical muon intensity at the ground. The solid line shows an analytic calculation [62]. The dotted line shows the spectrum in the absence of decay and energy loss, or equivalently the muon production spectrum integrated over the atmosphere.

Muon Spectrum at Sea Level

Muons from $\pi^\pm$ and $K^\pm$

$$\frac{dN_\mu}{dE_\mu} \simeq S_\mu(E_\mu) \frac{N_0(E_\mu)}{1 - Z_{NN}} \left\{ \mathcal{A}_{\pi\mu} \frac{1}{1 + \mathcal{B}_{\pi\mu} \cos \theta E_\mu / \epsilon_\pi} + 0.635 \mathcal{A}_{K\mu} \frac{1}{1 + \mathcal{B}_{K\mu} \cos \theta E_\mu / \epsilon_K} \right\} \quad (6.36)$$

larger path length
from production
to ground
 $\Rightarrow$ more muon
decay



— 0°
— 56°
— 75°

$d\tau/\mu < \lambda_{\pi/K}$ at
large zenith angle

(interaction at higher altitude
 $\Rightarrow$ thinner atmosphere
 $\Rightarrow$ longer path (in cm) to
reach grammage λ)

Atmospheric Neutrinos

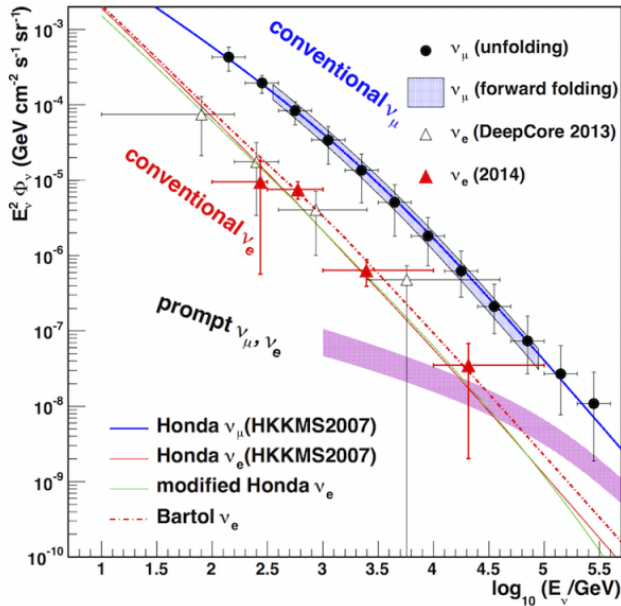
"Conventional ν_μ ": $\pi^\pm, K^\pm$

"Conventional ν_e " mostly

$K_L^0 \rightarrow \pi^\pm e^\mp \nu_e$ ($\approx 41\%$)

"prompt": charmed mesons $D^\pm, D^0$

$\rightarrow d \approx 0$



IceCube data

Electromagnetic Cascades

Bruno Rossi + Kenneth Greisen Rev. Mod. Phys. 1954

photon flux $\gamma(\omega, t)$ and e^+e^- flux $\pi(E, t)$, $t = X/X_0$

$$\frac{d\gamma}{dt} = -\frac{\gamma(\omega, t)}{\lambda_{\text{pair}}} + \int_{\omega}^{\infty} \pi(E', t) \frac{dne \rightarrow \gamma}{d\omega dt} dE'$$

$$\frac{d\pi}{dt} = -\frac{\pi(E, t)}{\lambda_{\text{brems}}} + \int_E^{\infty} \pi(E', t) \frac{dne \rightarrow e}{dE dt} dE' + 2 \int_E^{\infty} \gamma(\omega', t) \frac{dny \rightarrow e}{dE dt} d\omega'$$

"approximation A".

- $E \gg E_c \rightarrow$ neglect ionization losses
- scaling of bremsstrahlung + pair production

$$E' \frac{dne \rightarrow \gamma}{d\omega dt} = f\left(\frac{\omega}{E'}\right) \quad \text{Brems.}$$

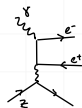
$$\omega' \frac{dny \rightarrow e}{dE dt} = g\left(\frac{E}{\omega'}\right) \quad \text{pair}$$

$$E' \frac{dne \rightarrow e}{dE dt} = f\left(1 - \frac{E}{E'}\right) \quad \text{Brems.}$$

interactions with nuclei of material (Z)



bremsstrahlung



pair production

boundary condition: photon induced shower

$$\pi(E, 0) = \delta(E, E_0)$$

Solution: "Greisen Profile"

$$N(t) = \frac{0.31}{\sqrt{\ln(E_0/E_c)}} e^{t(1 - \frac{3}{2})} \quad E_{\gamma, e^+e^-} > 20 \text{ MeV}$$

shower age s : $S = \frac{3t}{t + 2t_{\text{max}}}$

shower maximum at $t_{\text{max}} = \ln \frac{E_0}{E_c}$

$$N(t_{\text{max}}) = \frac{0.31}{\sqrt{\ln(E_0/E_c)}} \frac{E_0}{E_c}$$

γ -induced air shower profile

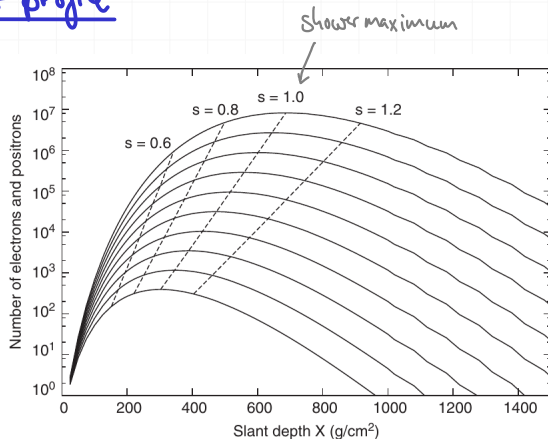


Figure 15.2 Shower size as a function of slant depth for photon-initiated showers in half-decade intervals of primary energy from 316 GeV (lowest curve) to 10^7 GeV (highest curve). The dashed lines trace the locus of size at specific shower ages across the same range of energies.