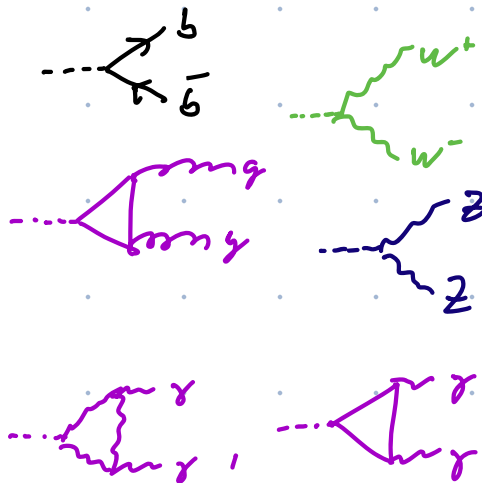
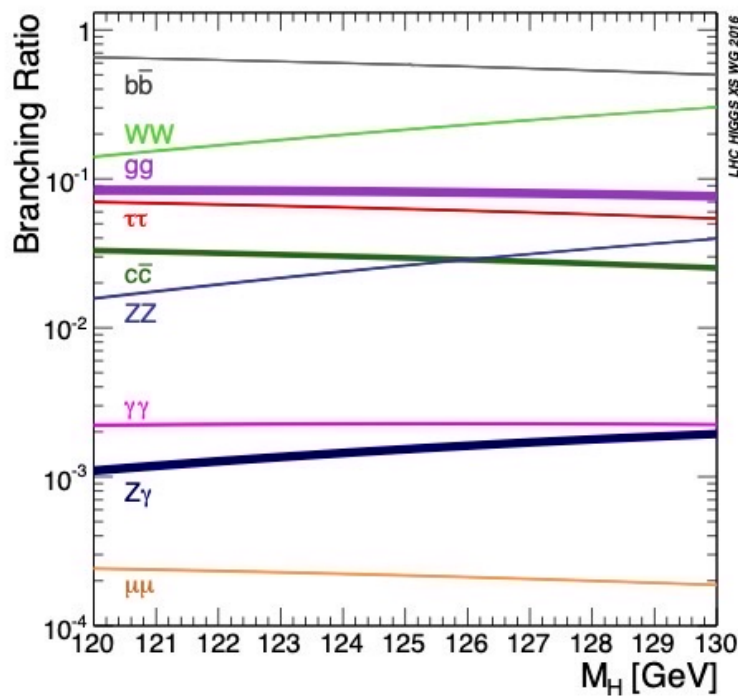


Lecture 29:

Higgs decays



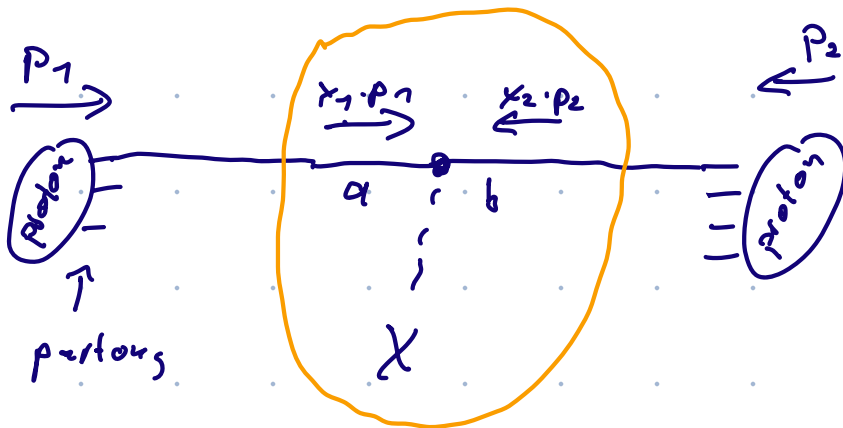
V.2 Higgs production at the LHC

The LHC collides protons with a center-of-mass energy of (currently) $\sqrt{s} = 13.6 \text{ TeV}$

Cross sections obtained via collinear factorization

$$d\sigma(PP \rightarrow X) = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 f_a(x_1, \mu_F) f_b(x_2, \mu_F) d\hat{\sigma}(a, b \rightarrow X)$$

from cross section of underlying partonic process



$$d\hat{\sigma}(a, b \rightarrow X)$$

$f_a(x, \mu_F)$: parton density function

$\rightarrow$ "probability" to find parton a with momentum $x \cdot p$ in the proton

