

①

a) $\vec{\tau}_e = \Psi_e^+ \vec{\tau} \Psi_e^-$ with $\Psi_e = \begin{pmatrix} C_{e\uparrow} \\ C_{-e\downarrow} \end{pmatrix}$, $\Psi_e^+ = (C_{e\uparrow}^+, C_{-e\downarrow})$, $\vec{\tau} = (\tau^x, \tau^y, \tau^z)$. ~~bla~~

$$\Rightarrow \tau_e^x = C_{e\uparrow}^+ C_{-e\downarrow}^+ + C_{-e\downarrow} C_{e\uparrow}$$

$$\tau_e^y = -i(C_{e\uparrow}^+ C_{-e\downarrow}^+ - C_{-e\downarrow} C_{e\uparrow})$$

$$\tau_e^z = C_{e\uparrow}^+ C_{e\uparrow} - C_{-e\downarrow} C_{-e\downarrow} = M_{e\uparrow} + \cancel{M_{-e\downarrow}} - 1 \text{ with } M_{e\uparrow} = C_{e\uparrow}^+ C_{e\uparrow}$$

Show:

$$[\tau_e^\alpha, \tau_e^\beta] = 2i \epsilon^{\alpha\beta\gamma} \tau_e^\gamma.$$

i) Show $[\tau_e^z, \tau_e^x] = 2i \tau_e^y$ on next page. Use

$$[AB, C] = A[B, C] - [A, C]B.$$

ii) $[\tau_e^x, \tau_e^y] = [C_{e\uparrow}^+ C_{e\uparrow} + C_{-e\downarrow} C_{-e\downarrow}, -i(C_{e\uparrow}^+ C_{-e\downarrow}^+ - C_{-e\downarrow} C_{e\uparrow})] =$
 $= -i \left\{ C_{e\uparrow}^+ [C_{e\uparrow}, C_{e\uparrow}^+ C_{-e\downarrow}^+] - \cancel{[C_{e\uparrow}^+, C_{-e\downarrow}^+ C_{-e\downarrow}]} C_{e\uparrow} + C_{-e\downarrow}^+ [\cancel{C_{-e\downarrow}}, C_{-e\downarrow}^+ C_{e\uparrow}] \right.$

~~see next
page~~

$$\left. - [C_{-e\downarrow}^+, C_{-e\downarrow} C_{e\uparrow}] C_{-e\downarrow} \right\} =$$

$$= -i \left\{ C_{e\uparrow}^+ C_{-e\downarrow}^+ + C_{-e\downarrow} C_{e\uparrow} - C_{-e\downarrow}^+ C_{e\uparrow}^+ - C_{e\uparrow} C_{-e\downarrow} \right\} = -2i \underline{\tau_e^x}. \quad \square$$

Number spinor: $\Psi_h = (c_{\alpha\uparrow}, c_{-\alpha\downarrow}^+)^T$

$$\Psi_h^+ = (c_{\alpha\uparrow}^+, c_{-\alpha\downarrow})$$

J so spin operators: $\bar{\tau}_h = \Psi_h^+ \bar{\tau} \Psi_h \Rightarrow \bar{\tau}_h^z = (c_{\alpha\uparrow}^+, c_{-\alpha\downarrow}) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} c_{\alpha\uparrow} \\ c_{-\alpha\downarrow}^+ \end{pmatrix} =$

$$\Rightarrow \boxed{\begin{aligned} \bar{\tau}_h^z &= c_{\alpha\uparrow}^+ c_{\alpha\uparrow} + c_{-\alpha\downarrow}^+ c_{-\alpha\downarrow} \cancel{- 1} \\ \bar{\tau}_h^x &= c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+ + c_{-\alpha\downarrow} c_{\alpha\uparrow} \\ \bar{\tau}_h^y &= -i(c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+ - c_{-\alpha\downarrow} c_{\alpha\uparrow}) \end{aligned}}$$

$$[\bar{\tau}_h^z, \bar{\tau}_h^x] =$$

$$= [c_{\alpha\uparrow}^+ c_{\alpha\uparrow} + c_{-\alpha\downarrow}^+ c_{-\alpha\downarrow}, c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+ + c_{-\alpha\downarrow} c_{\alpha\uparrow}] =$$

