

Problem set 5

Submission deadline: 10 July, 9:45
Discussion of solutions: 13 June, 11:30

Problem 13: Horizon exit

Consider a flat Λ CDM universe with

$$H = H_0 \sqrt{\Omega_{\text{rad}} \left(\frac{a_0}{a}\right)^4 + \Omega_{\text{m}} \left(\frac{a_0}{a}\right)^3 + \Omega_{\Lambda}} . \quad (1)$$

- a) Show that the conformal Hubble rate $\mathcal{H} = aH$ has a minimum for $a < a_0$. Find the corresponding value of a for $\Omega_{\text{m}} = 0.3$, $\Omega_{\Lambda} = 0.7$, $\Omega_{\text{rad}} \approx 0$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$.
- b) Argue that this means that there are perturbations that were subhorizon at some time in the past but are now superhorizon.
- c) Find the smallest conformal momentum k_{min} of any perturbation that was subhorizon at some point in the past. Calculate the corresponding physical size of this perturbation in the present universe.
- d) Compare your result to the physical size of perturbations that exit the horizon today.

Problem 14: Tensor perturbations

Let us consider tensor perturbations of the metric

$$ds^2 = a^2(\eta) \left[d\eta^2 - (\delta_{ij} + \hat{E}_{ij}) dx^i dx^j \right] , \quad (2)$$

where \hat{E}_{ij} is a traceless and transverse tensor. Any such tensor can be decomposed into a sum of two different polarization states ($\alpha = +, \times$):

$$\hat{E}_{ij}(\eta, \mathbf{x}) = \sum_{\alpha} h_{\alpha}(\eta, \mathbf{x}) e_{ij}^{\alpha} . \quad (3)$$

The linearised Einstein tensor is given by $\delta G_{00} = 0$, $\delta G_{0i} = 0$ and

$$\delta G_{ij} = \frac{1}{2a^2} \left(\partial_{\eta}^2 \hat{E}_{ij} + 2\mathcal{H} \partial_{\eta} \hat{E}_{ij} - \partial_k \partial_k \hat{E}_{ij} \right) . \quad (4)$$

- a) Assuming that there are no tensor perturbations in the matter, show that the Einstein equation implies

$$h_{\alpha}''(\eta, k) + 2\mathcal{H} h_{\alpha}'(\eta, k) + k^2 h_{\alpha}(\eta, k) = 0 . \quad (5)$$

- b) Argue that the third term is negligible in the superhorizon regime, leading to the solution $h_{\alpha} = h_{\alpha}^{(i)}$.

- c) Using the ansatz $h_\alpha(\eta, k) = f_\alpha(\eta, k)/a(\eta)$, show that an approximate solution in the subhorizon regime is given by

$$h_\alpha(\eta, k) = \frac{A}{a(\eta)} \cos(k\eta) + \frac{B}{a(\eta)} \sin(k\eta). \quad (6)$$

- d) Find the constants A and B by requiring that the superhorizon solution is recovered in the limit $\eta \rightarrow 0$ (you can assume radiation domination for small η).

Since tensor perturbations of the metric can exist in vacuum (i.e. for vanishing energy-momentum tensor), they are called **gravitational waves**. Possible evidence for gravitational waves from the early universe was reported last week by the NANOGrav collaboration, see <https://arxiv.org/pdf/2306.16213.pdf> and <https://arxiv.org/pdf/2306.16219.pdf>.

Problem 15: Evolution of dark matter perturbations

After photon decoupling, the equations of energy and momentum conservation for DM and baryons read

$$\delta'_b - k^2 v_b = 3\Phi' \quad \delta'_{\text{DM}} - k^2 v_{\text{DM}} = 3\Phi' \quad (7)$$

$$v'_b + \frac{2}{\eta} v_b = -\Phi \quad v'_{\text{DM}} + \frac{2}{\eta} v_{\text{DM}} = -\Phi \quad (8)$$

while the first Einstein equation gives

$$k^2 \Phi = -\frac{6}{\eta^2} \left(\frac{\Omega_{\text{DM}}}{\Omega_{\text{m}}} \delta_{\text{DM}} + \frac{\Omega_b}{\Omega_{\text{m}}} \delta_b \right). \quad (9)$$

- a) Show that for a subhorizon mode these equations imply

$$\delta''_{\text{DM}} + \frac{2}{\eta} \delta'_{\text{DM}} - \frac{6}{\eta^2} \delta_{\text{DM}} = -\frac{6}{\eta^2} \frac{\Omega_b}{\Omega_{\text{m}}} \Delta, \quad (10)$$

where $\Delta = \delta_{\text{DM}} - \delta_b$.

- b) Show that the general solution of this equation is given by

$$\delta_{\text{DM}} = \frac{\Omega_b}{\Omega_{\text{m}}} \Delta + \alpha \eta^2 + \frac{\beta}{\eta^3}, \quad (11)$$

where α and β are numerical constants.

- c) By matching to initial conditions (i.e. evaluating eqs. (7) and (8) at photon decoupling), show that

$$\alpha = \frac{1}{\eta_{\text{dec}}^2} \left[\frac{\Omega_{\text{DM}}}{\Omega_{\text{m}}} \delta_{\text{DM}}^{\text{dec}} + \frac{\Omega_b}{5\Omega_{\text{m}}} (3\delta_b^{\text{dec}} + k^2 \eta_{\text{dec}} v_b^{\text{dec}}) \right] \quad (12)$$

Hint: You can use the fact that dark matter perturbations are proportional to a during matter domination, such that

$$(\delta'_{\text{DM}})^{\text{dec}} = \left(\frac{a'}{a} \delta_{\text{DM}} \right)^{\text{dec}} = \frac{2}{\eta_{\text{dec}}} \delta_{\text{DM}}^{\text{dec}}. \quad (13)$$