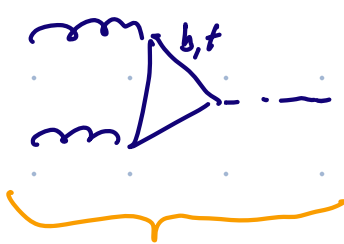
$$a \in \{g, u, \bar{u}, d, \bar{d}, s, \dots\}$$

μ_F : factorization scale

(similar to renormalization scale of running coupling $\rightarrow$ IV.2)

gluon fusion

dominant production mode for Higgs at LHC



$\mathcal{M}_{LO}$

$$\hat{\sigma}(gg \rightarrow H) = \frac{\alpha_F \alpha_s^2}{288 \sqrt{2} \pi} \left| \sum_Q A(\tau_Q) \right|^2$$

same as for $H \rightarrow gg$ decay

next-to-leading order corrections:

- virtual corrections

$$\mathcal{M}_V = \text{[diagram 1]} + \text{[diagram 2]} + \dots$$

- real corrections:

$$\mathcal{M}_R^{gg} = \text{[diagram 1]} + \text{[diagram 2]}$$

$$\mathcal{M}_R^{qg} = \text{[diagram 1]}$$

Need to be included, since

$h+g$ and $h+q$ final state indistinguishable with h final state if the q/g momentum is small or collinear to initial state.

→ "inclusive cross section"

$$\sigma(pp \rightarrow h + X)$$

$$\Rightarrow \hat{\sigma}_{NLO} = \hat{\sigma}_{LO} + \hat{\sigma}_V + \hat{\sigma}_R$$

$$= \int d\Phi_1 \underbrace{|\mathcal{M}_{LO}|^2}_{(\alpha_s)^2} + \int dd_1 \underbrace{(\mathcal{M}_V^* \cdot \mathcal{M}_{LO} + \mathcal{M}_{LO}^* \cdot \mathcal{M}_V)}_{\alpha_s^2 \cdot \alpha_s} + \int dd_2 \underbrace{(|\mathcal{M}_R^{gg}|^2 + |\mathcal{M}_R^{qg}|^2)}_{(\alpha_s^{3/2})^2}$$

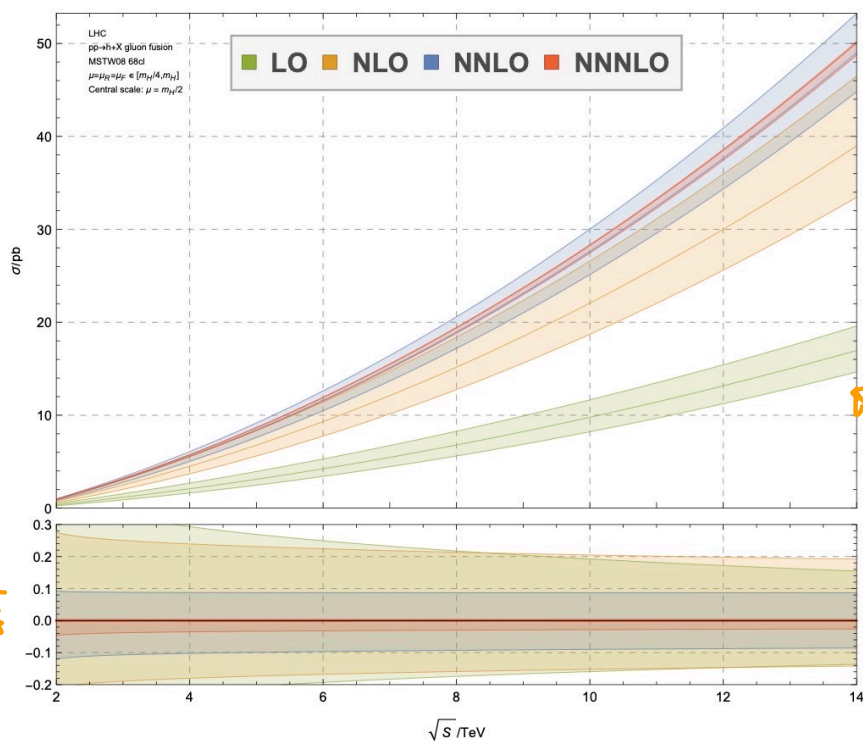
The NLO corrections increase the cross section by $\sim 700\%$

→ not sufficient for accurate description of process.

