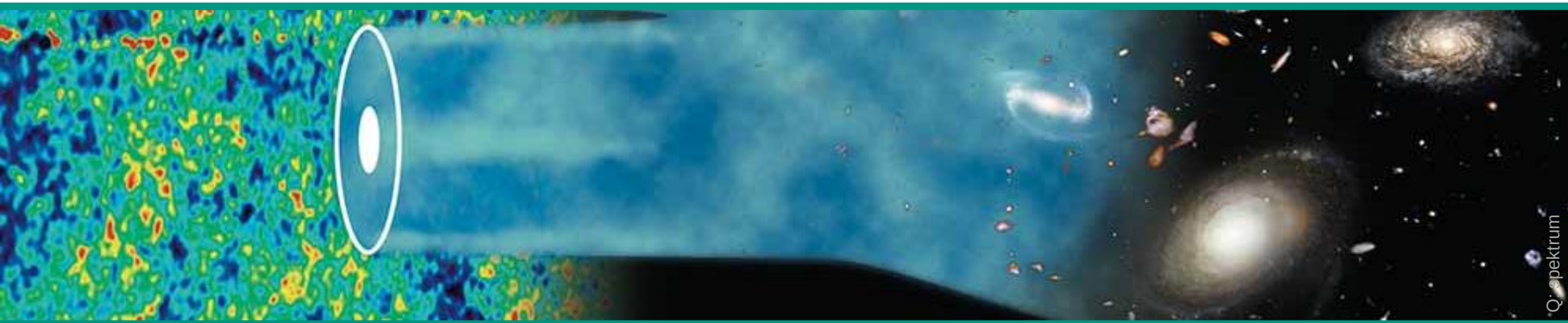


Introduction to Cosmology

Winter term 22/23

Lecture 6

Dec. 6, 2022



Recap of Lecture 5

■ Thermal universe: the radiation-dominated era with phase transitions

- temperature evolution: $T(t) \sim \sqrt{t}$ with fixpoint $T = 1 \text{ MeV}$ at $t = 1 \text{ s}$
- particles (energy E_J) in thermodynamical equilibrium with heat bath (photons)
- resulting Boltzmann distribution with $N_J(t) \sim e^{-E_J/k_B T(t)}$
- condition for **freeze-out**: $\Gamma(t) = H(t) \Leftrightarrow$ interaction rate = Hubble expansion
- photons from matter – anti-matter annihilation, tiny **baryon asymmetry η**
- *QCD* – phase transition: quark-gluon plasma $\rightarrow$ nucleons (p, n)

RECAP: neutrino freeze-out

■ during expansion of universe: **freeze-out** processes of ν 's are important

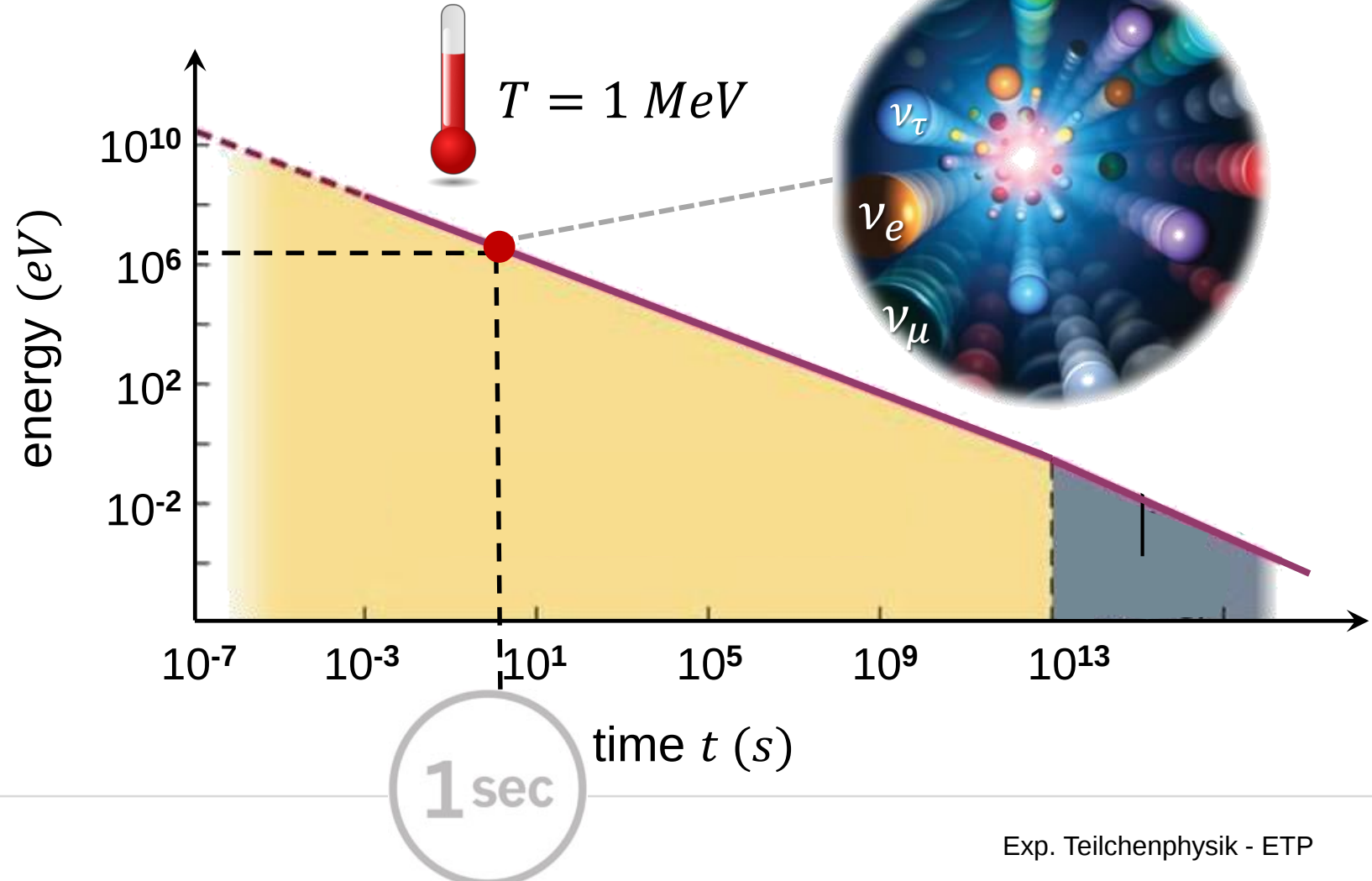
- since $t = 1\text{ s}$ neutrinos from the Big Bang are 'free-streaming'

$$T(t = 1\text{ s}) = 1\text{ MeV}$$

$$t = 1\text{ s}$$

$$\Gamma_\nu(t) = H(t)$$

weak interaction
cross section



neutrino freeze-out & # of relativistic particles

■ An important time stamp: ν 's decouple
at $t = 1\text{ s}$ & at $T = 1\text{ MeV}$

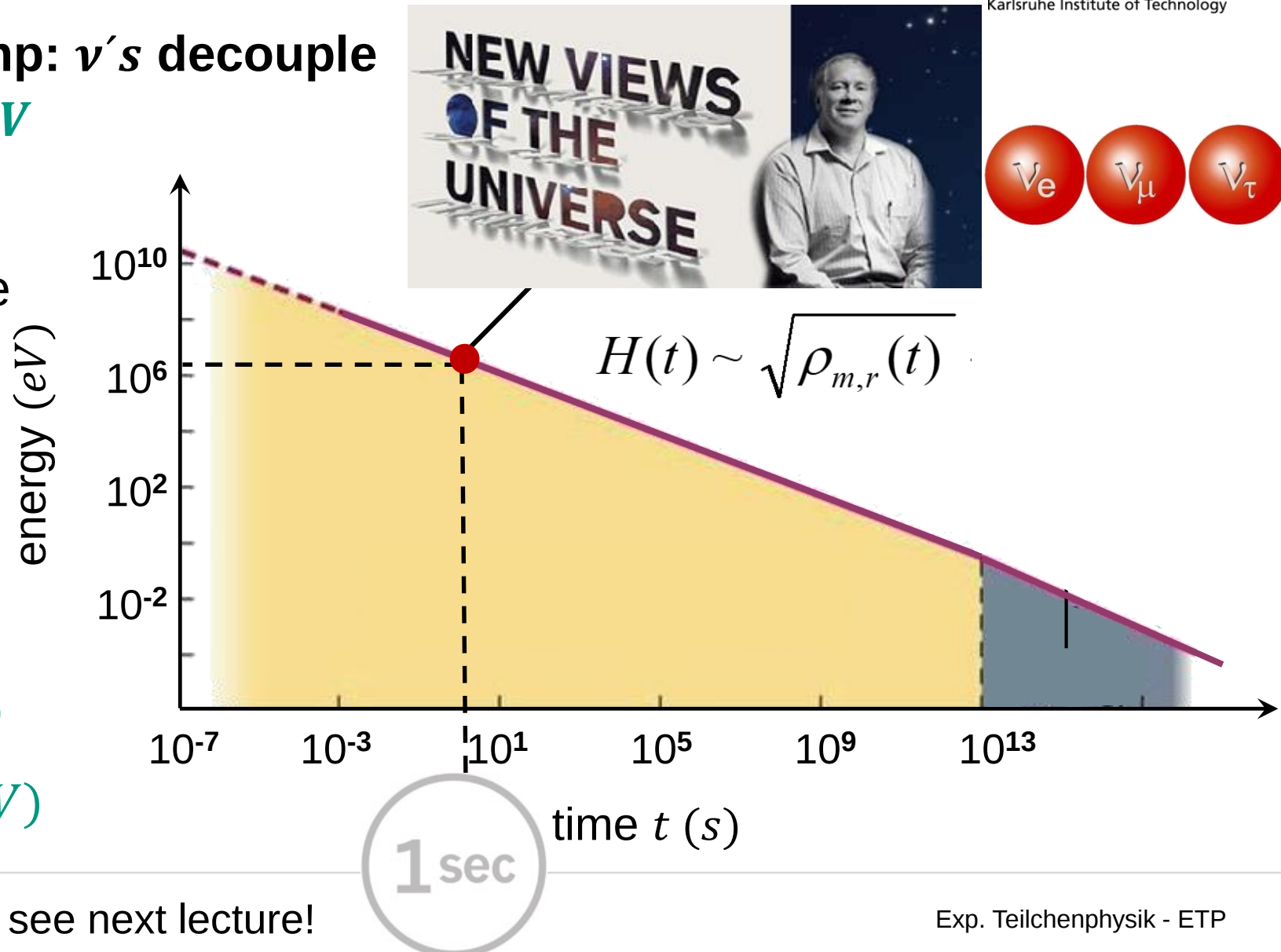
- how many relativistic
‘degrees of freedom’ were
present at this time?*

- preview: **relativistic
particles** at $E = 1\text{ MeV}$

γ : ☒ ($m_\gamma = 0$)

$\nu_e \nu_\mu \nu_\tau$: ☒ ($m_\nu < 0.8\text{ eV}$)

$e^- e^+$: ☒ ($m_e = 0.511\text{ MeV}$)



neutrino freeze-out & # of relativistic particles

■ An important time stamp: ν 's decouple
at $t = 1\text{ s}$ & at $T = 1\text{ MeV}$

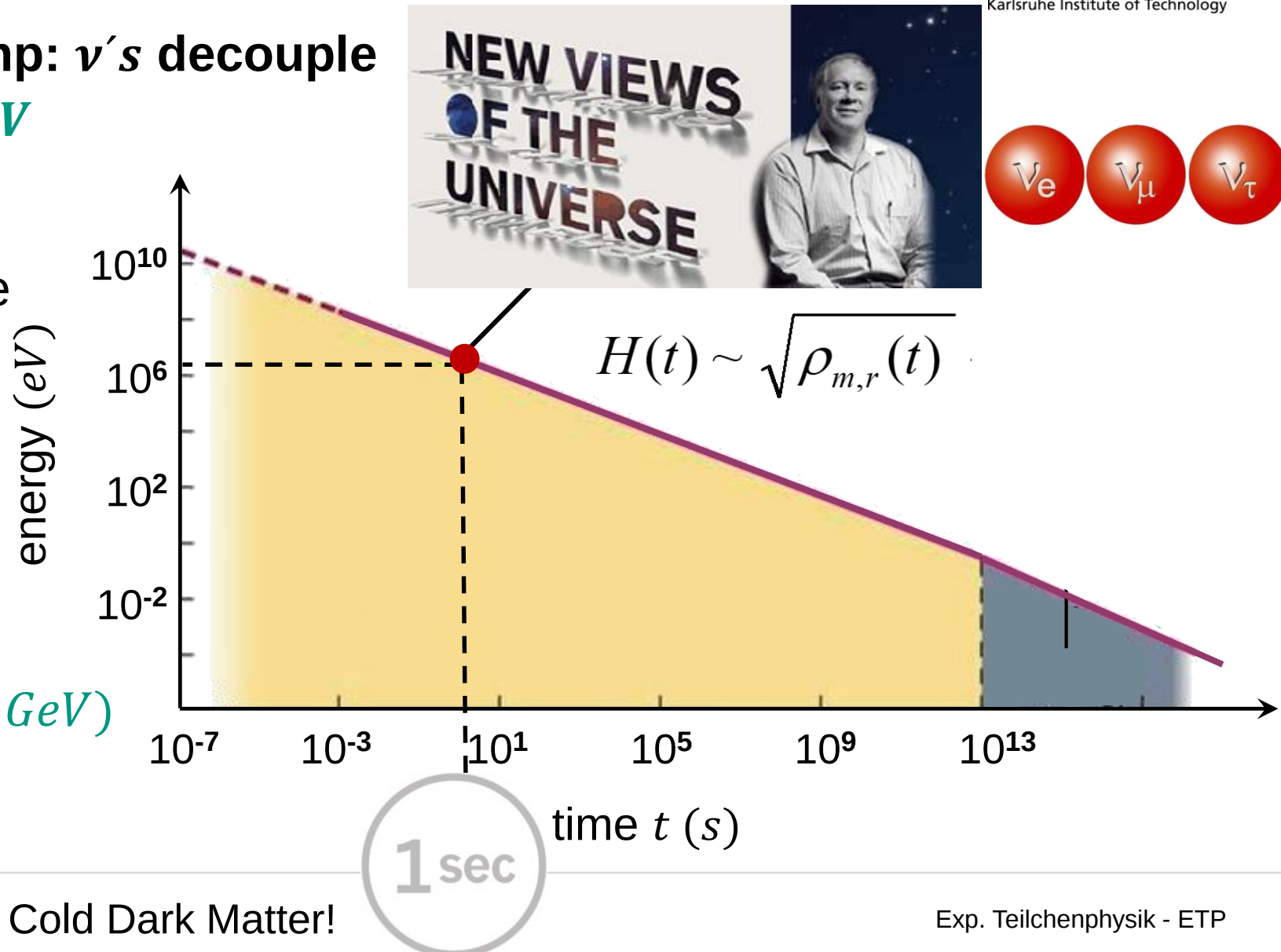
- how many relativistic
‘degrees of freedom’ were
present at this time?*

- preview: **relativistic
particles** at $E = 1\text{ MeV}$

p, n :  ($m_{p,n} \approx 1\text{ GeV}$)

χ (neutralino)*:  ($m_\chi \geq \text{GeV}$)

$\mu^- \mu^+$:  (decayed)



freeze-out of massless photons

■ Matter and radiation (CMB) decouple at later time $t = 3 \cdot 10^5 \text{ yr}$