$$= c_{\alpha\uparrow}^+ [c_{\alpha\uparrow}, c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+] + [c_{\alpha\uparrow}^+, c_{-\alpha\downarrow}^+ c_{\alpha\uparrow}] c_{\alpha\uparrow} + c_{-\alpha\downarrow}^+ [c_{-\alpha\downarrow}, c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}]$$

$$+ [c_{-\alpha\downarrow}^+, c_{-\alpha\downarrow}^+ c_{\alpha\uparrow}] c_{-\alpha\downarrow} =$$

$$[AB, C] = \{A\{B, C\} - \{A, C\}B$$

$$= ABC - CAB = ABC + ACB - ACB - CAB = ABC - CAB \checkmark$$

$$\cancel{[ABC] = B[A, C]} \quad \cancel{[ABC] = BAC} \quad \cancel{[BAC] = BAC} \quad \cancel{[BAC] = BAC}$$

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$$\begin{aligned}
& -C_{h\uparrow}^+ (-1) C_{-h\downarrow}^+ - C_{-h\downarrow} C_{h\uparrow} - C_{-h\downarrow}^+ C_{h\uparrow}^+ - (-1) C_{h\uparrow} C_{-h\downarrow} = \\
& = C_{h\uparrow}^+ C_{-h\downarrow}^+ + C_{h\uparrow}^+ C_{-h\downarrow} - C_{-h\downarrow} C_{h\uparrow} - C_{-h\downarrow}^+ C_{h\uparrow}^+ = \\
& = 2 (C_{h\uparrow}^+ C_{-h\downarrow}^+ - C_{-h\downarrow} C_{h\uparrow}) = \underline{2i \tau_a^y} = [\tau_a^z, \tau_a^x]. \quad \checkmark
\end{aligned}$$

(see p. 4. as well for $[\tau_a^x, \tau_a^y] = 2i \tau_a^z$)

$$[\tau_a^\alpha, \tau_b^\beta] = 2i \epsilon^{\alpha\beta\gamma} \tau_a^\gamma$$

⑤ BCS-Ham.:

$$\begin{aligned}
H_{BCS} &= \sum_{h,\sigma} \epsilon_{h\sigma} C_{h\sigma}^+ C_{h\sigma} + \sum_h [\Delta^* C_{-h\downarrow} \cancel{C_{h\uparrow}} + C_{h\uparrow}^+ C_{-h\downarrow}^* \Delta] + \sum_a \epsilon_a \\
&\quad + \frac{V}{g_0} |\Delta|^2 \cancel{\psi_h} = \sum_h \psi_h^+ (\vec{l}_{h\sigma} \cdot \vec{\tau}) \psi_h + \frac{V}{g_0} |\Delta|^2 \\
&= - \sum_h \vec{B}_h \cdot \vec{\tau}_h \stackrel{+ \text{const.}}{;} ; \vec{B}_h = -\vec{l}_h = -(\Delta_1, \Delta_2, \epsilon_h) \stackrel{-3-}{=} |B_h| \hat{m}_h . \quad |B_h| = \sqrt{\Delta_1^2 + \Delta_2^2 + \epsilon_h^2} = \\
&\quad = E_h .
\end{aligned}$$

$$\begin{aligned}
 \text{i)} [\tau_{\alpha}^x, \tau_{\alpha}^y] &= [c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+ + c_{-\alpha\downarrow} c_{\alpha\uparrow}, -i(c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+ - c_{-\alpha\downarrow} c_{\alpha\uparrow})] = \\
 &= -i \left\{ -c_{\alpha\uparrow}^+ [c_{-\alpha\downarrow}^+, c_{-\alpha\downarrow} c_{\alpha\uparrow}] - [c_{\alpha\uparrow}^+, c_{-\alpha\downarrow} c_{\alpha\uparrow}] c_{-\alpha\downarrow}^+ \right. \\
 &\quad \left. + c_{-\alpha\downarrow} [c_{\alpha\uparrow}, c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+] + [c_{-\alpha\downarrow}, c_{\alpha\uparrow}^+ c_{-\alpha\downarrow}^+] c_{\alpha\uparrow} \right\} = \\
 &= -i \left\{ -c_{\alpha\uparrow}^+ c_{\alpha\uparrow} + c_{-\alpha\downarrow}^+ c_{-\alpha\downarrow} + c_{-\alpha\downarrow}^+ c_{-\alpha\downarrow} - c_{\alpha\uparrow}^+ c_{\alpha\uparrow} \right\} = \\
 &= i \{ 2m_{\alpha\uparrow} + 2m_{-\alpha\downarrow} - 12 \} = \underline{\underline{2i\tau_{\alpha}^z}} \cdot \square
 \end{aligned}$$