The corrections up to N^3LO in α_s are known

(using the HTL)

asir: 1503.06056

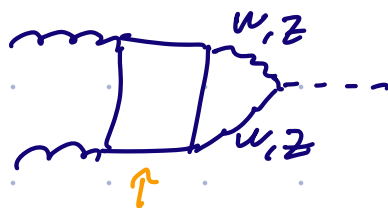
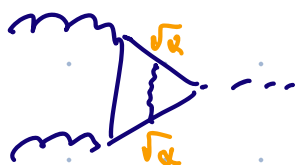


$N^3LO \sim 3\%$
 $NNLO \sim 20\%$
 $NLO \sim 100\%$ correction

LO $\sigma_{LO} \approx 18 \text{ pb}$
 uncertainty estimates, via variation of μ_R, μ_F , unreliable at LO due to new partonic channels appearing at NLO.

further corrections:

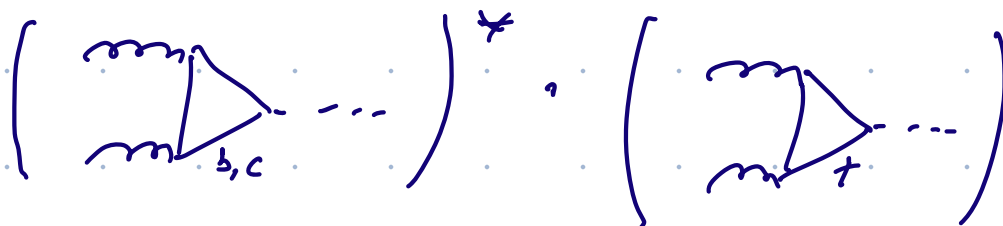
- NLO EW + 5%



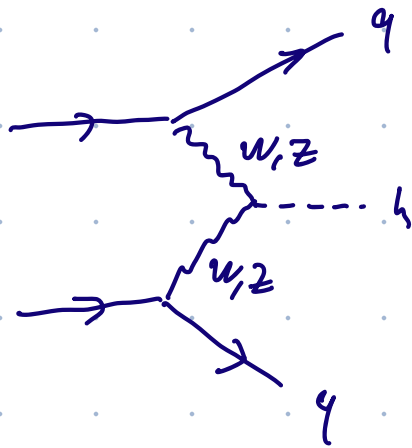
also light quarks contribute

- light quarks - 4%

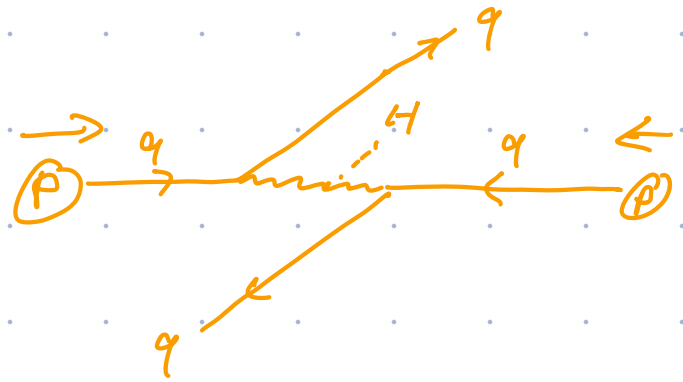
negative interference with top-quark loop



vector boson fusion



$$pp \rightarrow W^+W^+ / Z^+Z^+ \rightarrow H$$



$\Rightarrow$ characteristic detector signature:

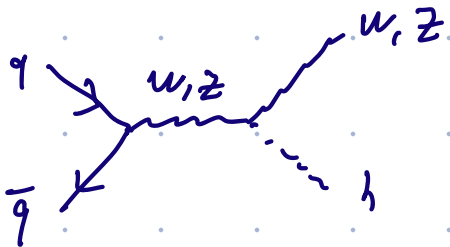
- 2 jets with large separation
- Higgs boson in central region

NLO QCD corrections: $\sim 1-10\%$ (depending on phase-space cuts)

also NNLO QCD and NLO EW corrections known

$\hookrightarrow$ N^3LO QCD in some approximation

Higgs strahlung



$$pp \rightarrow W^+/Z^+ \rightarrow W, Z + h$$

NLO QCD $+30\%$

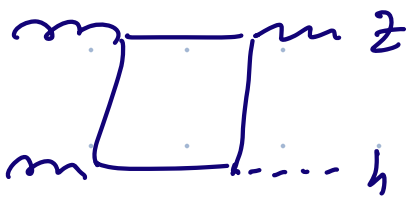
NNLO QCD $+3\%$

also N^3LO QCD

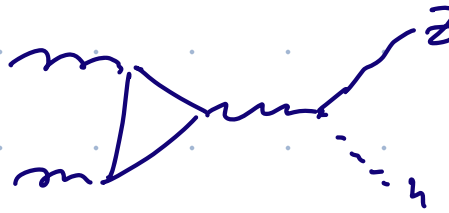
and NNLO EW known.