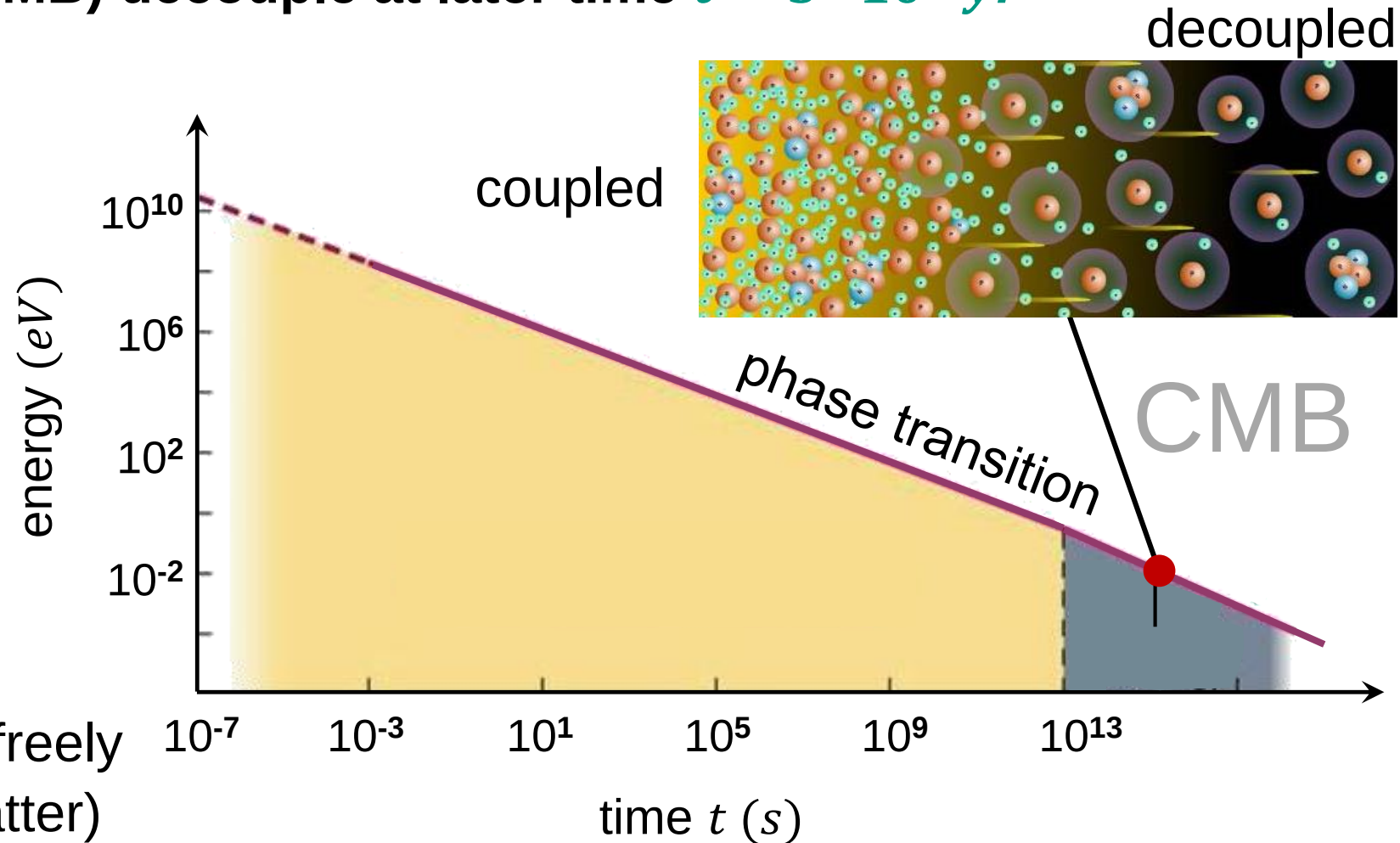
- freeze-out condition:

$$\Gamma(t) = H(t)$$

photon electromagnetic
interaction with charges



- photons (CMB) propagate freely
(without interaction with matter)



freeze-out of massless photons

■ Matter and radiation (CMB) decouple

- freeze-out condition:

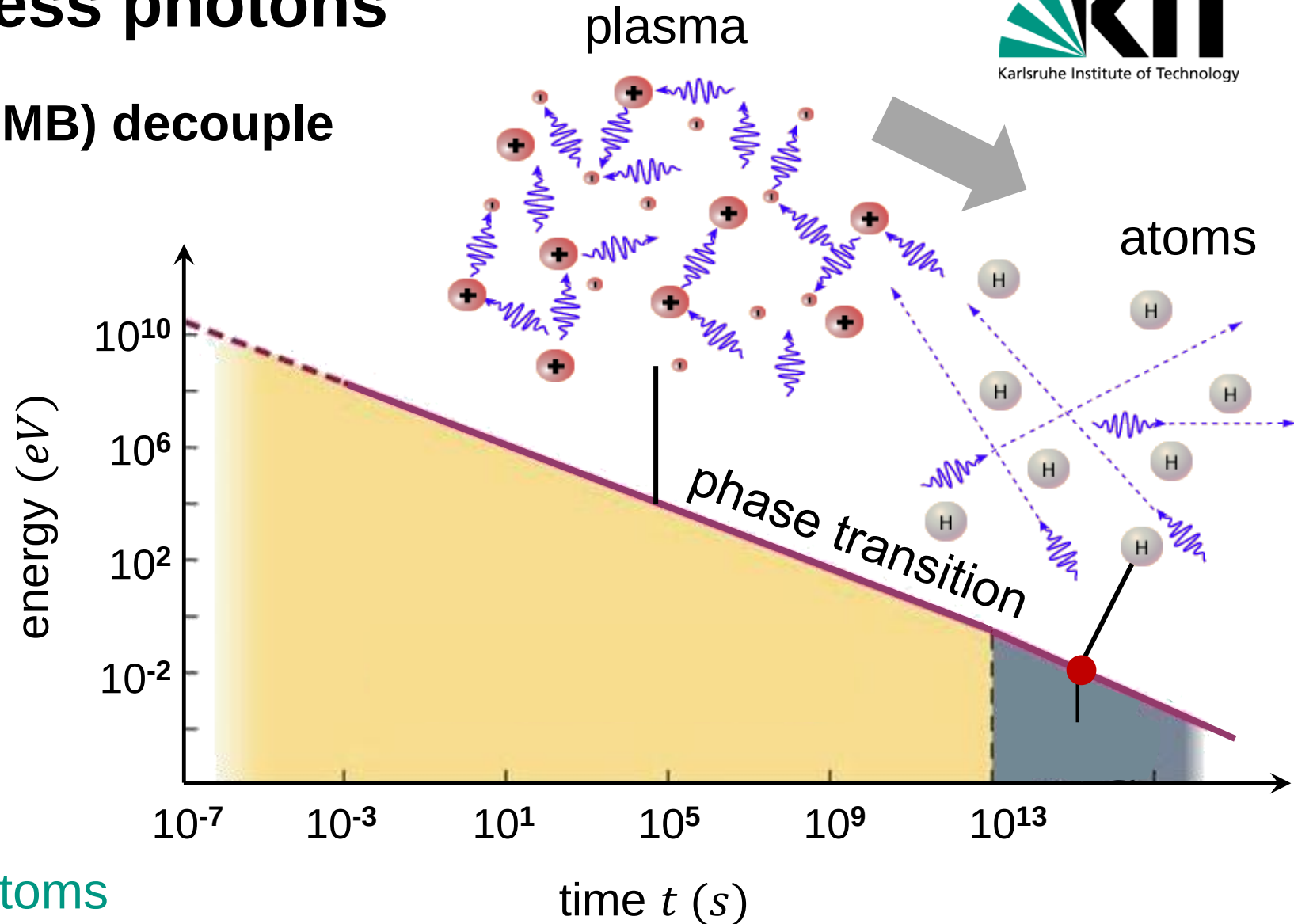
$$\Gamma(t) = H(t)$$

photon electromagnetic
interaction with charges

- phase transition:

$t < 3 \cdot 10^5 \text{ yr}$: plasma

$t > 3 \cdot 10^5 \text{ yr}$: neutral atoms



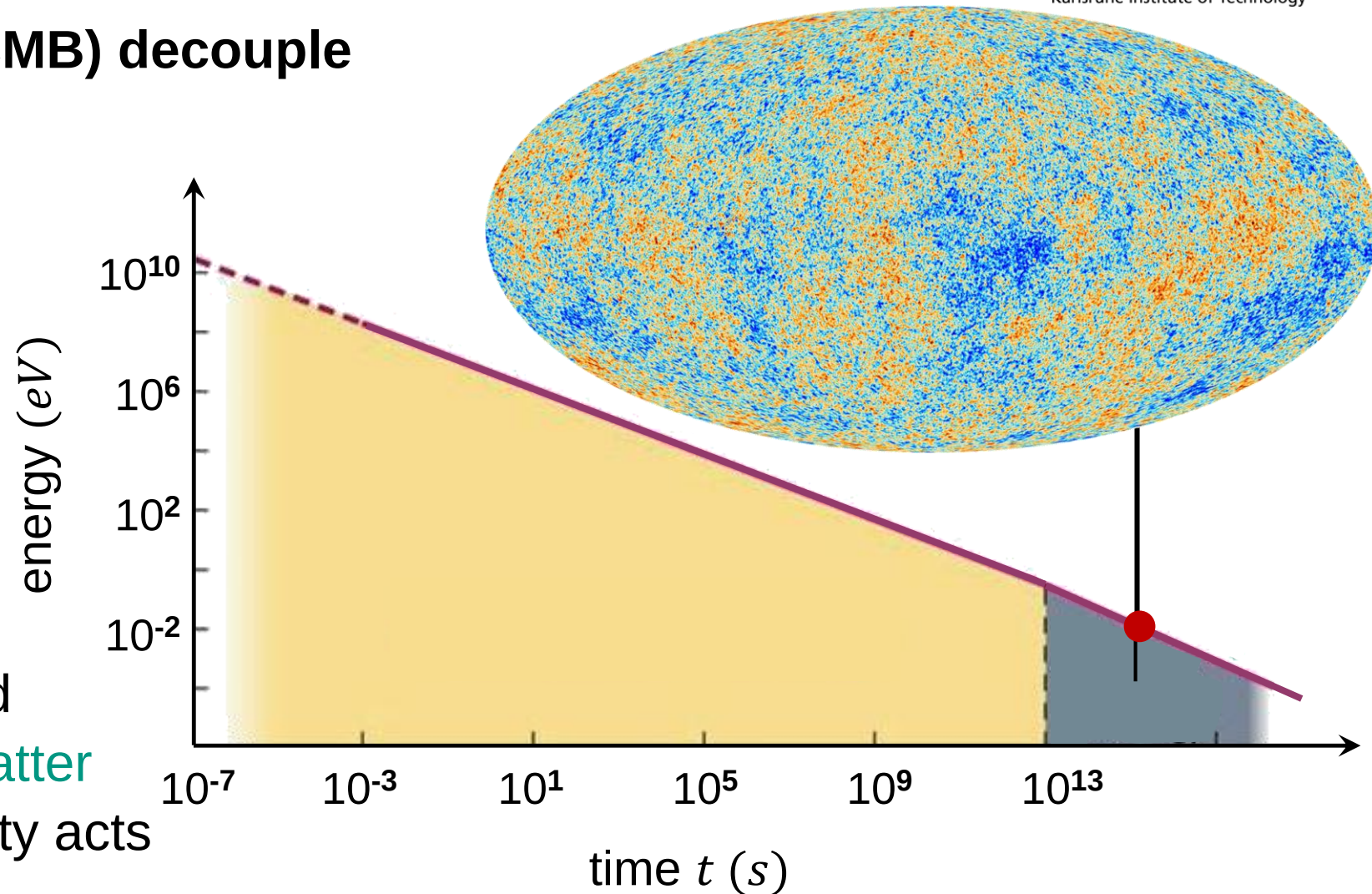
freeze-out of massless photons

■ Matter and radiation (CMB) decouple

- freeze-out condition:

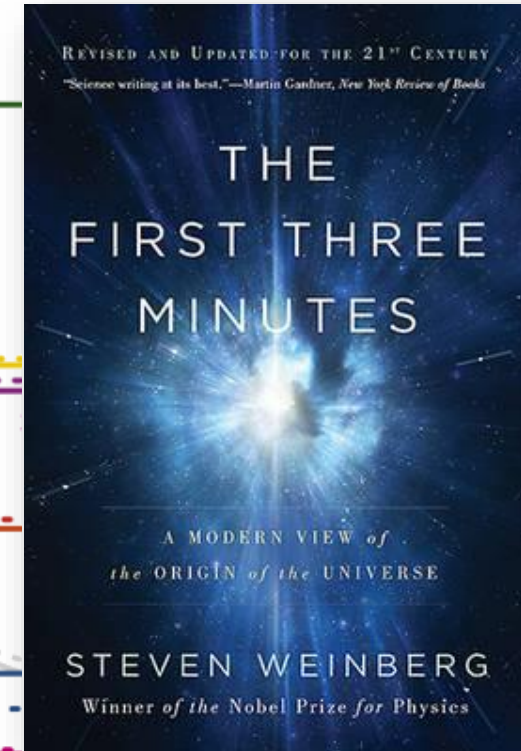
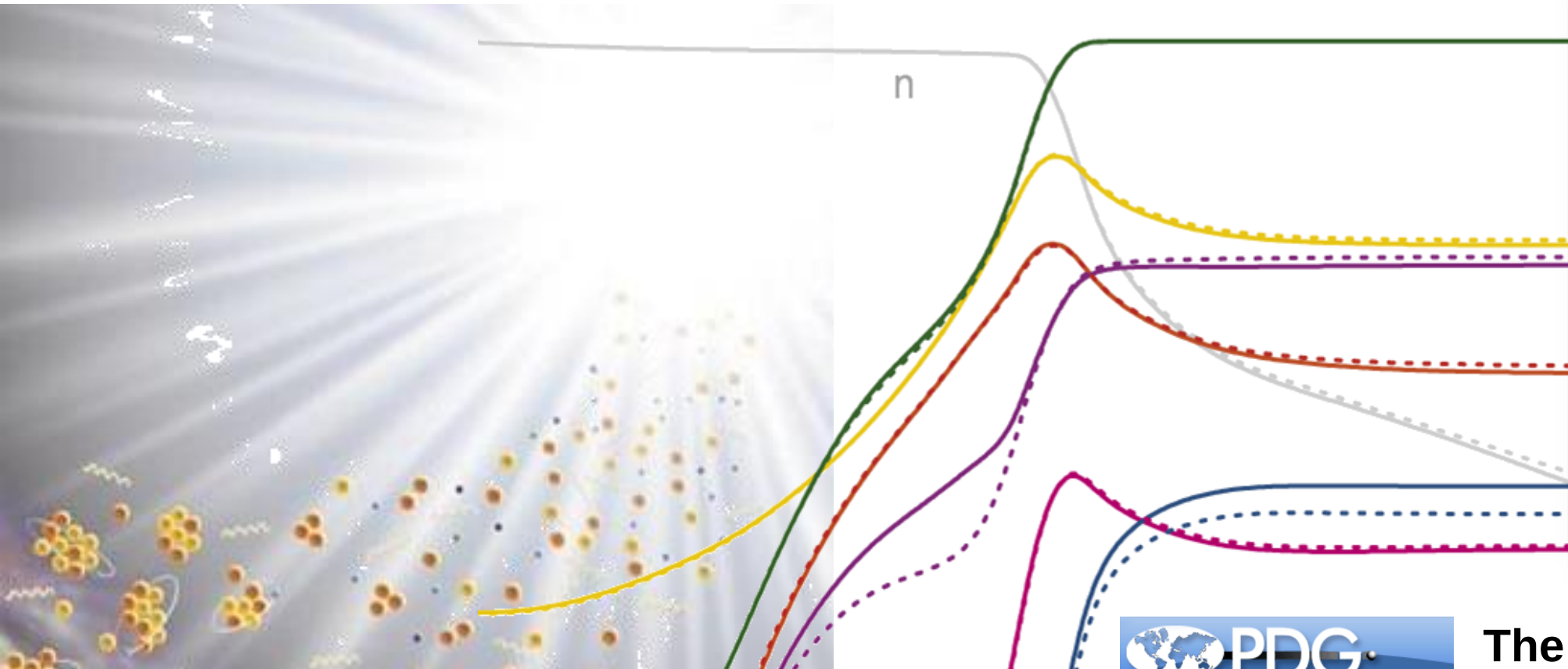
$$\Gamma(t) = H(t)$$