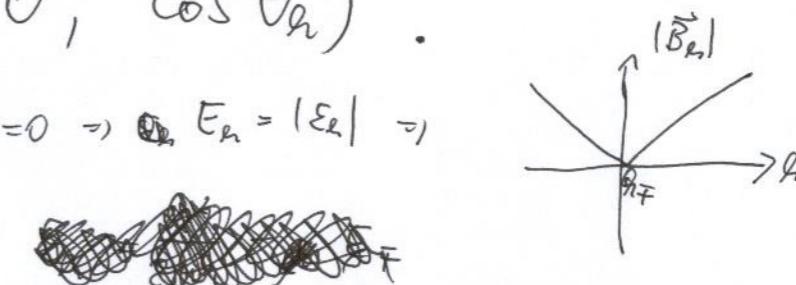
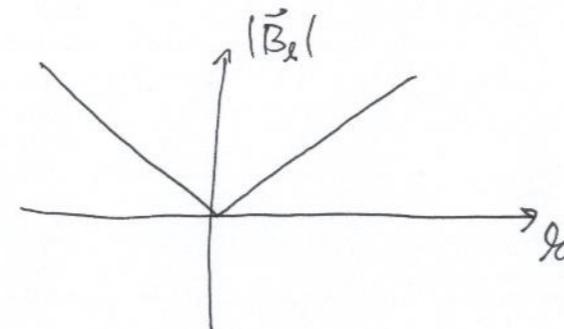
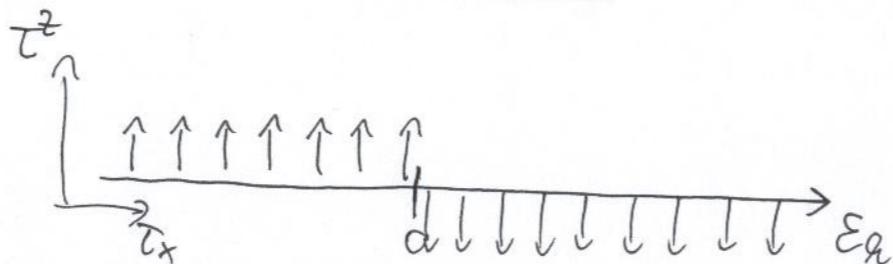
Magnitude: $|\vec{B}_n| = E_n = \sqrt{\Delta_1^2 + \Delta_2^2 + \varepsilon_n^2}$

Direction: $\hat{m}_n = \frac{\vec{B}_n}{|\vec{B}_n|} = -\left(\frac{\Delta_1}{E_n}, \frac{\Delta_2}{E_n}, \frac{\varepsilon_n}{E_n}\right)$.

c) Assume that gap is real $\Rightarrow \Delta_2 = 0$.

$$\Rightarrow \hat{m}_n = -\left(\frac{\Delta_1}{E_n}, 0, \frac{\varepsilon_n}{E_n}\right) = -\left(\sin \theta_n, 0, \cos \theta_n\right).$$

i) Spin configuration in normal state: (not asked for): $\Delta_1 = 0 \Rightarrow E_n = |\varepsilon_n| \Rightarrow$