additional contributions for ZH production

$$gg \rightarrow ZH \quad +10\%$$

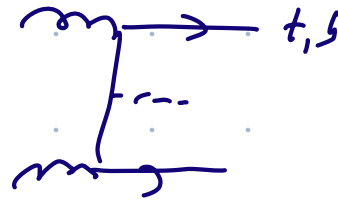
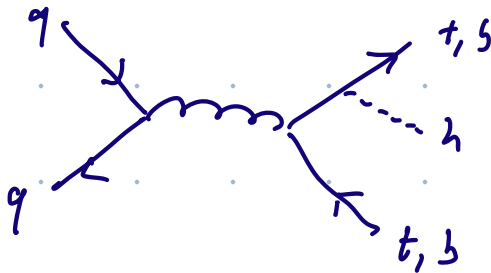


+



Associate production

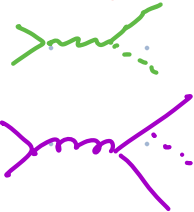
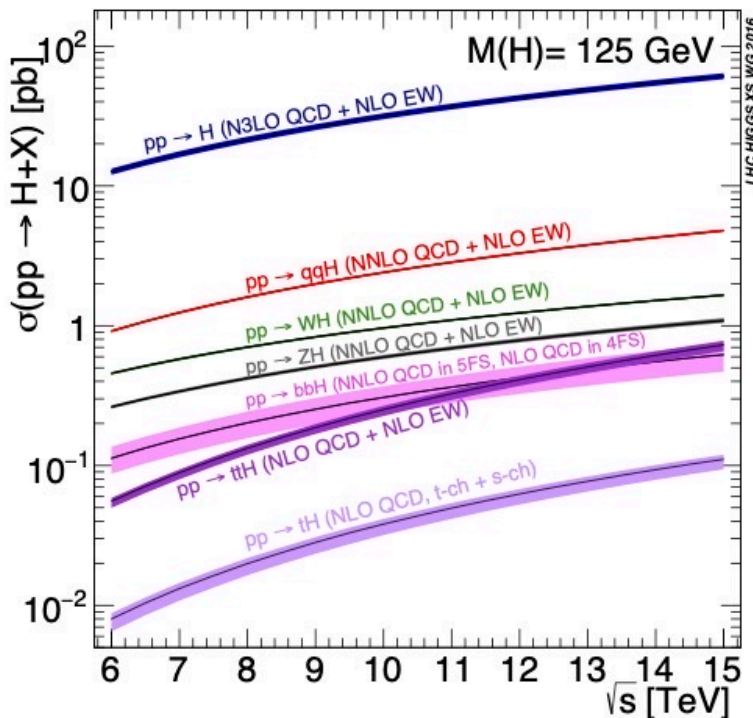
$$pp \rightarrow t\bar{t}H \quad / \quad b\bar{b}H$$



NLO QCD +20%

→ direct measurement of top-quark Yukawa coupling.
 complicated final state for $t\bar{t}H$ due to
 subsequent decays $t \rightarrow b + W^+$
 $\hookrightarrow q\bar{q}, \ell^+ \nu$

summary of production cross sections



V3. Higgs discovery and quantum numbers

discovery in July 2012 during LHC run 1:

2011: $\sqrt{s} = 7 \text{ TeV}$

$\int L = 5.6 \text{ fb}^{-1}$

2012: $\sqrt{s} = 8 \text{ TeV}$

$\int L = 6.6 \text{ fb}^{-1}$

c.f. run 2 2015-18 $\sqrt{s} = 13 \text{ TeV}$ $\int L \approx 130 \text{ fb}^{-1}$