- universe becomes
transparent for
radiation
- matter no longer coupled
to radiation fluid, thus matter
oscillations* stop & gravity acts



3.1 Primordial nucleosynthesis (BBN*)

■ The origin of the light elements: a cornerstone for the hot Big Bang



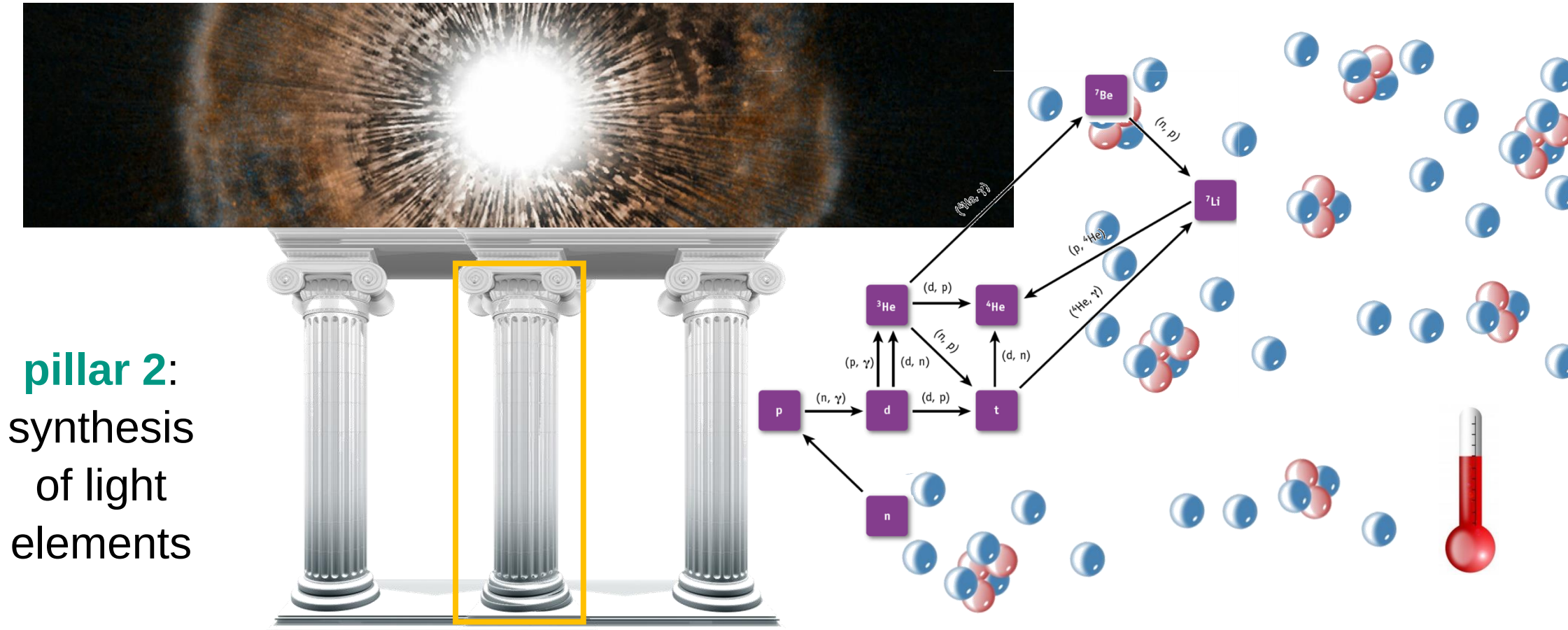
Steven Weinberg
The first three minutes
Spektrum, 2009

rpp2022-rev-bbang-nucleosynthesis.pdf (lbl.gov)



Primordial nucleosynthesis & hot Big Bang

■ Abundance of light elements as key evidence for a hot Big Bang



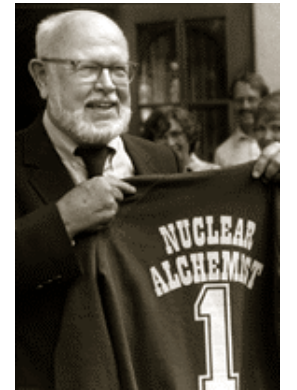
Primordial nucleosynthesis: history, breakthroughs

■ Phase-I: the 'classical' works & discoveries from 1948 - 1977

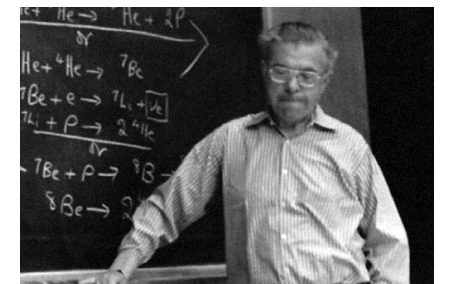
1948: **G. Gamov** & **R. Alpher** – all elements were forged in the early universe via neutron capture & β – decay



1957: **W. Fowler** *et al.* – heavy elements only made in stars (Nobel prize 1983)



1964: **F. Hoyle** *et al.* – He-4 production is 'primordial', production of heavy elements (no Nobel prize)



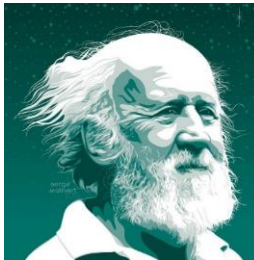
Primordial nucleosynthesis: history, breakthroughs

■ Phase-I: the 'classical' works & discoveries from 1948 - 1977

1965: **J. Peebles** – first 'modern' calculation of BBN reaction paths (Nobel prize 2019)



1970: **H. Reeves** – deuterium spectroscopy as observational means to measure the baryon density of the universe



1977: **D. Schramm et al.** – BBN limits the number of light ν -generations (link cosmology $\Leftrightarrow$ particle physics)



Primordial nucleosynthesis: fundamentals

■ The Big Bang Nucleosynthesis...

- ... allows detailed investigations of the **physics of the early universe**
- ... correctly predicts the **abundances of light elements** (despite a huge variation of more than 10 orders of magnitude!)
- ... allows an 'in-situ' determination of the **baryon density Ω_B** as well as of the **baryon asymmetry parameter η (ratio baryons / photons)**
- ... acts as an important test bed **for novel theories**: number of light ν -generations, sterile neutrinos, novel particles (such as gravitinos,...)
- ... is part of **nuclear astrophysics** (element synthesis in stars & elsewhere)

Primordial nucleosynthesis: the starting point

■ The world-famous ' $\alpha - \beta - \gamma$ ' paper from 1948 on nucleosynthesis

- all elements were forged in the very early universe via neutron capture & β - decay
- based on thesis R. Alpher, supervisor: G. Gamow
- name of 'author' H. Bethe was added '*in absentia*'
- publication date of PRL* print: April 1, 1948

INCORRECT

The Origin of Chemical Elements

R. A. ALPHER*

*Applied Physics Laboratory, The Johns Hopkins University,
Silver Spring, Maryland*

AND

H. BETHE

Cornell University, Ithaca, New York

AND

G. GAMOW

*The George Washington University, Washington, D. C.
February 18, 1948*

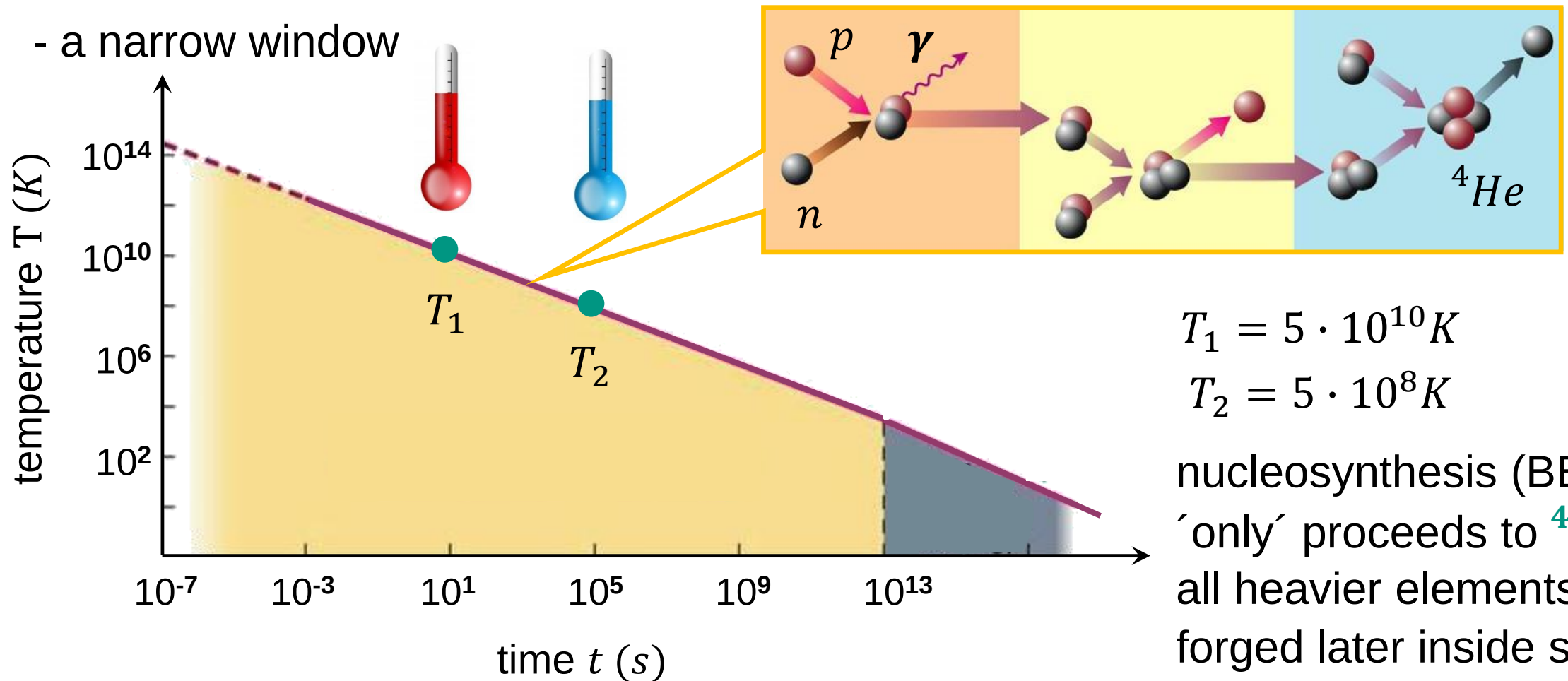
AS pointed out by one of us,¹ various nuclear species must have originated not as the result of an equilibrium corresponding to a certain temperature and density, but rather as a consequence of a continuous building-up process arrested by a rapid expansion and cooling of the primordial matter. According to this picture, we must imagine the early stage of matter as a highly compressed neutron gas (overheated neutral nuclear fluid) which started decaying into protons and electrons when the gas pressure fell down as the result of universal expansion. The radiative capture of the still remaining neutrons by the newly formed protons must have led first to the formation of deuterium nuclei, and the subsequent neutron captures resulted in the building up of heavier and heavier nuclei. It



Primordial nucleosynthesis: fundamentals

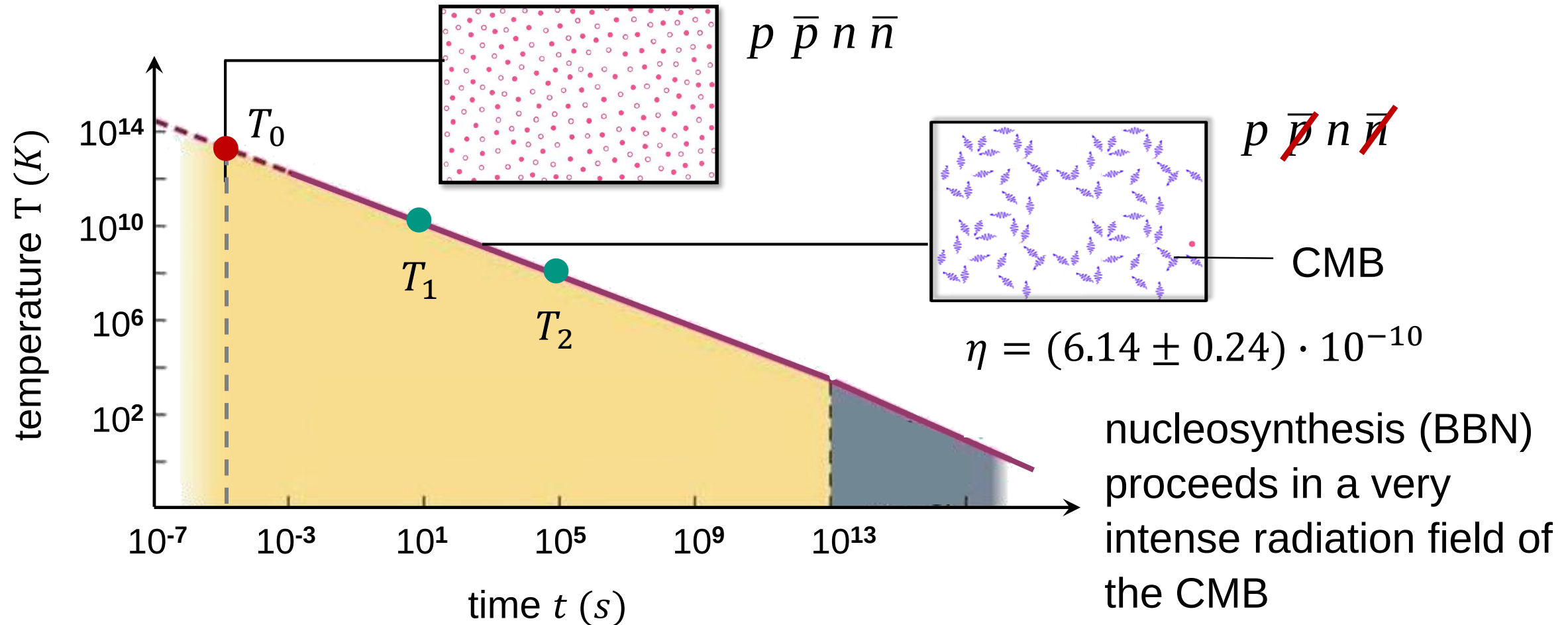
■ Modern description of processes during light element synthesis

- a narrow window



Primordial nucleosynthesis: specific phases

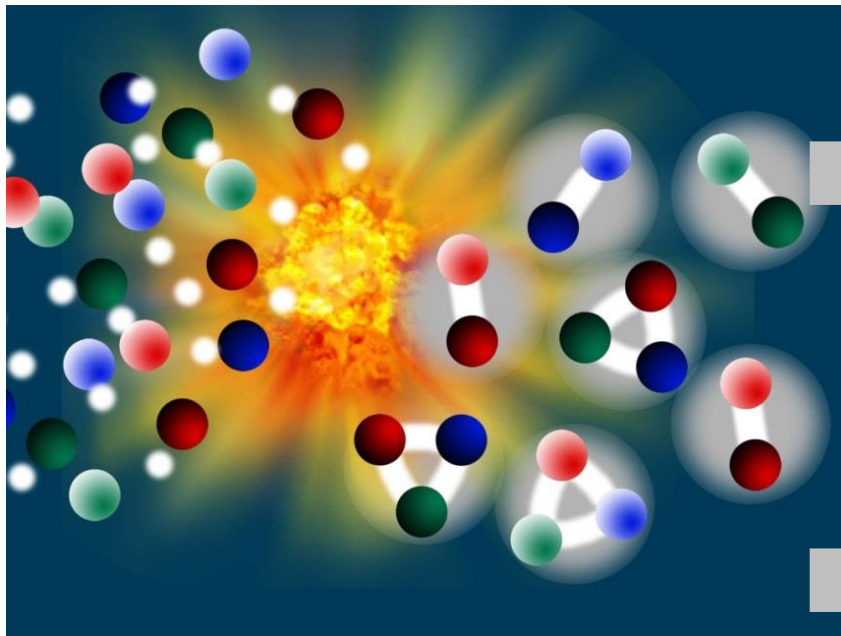
- Initial starting point $> T_0$: **annihilation of matter (p, n) & anti-matter ($\bar{p}, \bar{n}$)**



