"Nanotechnologie II"
Sommersemester 2016

C. Ausgewählte Kapitel zur Nanotechnologie

7. Nanostrukturen durch Selbstorganisation
   - 7.1 Voraussetzungen für Selbstanordnung
   - 7.2 Thermodynamische Aspekte der Selbstanordnung
   - 7.3 Weitere Beispiele
8. Partikuläre Nanostrukturen
   - 8.1 Festkörper in reduzierter Dimension
   - 8.2 Elektronische Eigenschaften eindimensionaler Strukturen
   - 8.3 Kohlenstoff-Nanoröhrchen
   - 8.4 Cluster und Kolloide
   - 8.5 Einzelladungseffekte
9. Nanoelektronik
   - 9.1 Der Einzelladungstransitor (SET)
   - 9.2 Quanten-Computing
   - 9.3 Molekulare Elektronik
10. Nanooptik
    - 10.1 Photonische Bandlückenmaterialien (PBG)
    - 10.2 Plasmonen
11. Nanotribologische Systeme
    - 11.1 Der Lotus-Effekt
    - 11.2 Der Gecko-Effekt
12. Biologische Nanostrukturen
    - 12.1 Abbildung und mechanische Eigenschaften lebender Zellen
    - 12.2 Biologische Nanostrukturen

Literatur zur Vorlesung "'Nanotechnologie II"'

Zur Vorbereitung der Vorlesung wurde verwendet:

- Rainer Waser (Ed.)
  Nanoelectronics and Information Technology
  2nd edition, John Wiley & Sons Ltd 2005
  ISBN 3-527-40363-9
  (Standort UB: nach 8.20, 2003 E 201(2)f)

- S. M. Lindsay
  Introduction to Nanoscience
  Oxford University Press 2010
  ISBN 978-019-954421-9

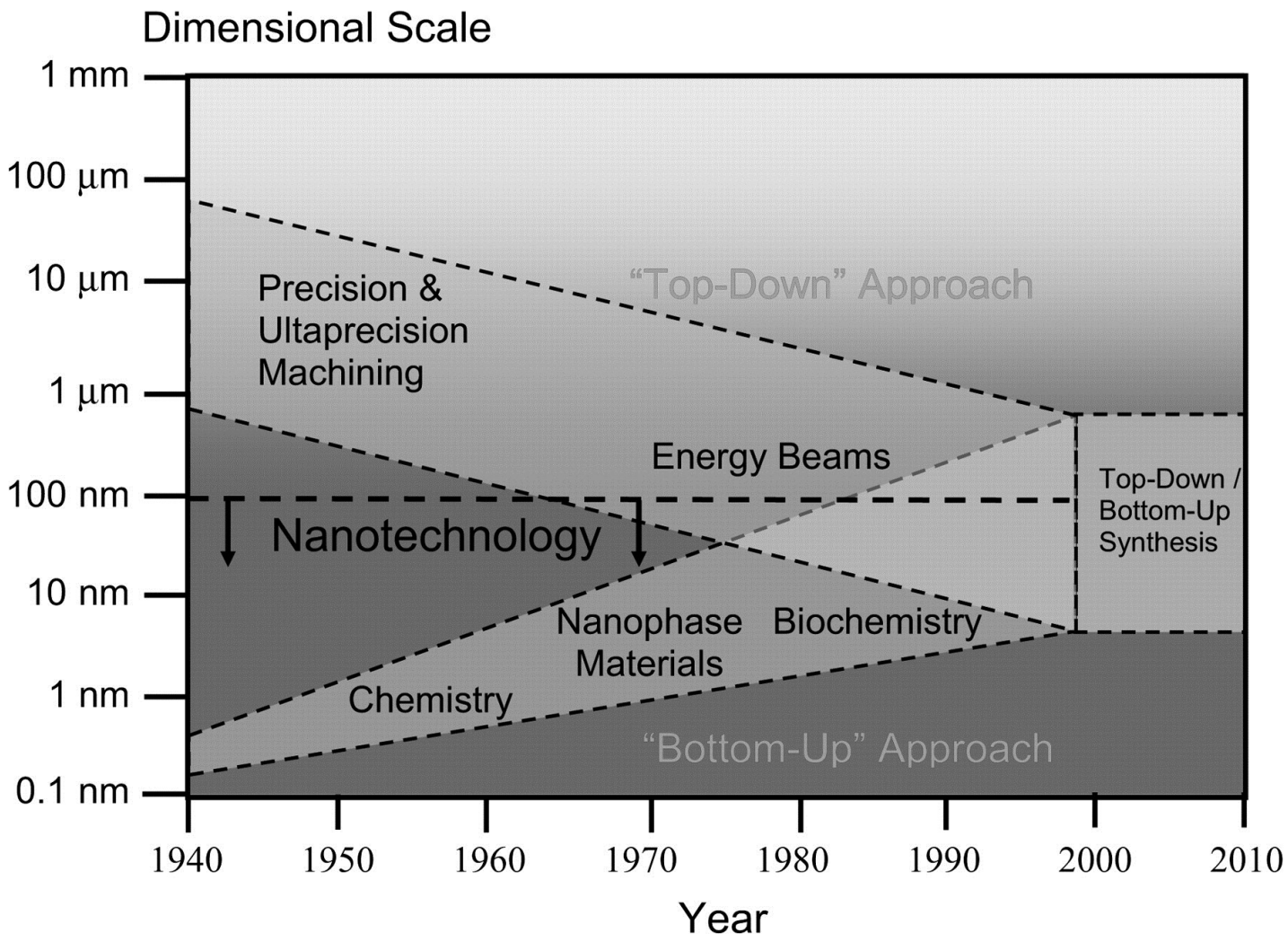
- Claire Dupas, Philippe Houdy, Marcel Lahmani (Eds.)
  Nanoscience
  Springer, Berlin 2004
  ISBN 3-540-28616-0
  (Standort UB: nano 0, 2007 A 4151)

- Günter Schmid (Ed.)
  Nanoparticles
  Wiley-VCH, Weinheim 2004
  ISBN 3-527-30507-6
  (Standort UB: nano 2, 2004 A 8753)

- Lukas Novotny und Bert Hecht
  Principles of Nano-Optics
  Cambridge University Press 2006
  ISBN 978-0-521-83224-3
  (Standort UB: phys 4.29, 2007 E 1379)

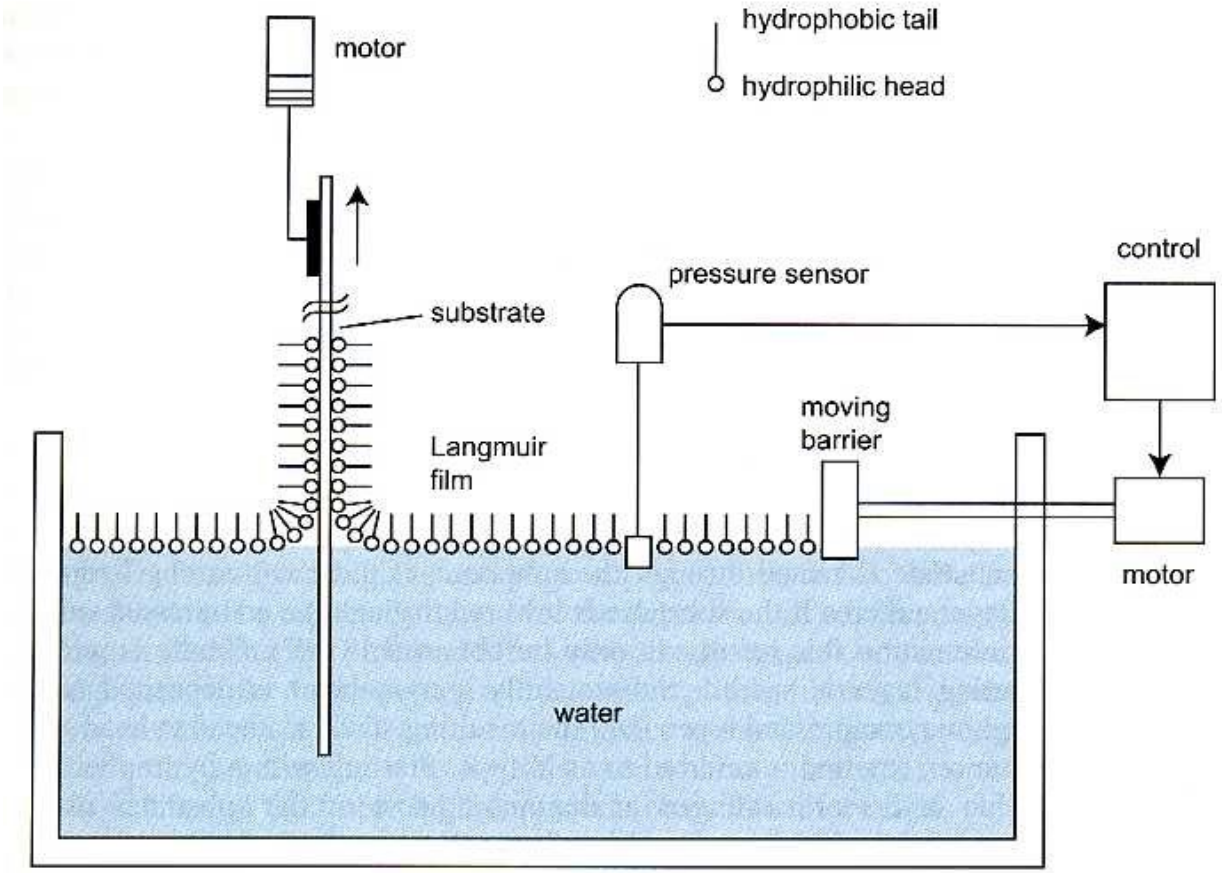
- Bharat Bushan (Ed.)
  Springer Handbook of Nanotechnology
  2nd edition, Springer Berlin 2007
  ISBN 3-540-29855-X

Konvergenz von top-down- und bottom-up-Nanotechnologien.



Whatmore R W Occup Med (Lond) 2006;56:295-299

Versuchsanordnung zur kontrollierten Abscheidung von Langmuir-Blodgett-Filmen



Versuchsführung bei
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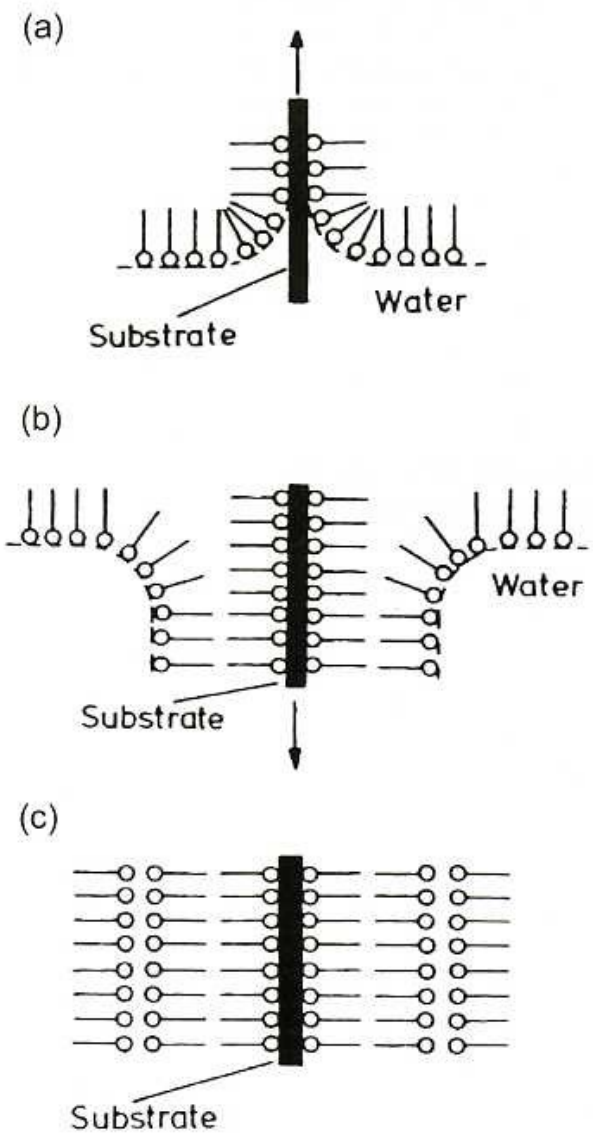
von Langmuir-Blodgett-Filmen