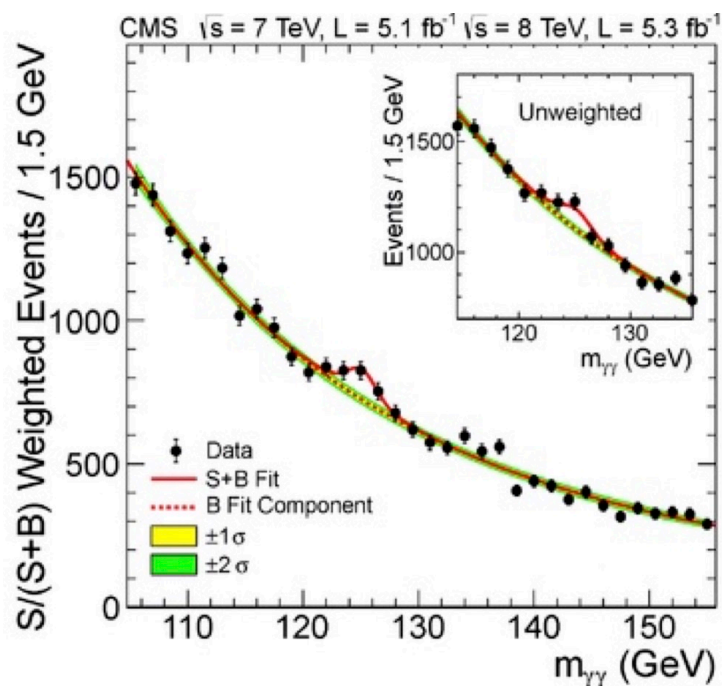
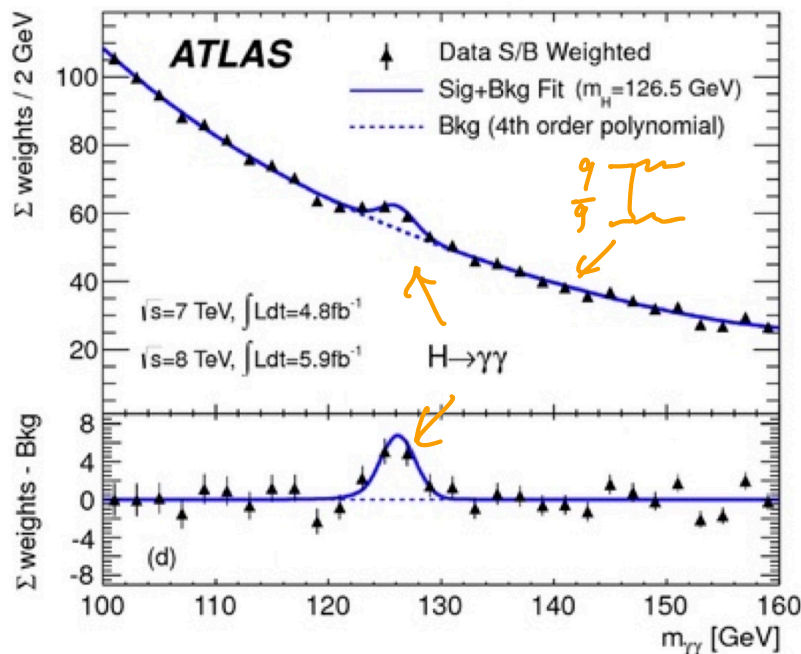
run 3 2022-24 $\sqrt{s} = 13.6 \text{ TeV}$ $\int L \approx 300 \text{ fb}^{-1}$

main decay channels for discovery:

