

Lecture 18 - WiSe2022/2023

Tuesday, December 27, 2022 10:47 AM

2.3. Atomic Paramagnetism

- e.g.
- insulator (e.g. MgO , $NaCl$) with a transition metal or a rare earth inside
 - Magnetic Molecules
 - Color Center in Diamond (Nitrogen-Vacancy-Center)

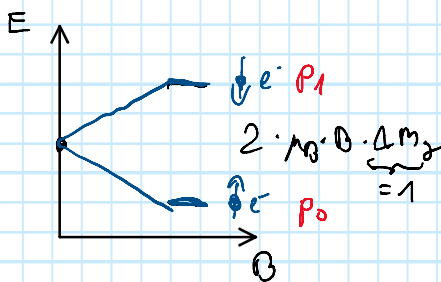
↳ no conduction electrons

$$E = \mu_B \cdot \vec{B} (\vec{L} + g \vec{S}) / \hbar = g_L \mu_B \cdot \frac{\vec{L}}{\hbar} \cdot \vec{B}$$

Landé factor total orbital moment

• example: Eigenstates for a spin- $\frac{1}{2}$ system

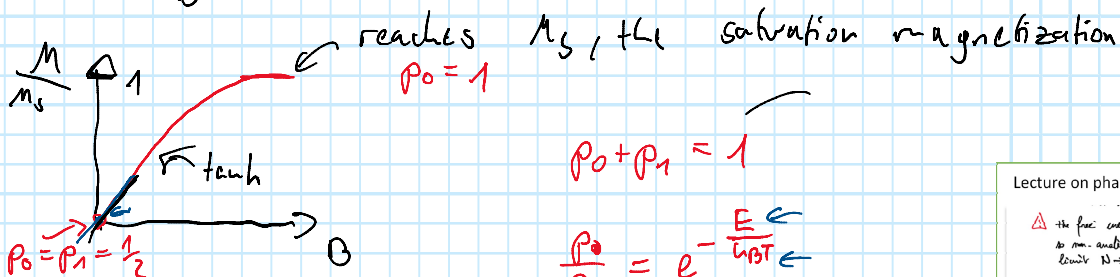
$$J = S = \frac{1}{2} \Rightarrow m_J = \pm \frac{1}{2} \quad g_J = 2$$



$$E = -\vec{\mu} \cdot \vec{B}$$

$$\left(\frac{\mu_B B}{k_B T} \right)$$

Q



$$P_0 + P_1 = 1$$

$$\frac{P_0}{P_1} = e^{-\frac{E}{k_B T}}$$

$$M = \langle m_J \rangle \propto m_0 \cdot P_0 + m_1 \cdot P_1$$

-1/2

Lecture on phase transitions

⚠ the free energy of a finite system is always analytical so no. singularity can only occur in the thermodynamic limit $N \rightarrow \infty$ → Coexistence of all degrees of freedom in the system.

Statistical physics $F = k_B T \langle e^{-\frac{E}{k_B T}} \rangle$

↳ average over all possible configurations of the system!

or fluctuations → new perspective of the quantum states and vacuum

Fluctuations are related to susceptibility

Degenerate case: $\chi = \left(\frac{\partial m}{\partial h} \right)_T$

↳ external field

Effect of the magnetic field?

} Z, the partition function

Free Energy $F = U - T \cdot S = \dots = -k_B \cdot T \cdot \ln(Z)$

$$M = - \frac{\partial F}{\partial B} = \frac{\mu_{1/2} \cdot e^x + \mu_{-1/2} \cdot e^{-x}}{e^x + e^{-x}} = \frac{e^x - e^{-x}}{e^x + e^{-x}}$$