Dendrimer mit eingeschlossenem Farbstoff-Radikal

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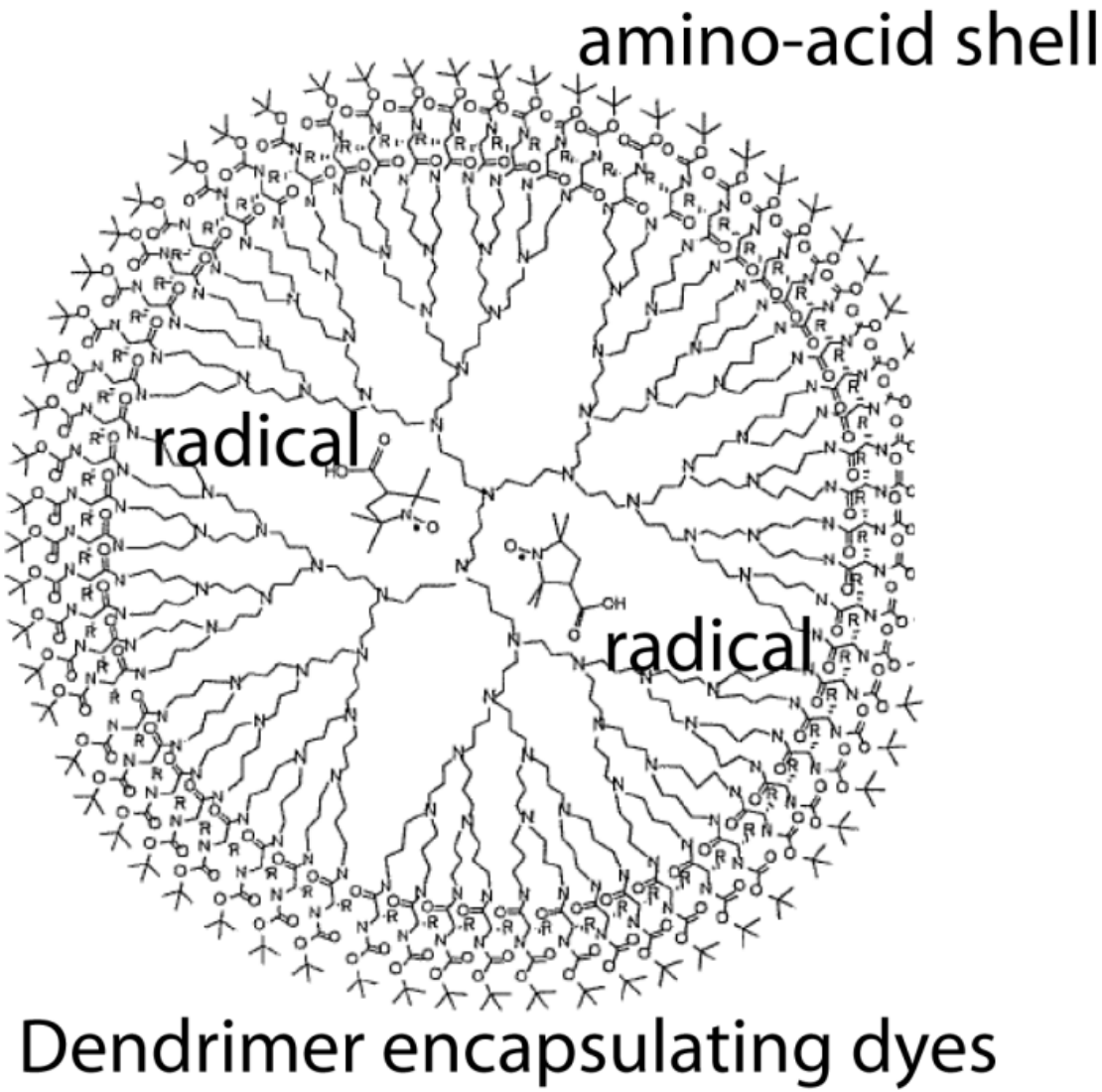
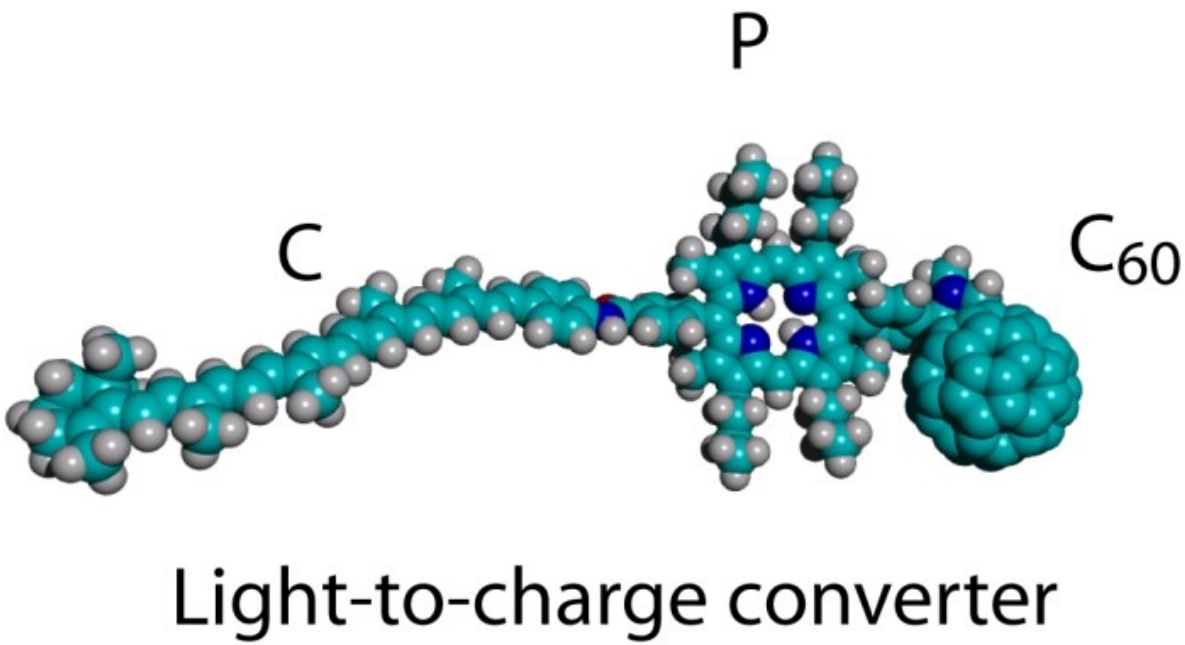


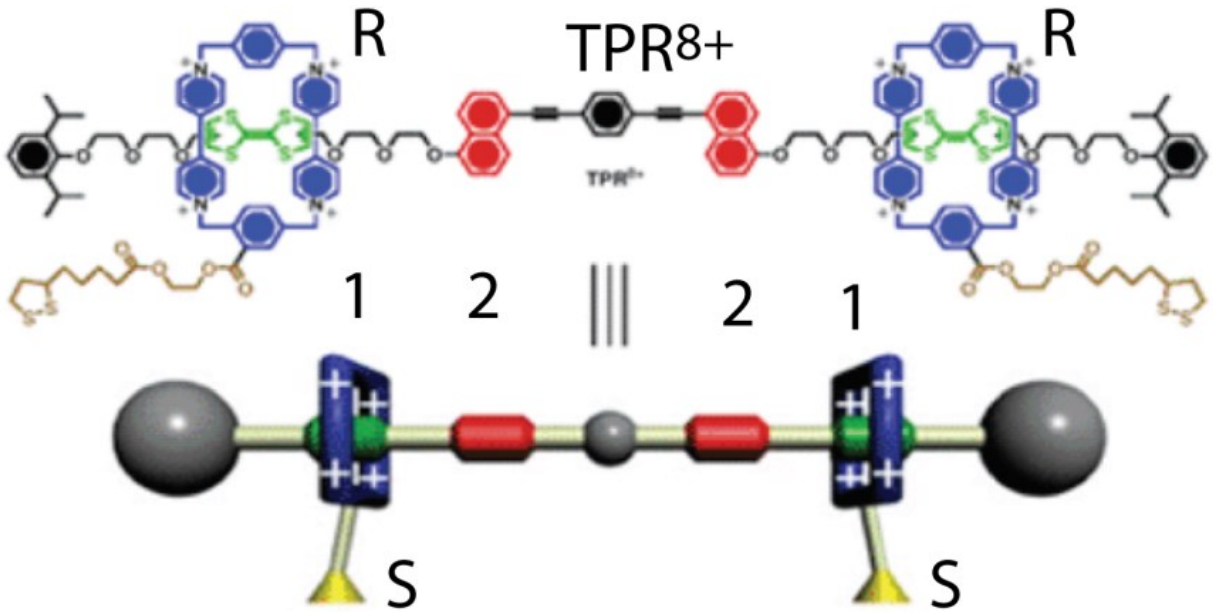
Licht-zu-Ladung-Konverter



(Courtesy of Professor Devens Gust of Arizona State University.)

Nano-Muskel: Aktuator auf Rotaxane-Basis

(Reprinted with permission from Y. Liu, A.H. Flood, P.A. Bonvallet, S.A. Vignon, B.H. Northrup, H-R Tseng, J.O. Jeppesen, T.J. Huang, B. Brough, M. Baller, S. Maganov, S.D. Solares, W.A. Goddard, C-M Ho, and J.F. Stoddart,. J. Am. Chem. Soc. 2005, **127**, 9745, Published 2005 by American Chemical Society).

