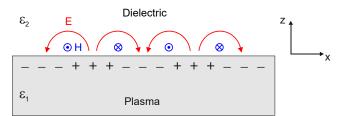


### Surface Plasmons

(SP - surface plasmon; SPP - surface plasmon polariton)



$$\vec{E}_{1} = (E_{x1}, 0, E_{z1}) \cdot e^{i(k_{x1}x - k_{z1}z - \omega t)}$$

$$\vec{H}_{1} = (0, H_{1}, 0) \cdot e^{i(k_{x1}x - k_{z1}z - \omega t)}$$

$$\vec{E}_{1} = (E_{x1}, 0, E_{z1}) \cdot e^{i(k_{x1}x - k_{z1}z - \omega t)}$$

$$\vec{H}_{1} = (0, H_{y1}, 0) \cdot e^{i(k_{x1}x - k_{z1}z - \omega t)}$$

$$\vec{H}_{2} = (0, H_{y2}, 0) \cdot e^{i(k_{x2}x + k_{z2}z - \omega t)}$$

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### 

#### SPP - dispersion relation

#### Dispersion relation:

$$\omega = \frac{c}{n} \cdot k$$
  $\Rightarrow$   $k^2 = \epsilon \cdot \left(\frac{\omega}{c}\right)^2$  (2)

Result: 
$$k_x(\omega) = \sqrt{\frac{\epsilon_1 \cdot \epsilon_2}{\epsilon_1 + \epsilon_2} \cdot \frac{\omega}{c}}$$

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## **SKIT**

#### SPP – dispersion relation

#### **Boundary conditions:**

$$\left. \begin{array}{ll} E_{x1} = E_{x2} & (=E_x) \\ H_{y1} = H_{y2} & (=H_y) \end{array} \right\} \quad \Rightarrow \quad k_{x1} = k_{x2} \quad (=k_x)$$

#### Maxwell equation:

$$\operatorname{curl} \vec{H} = \dot{\vec{D}} = \epsilon_0 \epsilon \dot{\vec{E}} \quad \Rightarrow \quad -\frac{\partial H_y}{\partial z} = \epsilon_0 \epsilon \dot{\vec{E}}_x \qquad \qquad \operatorname{curl} \vec{H} = \begin{pmatrix} \partial_y H_z - \partial_z H_y \\ \partial_z H_x - \partial_x H_z \\ \partial_x H_y - \partial_y H_x \end{pmatrix}$$

$$\frac{z>0:}{z<0:} -i k_{z2} H_y = -i\omega \epsilon_2 \epsilon_0 E_x$$

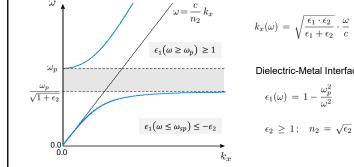
$$\frac{z<0:}{t} +i k_{z1} H_y = -i\omega \epsilon_1 \epsilon_0 E_x$$

$$\Rightarrow \begin{bmatrix} -\frac{k_{z2}}{k_{z1}} = \frac{\epsilon_2}{\epsilon_1} & (1) \end{bmatrix}$$

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## **SKIT**

### SPPs at a Dielectric-Metal Interface





#### Dielectric-Metal Interface:

$$\epsilon_1(\omega) = 1 - \frac{\omega_p^2}{\omega^2}$$

$$\epsilon_2 \ge 1; \quad n_2 = \sqrt{\epsilon_2}$$

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**SKIT** 

#### SPP - decay length

If  $k_{zi}$  is the z-component of the wavevector in media i with permittivity  $\epsilon_i$ , then:

$$k = \frac{\omega}{n}$$
  $\Longrightarrow$   $k_i^2 = \epsilon_i \cdot \left(\frac{\omega}{c}\right)^2 = k_x^2 + k_{zi}^2$   $\Longrightarrow$   $k_{zi} = \frac{\epsilon_i}{\sqrt{\epsilon_1 + \epsilon_2}} \cdot \frac{\epsilon_i}{c}$ 

Media 2 is a dielectric and media 1 is a metal. Then for the upper branch of the dispersion relation, we get:

 $\omega > \omega_p \implies \epsilon_1(\omega) > 0 \implies k_{z1}$  and  $k_{z2}$  are real (propagating waves)

For the lower branch of the dispersion relation, we have:

$$\omega < \frac{\omega_p}{\sqrt{1+\epsilon_2}} \quad \Longrightarrow \quad \epsilon_1(\omega) = 1 - \frac{\omega_p^2}{\omega^2} < -\epsilon_2 \quad \Longrightarrow \quad k_{z1} \text{ and } k_{z2} \text{ are imaginary!}$$

The decay lengths of the evanescent wave are: 
$$d_i = \frac{1}{|k_{zi}|} = \frac{\lambda}{2\pi} \sqrt{\left|\frac{\epsilon_1 + \epsilon_2}{\epsilon_i^2}\right|}$$

Silver @ 600 nm:  $d_1 = 24 \text{ nm}$   $d_2 = 390 \text{ nm}$ 

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# **SKIT** Excitation of SPPs by ATR (ATR = attenuated total reflection) metal Otto-Kretschmannconfiguration configuration H. Raether, Surface Plasmons; Novotny & Hecht, Nano-Optics Nanooptics 6/7

