

Computational Photonics

Finite-Difference Time-Domain

Boundary conditions

- most efficient boundary conditions are the PML
- more precise: material property, not an actual boundary
- published by Berenger in 1994 with reflections about 3000 times less than with 2nd order Mur boundary
- basic idea: constructing a media which absorbs light and at whose interfaces with the region of interest no reflections taking place

➡ impedance matching of the medium with the surrounding

- impedance of a medium with electric and magnetic conductivity

$$\eta = \sqrt{\frac{\tilde{\mu}}{\tilde{\epsilon}}} = \sqrt{\frac{\mu' - i\mu''}{\epsilon' - i\epsilon''}}$$

$$\epsilon'' = \frac{\sigma_e}{\omega} \quad (\text{electrical conductivity})$$
$$\mu'' = \frac{\sigma_m}{\omega} \quad (\text{magnetic conductivity})$$

Boundary conditions

- let the loss-less region being region 1, characterised by ϵ and μ

→ impedance: $\eta = \sqrt{\frac{\mu}{\epsilon}}$

- if the condition $\frac{\sigma_e}{\epsilon} = \frac{\sigma_m}{\mu}$ holds, the impedances are equal $\eta_l = \eta$

- plane waves propagating in each of the regions are characterised by their propagation constants, given as (by choosing $\epsilon = \epsilon'$ and $\mu = \mu'$)

$$k = \omega \sqrt{\epsilon \mu} \qquad k = \omega \sqrt{\epsilon \mu} \sqrt{\left(1 - i \frac{\sigma_e}{\omega \epsilon}\right) \left(1 - i \frac{\sigma_m}{\omega \mu}\right)}$$

- phase velocity given by

$$k = \omega \sqrt{\epsilon \mu} + i \eta \sigma_e$$

the same as in free space but light is additionally absorbed

Boundary conditions

- reflections coefficients of the interface are given generally by

$$R_{\perp} = \frac{E_{\perp}^r}{E_{\perp}^i} = \frac{\eta_2 \cos \theta_i - \eta_1 \cos \theta_t}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t} \quad R_{\parallel} = \frac{E_{\parallel}^r}{E_{\parallel}^i} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

(with $n_1 \sin \theta_i = n_2 \sin \theta_t$)

➡ reflection is 0, only if the angle of incidence = the angle of transmittance

➡ holds for interface between 2 media sustaining modes with same phase velocity

- following derivation exemplarily for 2D FDTD in TE polarisation

$$\begin{aligned} \epsilon \frac{\partial E_x}{\partial t} + \sigma_e E_x &= \frac{\partial H_z}{\partial y} \\ \epsilon \frac{\partial E_y}{\partial t} + \sigma_e E_y &= -\frac{\partial H_z}{\partial x} \end{aligned}$$

$$\mu \frac{\partial H_z}{\partial t} + \sigma_m H_z = - \left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right)$$

➡ absorb light only in one direction, fields shall propagate parallel to surface 4

Boundary conditions

- Berenger's idea: splitting the H field into a x and y component (x-derivative of the E field drives the H_{zx} component and vice versa)

$$H_z = H_{zx} + H_{zy}$$

- introduction of an anisotropy of all properties

$$\mu \frac{\partial H_{zx}}{\partial t} + \sigma_{mx} H_{zx} = -\frac{\partial E_y}{\partial x} \quad \epsilon \frac{\partial E_x}{\partial t} + \sigma_{ey} E_x = \frac{\partial (H_{zx} + H_{zy})}{\partial y}$$

$$\mu \frac{\partial H_{zy}}{\partial t} + \sigma_{my} H_{zy} = \frac{\partial E_x}{\partial y} \quad \epsilon \frac{\partial E_y}{\partial t} + \sigma_{ex} E_y = -\frac{\partial (H_{zx} + H_{zy})}{\partial x}$$

➡ can be shown, that the impedance of the Berenger medium equals the impedance of the free space, regardless of the angle of propagation

Boundary conditions

- problem that denies a sudden change in the material properties: electric and magnetic field are evaluated at different spatial positions (staggered grid)
- asymmetric absorption
- choosing appropriate absorption profile for σ_{ex} and σ_{mx}