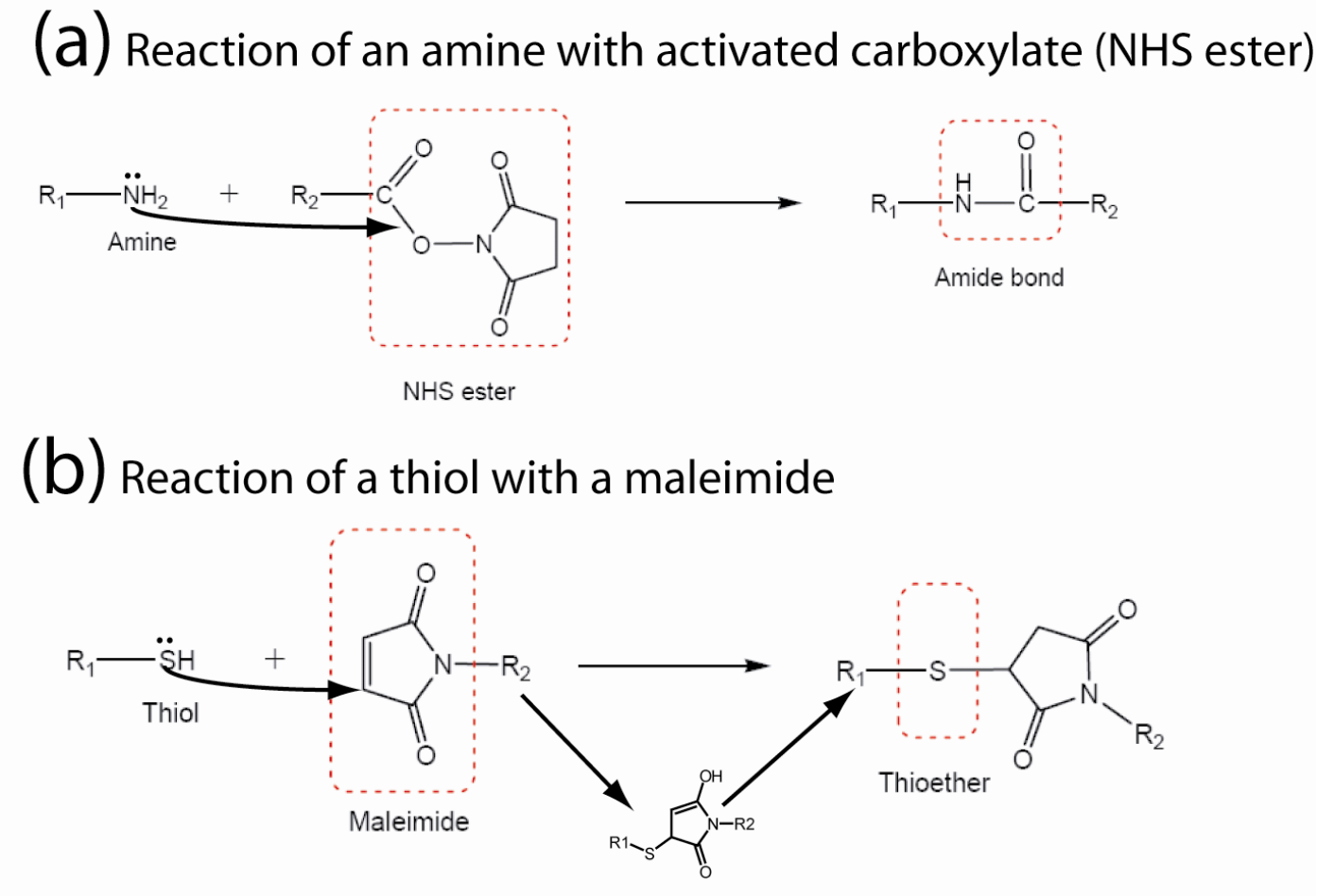


"Nano-muscle" actuator based on rotaxanes

Attacking agents

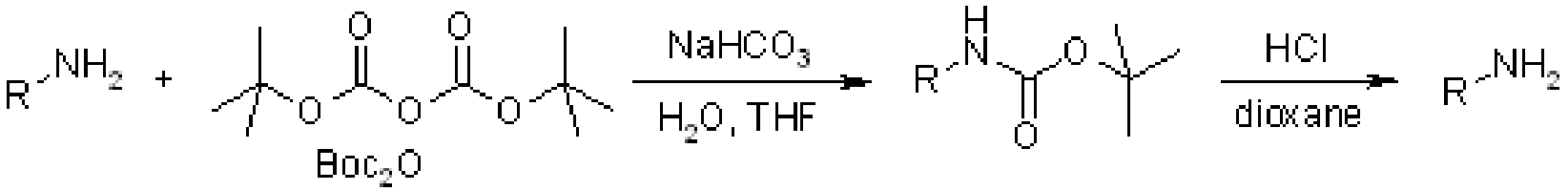
- Nucleophiles give 2 electrons to cause heterolytic cleavage
- Electrophiles take two electrons
- Free radical reactions give 1 electron to each section of a homolytic cleavage
- Note that electrons can transfer around organic rings (pericyclic), e.g. in benzene
- Acid-base reactions, e.g. HCl gas dissolving in water to form hydrochloric acid

Einige nützliche Reaktionen



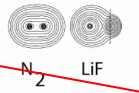
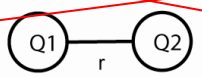
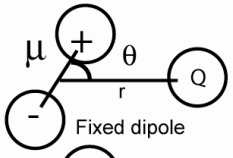
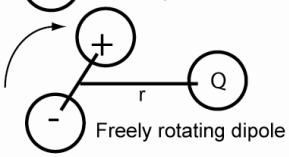
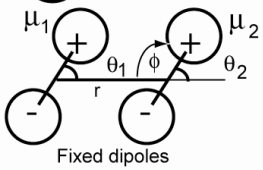
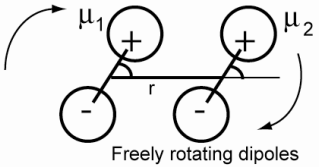
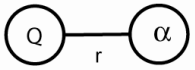
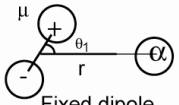
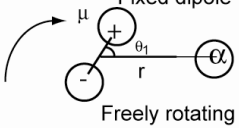
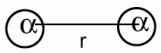
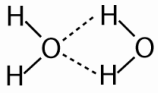
(Courtesy Professor Peiming Zhang, Arizona State University.)

Protection/deprotection

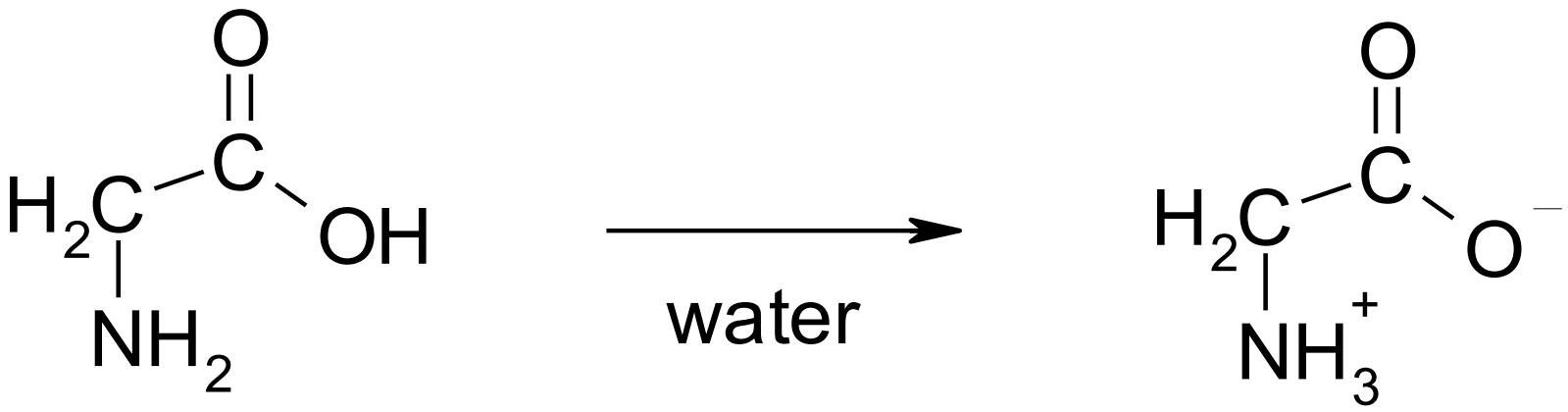


- BOC = *t*-ButylOxyCarbonyl
- Protects against nucleophiles
- Easily removed with acids

Beispiele für schwache Wechselwirkungen

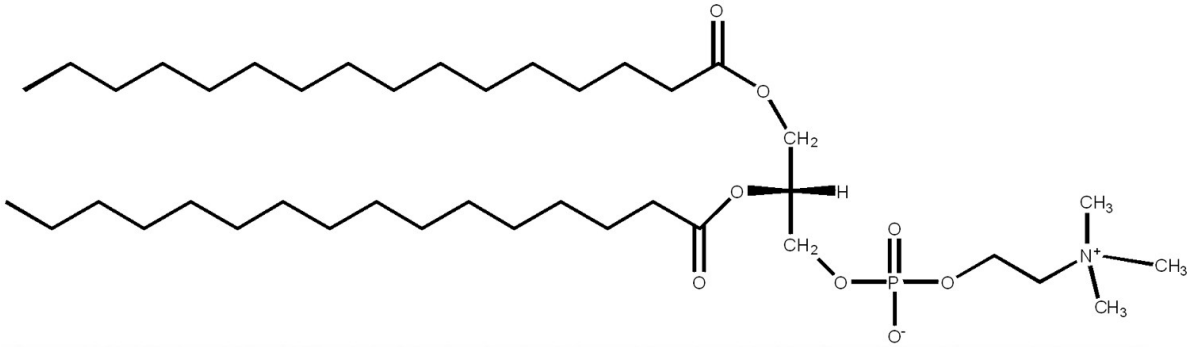
Interaction	Illustration	Formula for interaction energy
Covalent		Complex, short range - see Chapter 2
Charge-charge		Coulomb's law: $\frac{Q_1 Q_2}{4\pi\varepsilon_0 r}$
Charge-dipole		dipole moment is μ $-\frac{Q\mu\cos\theta}{4\pi\varepsilon_0 r^2}$
		$\frac{-Q\mu^2}{6(4\pi\varepsilon_0)k_B T r^4}$
Dipole-dipole		$-\frac{\mu_1\mu_2[2\cos\theta_1\cos\theta_2-\sin\theta_1\sin\theta_2\cos\phi]}{4\pi\varepsilon_0 r^3}$
		$\frac{-\mu_1^2\mu_2^2}{3(4\pi\varepsilon_0)^2 k_B T r^6}$
Charge-nonpolar		polarizability is α $\frac{-Q^2\alpha}{2(4\pi\varepsilon_0)^2 r^4}$
Dipole-nonpolar		$\frac{-\mu^2\alpha(1+3\cos^2\theta)}{3(4\pi\varepsilon_0)^2 r^6}$
		$\frac{-\mu^2\alpha}{(4\pi\varepsilon_0)^2 r^6}$
Nonpolar-nonpolar		$-\frac{3}{4}\frac{h\nu\alpha^2}{4\pi\varepsilon_0 r^6}$
Hydrogen bonds		Quantum-mechanical - roughly $\frac{1}{r^2}$

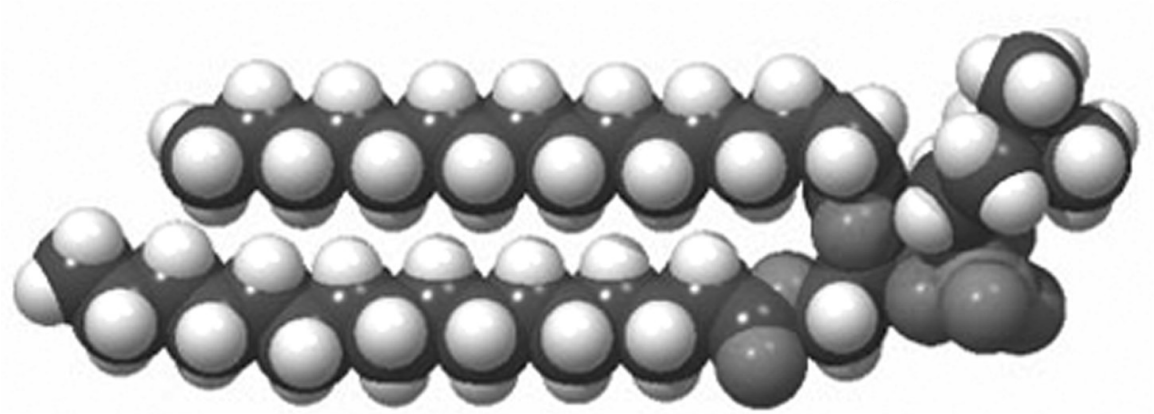
Ein „Zwitterion“: Unpolares Molekül wird in Wasser zum polaren Molekül



