

Exercise 3.1: $\mathfrak{SO}(N)$ and $\mathfrak{SU}(N)$

10P

Consider the Lie group $\mathrm{SO}(N)$ of orthogonal (i.e. $M^T M = \mathbb{I}$) matrices with unit determinant and the Lie group $\mathrm{SU}(N)$ of unitary (i.e. $M^\dagger M = \mathbb{I}$) matrices with unit determinant.¹

- Determine the dimension (i.e. the number of different matrices whose *real* linear combinations span the whole group) of $\mathrm{SO}(N)$ and of $\mathrm{SU}(N)$ for a generic $N \in \mathbb{N}$.
- Derive the conditions fulfilled by the corresponding Lie algebras $\mathfrak{so}(N)$ and $\mathfrak{su}(N)$. Remember that any element of a Lie group around the identity can be obtained through the exponentiation of an element of the corresponding Lie algebra: $G \ni M(\alpha) = \exp(\alpha N)$, $N \in \mathfrak{g}$.

Hint: Remember that $\det A = \exp \mathrm{Tr} A$.

Focus now on the $\mathfrak{so}(3)$ and $\mathfrak{su}(2)$ Lie algebras. In the *fundamental representation*, a basis for $\mathfrak{su}(2)$ is given by $T_F^a = \sigma^a/2$, where σ^a are the Pauli matrices, while a basis for $\mathfrak{so}(3)$ is given by $(T_F^i)_{jk} = \epsilon_{ijk}$, where ϵ_{ijk} is the three-dimensional Levi-Civita symbol, with $\epsilon_{123} = +1$.

- Calculate all the structure constants f^{ab}_c for both Lie algebras.
- The *Cartan–Killing form* of a Lie algebra is defined as

$$K^{ab} = -f^{am}_l f^{bl}_m \quad (1.1)$$

Prove that the Cartan–Killing form is symmetric. Calculate it explicitly for $\mathfrak{so}(3)$ and $\mathfrak{su}(2)$.

The *adjoint representation* of a Lie algebra is given by the structure constants themselves:

$$(T_A^a)^b_c = -i f^{ab}_c. \quad (1.2)$$

- Prove the Jacobi identity by using the definition of the structure constants, working in the adjoint representation:

$$f^{ab}_e f^{ec}_d + f^{bc}_e f^{ea}_d + f^{ca}_e f^{eb}_d = 0 \quad (1.3)$$

Operators that commute with every element of a Lie algebra are called *invariant Casimir operators*. It can be proved that there exist $(N-1)/2$ linearly-independent invariant Casimir operators for $\mathfrak{so}(N)$ if N is odd, $N/2$ if N is even, and $N-1$ for $\mathfrak{su}(N)$.

- For $\mathfrak{so}(3)$ and $\mathfrak{su}(2)$ compute the *quadratic Casimir operator*, defined as

$$C_2 = K_{ab} T^a T^b, \quad (1.4)$$

where K_{ab} is the inverse of K^{ab} (i.e. $K_{ab} K^{bc} = \delta_a^c$). Prove that it indeed commutes with every element of the corresponding Lie algebra.

¹In this exercise a superscript can only be contracted with a subscript, therefore we will pay great attention to the position of the indices. Superscripts and subscripts can be treated as equal for all practical purposes.

Exercise 3.2: Covariant derivative**5P**

In general, the covariant derivative

$$D^\mu = \partial^\mu - igA^\mu = \partial^\mu - igA_a^\mu T^a \quad (2.1)$$

explicitly depends on the choice of the representation of the gauge group generators T^a . You know that in the fundamental representation

$$D'^\mu = UD^\mu U^{-1} \quad \text{and} \quad A'^\mu = UA^\mu U^{-1} - \frac{i}{g}(\partial^\mu U)U^{-1}, \quad (2.2)$$

with $U = \exp(ig\vartheta_a T_F^a)$.

(a) Prove that the form of the transformations above is valid in any representation:

$$D'^\mu = VD^\mu V^{-1} \quad \text{and} \quad A'^\mu = VA^\mu V^{-1} - \frac{i}{g}(\partial^\mu V)V^{-1}, \quad (2.3)$$

where $V = \exp(ig\vartheta_a T)$, T being the generators in an arbitrary representation.

Hint: Start from the transformation law for A^μ in the fundamental representation and extract the exact form of the gauge transformation for A_a^μ , independent of the representation of the generators. To this end, make use of the Baker–Campbell–Hausdorff formula:

$$\exp(B)A\exp(-B) = \sum_{n=0}^{+\infty} \frac{1}{n!} A_n, \quad (2.4)$$

with $A_n = [B, A_{n-1}]$, $A_0 = A$, and remember that $f^{ab}_c = i(T_A^a)^b_c$.

(b) Show that

$$F^{\mu\nu} = \frac{i}{g}[D^\mu, D^\nu], \quad (2.5)$$

and is, therefore, representation-independent. Find how $F^{\mu\nu}$ transforms under a gauge transformation.

Exercise 3.3: Lorentz transformations**5P**

Taking $x \rightarrow x'^\mu = \Lambda^\mu_\nu x^\nu$ and the orthogonality condition, $\Lambda^\mu_\nu \Lambda^\rho_\sigma g_{\mu\rho} = g_{\nu\sigma}$,

(a) Show that $a^\mu b_\mu$ is invariant under Lorentz transformations.

Let $(\Lambda^{-1})^\mu_\nu$ be the inverse of Λ^μ_ν , meaning that

$$(\Lambda^{-1})^\mu_\nu \Lambda^\nu_\rho = \delta^\mu_\rho = g^\mu_\rho. \quad (3.1)$$

(b) Show that $(\Lambda^{-1})^\mu_\nu = \Lambda_\nu^\mu$.

(c) Show that $\partial'_\mu = (\Lambda^{-1})^\nu_\mu \partial_\nu$ where $\partial'_\mu = \partial/\partial x'^\mu$.

(d) Does $M_{\nu\mu} = \pm M_{\mu\nu}$ imply $M_\nu{}^\mu = \pm M_\mu{}^\nu$, $M_\nu{}^\mu = \pm M^\mu{}_\nu$, or $M^{\nu\mu} = \pm M^{\mu\nu}$? If not, how should the metric tensor be modified to allow this?