→ polynomial scaling

$$\sigma_{ex} = \left(\frac{x}{d}\right)^m \sigma_{emax}$$

practical
considerations

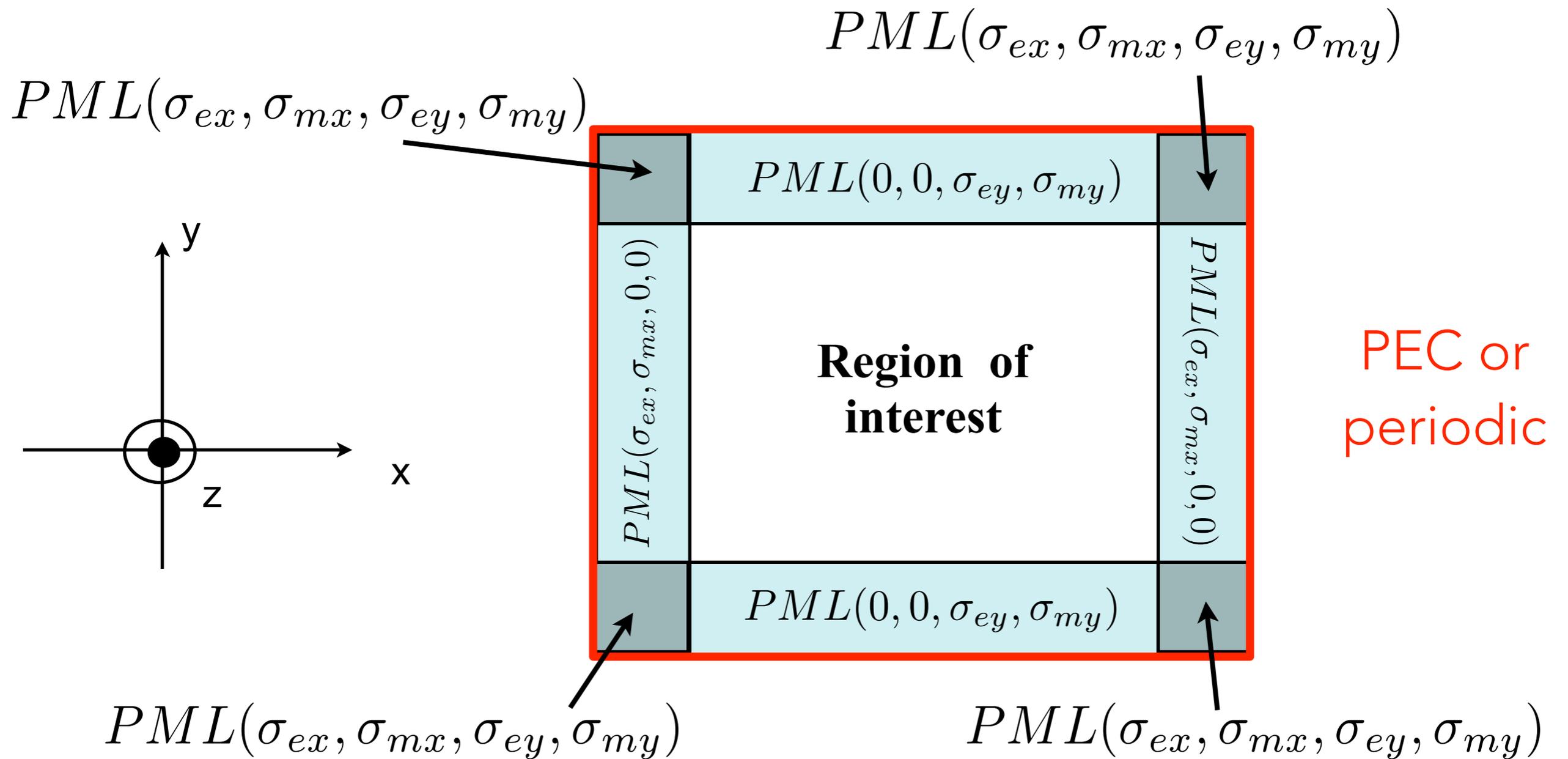
$$2 \ll m \ll 6$$

$$d = 10\Delta x$$

→ $R = 10^{-16}$

Boundary conditions

- important: waves propagating along the y-axis are not absorbed in the x-boundary