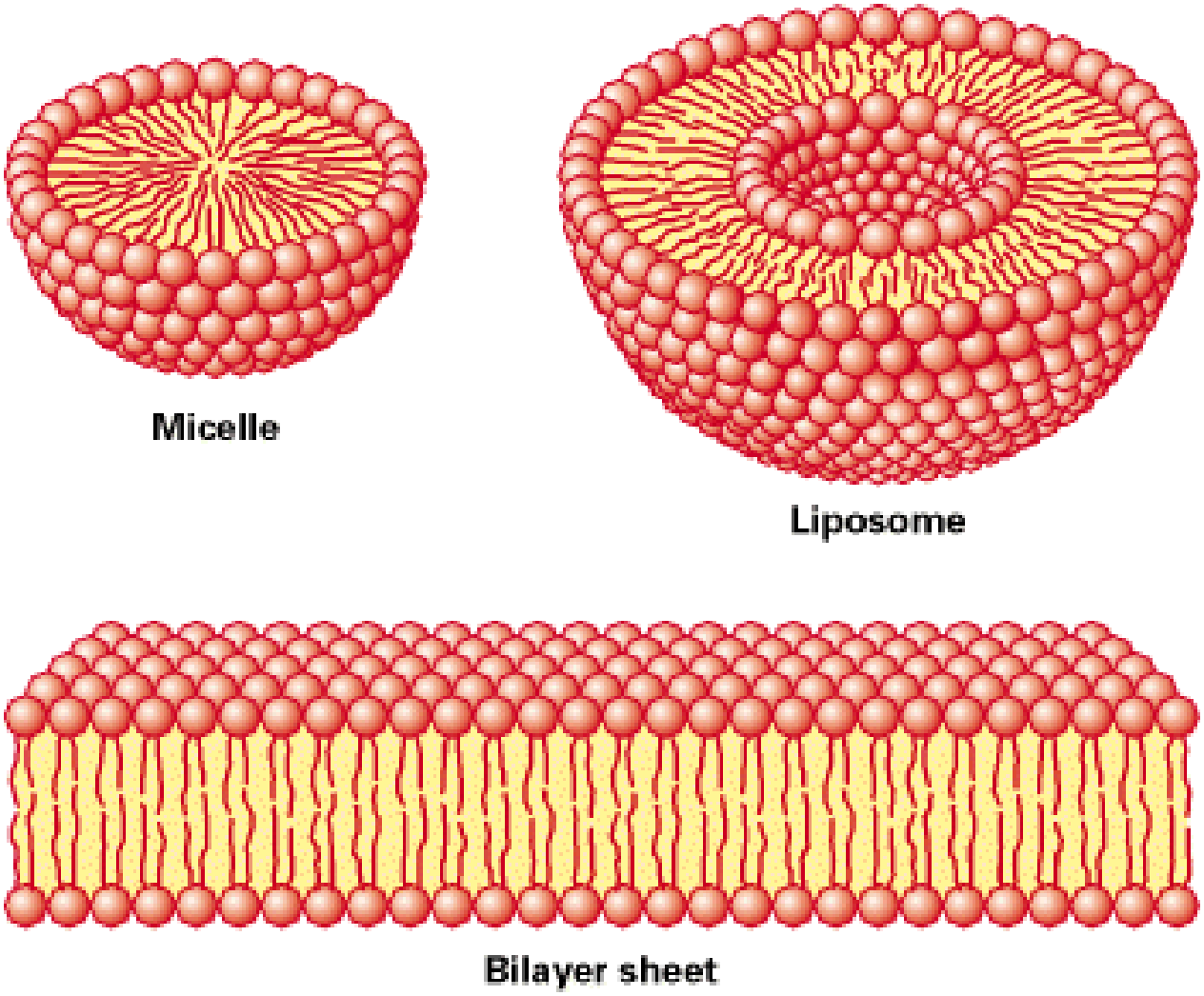
Amphiphile Moleküle: Moleküle mit hydrophoben und hydrophilem Ende

Bsp.: Phospholipidmoleküle mit hydrophiler Phosphat-Kopfgruppe und hydrophobem paarigen Schwanz aus Alkanketten (CH_2)



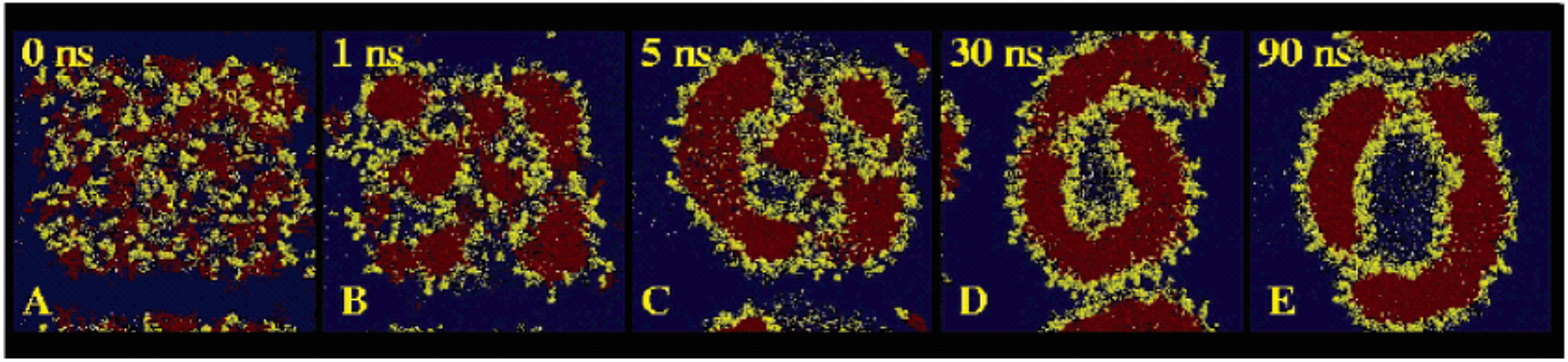


Selbstaufbauende amphiphile Strukturen



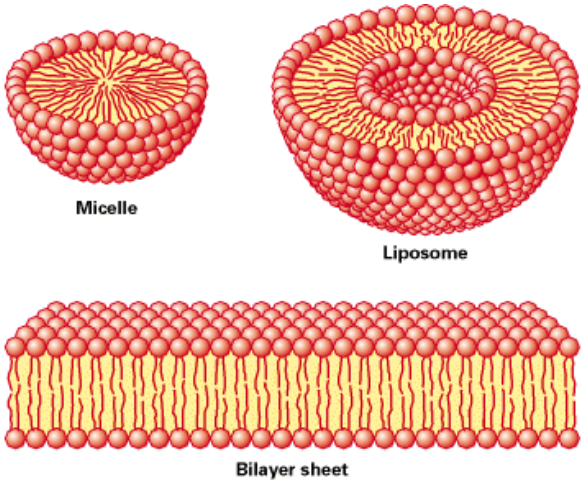
(From Molecular Cell Biology, 4th ed. By H. Lodish, A. Berk, S.L. Zipursky, P. Matsudara, D. Baltimore, J. Darnell.

Simulation der Formierung eines Vesikels



(Reprinted with permission from Molecular dynamics simulation of the spontaneous formation of a small DPPC vesicle in water in atomistic detail, A.H de Vries et al., A.E. Mark, and S.J. Marrink, J. Am. Chem. Soc. 2004 **126**: 4488. Published 2006 by American Chemical Society)

- Packing effects depend on geometry
- Non-spherical micelles $\frac{v}{a_0 \ell_c} < \frac{1}{3}$
- Spherical micelles $\frac{1}{3} < \frac{v}{a_0 \ell_c} < \frac{1}{2}$
- Vesicles/bilayers $\frac{1}{2} < \frac{v}{a_0 \ell_c} < 1$
- cones $\frac{v}{a_0 \ell_c} > 1$



(From Molecular Cell Biology, 4th ed. By H. Lodish, A. Berk, S.L. Zipursky, P. Matsudara, D. Baltimore, J. Darnell.
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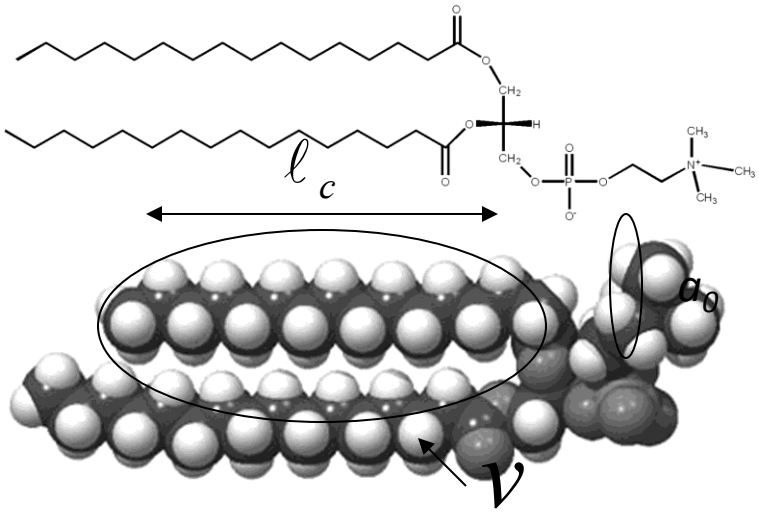


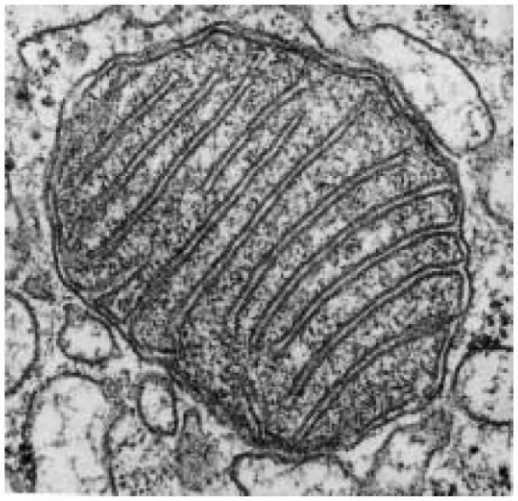
ℓ_c is the length of the hydrocarbon chain

