

## Einführung in Theoretische Teilchenphysik

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## Exercise Sheet 10

<u>Hand-in Deadline</u>: Friday 05.02.21, 14:00. <u>Discussion</u>: Tuesday 09.02.21, Thursday 11.02.21.

1. **[13 points] Higgs Production via** W **Fusion**: Higgs production via W fusion is one of the main Higgs production processes at electron-positron colliders. At leading order, this process is given by a single Feynman diagram:



The corresponding Feynman rules are



with the weak coupling constant  $g \simeq 0.65$  and the vacuum expectation value of the Higgs field  $v \simeq 246$  GeV.

(a) [2 points] Show the following properties of the chirality projection operators  $P_R = \frac{1 \pm \gamma^5}{2}$ :

$$\begin{pmatrix} P_L \\ R \end{pmatrix}^2 = P_L \\ R \end{pmatrix}, \qquad P_R + P_L = 1, \qquad P_R - P_L = \gamma^5,$$
$$\begin{pmatrix} P_L \\ R \end{pmatrix}^{\dagger} = P_L \\ R \end{pmatrix}, \qquad P_L \gamma^{\mu} = \gamma^{\mu} P_R \\ L \end{pmatrix}, \qquad \gamma^0 \begin{pmatrix} P_L \\ R \end{pmatrix} \gamma^0 = P_R \\ L \end{pmatrix}.$$

(b) [3 points] Consider first the following sub-diagram:

where the W boson is "amputated" in that the subdiagram is not connected to the rest of the diagram, i.e. the contraction of the W boson with its polarisation vector is omitted and instead the matrix element contains an open Lorentz index  $\mu$ .

- i. [1 point] Write down the corresponding matrix element  $\mathcal{M}_1^{\mu}$  and calculate  $q_{1\mu}\mathcal{M}_1^{\mu}$ , eliminating the explicit momentum dependence of the expression, assuming also that the neutrino has a mass  $m_{\nu}$ .
- ii. [2 points] What happens in the limit  $m_e = m_{\nu}$ ? Show that for  $m_e = m_{\nu} = 0$ ,  $q_{1\mu}\mathcal{M}_1^{\mu}$  vanishes.
- (c) [3 points] Considering only massless fermions,  $m_e = m_{\nu} = 0$ , show that one obtains, after spin summation, the following for the squared matrix element of the sub-diagram:

$$\sum_{s_1,s_3} |\mathcal{M}_1|^{2,\mu\nu} \equiv \sum_{s_1,s_3} \mathcal{M}_1^{\mu} \mathcal{M}_1^{\dagger,\nu} = g^2 \left( p_1^{\mu} p_3^{\nu} + p_3^{\mu} p_1^{\nu} - p_1 \cdot p_3 g^{\mu\nu} + i \epsilon^{\mu\nu\rho\sigma} p_{1,\rho} p_{3,\sigma} \right) \,.$$

(d) [5 points] Use this to calculate the spin-averaged squared matrix element of the whole Feynman diagram. You should obtain

$$\overline{\sum} |\mathcal{M}|^2 = \frac{g^8 v^2}{4} \frac{1}{(q_1^2 - M_W^2)^2 (q_2^2 - M_W^2)^2} \ p_1 \cdot p_4 \ p_2 \cdot p_3 \,.$$

*Hint:* Consider first the result of the "middle part" (W propagators, HWW vertex) separately. The limit  $\epsilon \to 0$  of the  $i\epsilon$  terms in the propagators can be taken immediately. Take care to contract the correct Lorentz indices.

2. [7 points] Weak Boson Decay: Consider a massive vector field  $Z^{\mu}$  (with mass M and momentum  $k = p_1 + p_2$ ) and a Dirac fermion field  $\Psi$ , interacting via

$$\mathcal{L}_{\text{int}} = Z^{\mu} \Psi (g_V - g_A \gamma_5) \Psi$$

The amplitude for the decay of a massive vector boson to two fermions is given by

$$\mathcal{M} = \epsilon^{*\mu} \bar{v}_2 \gamma_\mu (g_V - g_A \gamma_5) u_1 \,,$$

and this result holds both if  $\Psi$  and  $\bar{\Psi}$  are related, e.g. in the decay  $Z^0 \to e^+e^-$ , and if  $\bar{\Psi}$  is a different Dirac field than  $\Psi$ , e.g. in the decay  $W^+ \to e^+\bar{\nu}$ . Compute the rates for the decay processes,

(a)  $W^+ \to e^+ \bar{\nu}_e$ ,

(b)  $Z^0 \to e^+ e^-$ , (c)  $Z^0 \to \bar{\nu}_e \nu_e$ ,

neglecting the electron mass. Express your results in GeV.

*Hint:* start by computing the decay rate  $|\mathcal{M}|^2$  without specifying the quantities  $g_{V,A}$ . Sum over the final spins and average over the three initial polarisations. Find the total decay rate,  $\Gamma$ , remembering to divide by the symmetry factor. Generally you should find, for distinguishable massless outgoing particles, that

$$\Gamma = \frac{1}{12\pi} (g_V^2 + g_A^2) M \,.$$

Useful data:

$$\begin{array}{c|c} W^+ \to e^+ \bar{\nu}_e \\ Z^0 \to \bar{\nu}_e \nu_e \\ Z^0 \to e^+ e^- \end{array} \begin{array}{c} g_V = g_A = g_2/2\sqrt{2} \\ g_V = g_A = \frac{e}{4\sin\theta_W \cos\theta_W} \\ g_V = (-\frac{1}{4} + \sin^2\theta_W) \frac{e}{\sin\theta_W \cos\theta_W}, \qquad gA = -\frac{1}{4}\frac{e}{\sin\theta_W \cos\theta_W} \end{array}$$

with  $g_2 = \frac{e}{\sin \theta_W}$  and  $e^2 = 4\pi \alpha$ .