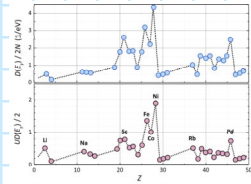


Lecture 24

Wednesday, January 19, 2022 1:13 PM

Review

$$\frac{1}{2} \mu \cdot D(E_F) > 1$$

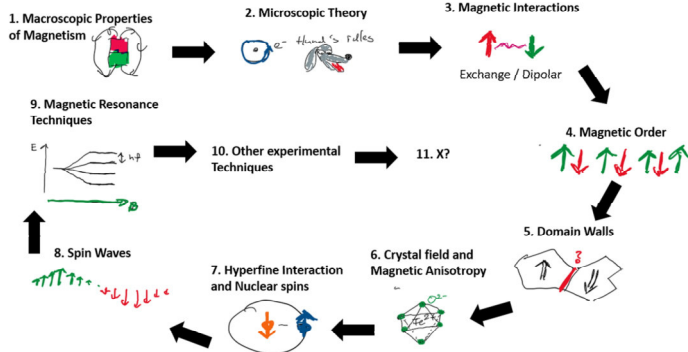


$$\chi = I \cdot \chi_p$$

χ_p : Pauli-susceptibility

$$I = \frac{1}{1 - \frac{1}{2} \mu \cdot D(E_F)}$$

Schedule

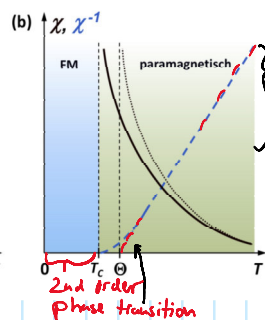
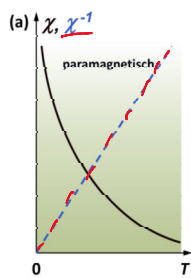


4. Magnetic Order

$$T_c = \gamma \cdot C$$

$$\chi = \frac{C}{T - T_c}$$

Curie-Weiss-law



works well for $T \gg T_c$

works not so well for $T \approx T_c$
due to magnetic interaction in the paramagnetic region

works not so well for $T \gg T_c$
due to magnetic interaction in the
paramagnetic region

- Introduction of a "paramagnetic" Curie-temperature
(Weiss-Constant) $\Theta = T_c + \text{a bit of interaction}$

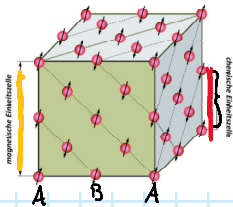
Material	T_c (K)	Θ (K)	C (K)	M_s (kA/m)	ρ
Fe	1044	1100	2.22	1750	2.22
Co	1360	1415	2.24	1450	1.72
Ni	629	649	0.588	510	0.60
Gd	289	302	5.00	2060	7.63
Dy	85	157	—	3050	10.4
EuO	69.4	78	4.68	1880	7.0
CrO ₂	396	—	—	325	2.0
MnAs	630	318	—	501	3.4

BA
1030 T
685 T

- MFT works also not so well $T \ll T_c$
(Spin-Waves introduce a T-dependence in
FM region)

2. Antiferromagnet

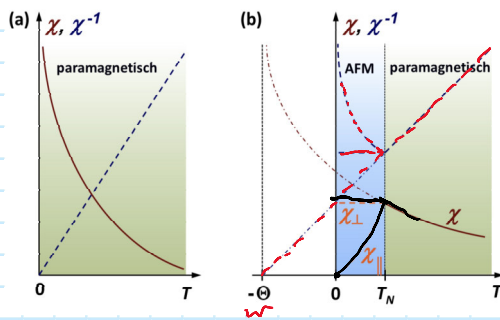
MnO:



X-ray: chemical unit cell
neutron: magnetic unit cell
scattering

$$T_N = |\delta_{AB} - \delta_{AA}| \cdot C$$

$$\chi = \frac{2C}{T + \Theta} \quad \Theta = |\delta_{AB} + \delta_{AA}| \cdot C$$