v is the volume occupied by the hydrocarbon chain

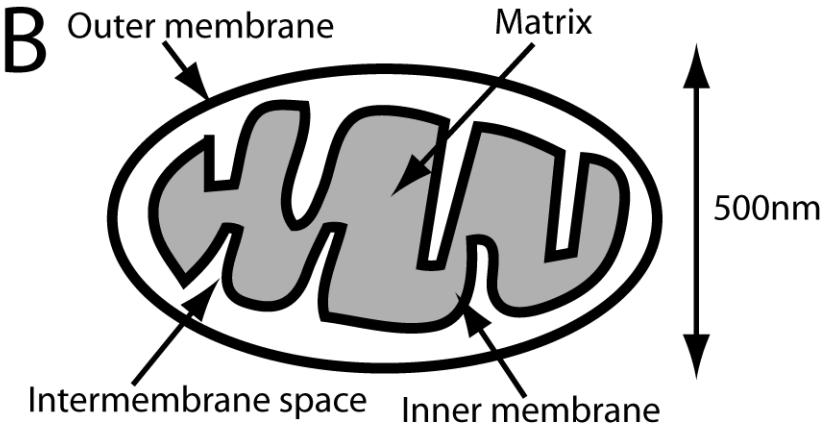
a_0 is the area of the head group

Beispiel für eine Lipid-Doppelschicht: Das Mitochondrium

A



B



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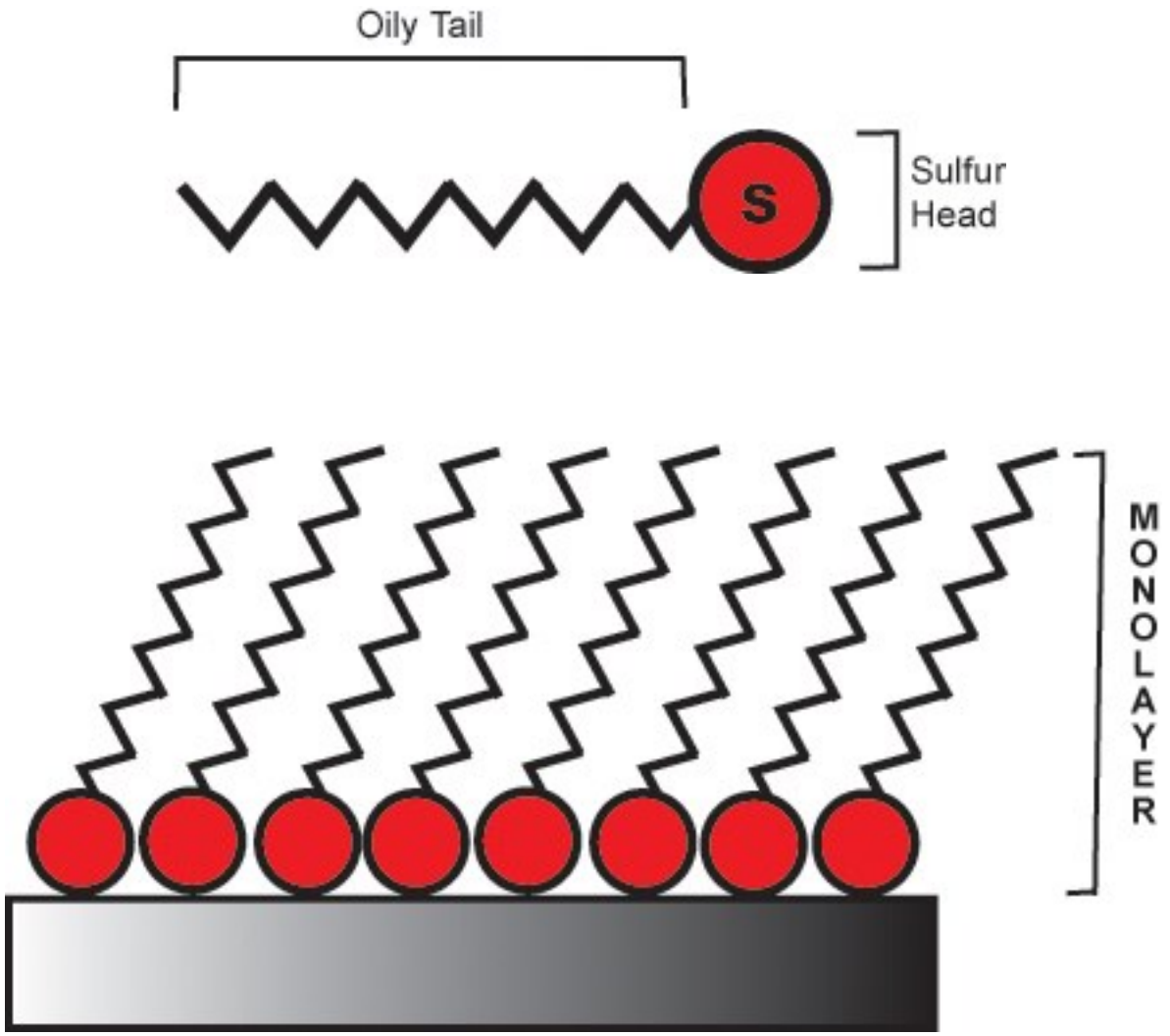
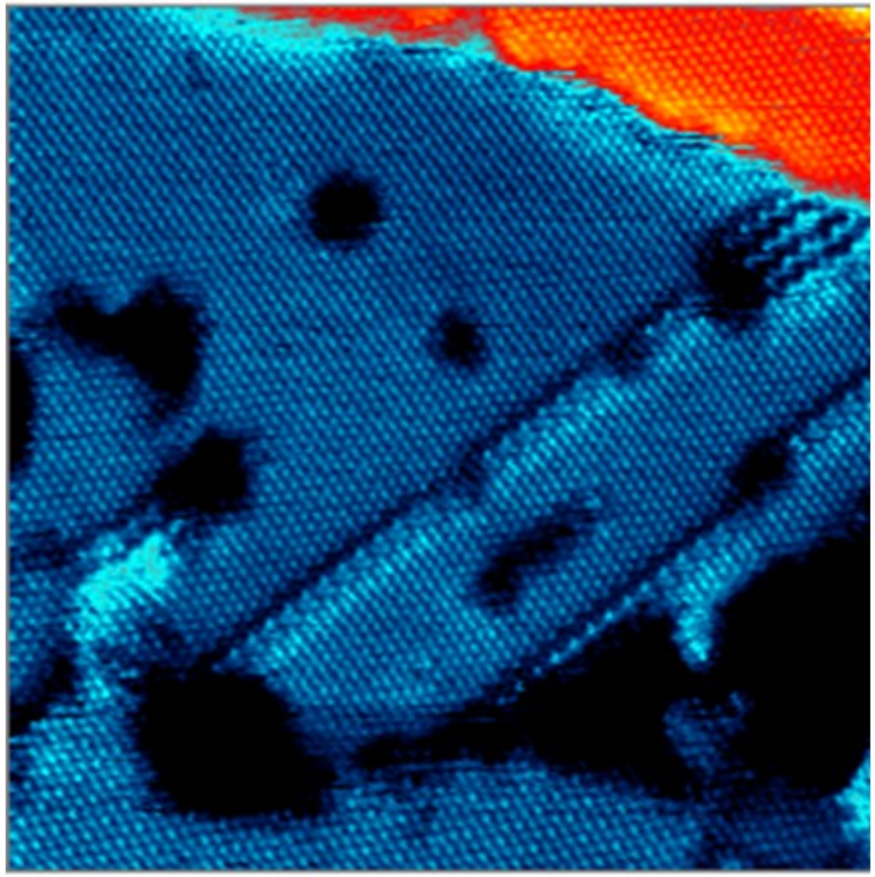
Pyruvate oxidation – 30 ATPs vs 2

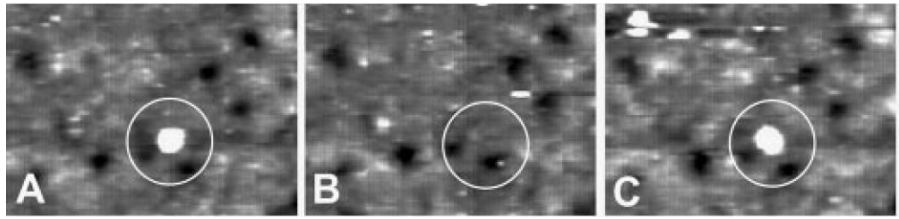
Beispiel für SAM



Self-assembled monolayers (SAM)



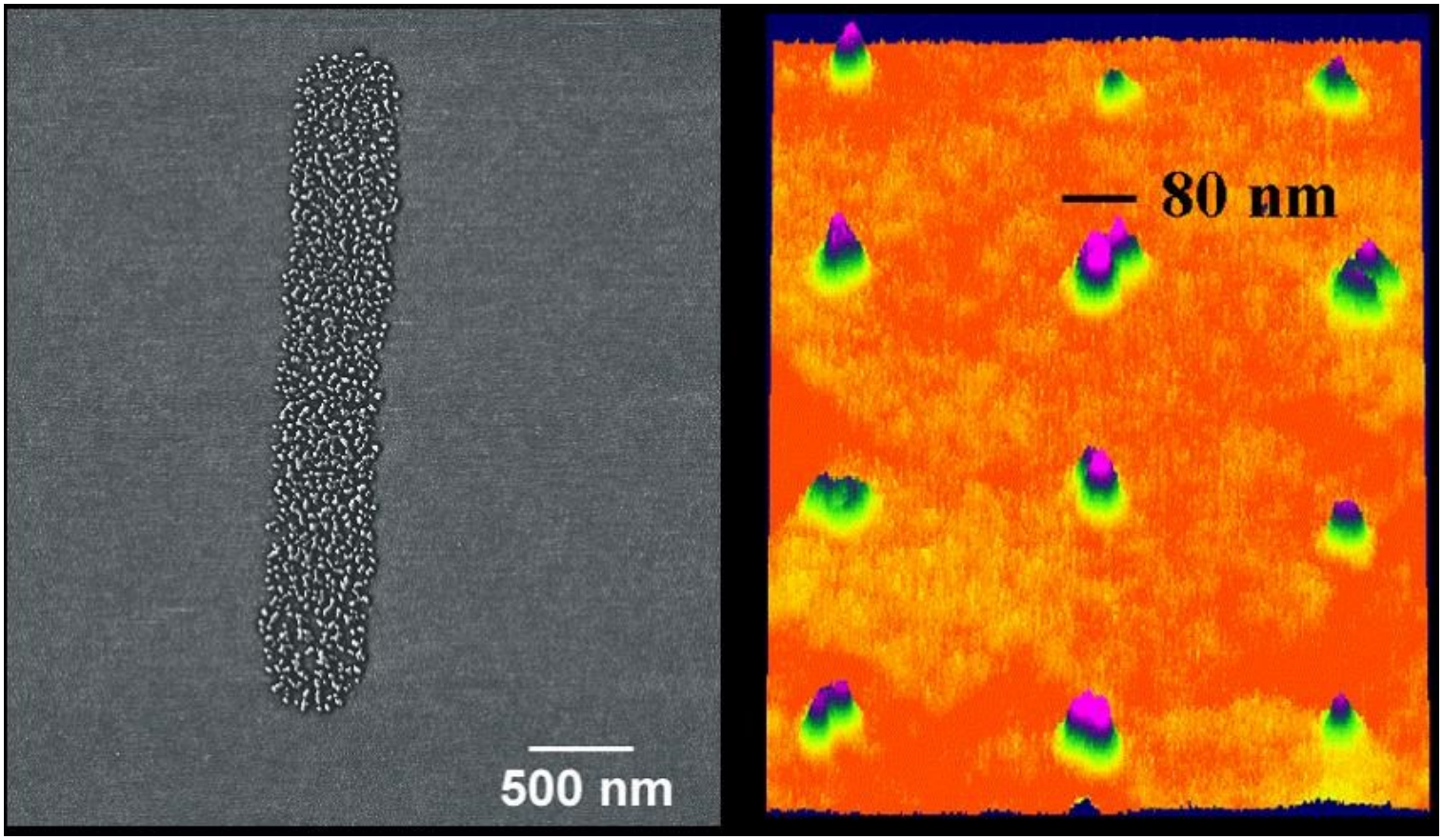
(Reproduced with permission from Functional molecules and assemblies in controlled environments, Weiss, P.S. published by Accounts of Chemical Research, 2008, courtesy of Professor Paul Weiss.)



Bond fluctuations

(From A bond-fluctuation mechanism for stochastic switching in wired molecules, G.K. Ramachandran, T.J. Hopson, A.M. Rawlett, L.A. Nagahara, A. Primak and S.M. Lindsay, Science 2003, 300, 3413. Reprinted with permission AAAS. Readers may view, browse and/or download material for temporary copying purposes only, provided that these uses are for noncommercial personal purposes. Except as provided by law, this material may not be further reproduced, distributed, transmitted, modified, adapted, performed, displayed, published or sold in whole or part without prior written permission from the publisher.)

Beispiel für die Anwendung von SAMs: Biopatterning von Proteinen
Array zur Untersuchung <1000 Antikörper (links) und einzelner Moleküle (rechts)



Dip-Pen-Nanolithografie

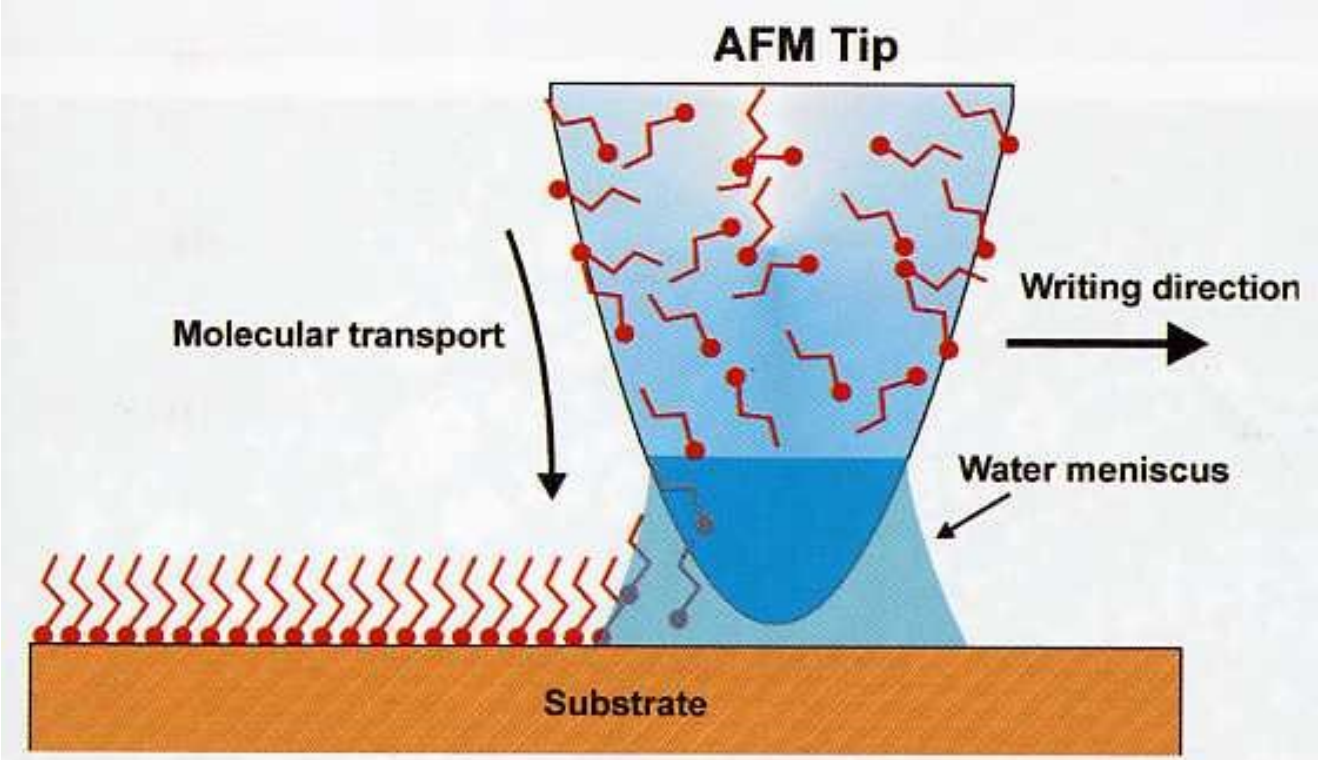
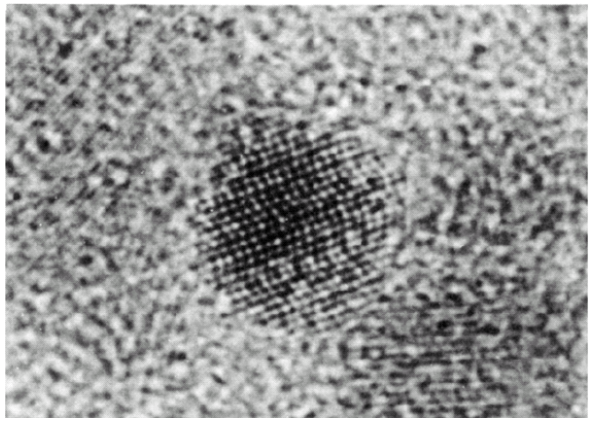


Figure 4.3

Schematic of the dip pen lithography process—the wiggly lines are molecular "ink."