Boundary conditions

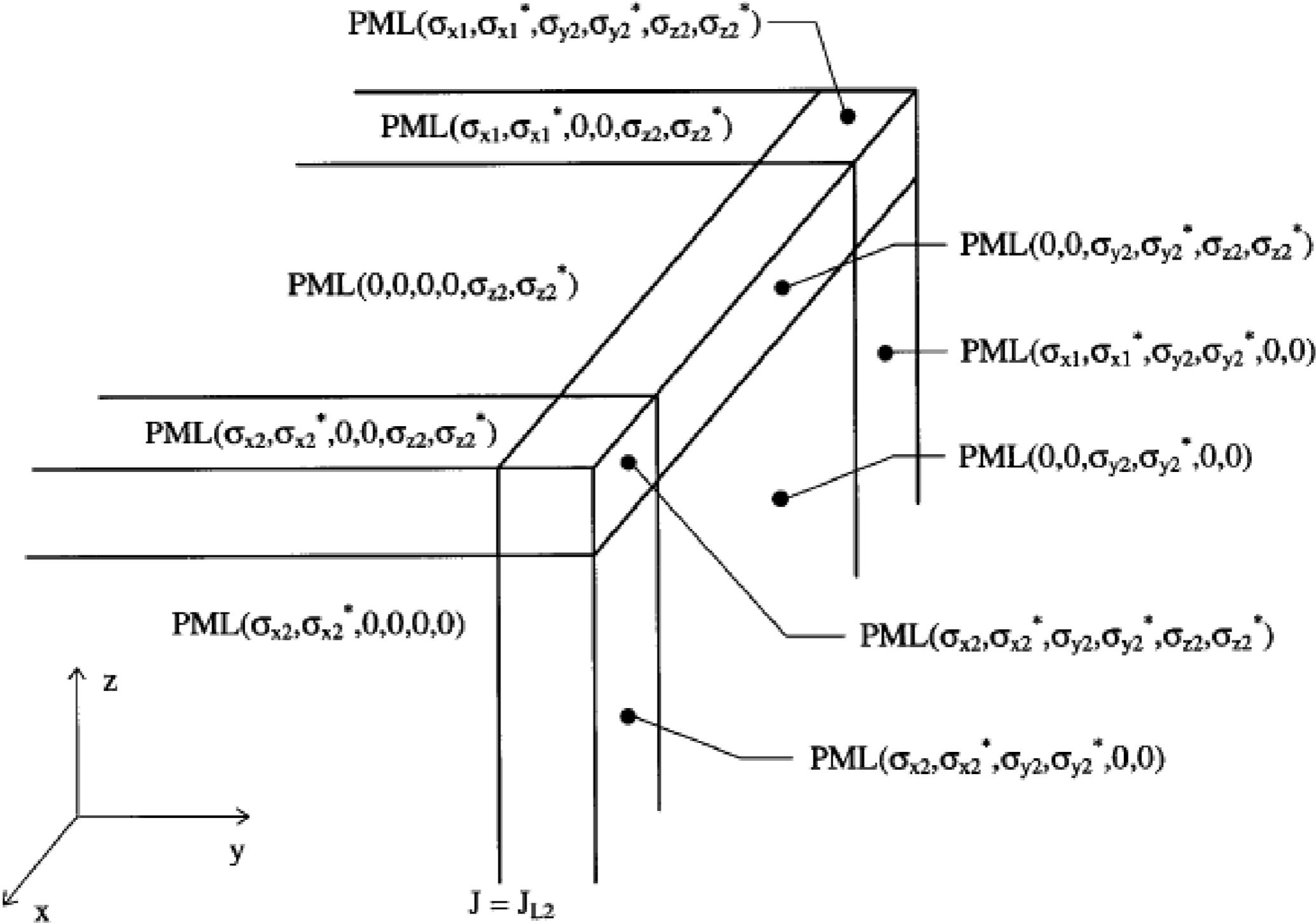
- splitting the field in the boundaries
- adding the electrical and magnetic conductivity to the equations as material parameters
- only the components normal to the boundary are absorbed, adjusting the proper absorption profile for each component
- electric and magnetic fields are evaluated a half discretisation step apart, hence the absorption profile is evaluated likewise at different spatial coordinates