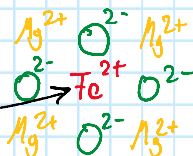


$\uparrow \downarrow \uparrow \downarrow \uparrow \Rightarrow \chi_{\perp} = \frac{1}{15\mu_B}$
 $\uparrow \downarrow \uparrow \downarrow \uparrow \uparrow \Rightarrow \chi_{\parallel} = 0 \quad @ T=0$
 $\chi_{\parallel} = \frac{1}{15\mu_B} @ T_c$

5. Crystal field and magnetic anisotropy

5.1. Crystal field

Transition metal ions:



e^- in 3d of Fe are
exposed to an inhomogeneous
electric field induced by neighboring ions

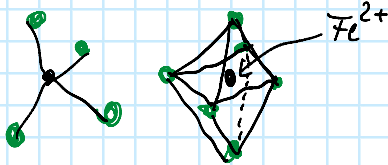
e in 3d of Fe are exposed to an inhomogeneous electric field induced by neighboring ions

⇒ Crystal field

- break spherical symmetry
- $\approx 1\text{eV}$ for 3d
- stronger than spin-orbit interaction $\lambda \cdot \vec{S} \cdot \vec{L}$

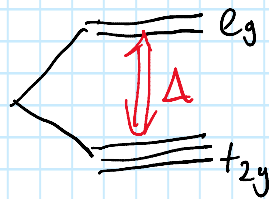
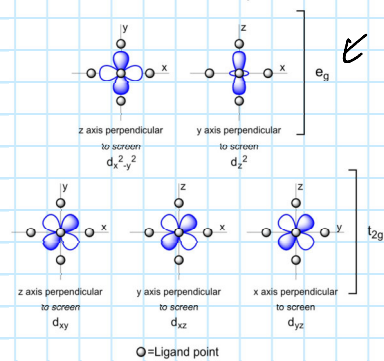
↳ Violation for 3rd Hund's rule (J is not a good quantum number anymore)

Tetrahedral



• Imagine you are a d-shell electron
→ Where would you like to sit?

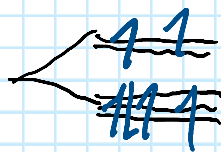
Metal d orbitals in octahedral crystal field



distance Δ between the levels now will decide the spin state

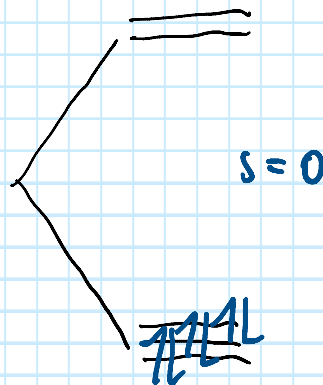
[https://chem.libretexts.org/Bookshelves/Inorganic_Chemistry/Modules and Websites \(Inorganic Chemistry\)/Crystal Field Theory/Crystal Field Theory](https://chem.libretexts.org/Bookshelves/Inorganic_Chemistry/Modules_and_Websites_(Inorganic_Chemistry)/Crystal_Field_Theory/Crystal_Field_Theory)

again for Fe^{2+} (d^6)



$\Delta \ll 1\text{eV}$

(low spin state)



$\Delta \gg 1\text{eV}$

(high spin state)

Orbital quenching

$$g_2(\lambda(\lambda+1))^{1/2} \quad g_s(S(S+1))^{1/2}$$

$g_L(\frac{L(L+1)}{2})^{1/2}$ $g_S(\frac{S(S+1)}{2})^{1/2}$

ion	shell	S	L	J	term	P_1	P_{exp}	P_2
Ti^{3+}, V^{4+}	$3d^1$	$\frac{1}{2}$	2	$\frac{3}{2}$	${}^2D_{3/2}$	1.55	1.70	1.73
V^{3+}	$3d^2$	1	3	2	3F_2	1.63	2.61	2.83
Cr^{3+}, V^{2+}	$3d^3$	$\frac{3}{2}$	3	$\frac{3}{2}$	${}^4F_{3/2}$	0.77	3.85	3.87
Mn^{3+}, Cr^{2+}	$3d^4$	2	2	0	5D_0	0	4.82	4.90
Fe^{3+}, Mn^{2+}	$3d^5$	$\frac{5}{2}$	0	$\frac{5}{2}$	${}^6S_{5/2}$	5.92	5.82	5.92
Fe^{2+}	$3d^6$	2	2	4	5D_4	6.70	5.36	4.90
Co^{2+}	$3d^7$	$\frac{3}{2}$	3	$\frac{9}{2}$	${}^4F_{9/2}$	6.63	4.90	3.87
Ni^{2+}	$3d^8$	1	3	4	3F_4	5.59	3.12	2.83
Cu^{2+}	$3d^9$	$\frac{1}{2}$	2	$\frac{5}{2}$	${}^2D_{5/2}$	3.55	1.83	1.73
Zn^{2+}	$3d^{10}$	0	0	0	1S_0	0	0	0