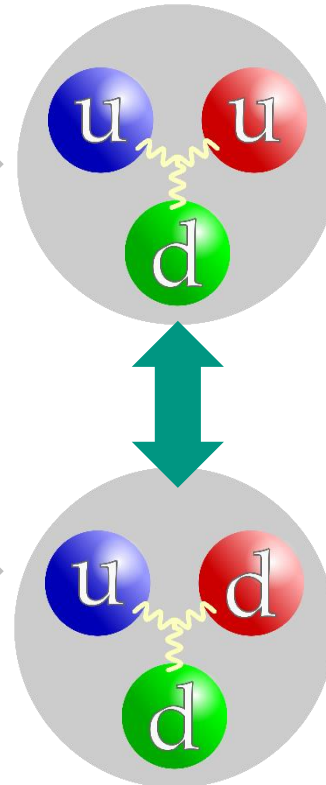
Primordial nucleosynthesis: phase 1

■ Thermodynamical equilibrium

- nucleons p, n condense out of initial quark-gluon plasma (hadronisation)



QCD phase transition
at $T \sim 150 \text{ MeV}$



p (proton)
 $m = 938.3 \text{ MeV}$

$W^\pm$ – bosons

n (neutron)
 $m = 939.6 \text{ MeV}$

$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$

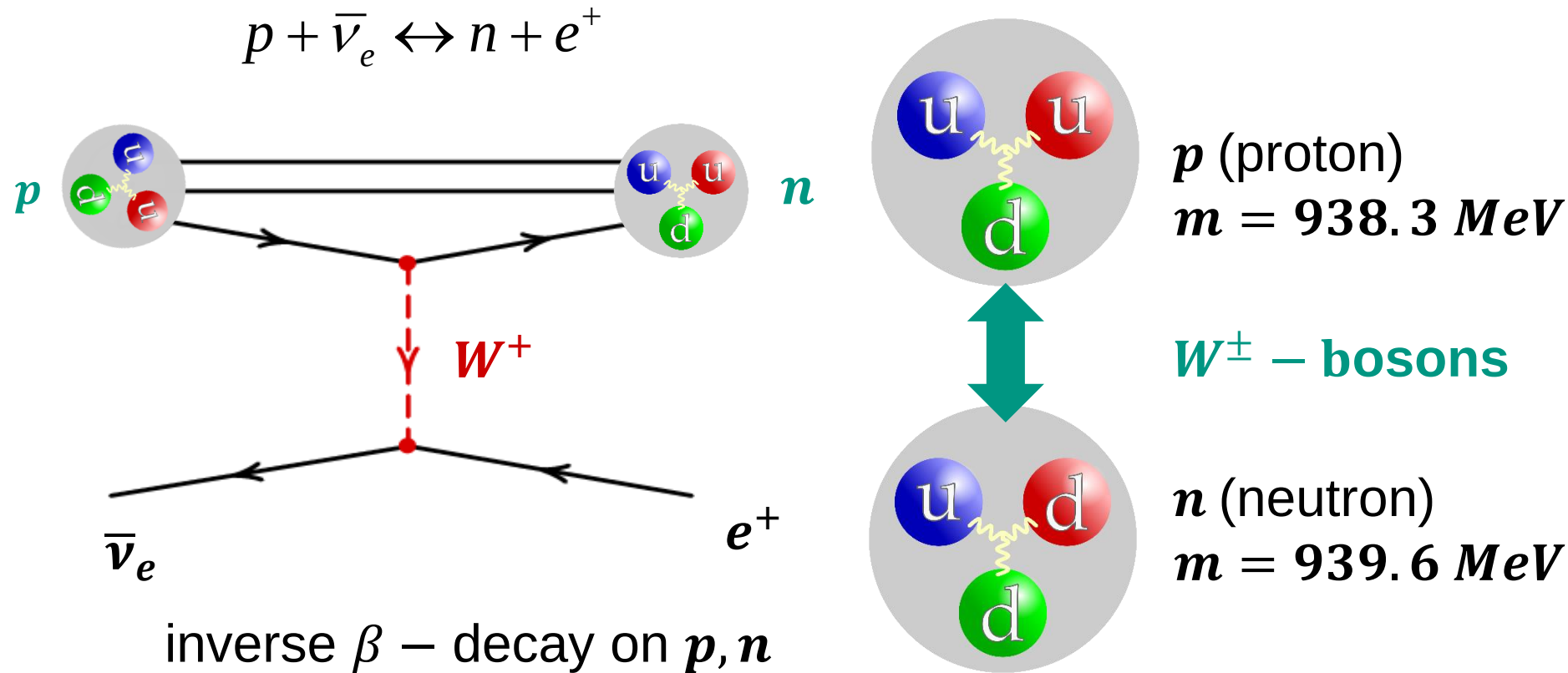


$$t = 0.01 \text{ s}$$

Primordial nucleosynthesis: phase 1

■ Thermodynamical equilibrium: example inverse β – decay

- nucleons p, n transform into each other due to weak interaction processes



$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$

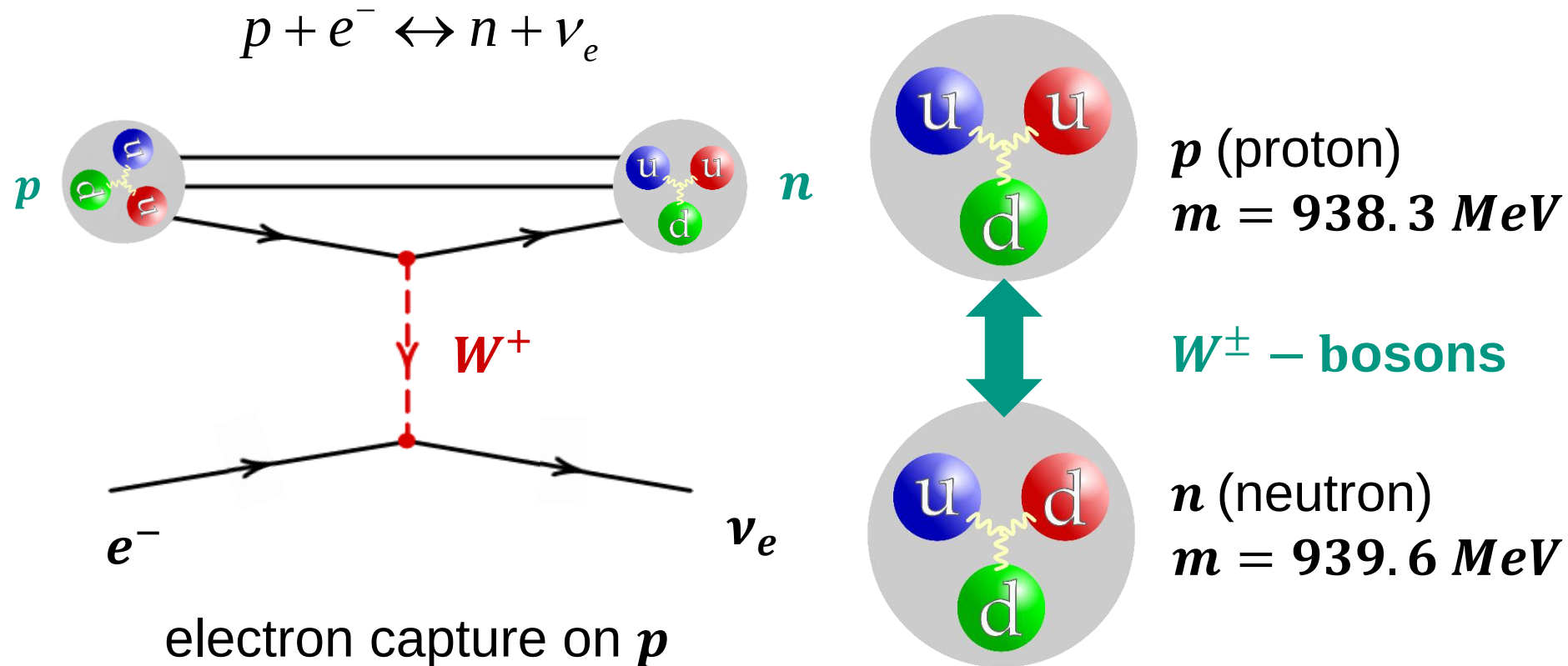


$$t = 0.01 \text{ s}$$

Primordial nucleosynthesis: phase 1

■ Thermodynamical equilibrium: example electron capture process

- nucleons p, n transform into each other due to weak interaction processes



$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$

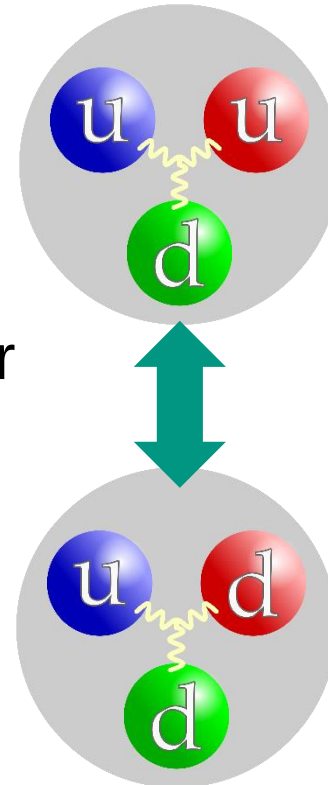


$$t = 0.01 \text{ s}$$

Primordial nucleosynthesis: phase 1

■ Thermodynamical equilibrium: example electron capture process

- nucleons p, n transform into each other due to weak interaction processes
- no change of (p, n) ratio via strong/electromagnetic interactions (strong isospin), **only via weak interactions**, i.e. ν –processes (before their decoupling!)
- early universe ($t < 1$ s): extremely high ν –densities*:
 $N_\nu \approx 10^{31} \dots 10^{32} \text{ cm}^{-3}$



p (proton)
 $m = 938.3 \text{ MeV}$

$W^\pm$ – bosons

n (neutron)
 $m = 939.6 \text{ MeV}$

$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$



$$t = 0.01 \text{ s}$$

Primordial nucleosynthesis: phase 1

■ Thermodynamical equilibrium: calculating initial (p, n) – ratio at high T

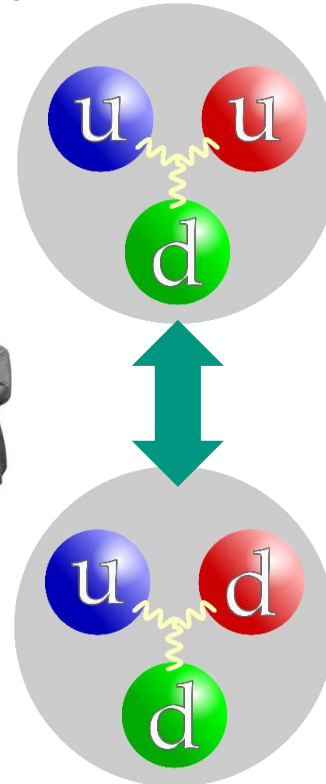
- (p, n) – ratio follows a **Boltzmann distribution**
different masses $\Delta m = 1.3 \text{ MeV}$!

- for **high temperatures**
 $T(10 \text{ MeV}) \gg \Delta m (1.3 \text{ MeV})$
we have:

$$n/p = e^{-\Delta m \cdot c^2 / k_B \cdot T}$$

$$n/p = e^{-1.3/10} = 0.88$$

⇒ almost equal numbers



p (proton)
 $m = 938.3 \text{ MeV}$

$W^\pm$ – bosons

n (neutron)
 $m = 939.6 \text{ MeV}$

$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$



$$t = 0.01 \text{ s}$$

Primordial nucleosynthesis: phase 1

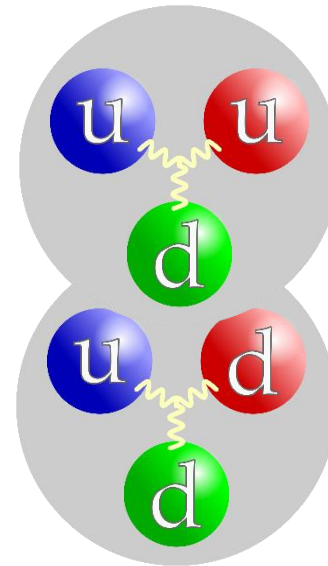
■ Thermodynamical equilibrium: fusion reactions to deuterium at high T ?

- can nucleons (p, n) already undergo fusion processes to deuterium?

- for **high temperatures**

$T(10 \text{ MeV})$

we have to consider
whether deuterium is
stable under these
conditions



d (deuterium)
 $E_B = 2.2 \text{ MeV}$

strong interaction

$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$



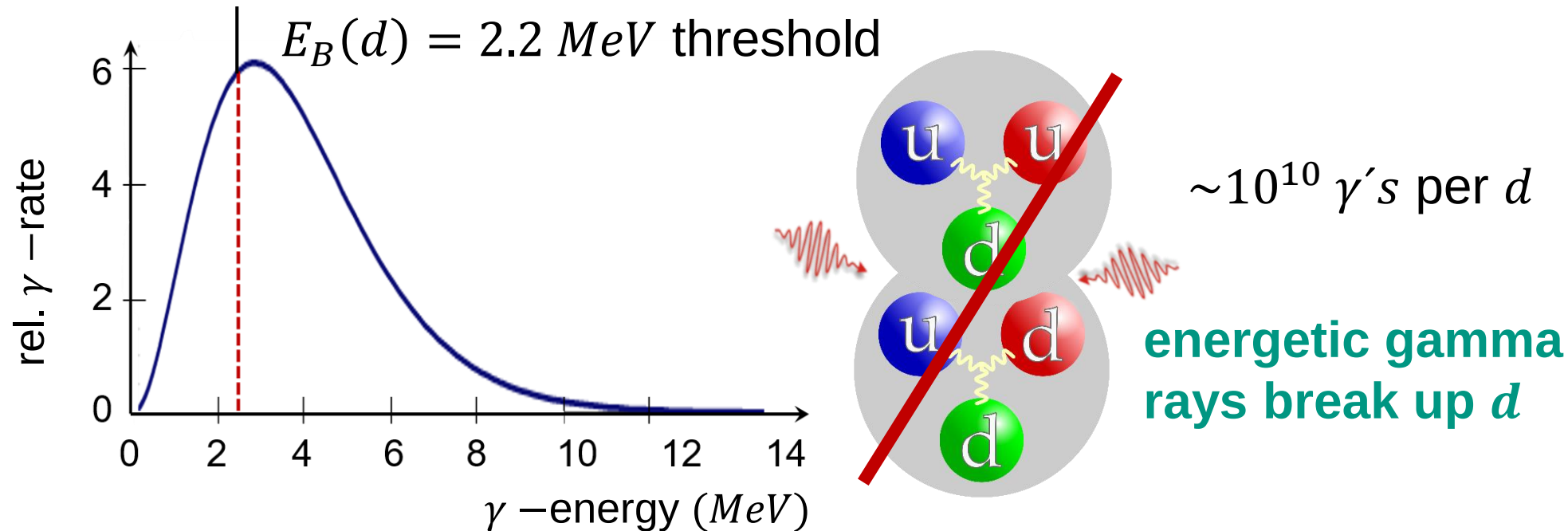
$$t = 0.01 \text{ s}$$

⇒ fusion reactions (?) in an **intense heat bath**

Primordial nucleosynthesis: phase 1

■ Thermodynamical equilibrium: fusion reactions to deuterium at high T ?