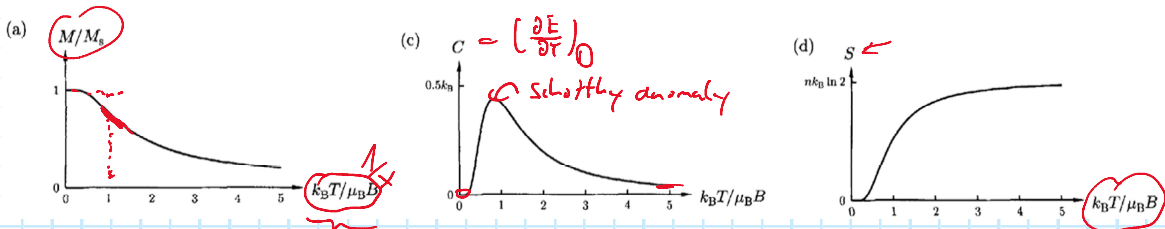
$$M := \langle m \rangle = - \frac{\partial F}{\partial B} = \frac{\overset{\downarrow \mu_B}{m/2} \cdot e^x + \overset{\downarrow -\mu_B}{m/2} \cdot e^{-x}}{e^x + e^{-x}} \quad \left. \vphantom{\frac{\partial F}{\partial B}} \right\} \frac{e^x - e^{-x}}{e^x + e^{-x}} = \tanh$$

$$x = \frac{\mu_B \cdot B}{k_B T}$$

$$\frac{M}{M_s} = \tanh \left(\frac{\mu_0 \cdot B}{k_B \cdot T} \right) \quad M_s = \max(\langle m \rangle) = \mu_0$$

$$\text{Susceptibility} \quad \chi = \mu_0 \cdot \left(\frac{\partial M}{\partial B} \right)_{T,V} \stackrel{x \ll 1 \rightarrow \tanh(x) \approx x}{\approx} \frac{\mu_0 \cdot \mu_0^2}{k_B T} = \frac{C}{T}$$

$$k_B = 1.38 \cdot 10^{-23} \frac{J}{K} = 86 \frac{meV}{K}$$



$$\langle m_j \rangle = \frac{\sum_{m_j = -1/2}^{+1/2} m_j \cdot e^{-\frac{m_j \cdot g_j \cdot \mu_B \cdot B}{k_B T}}}{\sum_{m_j = -1/2}^{+1/2} e^{-\frac{m_j \cdot g_j \cdot \mu_B \cdot B}{k_B T}}} = - \frac{1}{Z} \frac{\partial Z}{\partial x}$$

$$M = - \frac{N}{V} \cdot g_j \cdot \mu_B \cdot \langle m_j \rangle = n \cdot \frac{g_j \cdot \mu_B}{Z} \frac{\partial Z}{\partial B} \cdot \frac{\partial B}{\partial x}$$

$$= n \cdot k_B \cdot T \cdot \frac{\partial \ln Z}{\partial B} \quad Z = \frac{\sinh[(2j+1) \cdot \frac{x}{2}]}{\sinh(\frac{x}{2})}$$

$$\frac{M}{M_s} = B_2(y) \stackrel{\text{Brillouin function}}{=} \frac{2j+1}{2j} \cdot \coth\left(\frac{2j+1}{2j} \cdot y\right) - \frac{1}{2j} \cdot \coth\left(\frac{1}{2j} \cdot y\right)$$

$$y = x \cdot j = \frac{g_j \cdot \mu_B \cdot j \cdot B}{k_B \cdot T}$$

$$M_s = n \cdot g_j \cdot \mu_B \cdot j$$

$$B_2(\gamma) \approx \frac{\gamma+1}{3\gamma} \cdot \gamma = \frac{\gamma+1}{3} \cdot x$$

$$\chi = \mu_0 \cdot \left(\frac{\partial M}{\partial B} \right)_{T,V} = \frac{\mu_0 \cdot n \cdot \gamma (\gamma+1) g_J^2 \cdot \mu_B^2}{3 k_B T} = \frac{C}{T}$$

$$C = \frac{\mu_0 \cdot n \cdot (\rho \cdot \mu_0)^2}{3 k_B} = \frac{\mu_0 \cdot n \cdot \mu_{\text{eff}}^2}{3 k_B}$$

effective magneton number $\rho = g_J \sqrt{\gamma(\gamma+1)}$

$$\mu_{\text{eff}} = \rho \cdot \mu_0$$