- For 3d elements the predicted value of the total magnetic moment leads to poor results

⇒ crystal field > spin-orbit (3rd Hund's rule)

- instead, often $L=0$, $J=S$

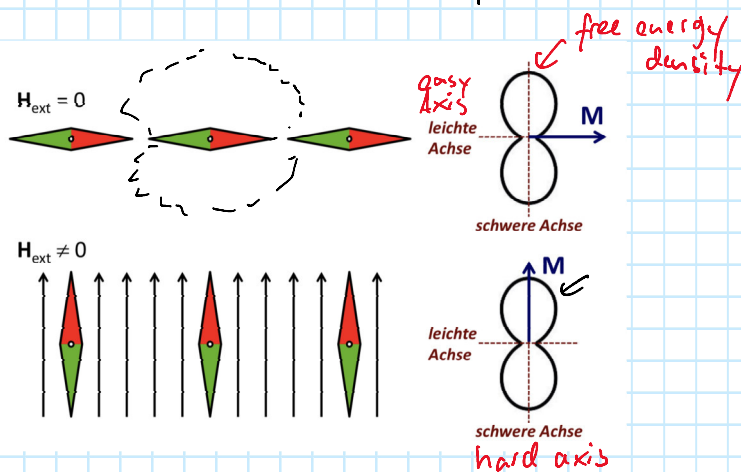
⇒ orbital moment is quenched

Semiclassical: orbital angular momentum precesses in the crystal field with unchanged magnitude, but components which average to zero

- Often orbital angular moment is not fully quenched → spin-orbit is not fully negligible

5.2. Magnetic Anisotropy

- In general: Different spatial directions do not have the same magnetic properties



- Three main sources for magnetic sources

$$E_{ani} = E_{MC} + E_{shape} + E_{ind}$$

$\uparrow$ $\uparrow$

$$E_{ani} = E_{MC} + E_{shape} + E_{ind}$$

$\uparrow$
Magnetocrystalline
 $\uparrow$
Induced

Magnetocrystalline Anisotropy

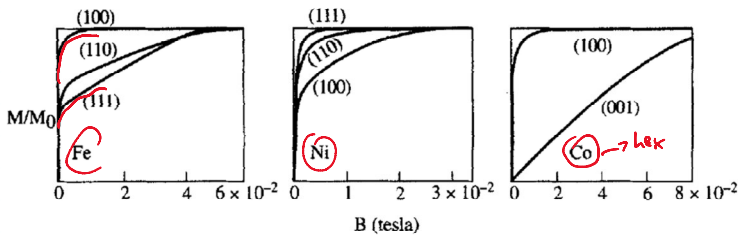


Fig. 6.22 Magnetization in Fe, Co and Ni for applied fields in different directions showing anisotropy. After Honda and Kaya 1926, Kaya 1928.

Phenomenological description:

magnetic free energy (density) $f_{ani} = \frac{E_{ani}}{V}$

as a function of $\vec{m} = \frac{\vec{M}}{|\vec{M}|} = \begin{pmatrix} m_x \\ m_y \\ m_z \end{pmatrix} = \begin{pmatrix} \sin\theta \cdot \cos\phi \\ \sin\theta \cdot \sin\phi \\ \cos\theta \end{pmatrix}$

Uniaxial anisotropy

$$f_{uni} \approx K_{u1} \cdot \cos^2\theta$$

general:

$$f_{uni} = K_{u1} |\vec{m} \cdot \vec{u}|^2 + K_{u2} |\vec{m} \cdot \vec{u}|^4 + \dots$$