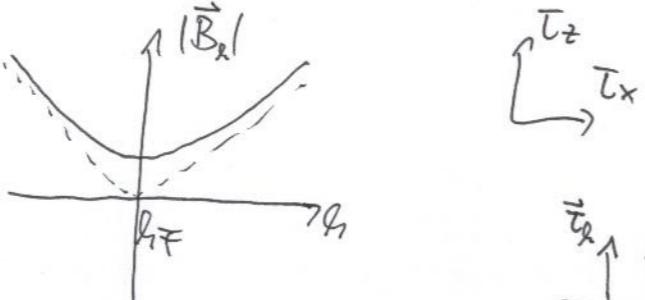
$$\varepsilon_F = v_F(k - k_F) \quad (\text{linearized around FS}).$$

$$H = - \sum_n \vec{B}_n \cdot \vec{t}_n = - \sum_n |\vec{B}_n| \hat{m}_n \cdot \vec{t}_n.$$

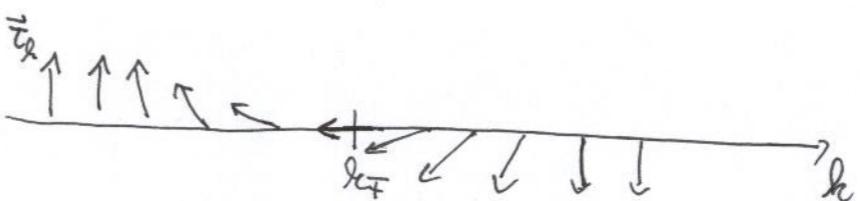
\vec{t}_n points ~~along~~ ^{along} \vec{m}_n in GS.

ii) Spin configuration in sc. state :

$$\hat{m}_h = \left(-\frac{\Delta_1}{E_h}, 0, \frac{E_h}{E_h} \right) = \frac{-1}{E_h} (\Delta_1, 0, V_F(h-h_F)) = -(\sin \theta_h, 0, \cos \theta_h).$$



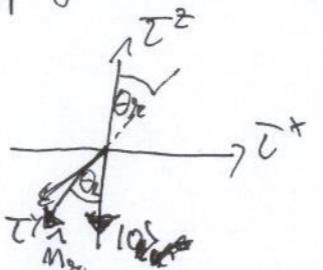
$$\sin \theta_h = \frac{\Delta_1}{E_h}; \quad \cos \theta_h = \frac{E_h}{E_h}.$$



$$H = - \sum_h \vec{B}_h \cdot \vec{\tau}_h.$$

(d) Isospin ~~state~~ state that corresponds to electron vacuum reads : $| \uparrow \downarrow \uparrow \downarrow \dots \rangle =$ ~~all~~ spins pointing down $\forall h$ (i.e. along $-z$) .

Now apply rotation operator $\exp \left[-\frac{i}{2} \theta_h \vec{\tau}_h^y \right] |0\rangle = |\theta_h\rangle$



that rotates each spin around $\vec{\tau}_h^y$ axis to be in the configuration to be aligned with direction of \vec{B}_h (or \hat{m}_h). This is the BCS ground state.

$$\begin{aligned}
 |\theta_n\rangle &= \exp\left[-\frac{i}{2}\theta_n \psi_n^+ \tau^+ \psi_n\right] |1\rangle = \left(\cos\left(\frac{\theta_n}{2}\right) - i \sin\left(\frac{\theta_n}{2}\right) \psi_n^+ \tau^+ \psi_n\right) |1\rangle = \\
 &= \cos\left(\frac{\theta_n}{2}\right) |1\rangle - \sin\left(\frac{\theta_n}{2}\right) \underbrace{\psi_{n\uparrow}^+ \psi_{n\downarrow}^+}_{=|\uparrow\rangle} |1\rangle
 \end{aligned}$$

BCS ground state.

$$|\text{BCS}\rangle = \prod_n |\theta_n\rangle = \prod_n \left(\underbrace{\cos\left(\frac{\theta_n}{2}\right)}_{=u_n} + \underbrace{\sin\left(\frac{\theta_n}{2}\right)}_{=v_n} \psi_{n\downarrow}^+ \psi_{n\uparrow}^+ \right) |0\rangle.$$

$$u_n = \cos \frac{\theta_n}{2} = \sqrt{\frac{1}{2}(1 + \cos \theta_n)} = \sqrt{\frac{1}{2}\left(1 + \frac{\epsilon_n}{E_n}\right)}.$$

$$v_n = \sin \frac{\theta_n}{2} = \sqrt{\frac{1}{2}(1 - \cos \theta_n)} = \sqrt{\frac{1}{2}\left(1 - \frac{\epsilon_n}{E_n}\right)}.$$

