

# Dynamics of electrons in Solids (= Crystals)

- \* Drude model / free electrons
- \* Add Quantum mechanics (Sommerfeld Model)
- \* Add Lattice (Bloch)

## ① Why is the transport / dynamics of electrons interesting.

Electrons are charged so they are put in motion by  $E$  field and/or by thermal gradient. This results in a "current" of electricity or heat that is related to the applied field but also to the intrinsic properties of materials

$$U = R \cdot I$$

voltage  $\rightarrow$   $U$   $\leftarrow$  flow of electrons  
 $\downarrow$   
 "transport coefficient"

	$B = 0, M = 0$	
	$\vec{E} = -\nabla\phi_{el}$	$-\nabla T$
$J_q$	Electrical Conductivity ( $\sigma$ ) or resistivity ( $\rho$ )	Seebeck-Effect (S)
$J_H$	Peltier-Effect (P)	Heat Conductivity ( $\kappa$ )
	$B \neq 0$	
	$-\nabla\phi_{el} \times \vec{B}$	$-\nabla T \times \vec{B}$
$J_q$	Magnetoresistance $\sigma_{xx}(B)$ Hall Effect ( $R_H$ )	Nernst Effect (N)
$J_H$	Ettinghausen-Effect	Thermal Hall Effect

## ② A brief Recollection on Drude Model 1897

Key idea: Classical gas of electrons which follow Maxwell-Boltzmann statistics.

Probability of having an electron with a velocity  $\vec{v}$  between  $\vec{v}$  and  $\vec{v} + d\vec{v}$  is given by:

$$f(\vec{v}) = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi v^2 \exp\left(-\frac{mv^2}{2k_B T}\right)$$

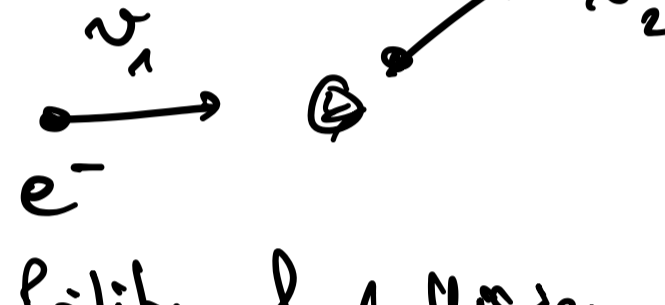
or equivalently since  $E = \frac{1}{2} m v^2$

$$f(E) = 2 \left(\frac{E}{\pi(k_B T)^3}\right)^{1/2} e^{-E/k_B T}$$

$$\text{Average electronic energy } \langle E \rangle = \int_0^\infty f(E) \cdot E \cdot dE = \frac{3}{2} k_B T = U$$

Assumption = thermal equilibrium is achieved through collisions

- ① No interactions between the collisions
- ② Instantaneous change of the velocity after the collision



- ③ Probability of collision is given by  $\tau$ ,  $1/\tau =$  scattering rate [given in  $s^{-1}$ ]  
 $\Rightarrow$  electronic mean free path = average distance between two collisions  
 $l = \langle v \rangle \cdot \tau$