Courtesy of the Mirkin Group, Northwestern University.

Kinetisch getriebene Prozesse: Ausfällung von Nanopartikeln aus Lösungen



Quantum dots from 2 phase synthesis with Ostwald ripening

(Reprinted with permission from Synthesis and characterization of nearly monodisperse CdE (E=S,Se,Te) semiconductor nanocrystallites, C.B. Murray, D.J. Noms and M.G. Bawendi, J. Am. Chem. Soc. 115 8706 Published 1993 by American Chemical Society).

Kinetisch getriebene Prozesse: Ausfällung von Nanopartikeln aus der Gasphase

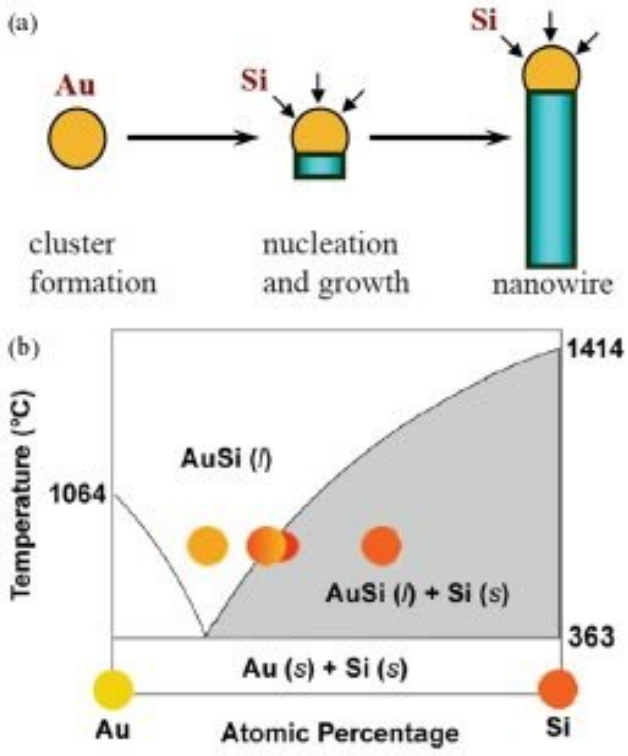
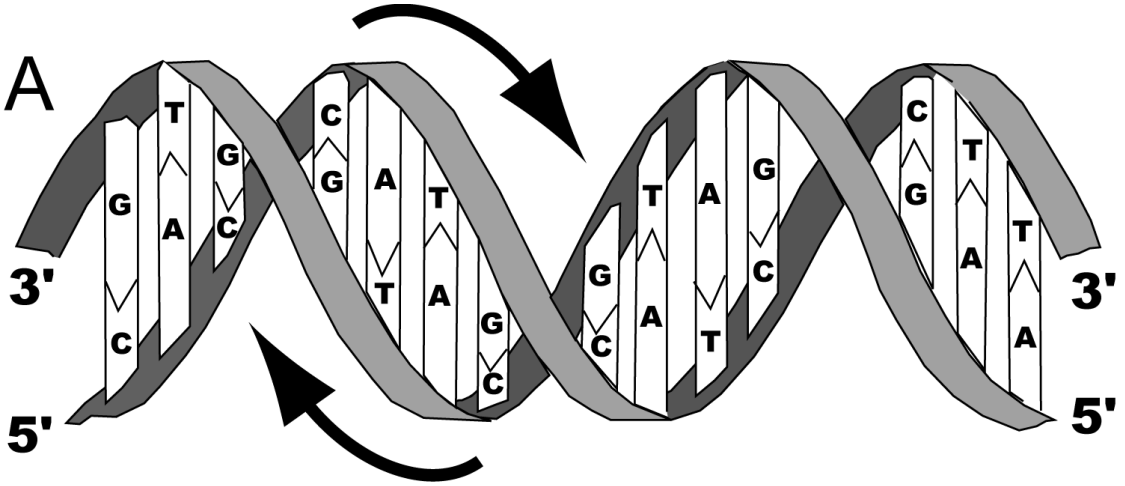


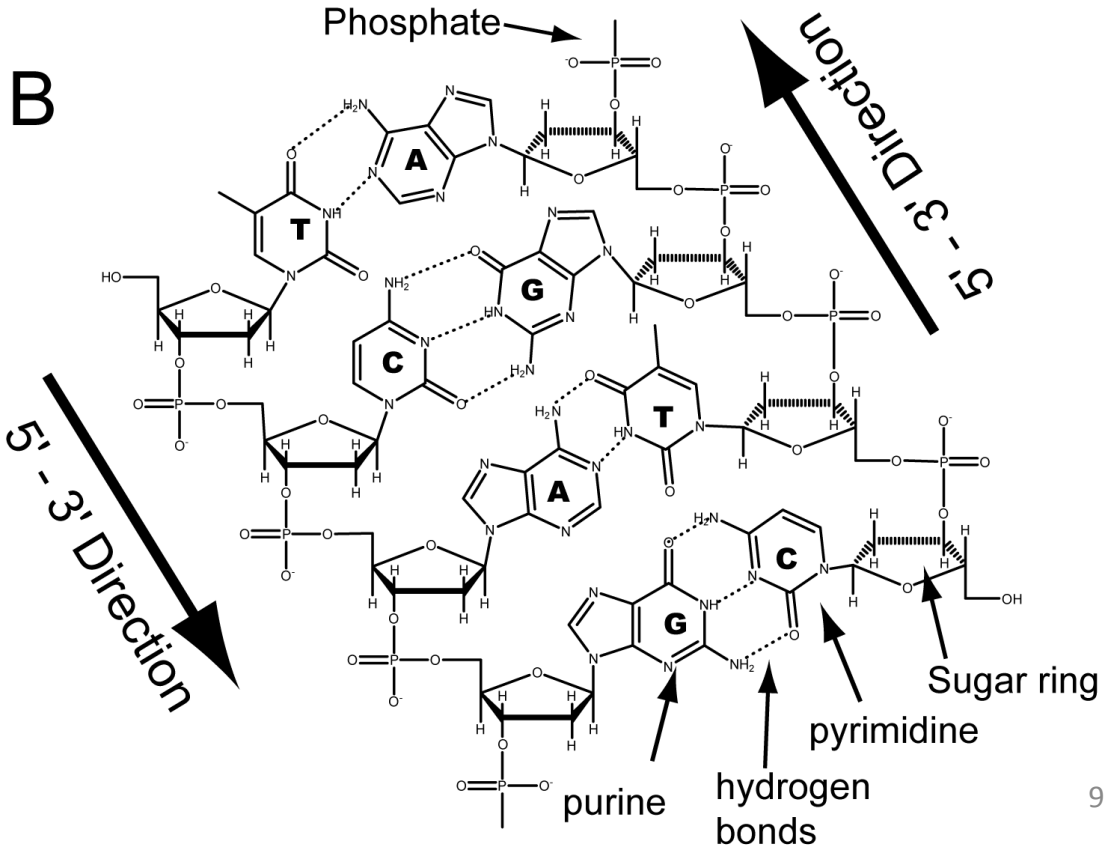
Figure 1. Schematic of VLS growth of Si nanowires (SiNWs). (*a*) A liquid alloy droplet AuSi is first formed above the eutectic temperature (363 °C) of Au and Si. The continued feeding of Si in the vapour phase into the liquid alloy causes oversaturation of the liquid alloy, resulting in nucleation and directional nanowire growth. (*b*) Binary phase diagram for Au and Si illustrating the thermodynamics of VLS growth.

Si Nanowires from Au/Si eutectic seeded on Au NanoParticles

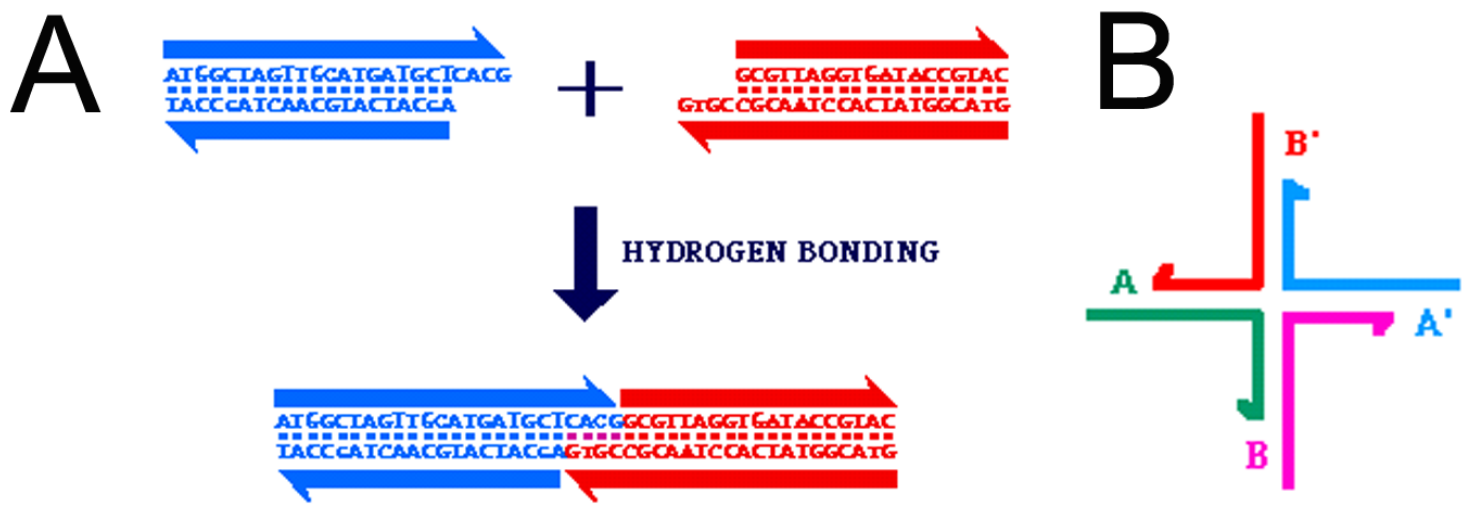
(Reproduced with permission from Semiconductor nanowires, W. Lu and C.M. Lieber J. Phys. D: Applied Physics 2006 with permission from IOP publishing and courtesy Wei Lu.)