- all deuterium nuclei forged at $T = \text{MeV} - \text{scale}$ are immediately destroyed



$$T = 10^{11} \text{ K}$$

$$E = 10 \text{ MeV}$$



$$t = 0.01 \text{ s}$$

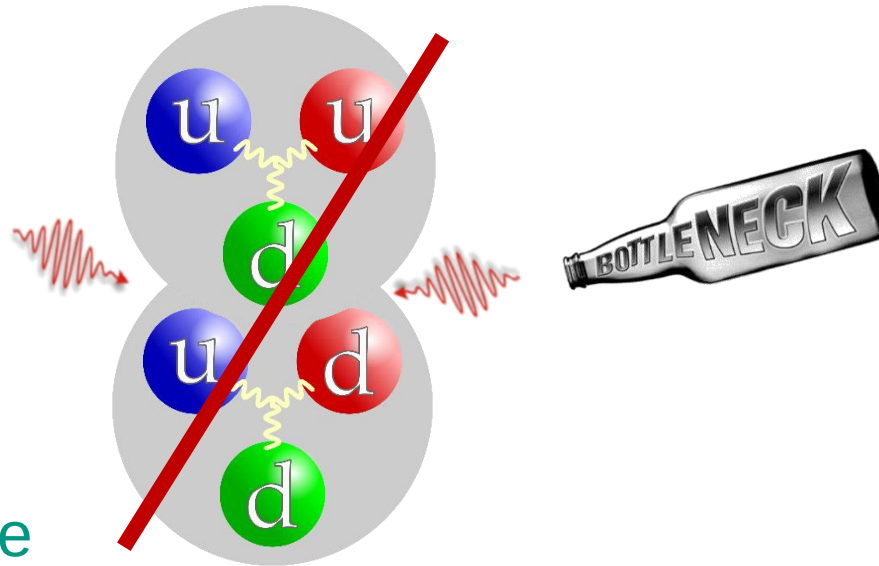
⇒ fusion reactions cannot proceed in an **intense heat bath**
nucleons (p, n) remain unpaired

Primordial nucleosynthesis: phase 2

■ Universe has to expand & cool before fusion reactions can proceed

- d is **bottleneck** for fusion reactions due to *energetic γ 's at MeV – scale*

- only 'solution':



- waiting for temperature T
to drop: *neutrinos decouple
after $t = 1\text{ s}$ (1 MeV)*

$\Rightarrow$ no further (n, p) transformations $\Rightarrow (n, p)$ – ratio is fixed

$$T \ll 10^{11}\text{ K}$$

$$E \ll 10\text{ MeV}$$



$$t \gg 0.01\text{ s}$$

Primordial nucleosynthesis: phase 2

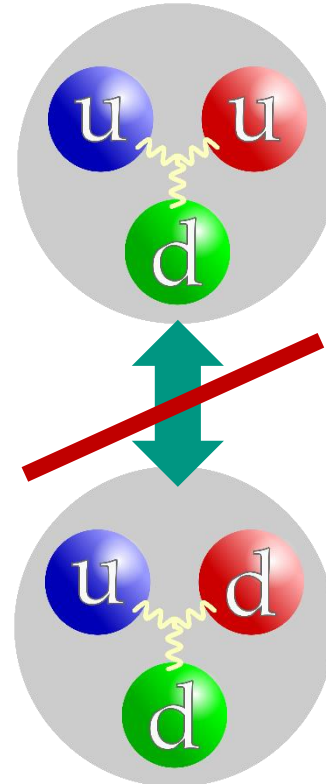
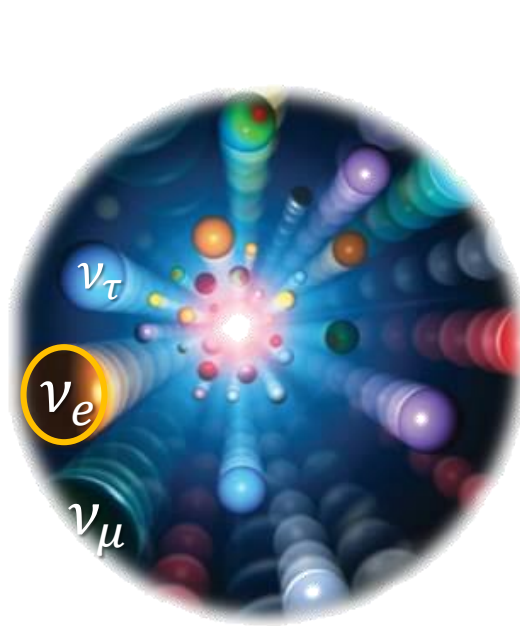
■ Freeze-out of neutrinos at $t \approx 1\text{ s}$ of key importance for (n, p) –ratio

- detailed calculations for neutrino freeze-out energy for ν_e

$$\Gamma_\nu(t) = H(t)$$

$$\sim G_F \cdot T^5 \quad \sim T^2$$

$$T_{fr} = 0.7\text{ MeV}$$



p (proton)
 $m = 938.3\text{ MeV}$

~~$W^\pm$ - bosons~~

n (neutron)
 $m = 939.6\text{ MeV}$

$$T \approx 10^{10}\text{ K}$$

$$E = 0.7\text{ MeV}$$



$$t \approx 1\text{ s}$$

Primordial nucleosynthesis: phase 2

■ Freeze-out of neutrinos at $t \approx 1\text{ s}$ of key importance for (n, p) –ratio

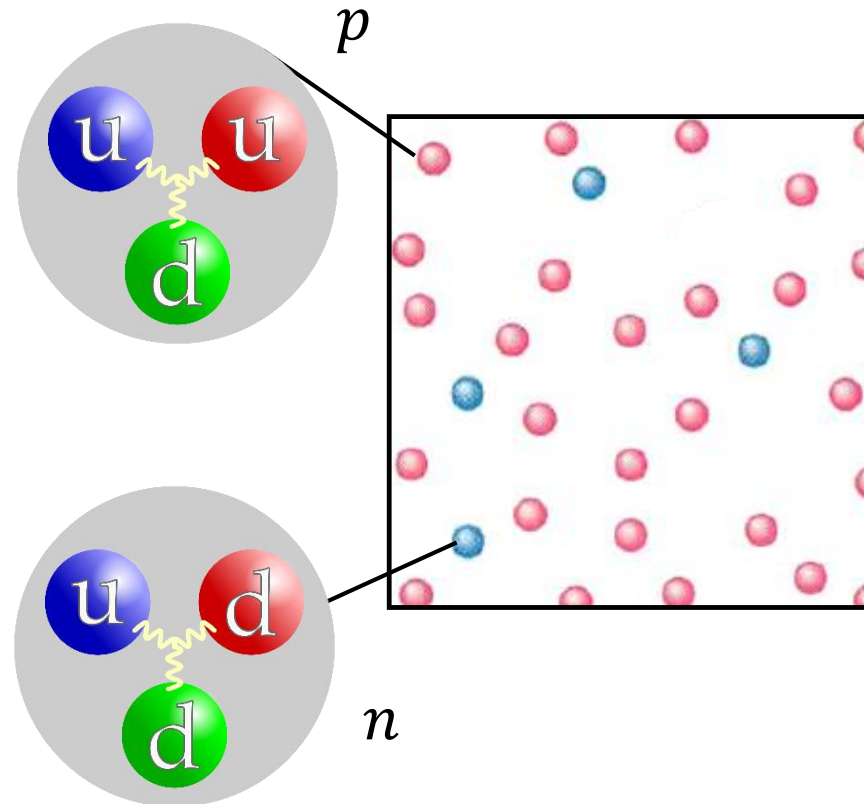
- calculated (n, p) – ratio at $T = 0.7\text{ MeV}$

$$n/p = e^{-\Delta m \cdot c^2 / k_B \cdot T}$$

$$n/p = e^{-1.3/0.7} = 0.15 \approx \frac{1}{6}$$

- (n, p) – ratio not further modified by thermodynamics

FIXED



$$T \approx 10^{10}\text{ K}$$

$$E = 0.7\text{ MeV}$$



$$t \approx 1\text{ s}$$

INSERTION: Primordial nucleosynthesis'

- We consider a different universe (#42) with larger freeze-out temperature

- in case of significantly changed physical constants we have for example:

$$T_{fr} = 70 \text{ MeV}$$

$$n/p = e^{-1.3/70} = 0.981 \approx 1$$

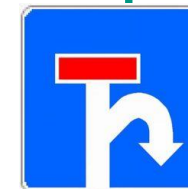
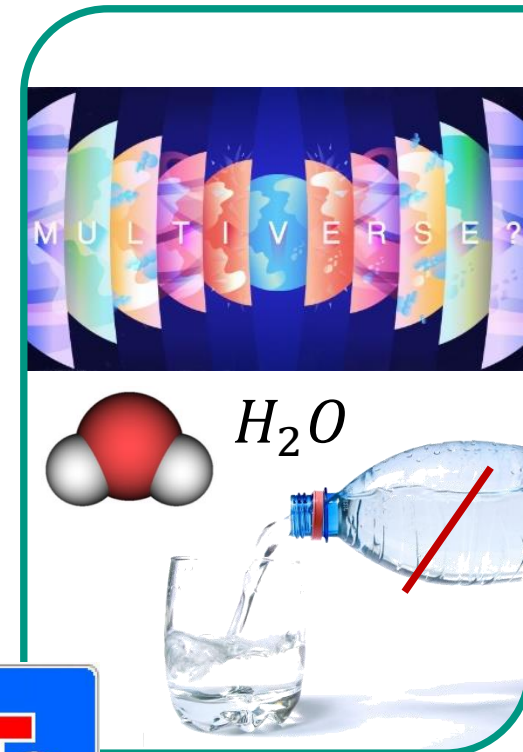
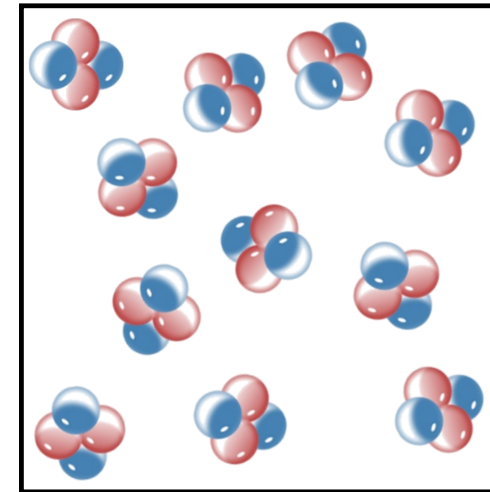


formation of ${}^4\text{He}$

- no free protons p remain to form

molecules: H_2O and CH_2

molecules: H_2 (only fast-burning He-stars...)

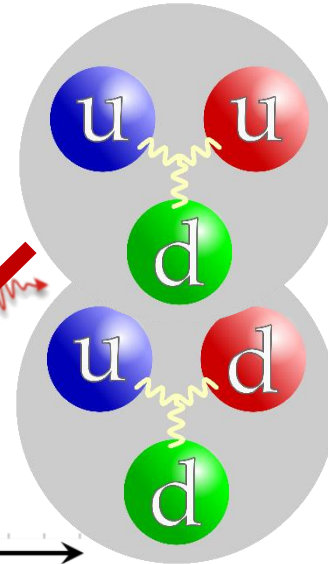
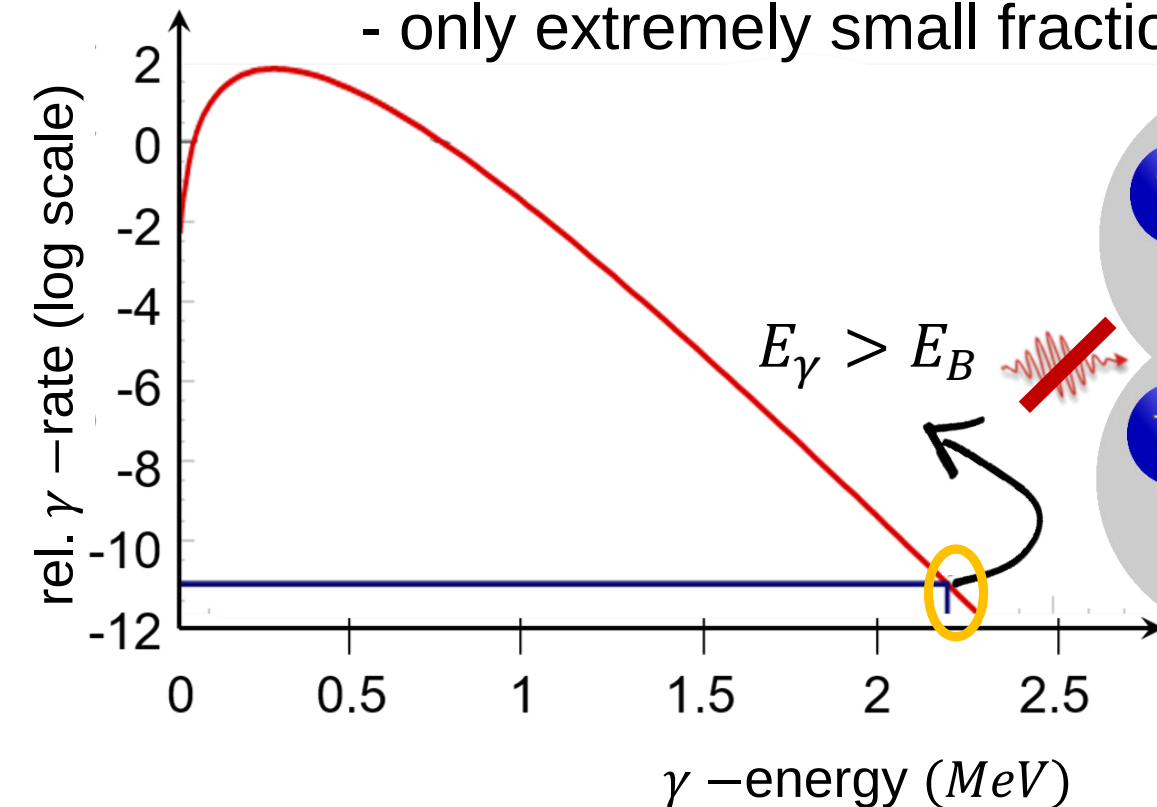


Primordial nucleosynthesis: phase 2 ends

■ We need to wait for the universe to cool and expand, until d can be formed

- only after $t = 60\text{ s}$ the universe is cool enough ($T < 0.1\text{ MeV}$) so that d is stable

- only extremely small fraction of γ 's with $E > 2.2\text{ MeV}$!



d (deuterium)
 $E_B = 2.2\text{ MeV}$
is finally stable!



$$T < 10^9\text{ K}$$

$$E \approx 0.1\text{ MeV}$$

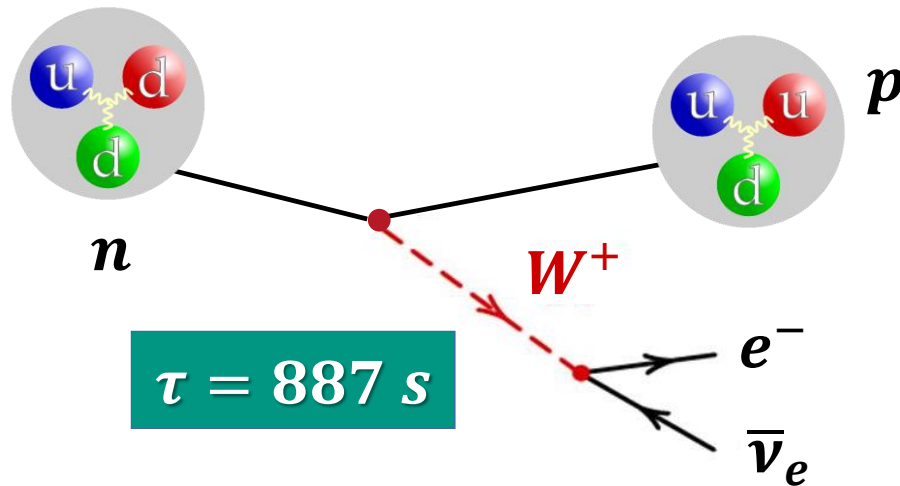


$$t \approx 60\text{ s}$$

Primordial nucleosynthesis: phase 2 ends

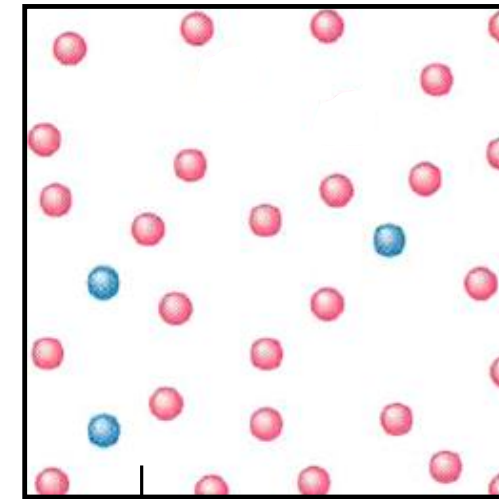
■ We now need to take into account the finite life-time of neutrons