Main Result: estimate of the electrical conductivity

$$\vec{J}_j = \vec{\sigma} \cdot \vec{E}$$

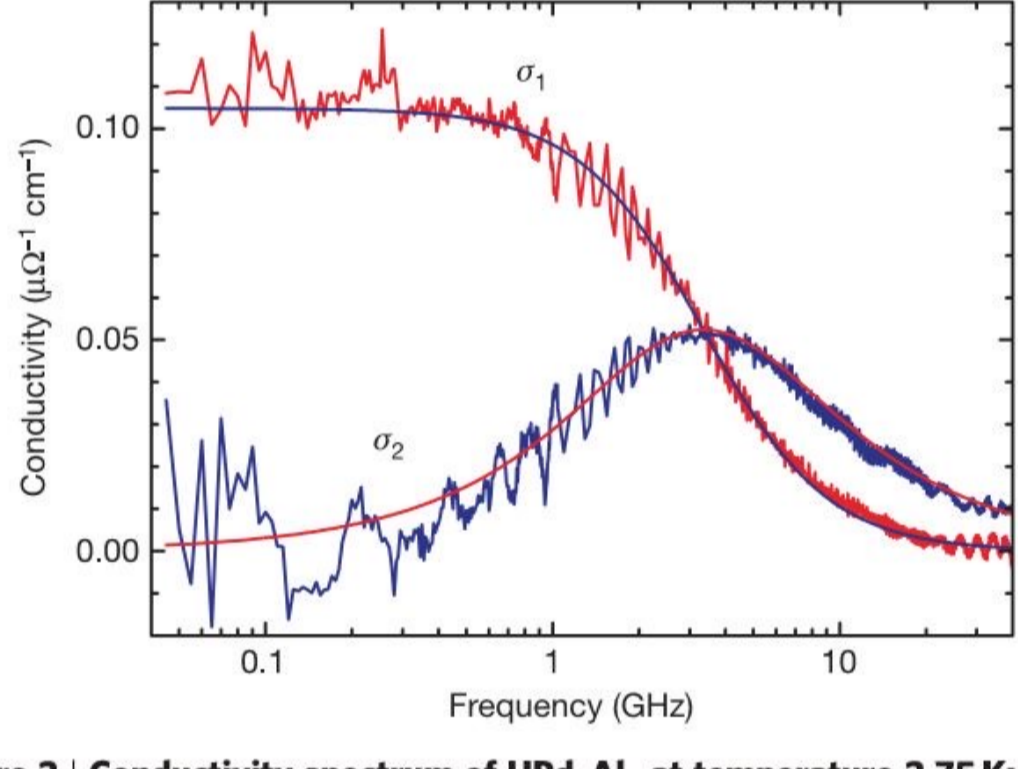
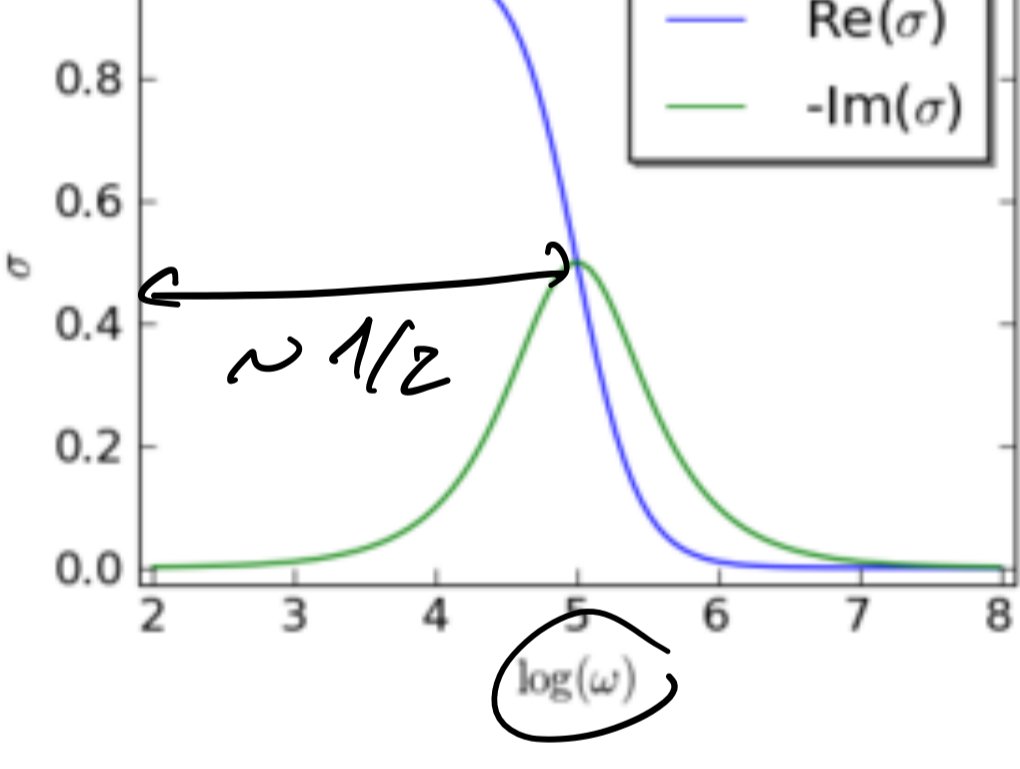
$\hookrightarrow$  conductivity tensor

- DC-conductivity  $\sigma_0 = \frac{n e^2 \tau}{m}$ 
  - $n$  = charge density
  - $m$  = mass of the charge carriers
  - $e$  = charge.

• AC / optical conductivity = conductivity at finite frequency  $\omega$

$$\sigma(\omega) = \frac{\sigma_0}{1 - i\omega\tau} = \frac{\sigma_0}{1 + \omega^2\tau^2} + i \frac{\omega\tau\sigma_0}{1 + \omega^2\tau^2}$$

$\underbrace{\hspace{10em}}_{\text{re part}} \quad \underbrace{\hspace{10em}}_{\text{imaginary part}}$



(Source: Wikipedia)

Source: Scheffler et al. Nature 438, 1135 (2005)

Main failures: + electronic heat capacity: energy to provide to the system to increase its temperature by  $1K$

$$C_{el} = \frac{dU}{dT} \quad U = \text{external energy}$$

$$= \frac{3}{2} n k_B \rightarrow T\text{-independent}$$

$$\text{Experimentally: } C = \gamma T + \beta T^3$$

$\gamma$  is  $T$ -independent,  $\beta$  is  $T$ -dependent

Overall the experimental value is 2 orders of magnitude smaller than  $\frac{3}{2} n k_B$

\* thermal conductivity  $\kappa$

$$\vec{J}_H = -\vec{\kappa} \cdot \nabla T$$

$$\kappa = \frac{1}{3} v^2 \tau C_{el} = \text{also off by two orders of magnitude}$$

\* Wiedemann-Franz law:  $\frac{\kappa}{\sigma T} = d = \text{constant}$   
OK (by chance)

\* Hall effect:  $R_H = -\frac{1}{ne} = T\text{-independent}$   
 $\hookrightarrow$  contradiction with experiment

## ③ Drude - Sommerfeld Model

Replace Boltzmann distribution with the Fermi Dirac one

$$f_D(E) = \frac{1}{\exp\left(\frac{E-\mu}{k_B T}\right) + 1}$$

$$\mu = \text{chemical potential} = E_F \left(1 - \frac{\pi^2}{12} \left(\frac{T}{T_F}\right)^2\right)$$