$$\begin{aligned}
 \textcircled{1} \quad & \cos \theta = \cos^2 \frac{\theta}{2} - \sin^2 \frac{\theta}{2} = \\
 & = \cos^2 \frac{\theta}{2} - 1 + \cos^2 \frac{\theta}{2} = \\
 & = -1 + 2 \cos^2 \frac{\theta}{2} \\
 \Leftrightarrow & \cos^2 \frac{\theta}{2} = \frac{1}{2}(1 + \cos \theta).
 \end{aligned}$$

$$\begin{aligned}
 \textcircled{2} \quad & \sin \theta = 1 - 2 \sin^2 \frac{\theta}{2} \\
 \Leftrightarrow & \sin^2 \frac{\theta}{2} = \frac{1}{2}(1 - \cos \theta).
 \end{aligned}$$

e) Find two eigenvectors with isospin parallel & anti-parallel to \vec{B}_e .

$$i) \hat{M}_e \cdot \vec{\tau} \begin{pmatrix} u_e \\ v_e \end{pmatrix} = \begin{pmatrix} -\cos \theta_e & -\sin \theta_e \\ -\sin \theta_e & \cos \theta_e \end{pmatrix} \begin{pmatrix} \cos \frac{\theta_e}{2} \\ \sin \frac{\theta_e}{2} \end{pmatrix} = \begin{pmatrix} -\cos \theta_e \cos \frac{\theta_e}{2} - \sin \theta_e \sin \frac{\theta_e}{2} \\ -\sin \theta_e \cos \frac{\theta_e}{2} + \cos \theta_e \sin \frac{\theta_e}{2} \end{pmatrix} =$$

$$\hat{M}_e = (\sin \theta_e, 0, \cos \theta_e) = - \begin{pmatrix} \cos \frac{\theta_e}{2} \\ \sin \frac{\theta_e}{2} \end{pmatrix}$$

The state $\begin{pmatrix} u_e \\ v_e \end{pmatrix}$ is aligned against anti-parallel to the field $\vec{B}_e \sim \hat{M}_e$.

It has energy $+E_h$.

$$\Rightarrow H = -\sum_e \vec{B}_e \cdot \vec{\tau}_e = -\sum_e E_h \hat{M}_e \cdot \vec{\tau}_e$$

\vec{B}_e

\vec{B}_e !!

$$\textcircled{1} \quad (-\cos^2 \frac{\theta_e}{2} + \sin^2 \frac{\theta_e}{2}) \cos \frac{\theta_e}{2} - 2 \sin^2 \frac{\theta_e}{2} \cos \frac{\theta_e}{2} = - \underbrace{(\cos^2 \frac{\theta_e}{2} + \sin^2 \frac{\theta_e}{2})}_{=1} \cos \frac{\theta_e}{2}$$

$$\textcircled{2} \quad (-2 \sin \frac{\theta_e}{2} \cos^2 \frac{\theta_e}{2} + \cos^2 \frac{\theta_e}{2} \sin \frac{\theta_e}{2} - \sin^2 \frac{\theta_e}{2} \sin \frac{\theta_e}{2}) = - \underbrace{(\cos^2 \frac{\theta_e}{2} + \sin^2 \frac{\theta_e}{2})}_{=1} \sin \frac{\theta_e}{2}$$

$$-\vec{B}_e \cdot \vec{\tau}_e \begin{pmatrix} u_e \\ v_e \end{pmatrix} = +E_h \begin{pmatrix} u_e \\ v_e \end{pmatrix}$$

$$ii) \text{ the orthonormal state } \begin{pmatrix} -v_e^* \\ u_e^* \end{pmatrix} = \begin{pmatrix} -v_e \\ u_e \end{pmatrix} \begin{pmatrix} u_e \\ v_e \end{pmatrix} = 0.$$

(sesquilinear product $v_2^* \cdot v_1 = 0$)

Can also be explicitly checked: $\hat{M}_e \cdot \vec{\tau} \begin{pmatrix} -v_e^* \\ u_e^* \end{pmatrix} = + \begin{pmatrix} -v_e^* \\ u_e^* \end{pmatrix}$.