- after $t = 60\text{ s}$ a fraction of neutrons will have decayed (β – decay)



- (n, p) – ratio after $t = 60\text{ s}$ decreases to

$$\frac{n}{p} = \frac{1}{6} \cdot e^{-60/887} \approx \frac{1}{7}$$



(n, p) – ratio now at 1/7

$$T < 10^9\text{ K}$$

$$E \approx 0.1\text{ MeV}$$

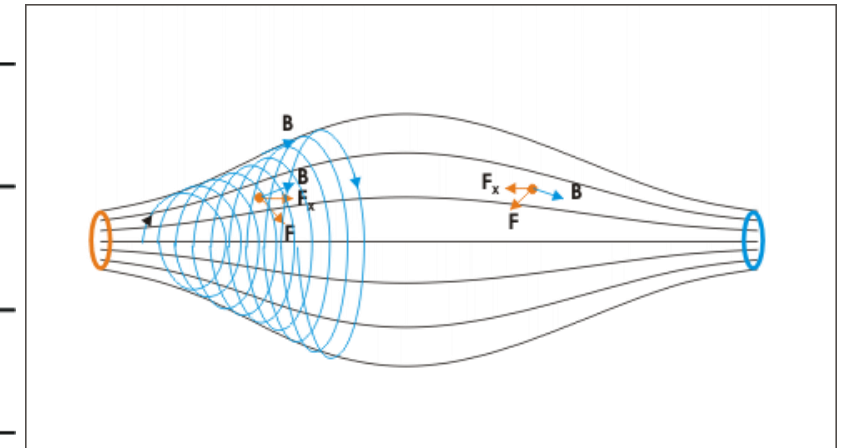
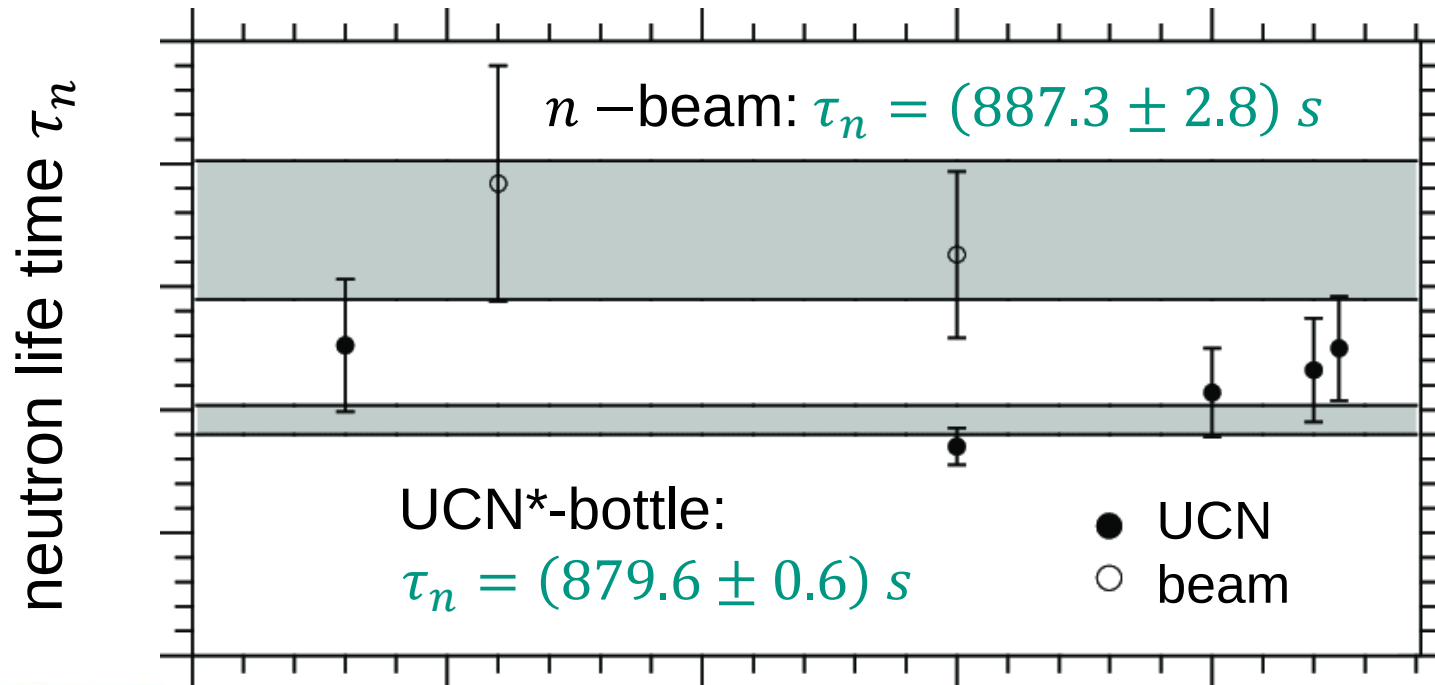


$$t \approx 60\text{ s}$$

INSERTION: Primordial nucleosynthesis

■ finite life-time of neutrons: overview of experimental results

- systematic effects in measurements of n – life times τ_n over past decades



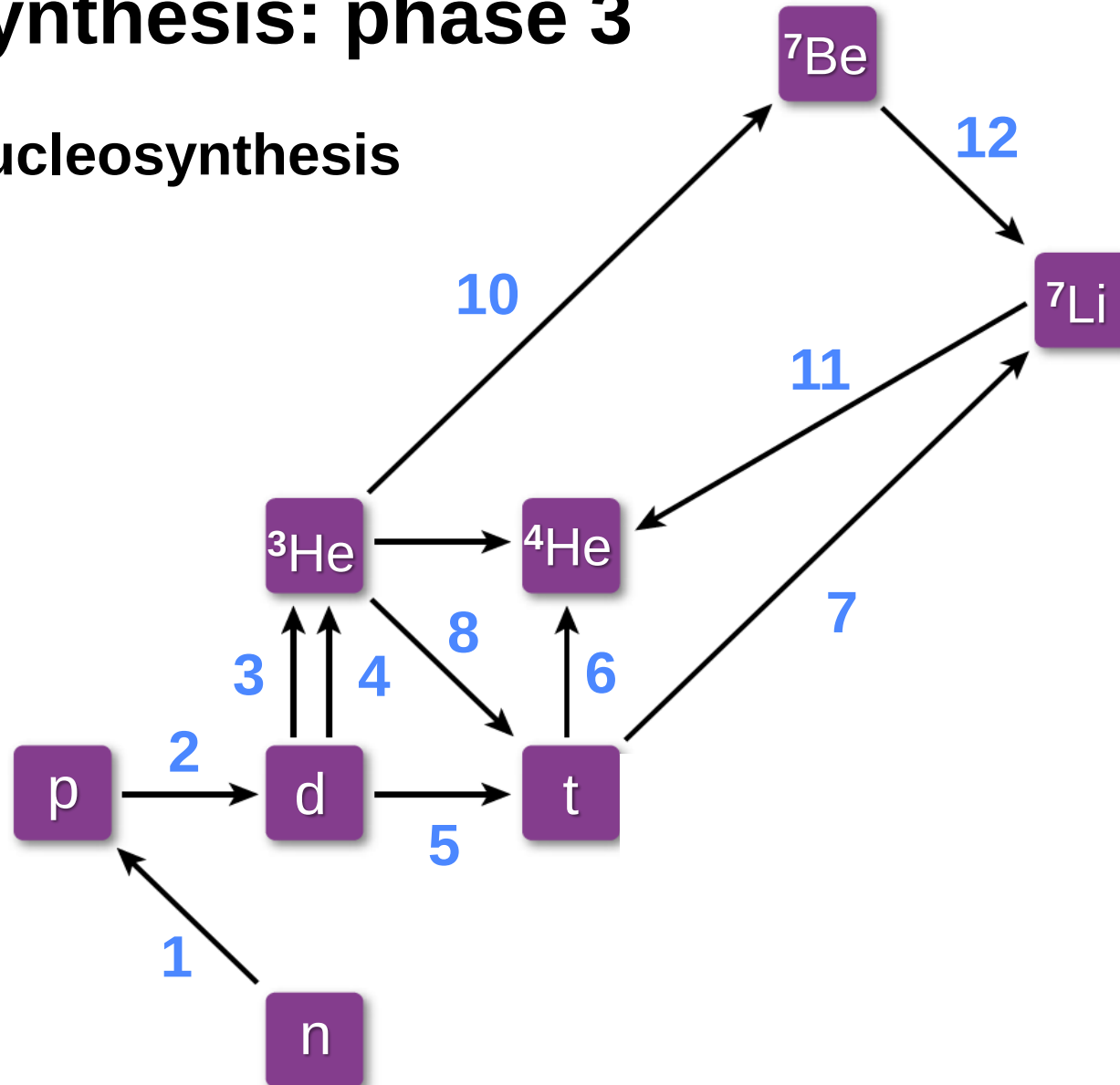
magnetic bottle



Primordial nucleosynthesis: phase 3

■ 12 main pathways of nucleosynthesis

- 1 $n \rightarrow {}^1\text{H} + e^- + \bar{\nu}$
- 2 ${}^1\text{H} + n \rightarrow {}^2\text{H} + \gamma$
- 3 ${}^2\text{H} + {}^1\text{H} \rightarrow {}^3\text{He} + \gamma$
- 4 ${}^2\text{H} + {}^2\text{H} \rightarrow {}^3\text{He} + n$
- 5 ${}^2\text{H} + {}^2\text{H} \rightarrow {}^3\text{H} + {}^1\text{H}$
- 6 ${}^2\text{H} + {}^3\text{H} \rightarrow {}^4\text{He} + n$
- 7 ${}^3\text{H} + {}^4\text{He} \rightarrow {}^7\text{Li} + \gamma$
- 8 ${}^3\text{He} + n \rightarrow {}^3\text{H} + {}^1\text{H}$
- 9 ${}^3\text{He} + {}^2\text{H} \rightarrow {}^4\text{He} + {}^1\text{H}$
- 10 ${}^3\text{He} + {}^4\text{He} \rightarrow {}^7\text{Be} + \gamma$
- 11 ${}^7\text{Li} + {}^1\text{H} \rightarrow {}^4\text{He} + {}^4\text{He}$
- 12 ${}^7\text{Be} + n \rightarrow {}^7\text{Li} + {}^1\text{H}$



Primordial nucleosynthesis: phase 3

■ all pathways of nucleosynthesis



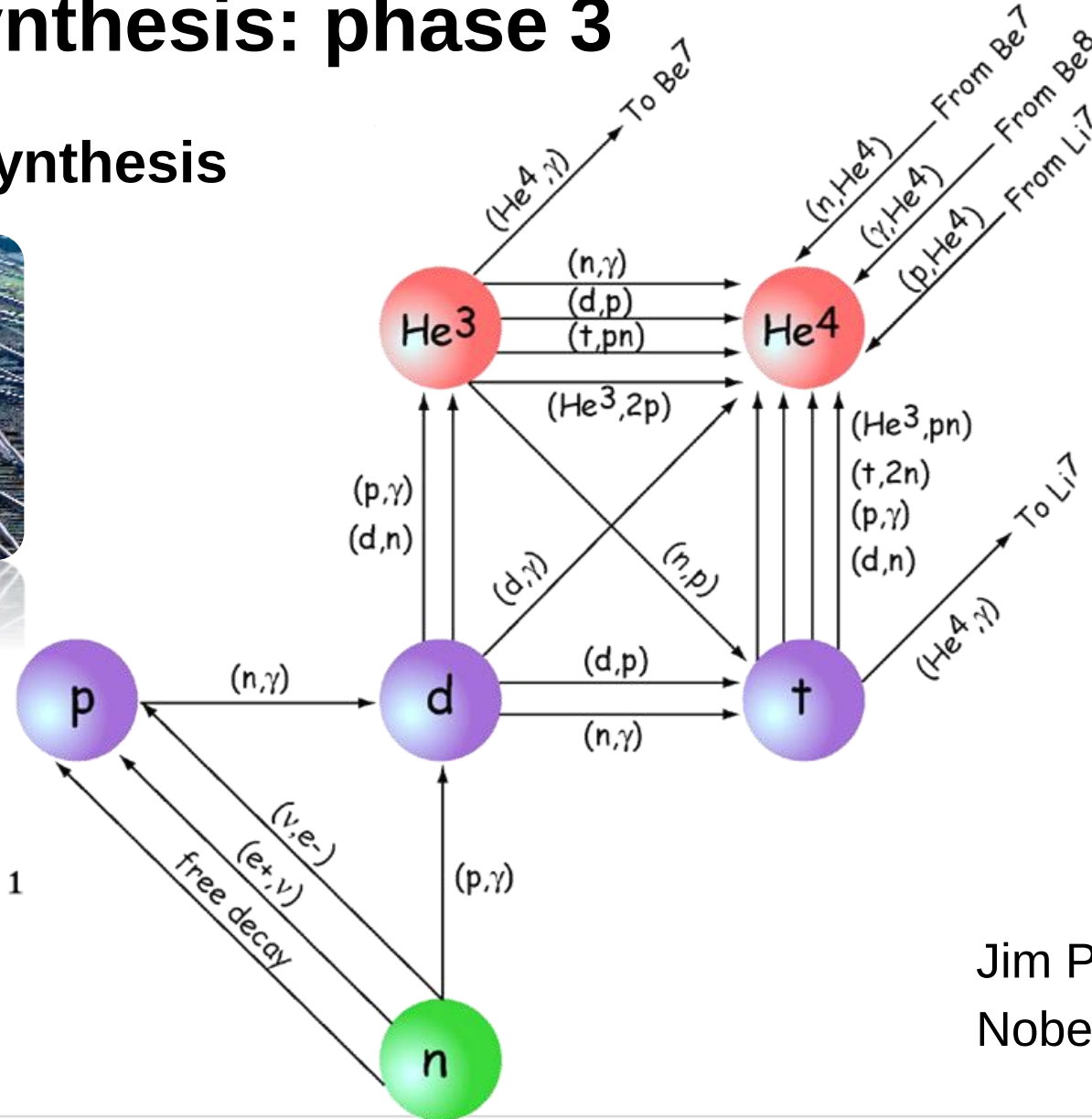
*system of
coupled
differential
equations*



Jim Peebles
Nobel prize 2019

$$N_A \sim (\sigma_A + \tau_0^{-1})^{-1}$$

$$N_{A'} \sim \left(\sigma_{A'} + \frac{\lambda_\beta}{\sigma_{A'} \Phi} + \tau_0^{-1} \right)^{-1}$$



Primordial nucleosynthesis: phase 3 channels

■ Grouping the different reaction pathways into channels (reaction ‘ladder’)

