

Problem set 1

Submission deadline: 1 May, 09:45
Discussion of solutions: 4 May, 11:30

Problem 1: Symmetries of the Riemann tensor

- a) Using the explicit form of the Christoffel symbol in terms of the metric, show that $\Gamma_{\nu\rho}^{\mu}$ is symmetric in the two lower indices: $\Gamma_{\nu\rho}^{\mu} = \Gamma_{\rho\nu}^{\mu}$.
- b) Using the explicit form of $R_{\nu\rho\sigma}^{\lambda}$ in terms of Christoffel symbols, show that

$$R_{\mu\nu\rho\sigma} = -R_{\mu\nu\sigma\rho} . \quad (1)$$

Let us define the fully covariant Riemann tensor

$$R_{\mu\nu\rho\sigma} = g_{\mu\lambda} R_{\nu\rho\sigma}^{\lambda} . \quad (2)$$

Using a so-called *locally inertial frame*, it can be shown that this tensor satisfies the block interchange symmetry $R_{\mu\nu\rho\sigma} = R_{\rho\sigma\mu\nu}$.

- c) Show that $R_{\mu\nu\rho\sigma} = -R_{\nu\mu\rho\sigma}$.
- d) Conclude that in two dimensions (i.e. for $\mu, \nu, \rho, \sigma \in \{0, 1\}$) the Riemann tensor is fully determined by the component R_{1010} .
- e) How many independent components are there in three dimensions?

Problem 2: Examples of curvature

Consider flat three-dimensional space. In cartesian coordinates $(y_1, y_2, y_3) = (x, y, z)$, the metric is simply given by

$$\eta_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} , \quad (3)$$

for which all Christoffel symbols vanish. Let us instead consider spherical coordinates $(x_1, x_2, x_3) = (r, \theta, \phi)$ with

$$x = r \sin \theta \cos \phi , \quad (4)$$

$$y = r \sin \theta \sin \phi , \quad (5)$$

$$z = r \cos \theta . \quad (6)$$

- a) Use the relation $g_{\mu\nu} dx^{\mu} dx^{\nu} = \eta_{\mu\nu} dy^{\mu} dy^{\nu}$ to show that the metric in spherical coordinates is given by

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & r^2 \sin^2 \theta \end{pmatrix} . \quad (7)$$

b) Calculate all non-vanishing Christoffel symbols. Show that $R_{\theta\phi\theta\phi} = 0$.

In fact, all components of the Riemann tensor vanish, as expected in flat space.

Let us now instead consider the two-dimensional surface of a sphere with constant radius R , i.e. $(x_1, x_2) = (\theta, \phi)$. By setting $dr = 0$ in the calculation above we obtain the metric as

$$g_{\mu\nu} = \begin{pmatrix} R^2 & 0 \\ 0 & R^2 \sin^2 \theta \end{pmatrix}. \quad (8)$$

c) Show that in this case $R_{\theta\phi\theta\phi} \neq 0$, representing the fact that the surface of the sphere is not flat.

d) Calculate the resulting Ricci tensor and Ricci scalar.

Problem 3: Geodesics

Consider a particle moving along a path from X_{start}^μ to X_{end}^μ . The path can be parametrised by a function $X^\mu(\lambda)$ with $0 \leq \lambda \leq 1$ such that $X^\mu(0) = X_{\text{start}}^\mu$ and $X^\mu(1) = X_{\text{end}}^\mu$. The length of the path (i.e. the proper time that passes for the particle) is given by

$$\tau = \int_0^1 d\lambda \sqrt{g_{\mu\nu} \dot{X}^\mu \dot{X}^\nu} \equiv \int_0^1 d\lambda L[X^\mu, \dot{X}^\nu] \quad (9)$$

with $\dot{X}^\mu = dX^\mu/d\lambda$.

a) Make use of the relation $d\tau/d\lambda = L$ to show that

$$g_{\mu\nu} U^\mu U^\nu = 1, \quad (10)$$

where $U^\mu = dX^\mu/d\tau$.

Out of all possible paths, the one that requires the shortest proper time must satisfy the Euler-Lagrange equation

$$\frac{d}{d\lambda} \left(\frac{\partial L}{\partial \dot{X}^\mu} \right) - \frac{\partial L}{\partial X^\mu} = 0. \quad (11)$$

b) Show that the Euler-Lagrange equation expressed in terms of U^μ is just the geodesic equation.