$$\rightarrow \text{has energy } -\vec{B}_e \cdot \vec{\tau}_e \begin{pmatrix} -v_e^* \\ u_e^* \end{pmatrix} = -E_h \begin{pmatrix} -v_e^* \\ u_e^* \end{pmatrix}$$

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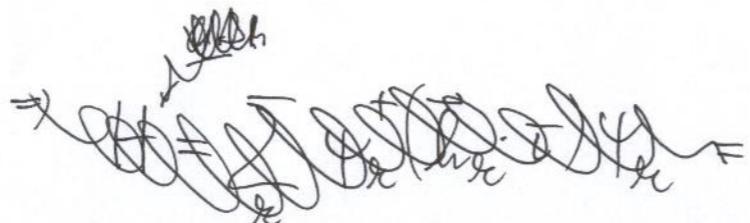
Projecting Nambu spinors on these states yields

$$a_{\alpha\uparrow}^+ := \Psi_e^+ \cdot \begin{pmatrix} u_e \\ v_e \end{pmatrix} = (c_{e\uparrow}^+, c_{-e\downarrow}) \begin{pmatrix} u_e \\ v_e \end{pmatrix} = u_e c_{e\uparrow}^+ + v_e c_{-e\downarrow}^+.$$

$$a_{-e\downarrow} := \Psi_e^+ \cdot \begin{pmatrix} -v_e^* \\ u_e^* \end{pmatrix} = -v_e^* c_{e\uparrow}^+ + u_e^* c_{-e\downarrow}^+.$$

or

$$\alpha_e^+ = (a_{e\uparrow}^+, a_{-e\downarrow}) = \Psi_e^+ \begin{pmatrix} u_e & -v_e^* \\ v_e & u_e^* \end{pmatrix} = \Psi_e^+ U_e$$



$$= -U_e \vec{\tau}^2$$

$$\begin{aligned} \Rightarrow H - \sum_h \frac{(\Delta_h)^2 - \sum E_h}{V_0} \sum_e \Psi_e^+ (-E_e) \hat{n}_e \cdot \vec{\tau}^2 \Psi_e &= \sum_h \underbrace{\Psi_e^+}_{=a_e^+} \underbrace{U_e}_{=a_e^+} \underbrace{U_e^+}_{=+E_e \vec{\tau}^2} \underbrace{(-E_e)}_{=+E_e \vec{\tau}^2} \underbrace{\hat{n}_e}_{=a_e} \underbrace{\vec{\tau}^2}_{=a_e} = \\ = - \sum_e \vec{B}_e \cdot \vec{\tau}_e &= - \sum_e E_e \hat{n}_e \cdot \vec{\tau}_e \quad \boxed{\sum_h a_e^+ E_e \vec{\tau}^2 a_h}, \end{aligned}$$

Explains.

$$H - \sum_h \left(\frac{|\Delta_{eh}|^2 + \varepsilon_{eh}}{V_0} \right) \sum_h (E_{eh} a_{eh\uparrow}^\dagger a_{eh\uparrow} - E_{eh} a_{eh\downarrow}^\dagger a_{eh\downarrow}) =$$

$$= \sum_h E_{eh} (a_{eh\uparrow}^\dagger a_{eh\uparrow} + a_{eh\downarrow}^\dagger a_{eh\downarrow}) - \sum_h E_{eh}$$

$$= \sum_{eh} E_{eh} (a_{eh\uparrow}^\dagger a_{eh\uparrow} + \underbrace{a_{eh\downarrow}^\dagger a_{eh\downarrow}}_{\text{assume } E_{eh\downarrow} = E_{eh}}) - \sum_{eh} E_{eh}$$

→ Ground state energy: $E_g = - \sum_h E_{eh} \times \sum_{eh} (|\Delta_{eh}|^2 + \varepsilon_{eh})$ (all neg. energies filled)

Excitation spectrum: $E_{ex} = \sqrt{|\Delta_{eh}|^2 + \varepsilon_{eh}^2}$ (spin flips)

$$E_g = \sum_{eh} \left(\frac{|\Delta_{eh}|^2}{V_0} + \varepsilon_{eh} - E_{ee} \right)$$