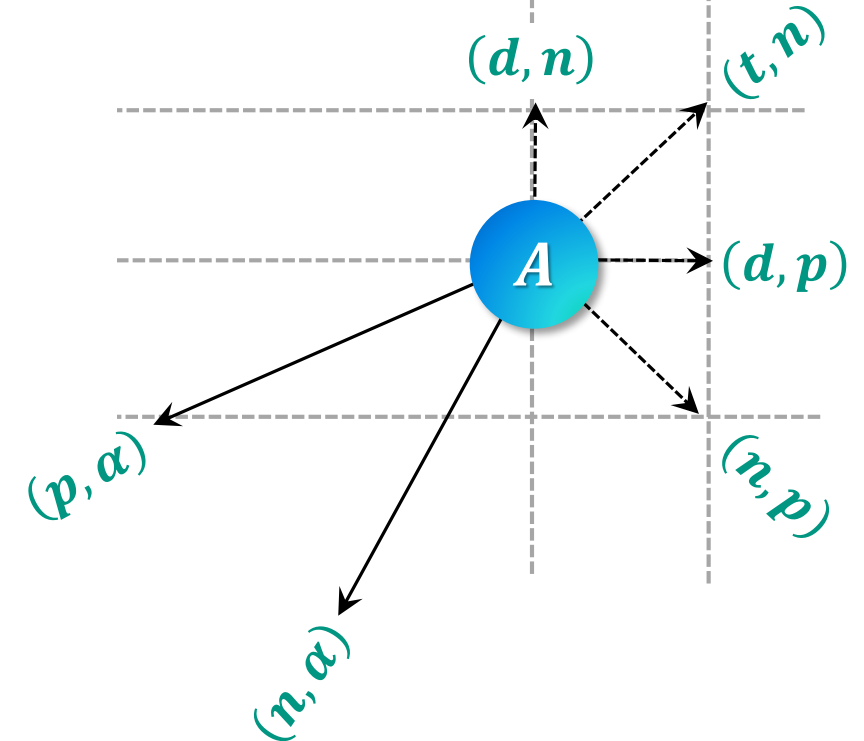
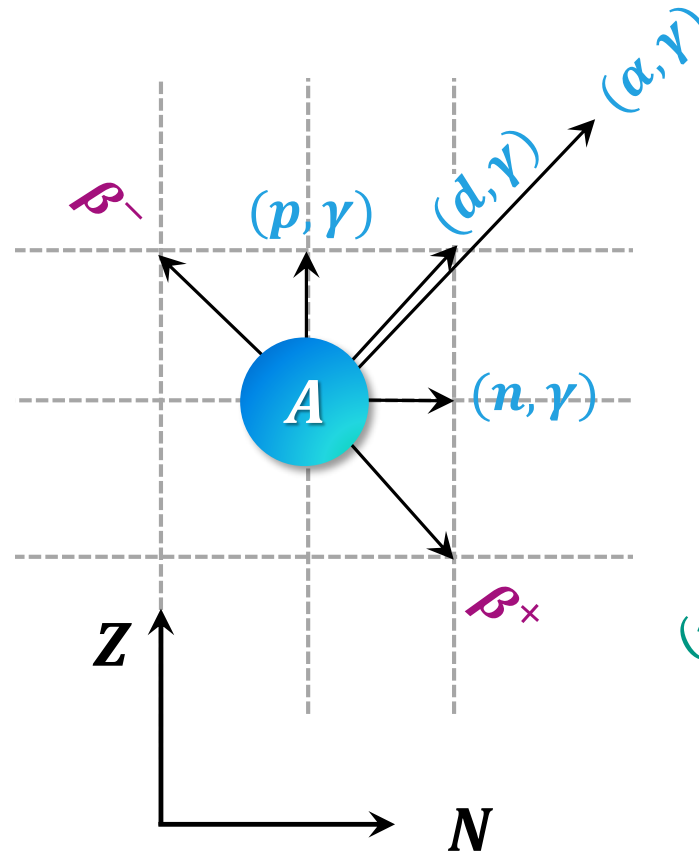
$\beta^\pm$ –decays

capture reactions:

$(p, \gamma), (n, \gamma),$
 $(d, \gamma), (\alpha, \gamma)$

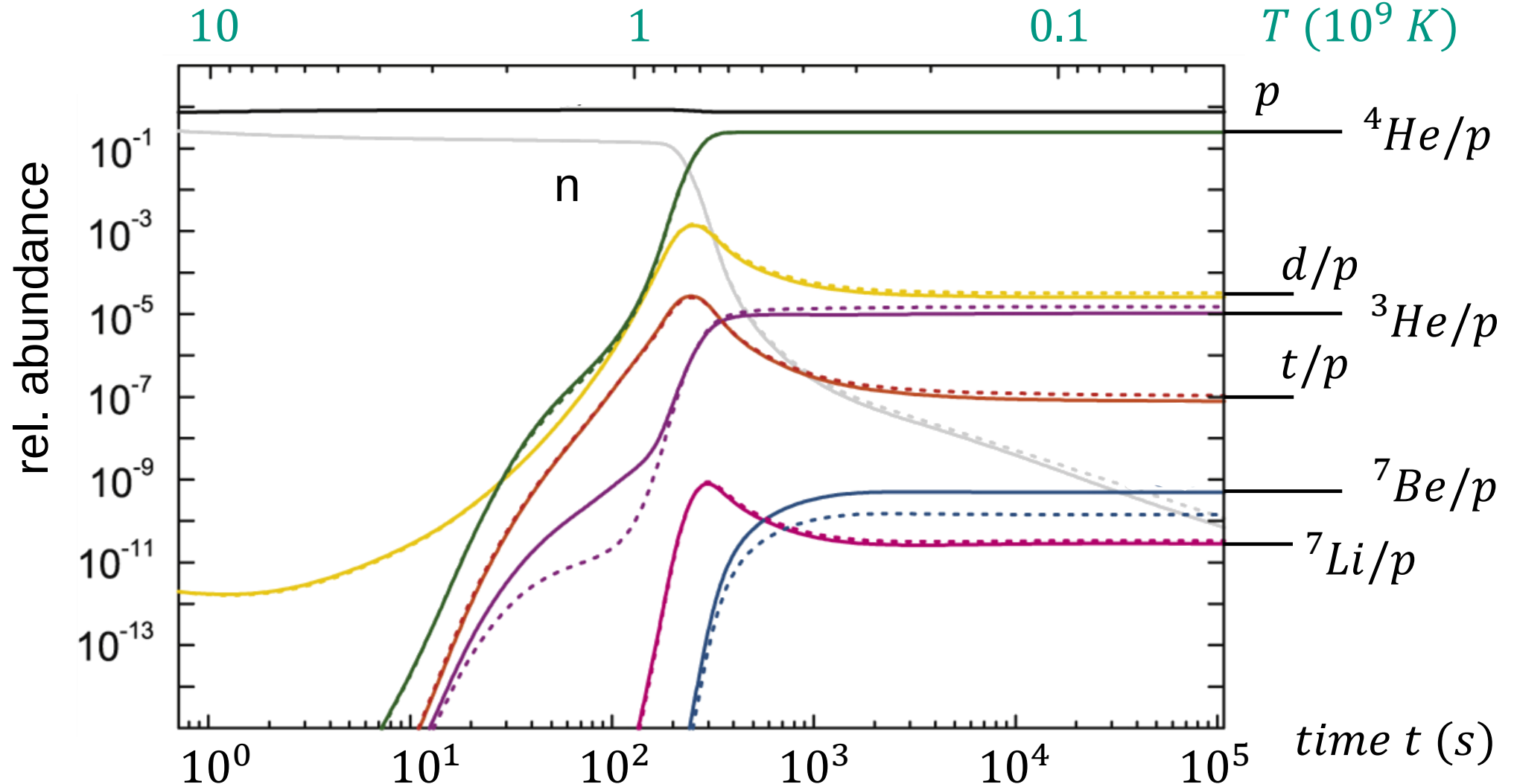
transfer processes:

$(n, p), (d, p),$
 $(d, n), (t, n), \dots$



Primordial nucleosynthesis: fusion yields

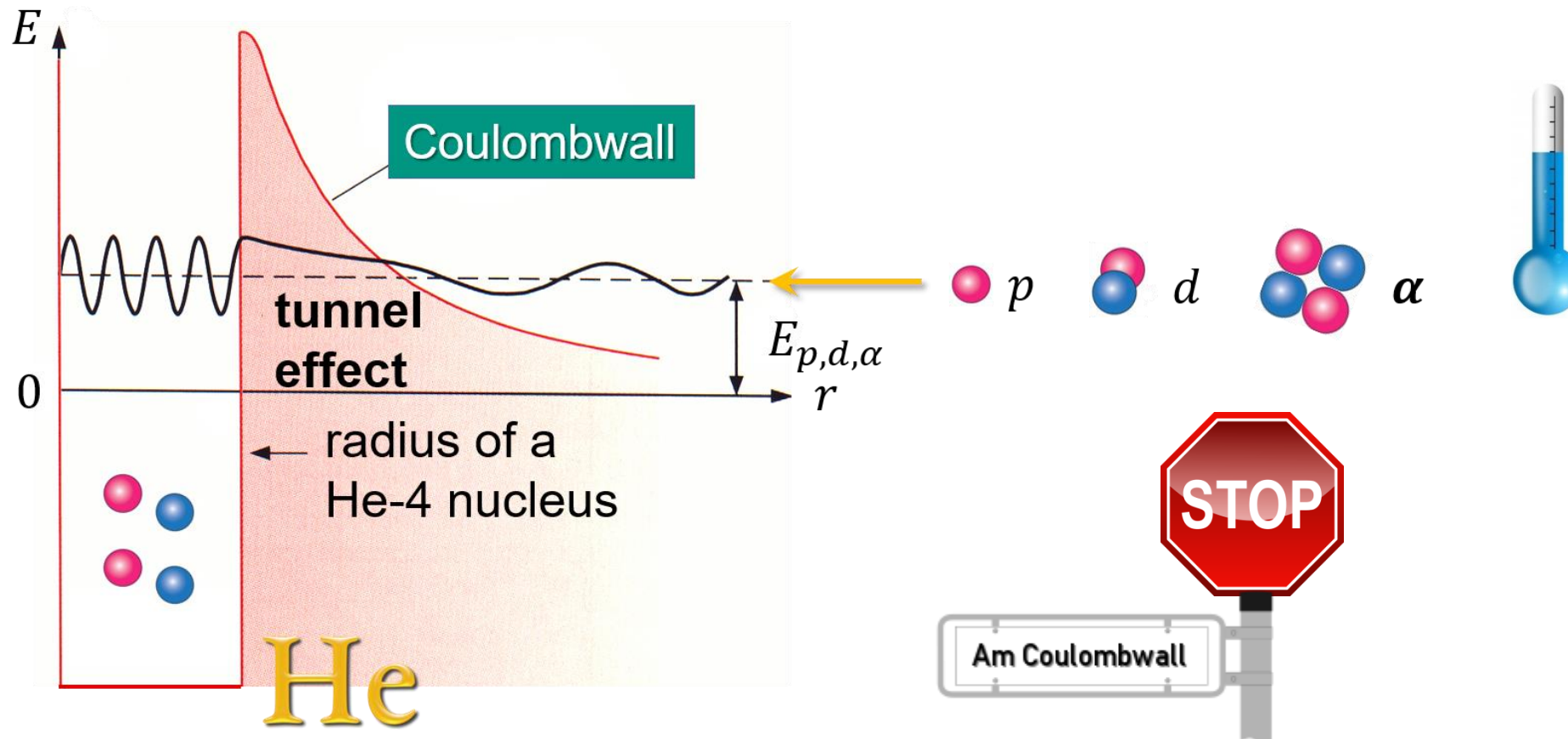
■ Results



Primordial nucleosynthesis: fusion stops

■ Synthesis of light elements at low temperature ends due to Coulomb barrier

- energy $E_{p,d,\alpha}$ of **charged particles** no longer sufficient to sustain fusion reactions



$$T < 5 \cdot 10^8 K$$

$$E \ll 0.1 \text{ MeV}$$



$$t \approx 300 \text{ s}$$

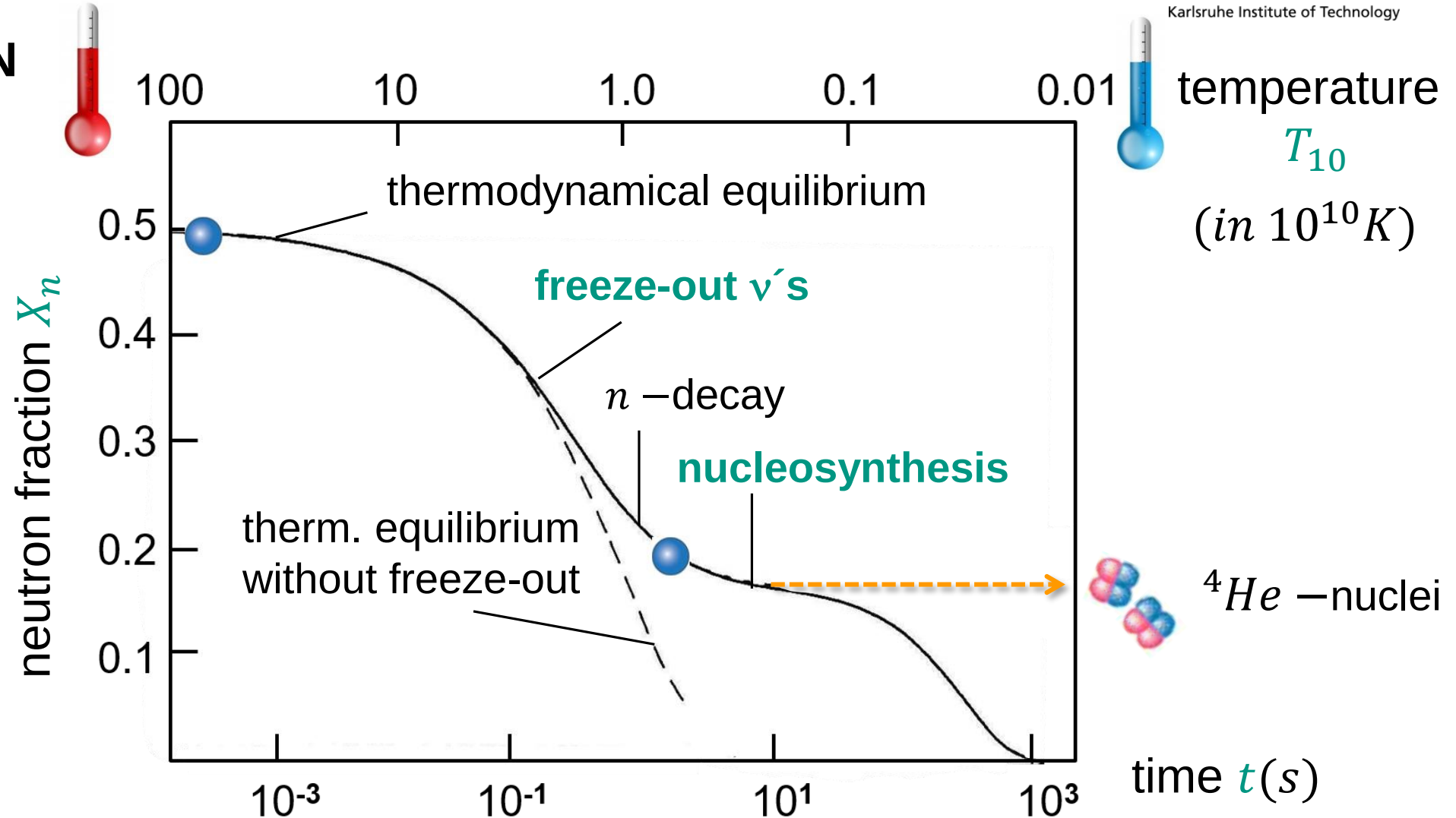
Primordial nucleosynthesis: all steps

■ Overview BBN

- fraction X_n of free n

$$X_n = \frac{N(n)}{N(p) + N(n)}$$

- all neutrons after $t > 60\text{ s}$ end up in ${}^4\text{He}$



Primordial Nucleosynthesis & parameter η

■ Fusion yield of light elements depends on the size of baryon asymmetry η

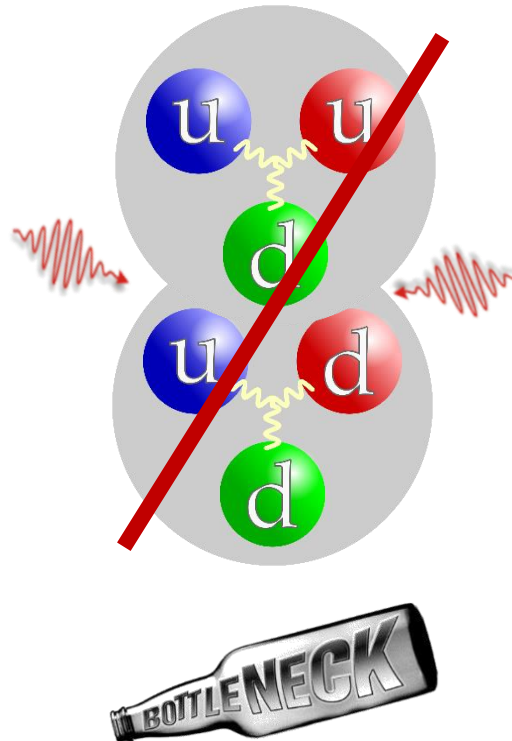
- theoretical estimates of the light element fusion yields have been performed for different values of η :