Boundary conditions

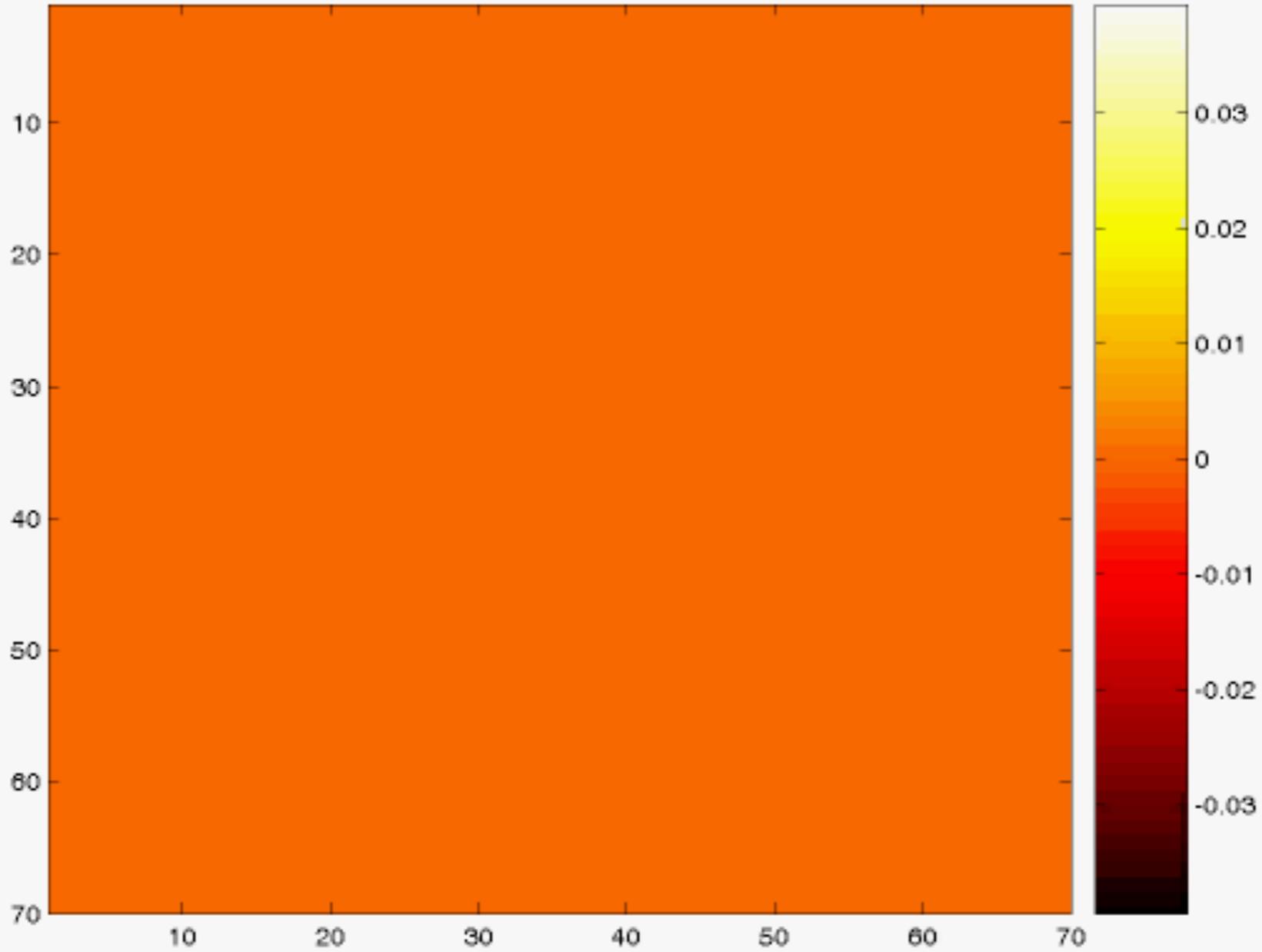
→ extensions for TM polarisation and 3D are straight forward

$$\begin{aligned}
 \left(\epsilon \frac{\partial}{\partial t} + \sigma_{ey} \right) E_{xy} &= \frac{\partial}{\partial y} (H_{zx} + H_{zy}) & \left(\epsilon \frac{\partial}{\partial t} + \sigma_{my} \right) H_{xy} &= -\frac{\partial}{\partial y} (E_{zx} + E_{zy}) \\
 \left(\epsilon \frac{\partial}{\partial t} + \sigma_{ez} \right) E_{xz} &= -\frac{\partial}{\partial z} (H_{yx} + H_{yz}) & \left(\epsilon \frac{\partial}{\partial t} + \sigma_{mz} \right) H_{xz} &= \frac{\partial}{\partial z} (E_{yx} + E_{yz}) \\
 \left(\epsilon \frac{\partial}{\partial t} + \sigma_{ex} \right) E_{yx} &= -\frac{\partial}{\partial x} (H_{zx} + H_{zy}) & \left(\epsilon \frac{\partial}{\partial t} + \sigma_{mx} \right) H_{yx} &= \frac{\partial}{\partial x} (E_{zx} + E_{zy}) \\
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 \left(\epsilon \frac{\partial}{\partial t} + \sigma_{ey} \right) E_{zy} &= -\frac{\partial}{\partial y} (H_{xy} + H_{xz}) & \left(\epsilon \frac{\partial}{\partial t} + \sigma_{my} \right) H_{zy} &= \frac{\partial}{\partial y} (E_{xy} + E_{xz})
 \end{aligned}$$