$H \rightarrow \gamma\gamma$, $H \rightarrow ZZ^* \rightarrow 4\ell$

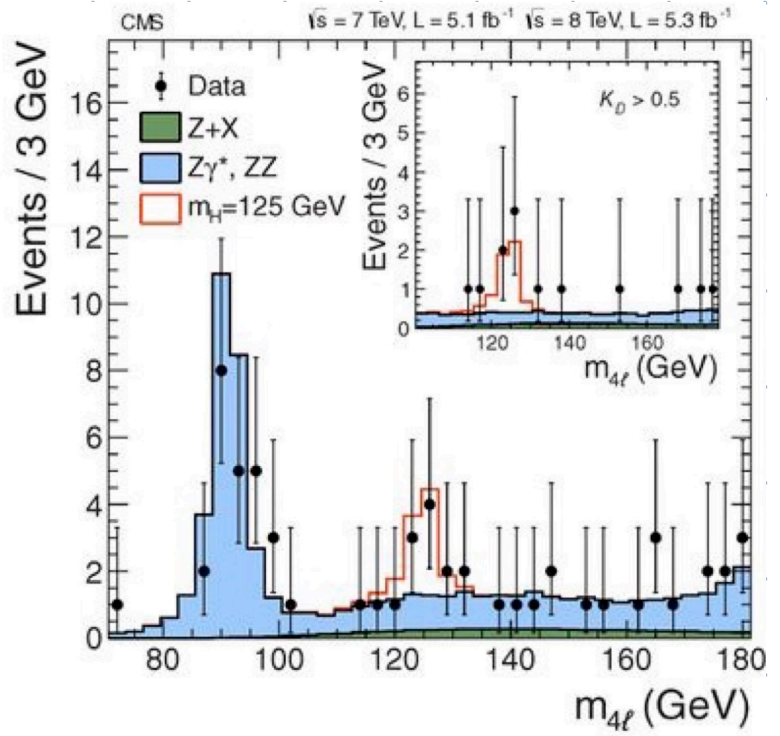
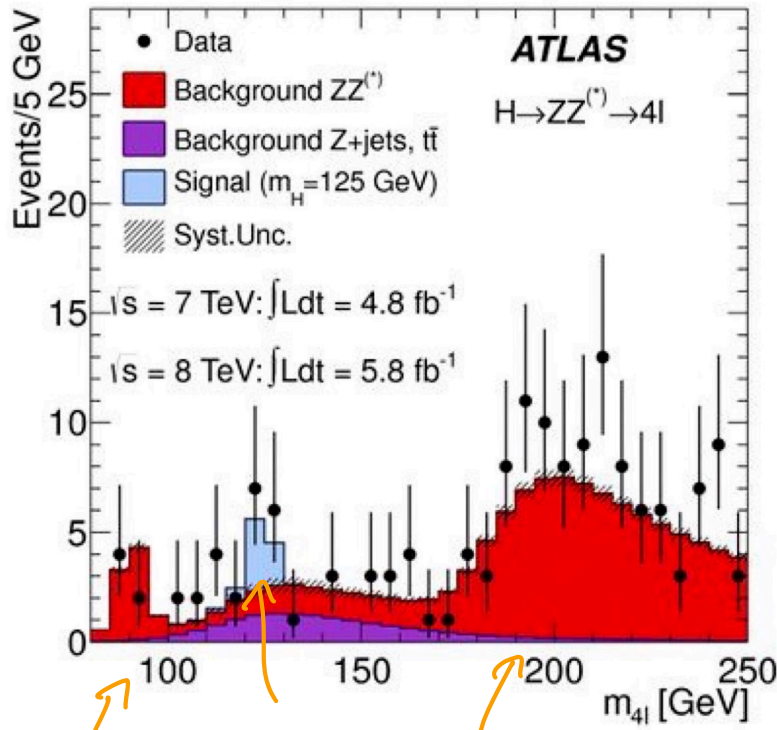
$\Rightarrow$ small BR, but distinct detector signal

$H \rightarrow \gamma\gamma$:



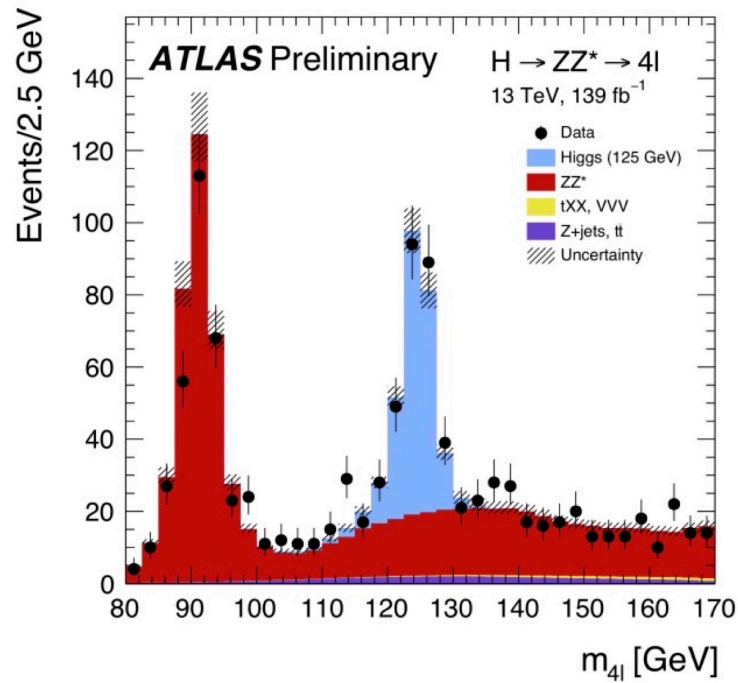
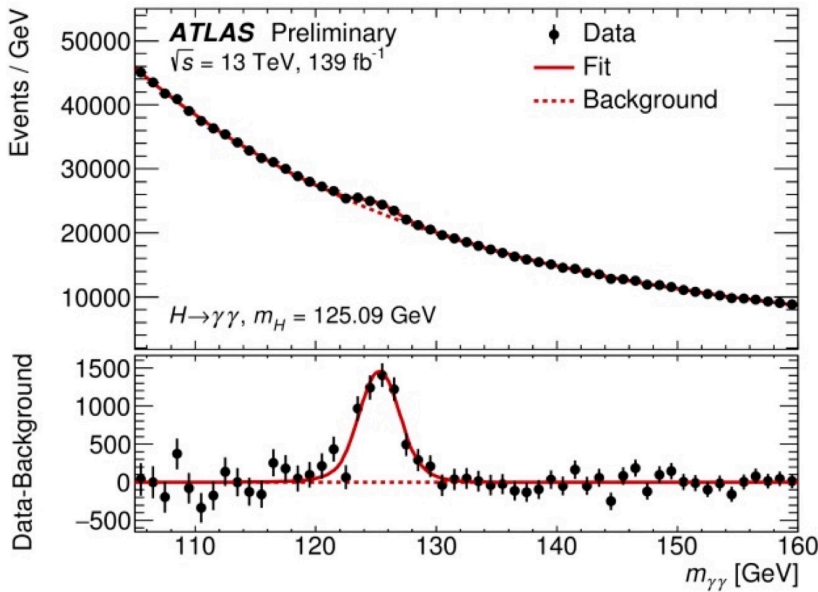
$$\hookrightarrow m_{\gamma\gamma}^2 = (p_{\gamma 1} + p_{\gamma 2})^2$$