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma}$$

$$\eta = 10^{-6} \dots 10^{-12}$$

η

different
fusion
yields



observed precision
fusion yields of light
elements in BBN

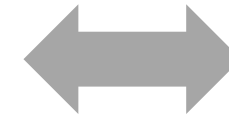
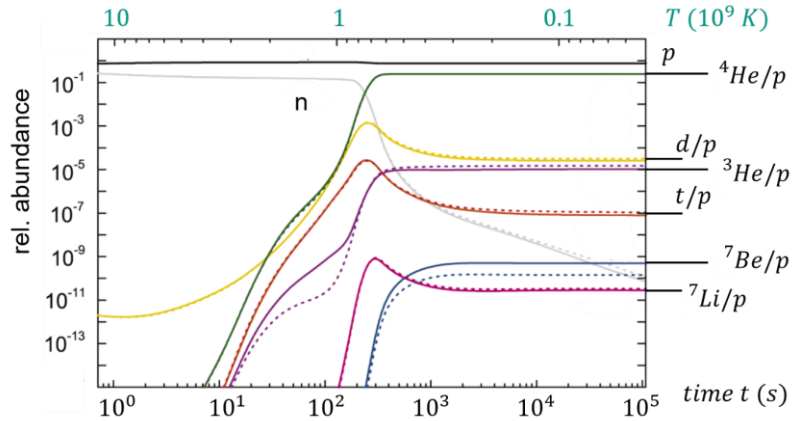


estimates of
baryon asymmetry
parameter η

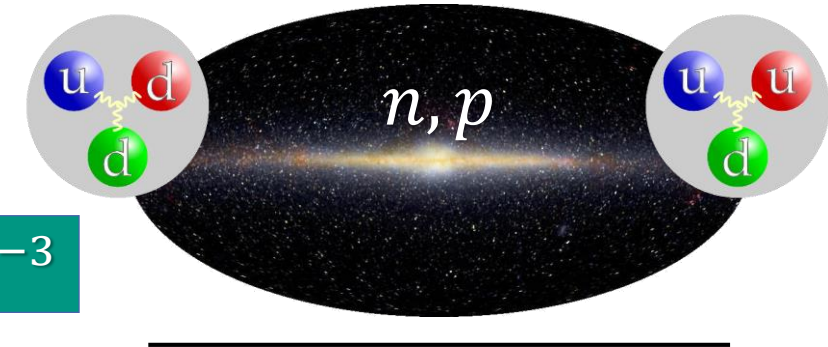
Primordial Nucleosynthesis & baryon density ρ_B

■ Fusion yield of light elements depends on the value of the **baryon density ρ_B**

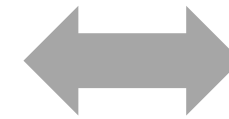
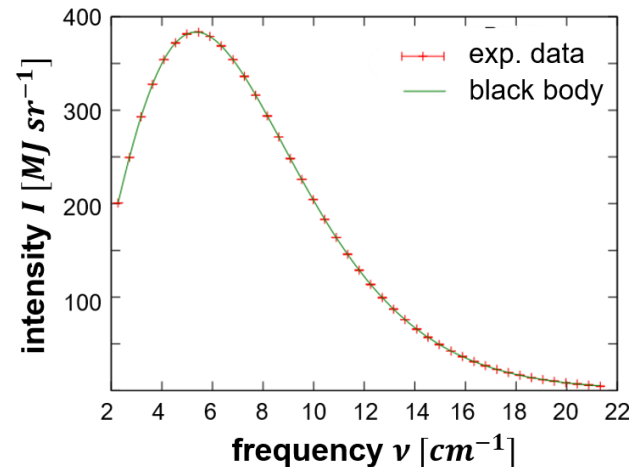
BBN
calculation



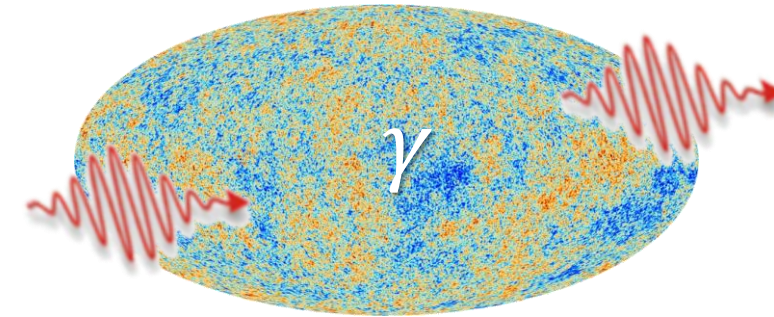
$$\rho_B = ? \text{ cm}^{-3}$$



Planck
distribution



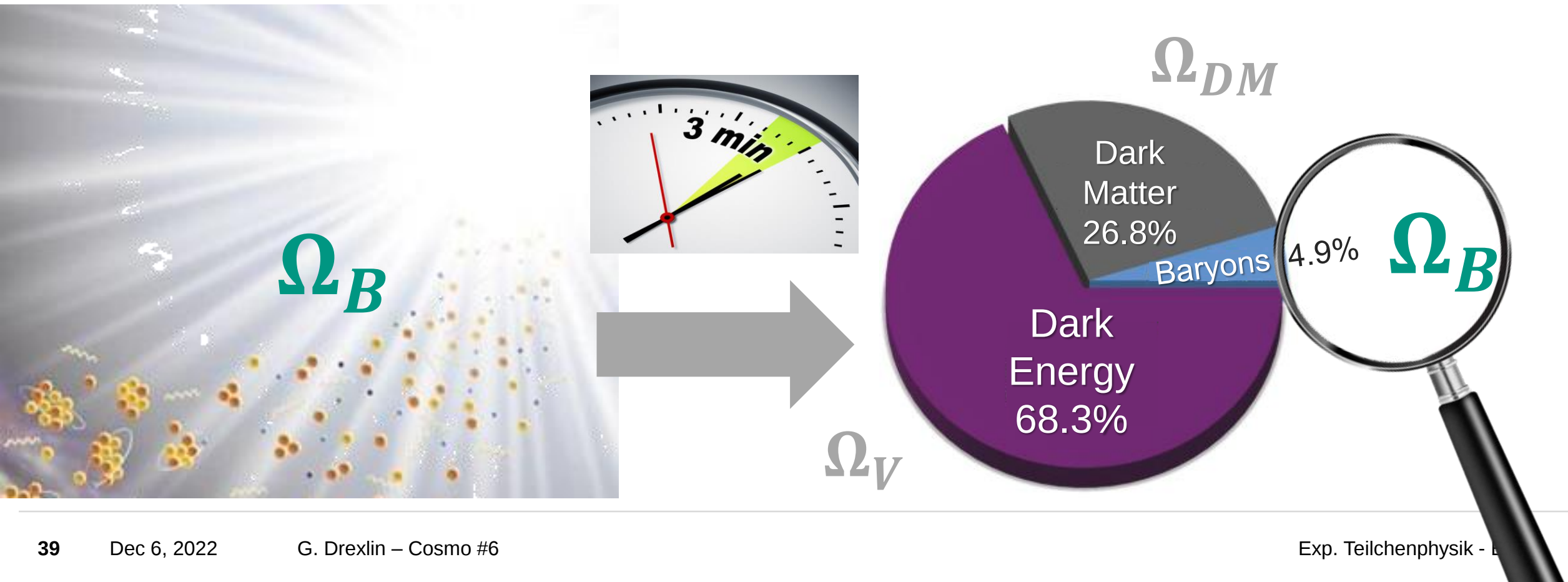
$$n_{\text{photons}} = (412 \pm 2) \text{ cm}^{-3}$$



Primordial Nucleosynthesis & baryon density Ω_B

■ Extracting the baryon density parameter Ω_B from BBN synthesis of elements

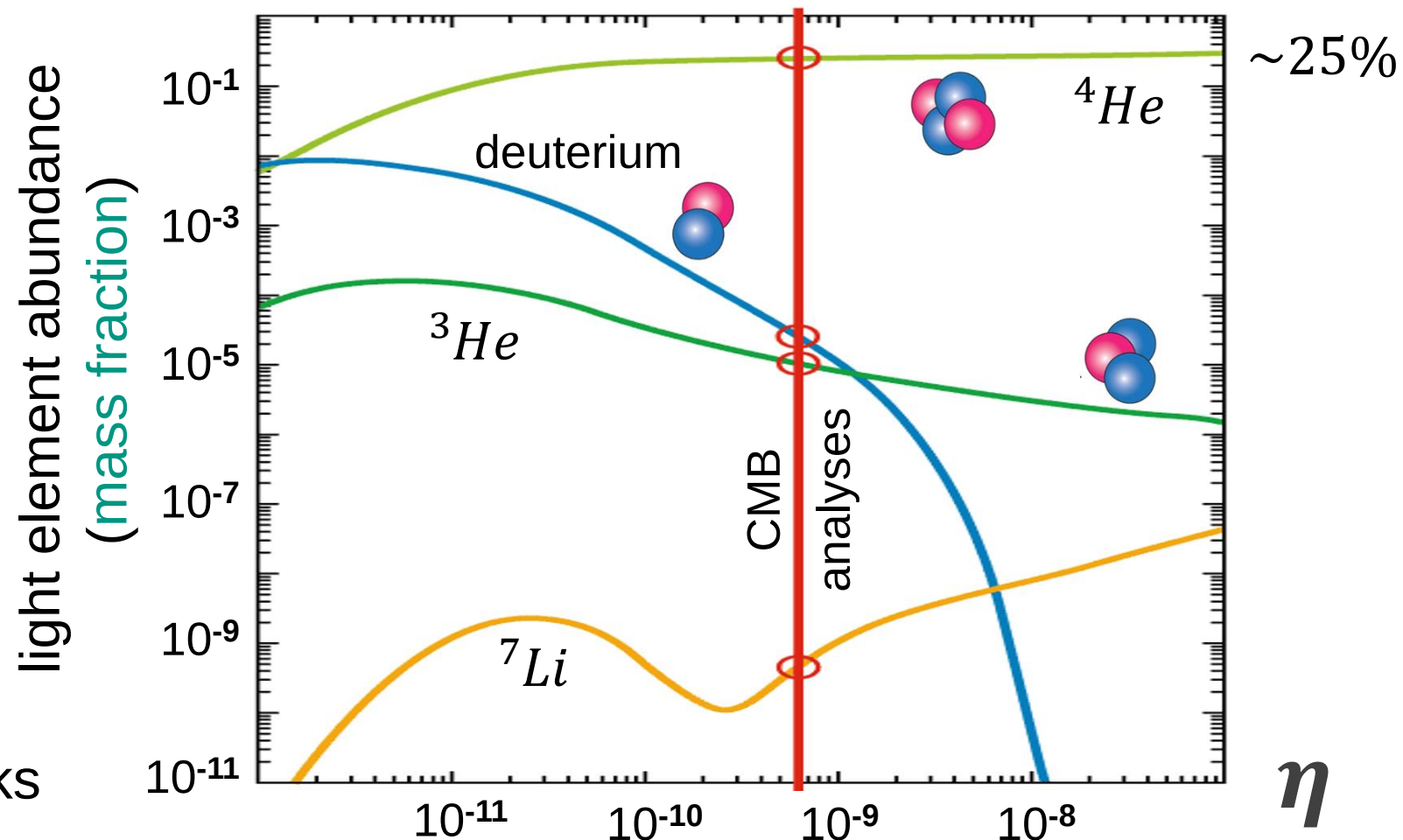
- a spectroscopic determination of light elements d , ${}^3\text{He}$, ${}^{6,7}\text{Li}$ determines Ω_B



Primordial nucleosynthesis - calculations

■ Results of fusion yields for light elements as function of parameter η

- fusion yield of ${}^4\text{He}$ is **dominating** & increases for a larger baryon asymmetry η
- yield of **other isotopes** is much smaller & covers many orders of magnitude
- yield of **each isotope as $f(\eta)$ is unique** and thus allows important cross-checks

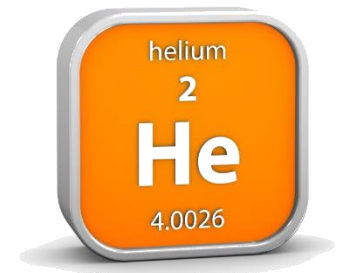
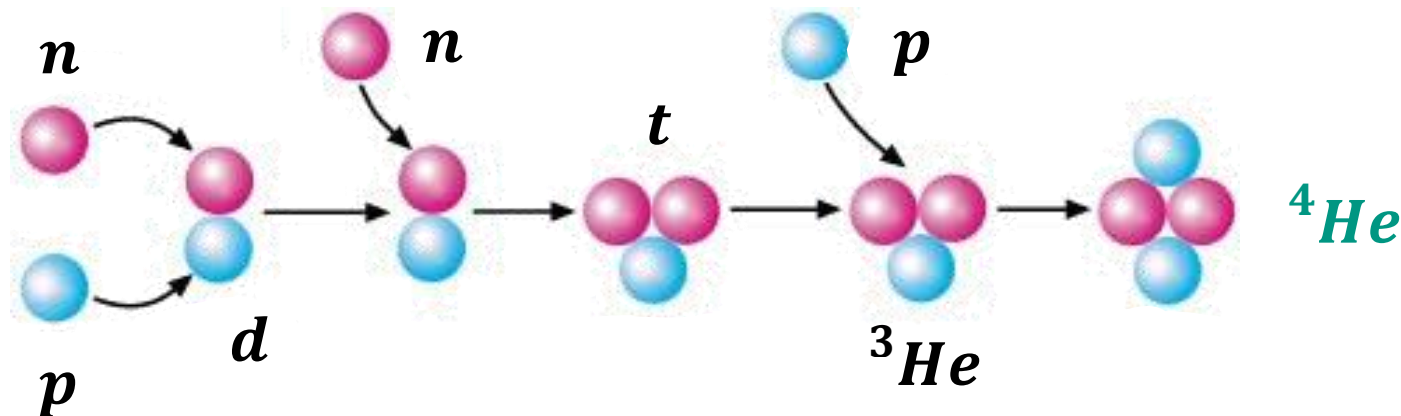


Primordial nucleosynthesis & yield of ${}^4\text{He}$

■ Synthesis of light elements essentially stops at ${}^4\text{He}$

- fusion beyond ${}^4\text{He}$ difficult, as this is a **very stable** (doubly magical) **nucleus**
- ${}^4\text{He}$ is **difficult to destroy**, also tiny cross-sections for capture of p, n, d
- almost **all neutrons are incorporated into ${}^4\text{He}$** (n limit its formation)
- in literature: primordial ${}^4\text{He}$ – **mass fraction Y_p**

one of many
BBN reaction
pathways to
 ${}^4\text{He}$ -formation



Primordial nucleosynthesis & yield of ${}^4\text{He}$

■ A quick estimate on the ${}^4\text{He}$ –mass fraction Y_p as function of (n, p) ratio

- we start with the calculated (n, p) –ratio of 1 : 7 at the onset of BBN
- we then form ${}^4\text{He}$ via the initial ratio out of $(2n + 2p) + 12 p = {}^4\text{He} + 12 p$

$$2 n : 14 p$$

- we now calculate the ${}^4\text{He}$ mass fraction $Y_p = {}^4\text{He} / (2 n + 14 p)$

$$\text{with } m_n = m_p = m_{nucl} \text{ we have} \quad = 4 m_{nucl} / 16 m_{nucl} = 0.25$$

- for an (n, p) –ratio of 1 : 7 we expect a ${}^4\text{He}$ mass fraction $Y_p = 0.25$