Boundary conditions

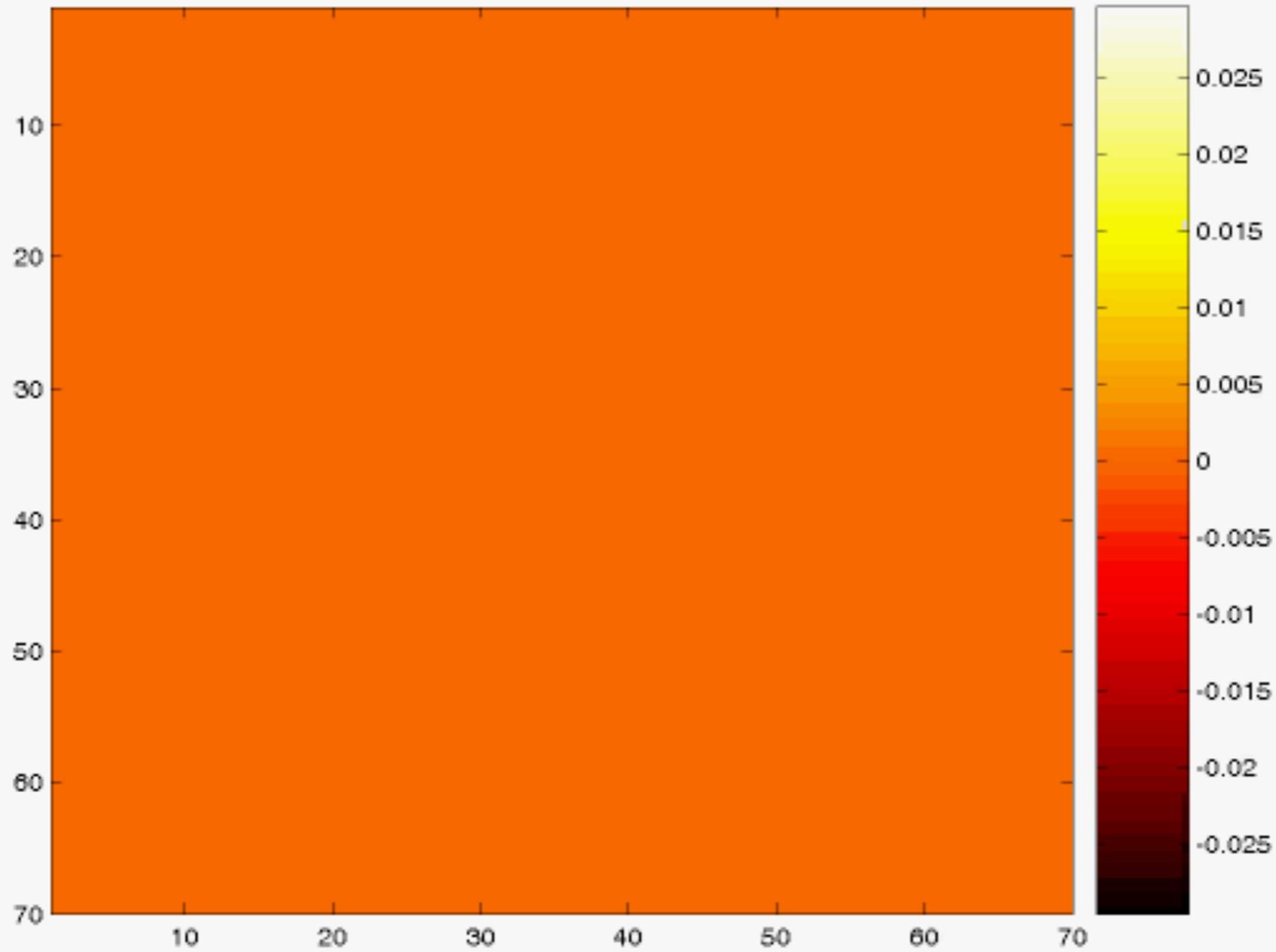


Boundary conditions



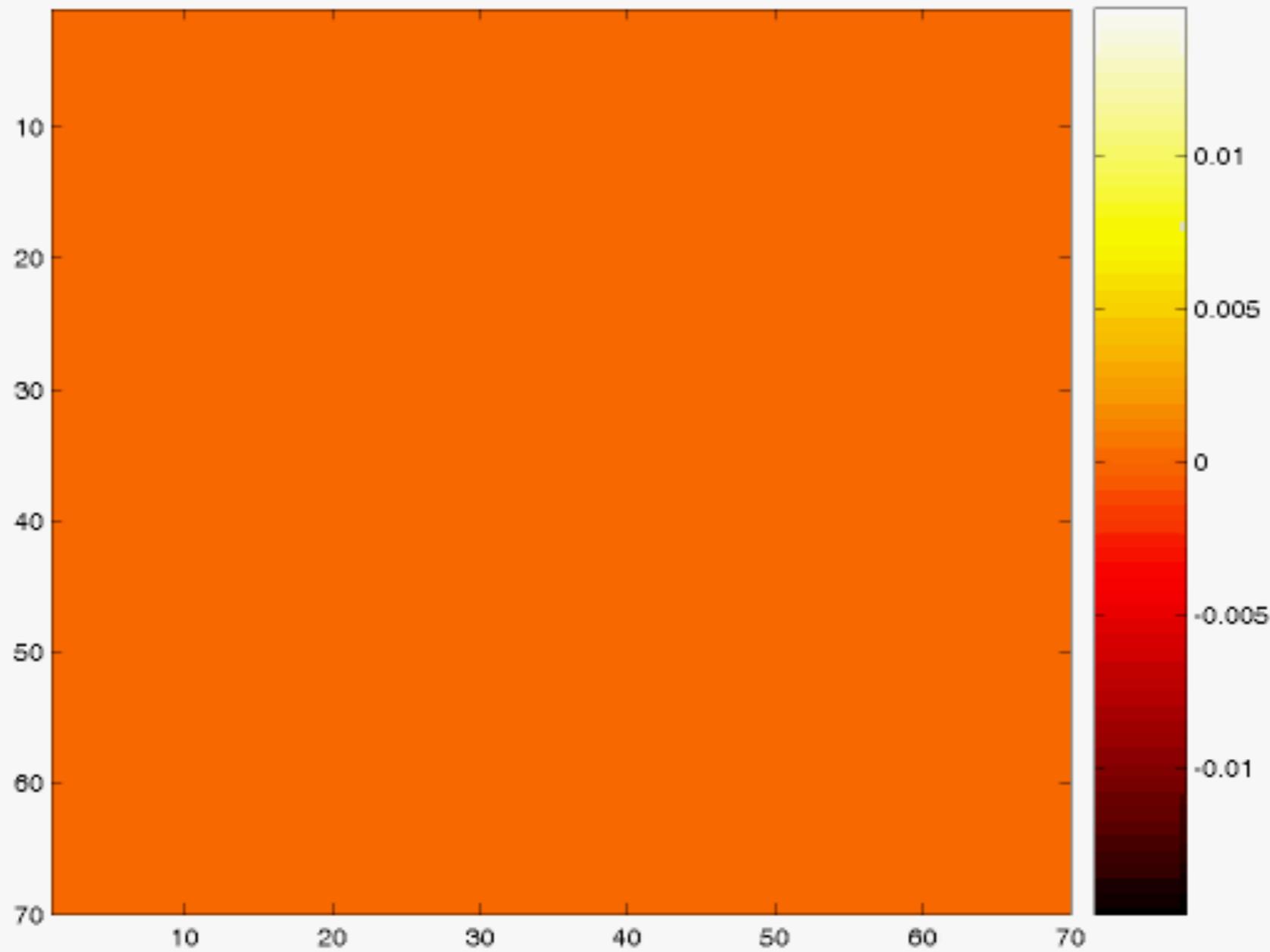
perfectly
conducting
boundary

Boundary conditions



2nd order
Mur
boundary

Boundary conditions



PML layers
with

$$m = 3$$

$$d = 16\Delta x$$

Computational Photonics

Finite-Difference Time-Domain