1-7 22



Handwritten notes: $\sqrt{s} \leftarrow m \leftarrow$ and $\int \mathcal{L} dt$

updated results with full run 2 statistics



spin and parity quantum numbers

The Higgs boson in the SM is CP even, with couplings

$$\mathcal{L} \supset \frac{1}{v} \left(m_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} m_Z^2 Z_\mu Z^\mu + m_f \bar{\psi} \psi \right)$$

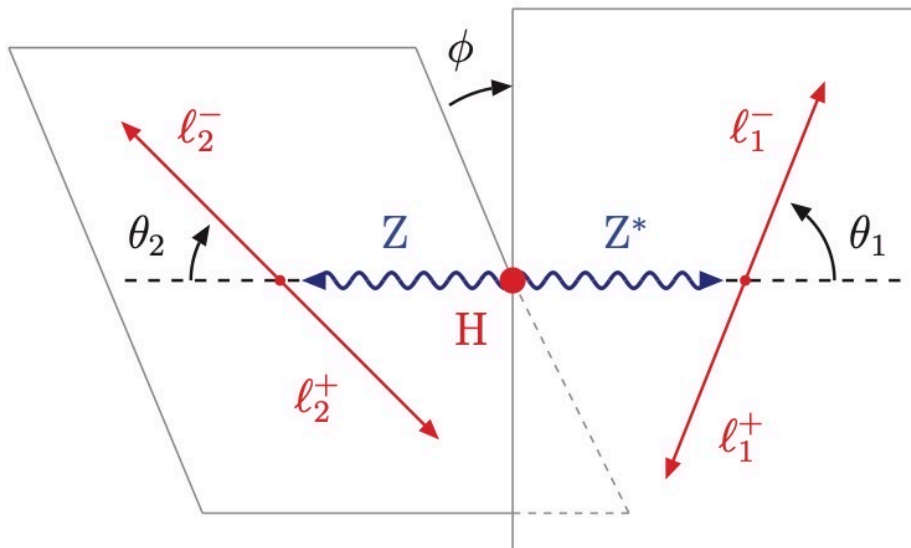
The corresponding CP-odd coupling structures would be

$$\sim h W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} , \quad \frac{1}{2} h Z_{\mu\nu} \tilde{Z}^{\mu\nu} , \quad h \bar{\psi} \gamma_5 \psi$$

$$\text{with } \tilde{F}_{\mu\nu} = \frac{1}{2} \epsilon_{\mu\nu\sigma\tau} F^{\sigma\tau}$$