DNA Nanotechnology



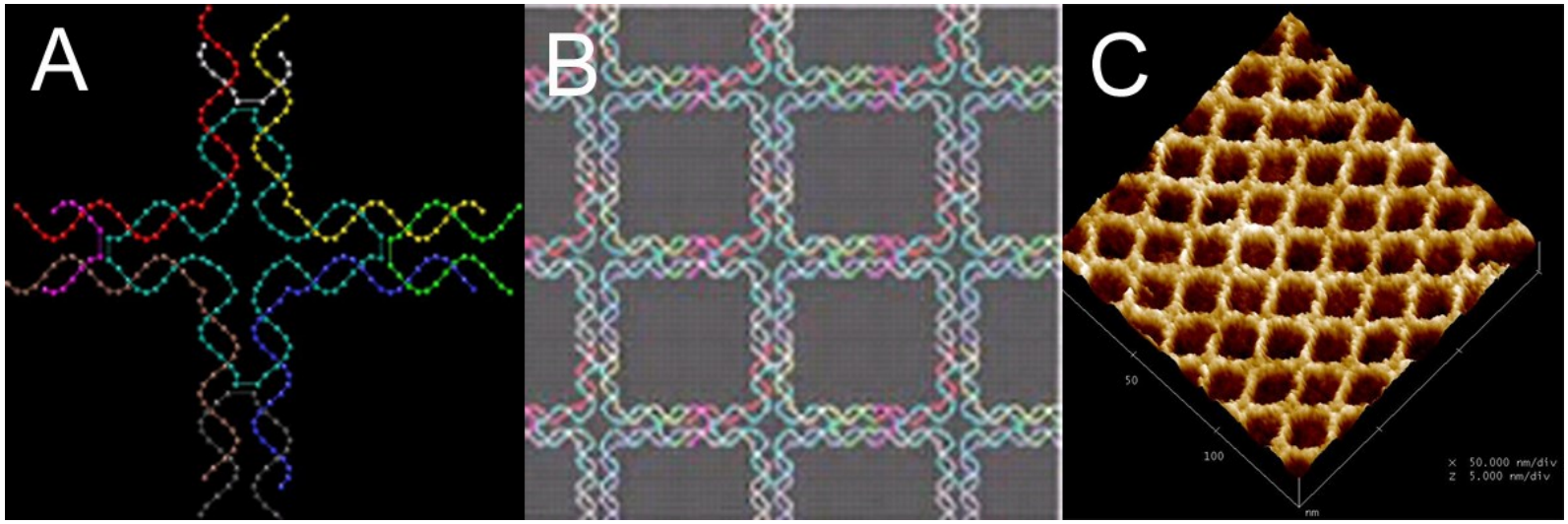


DNA Nanotechnology



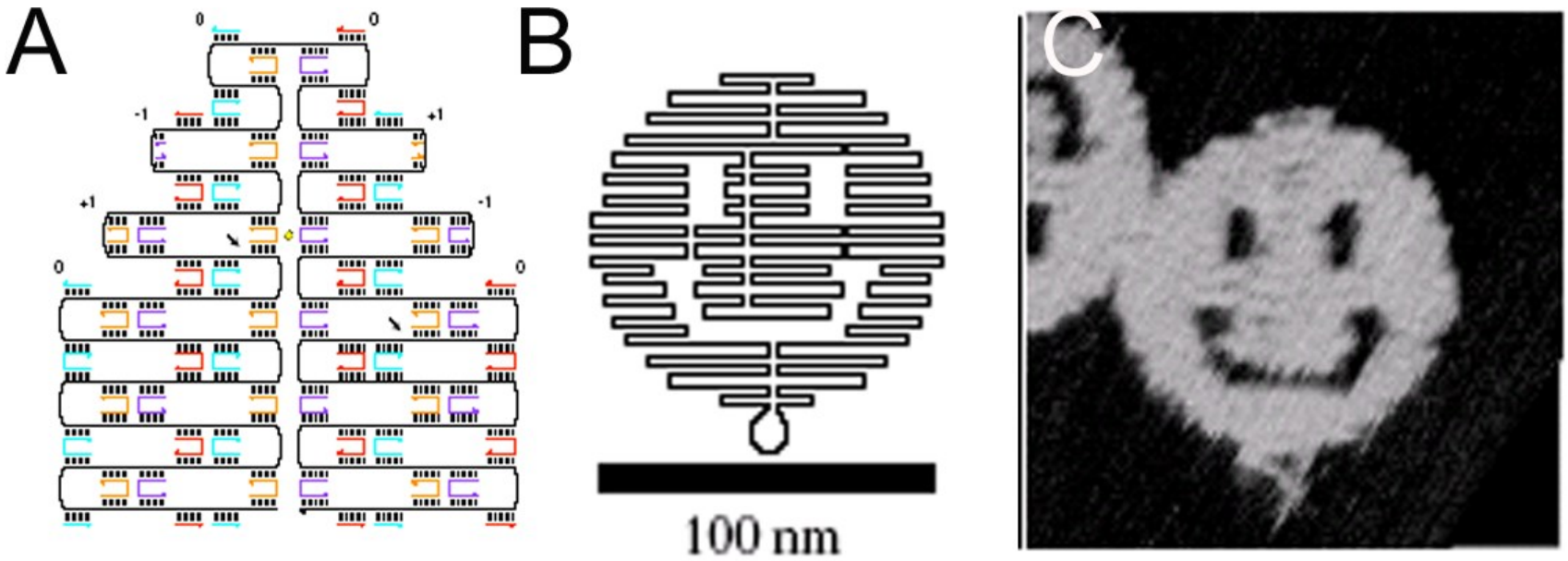
(Courtesy of Professor Hao Yan, Arizona State University)

DNA Nanotechnology



(Courtesy of Professor Hao Yan, Arizona State University)

DNA Nanotechnology



(Reprinted by permission from McMillan Publishers Ltd.: Nature Publishing Group, Folding DNA to create nanoscale shapes and patterns, P. Rothmunde, Nature 2006, **440**, 297.)

DNA Nanotechnology

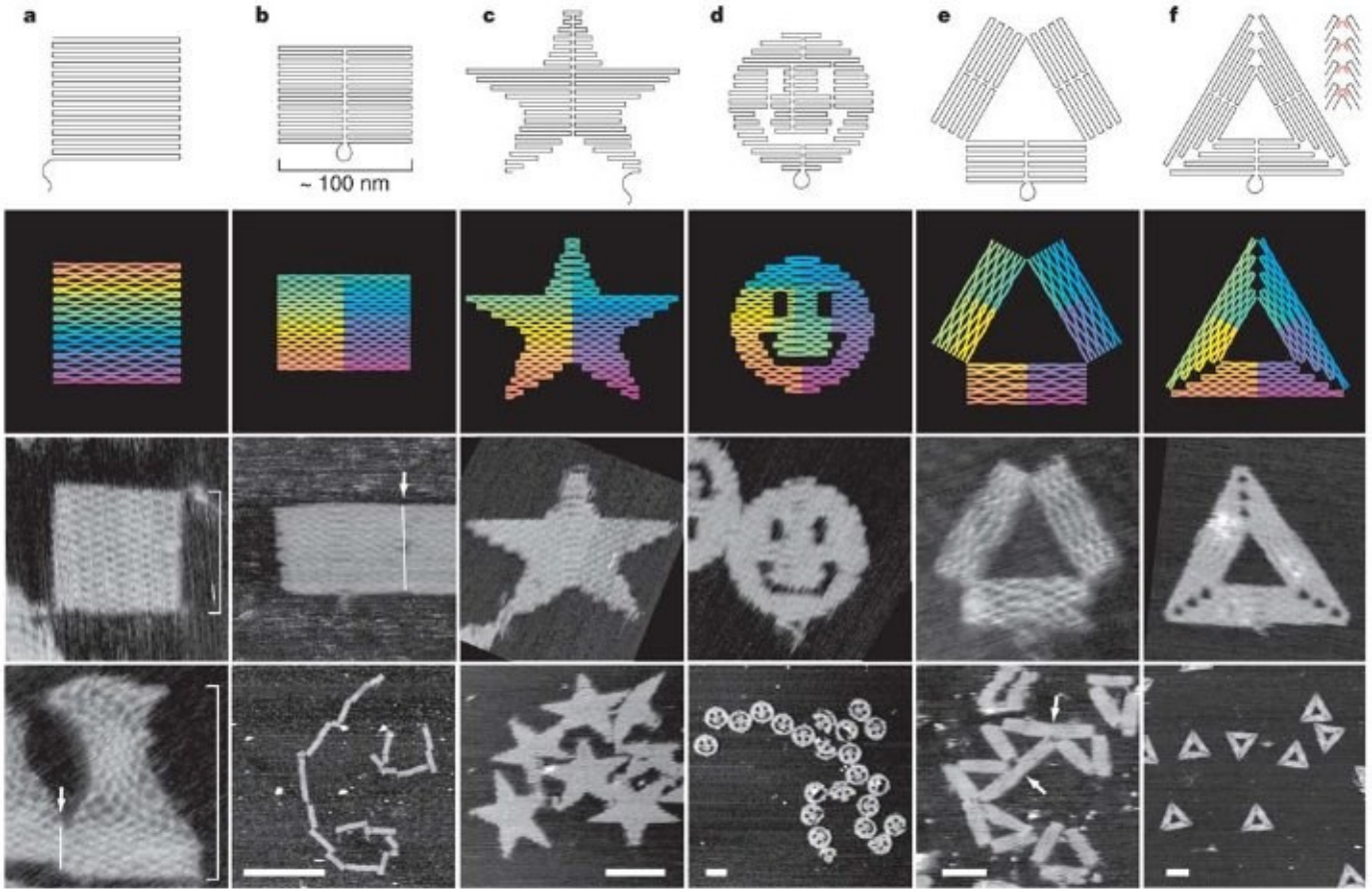


Figure 2 | DNA origami shapes. Top row, folding paths. **a**, square; **b**, rectangle; **c**, star; **d**, disk with three holes; **e**, triangle with rectangular domains; **f**, sharp triangle with trapezoidal domains and bridges between them (red lines in inset). Dangling curves and loops represent unfolded sequence. Second row from top, diagrams showing the bend of helices at crossovers (where helices touch) and away from crossovers (where helices bend apart). Colour indicates the base-pair index along the folding path; red is the 1st base, purple the 7,000th. Bottom two rows, AFM images. White lines and arrows indicate blunt-end stacking. White brackets in **a** mark the height of an unstretched square and that of a square stretched vertically (by a factor >1.5) into an hourglass. White features in **f** are hairpins; the triangle is labelled as in Fig. 3k but lies face down. All images and panels without scale bars are the same size, 165 nm × 165 nm. Scale bars for lower AFM images: **b**, 1 μm; **c**–**f**, 100 nm.

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Partikuläre Nanostrukturen

Abb. 2 Der „Lycurgus-Kelch" aus dem 4. Jahrhundert erscheint im reflektierten Licht grün. Wird er jedoch von innen beleuchtet, dann erstrahlt er in einem satten Rot.

Abb. 1 Je nach der Größe der sphärischen Gold-Cluster im Lösungsmittel variiert der Farbeindruck.

Eines der farbenprächtigen Fenster der St. Nicolai-Kirche in Kalkar. Bei seinen Entwürfen ließ sich der Künstler Karl-Martin-Hartmann von physikalischen Motiven inspirieren.

Physik Journal 5 (2006) 52

Partikuläre Nanostrukturen

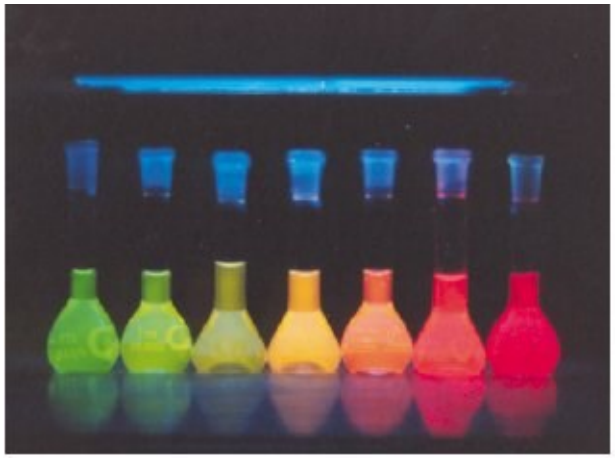


Abb. 1:
Unterschiedlich große CdSe-Nanopartikel in Lösung lassen sich mit UV-Licht zur Emission in den verschiedensten Farben anregen.