$\Rightarrow$ modification of angular distributions, e.g. in

$H \rightarrow Z Z \rightarrow 4\ell$ decay

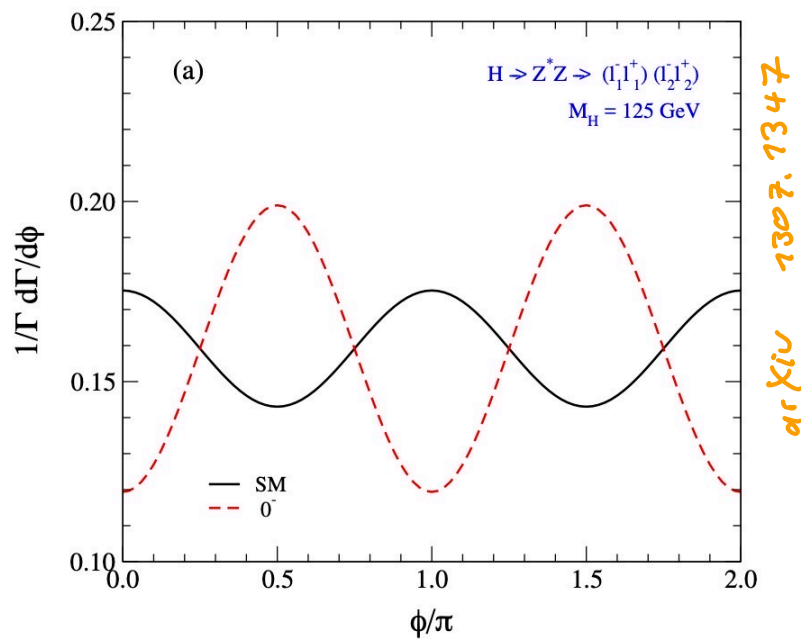


arXiv:1307.1347

for CP-even: $\frac{d\Gamma}{d\phi} \sim 1 + \frac{1}{2g^4+4} \cos 2\phi$

CP-odd: $\frac{d\Gamma}{d\phi} \sim 1 - \frac{1}{4} \cdot \cos 2\phi$

$$\gamma^2 = \frac{m_h^2 - m_Z^2 - m_{Z^*}^2}{2 m_{Z^*} m_Z}$$



→ Experimental confirmation that Higgs is CP even by ATLAS, CMS in 2012/13

Also confirmed spin = 0:

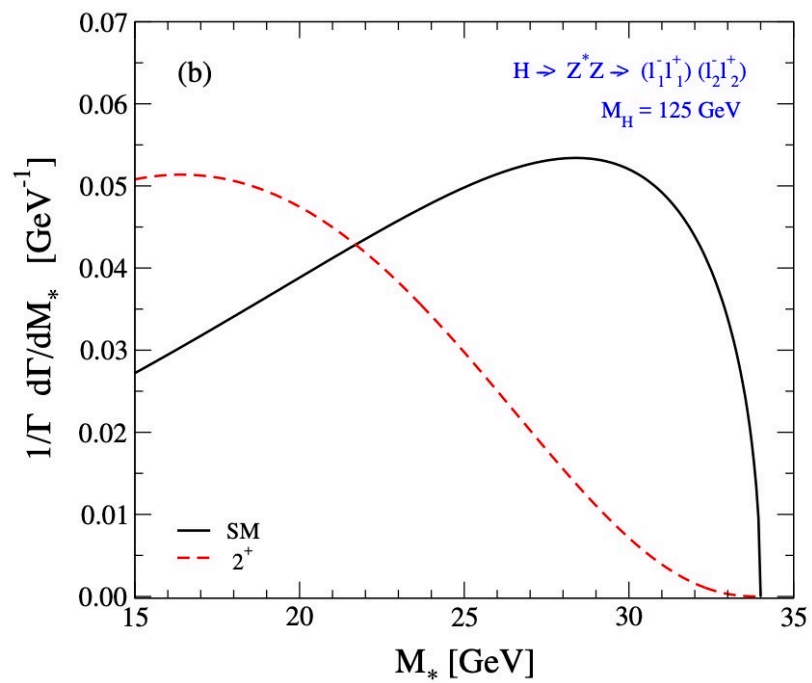
• spin = 1 excluded, since decay $h \rightarrow \gamma\gamma$ would be forbidden by angular momentum conservation and parity

→ Landau, Yang theorem, 1948/50

• spin = 2 would lead to different partial decay width for $h \rightarrow ZZ^*$

spin 0 : $\frac{d\Gamma}{dm_{Z^*}} \sim \beta = \sqrt{(m_h - m_Z)^2 - m_{Z^*}^2} / m_h$

spin 2 : $\frac{d\Gamma}{dm_{Z^*}} \sim \beta^3$



arXiv:1307.1347