Physik Journal 3 (2004) 118

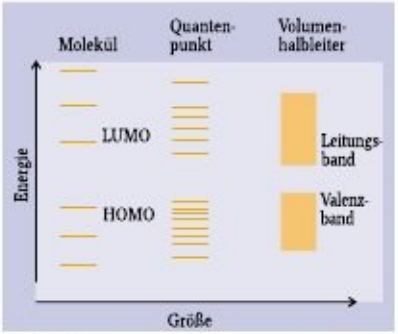


Abb. 2:
Das Energiespektrum eines Quantenpunkts (Mitte) liegt zwischen dem eines Moleküls (links) und dem eines Festkörpers (rechts). Die Energiezustände sind diskret und man spricht vom niedrigsten unbesetzten bzw. höchsten besetzten Molekül-Orbital (LUMO bzw. HOMO)

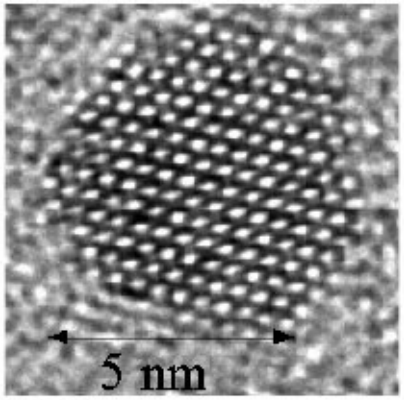
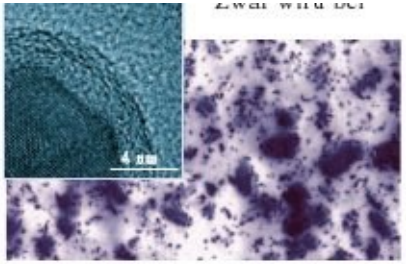


Abb. 3:
Durch nasschemische Verfahren lassen sich Halbleiter-Nanopartikel in Gramm-Mengen herstellen. Die Elektronenmikroskop-Aufnahme zeigt einen 6nm großen, einkristallinen CdSe-Nanokristall mit einer facettierten Oberfläche.



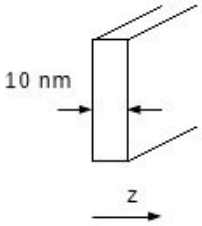
TEM-Aufnahme von kristallinen TiO_2-Teilchen (dunkel) mit einem Diketopyrrolopyrrol-Pigment (hell, größer) in einem Bindemittel. In höherer Vergrößerung (Inset) lässt sich sogar die Kristallstruktur und die Beschichtung eines TiO_2-Nanopartikels erkennen. (Foto: Sachtleben)

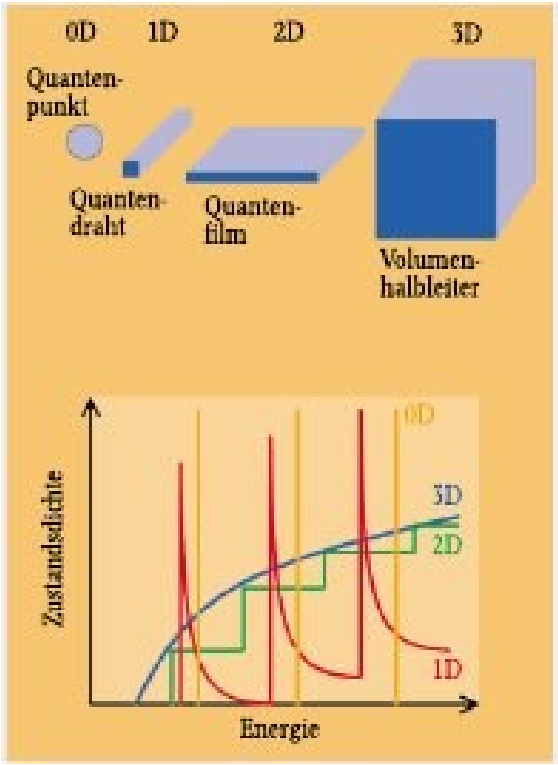
Physik Journal 3 (2004) 118

Physik Journal 1 (2002) 47

Festkörper in reduzierter Dimension

	$d = 3$	$d = 2$	$d = 1$	d=0
$\epsilon(\vec{k})=$	$\frac{\hbar^2(k_x^2+k_y^2+k_z^2)}{2m^*}$	$\frac{\hbar^2(k_x^2+k_y^2)}{2m^*}+E_m$	$\frac{\hbar^2 k_x^2}{2m^*}+E_m+E_n$	$E_m+E_n+E_l$
		„Quantenfilm“	„Quantendraht“	„Quantenpunkt“

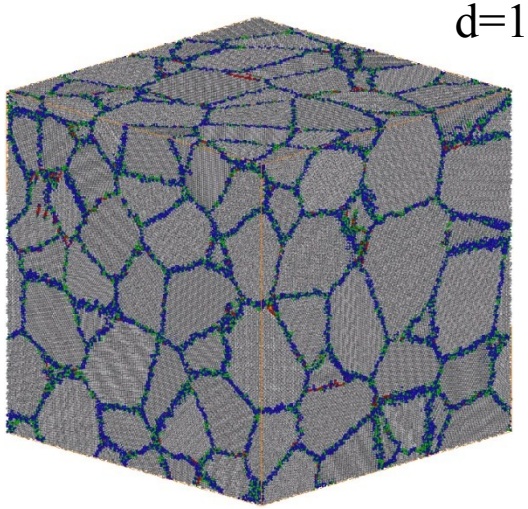
	Metall	2-dimensionales Elektronengas (z.B. GaAs)	Quantenpunkt (z.B. GaAs) „künstliches Atom“	Atom
Ausdehnung	1 cm^3		$(10\ nm)^3$ 	0.1 nm
charakteristische Wellenlänge	$\lambda_F \sim 0.5$ nm	$\lambda_F \sim 10\text{-}50$ nm	$a^* = 10$ nm	$a_B = 0.05$ nm
Energieskala	$\varepsilon_F \sim 5$ eV	$\varepsilon_F \sim$ meV	$\Delta\varepsilon \sim$ meV	$\Delta\varepsilon \sim 10$ eV
Anregungsenergie	10^{-10} eV	~ meV z-Quantisierung	~ meV	~ 10 eV

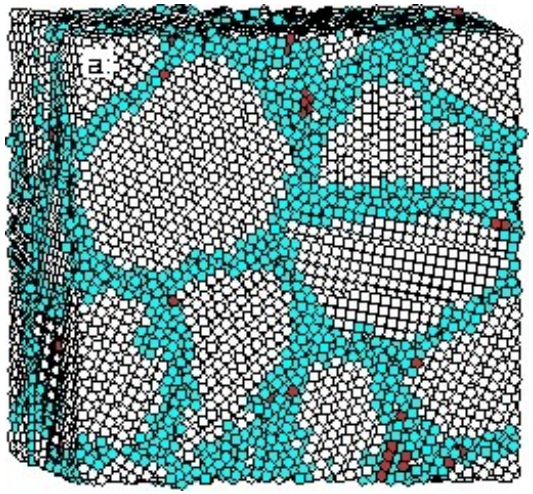


Physik Journal 1 (2002) 47

Nanokristalline Materialien

d=10nm





d=5.2nm

http://sic.epfl.ch/publications/SCR02/scr13_page9e.html

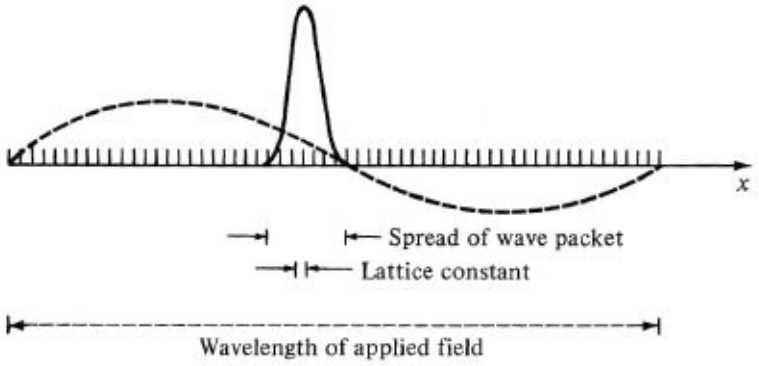
blau/cyan: Atome an Korngrenzen

grau: Atome im Volumen (fcc)

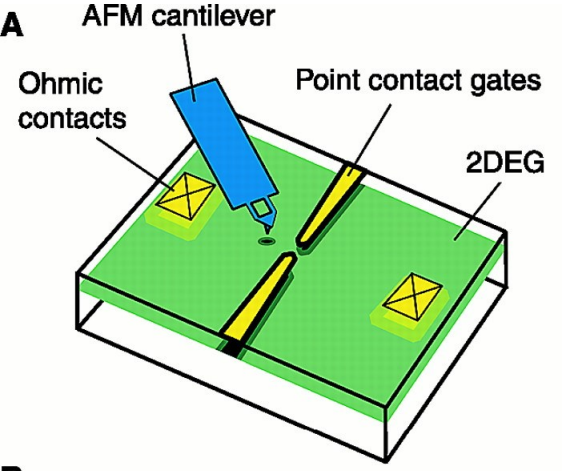
rot: Atome in Stapelfehlern (hcp)

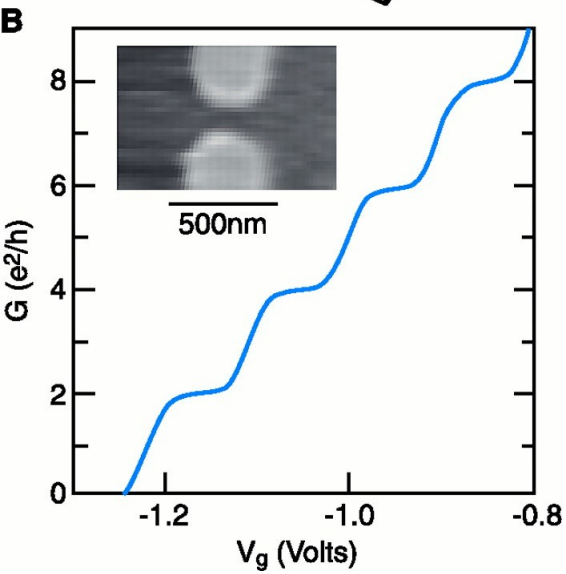
	bulk Ni	n-Ni d=100nm	n-Ni d=10nm
Fließfestigkeit/ Streckfestigkeit [MPa]	103	690	>900
Bruchfestigkeit [MPa]	403	1100	>2000
Vickers Härte [kg/mm^2]	140	300	650

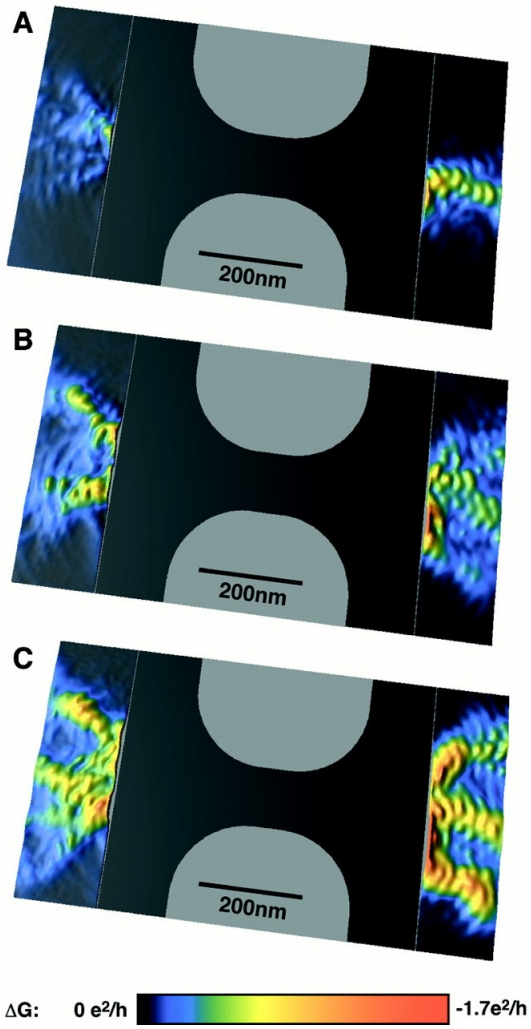
Elektronischer Transport im ausgedehnten Festkörper: semiklass. Beschreibung

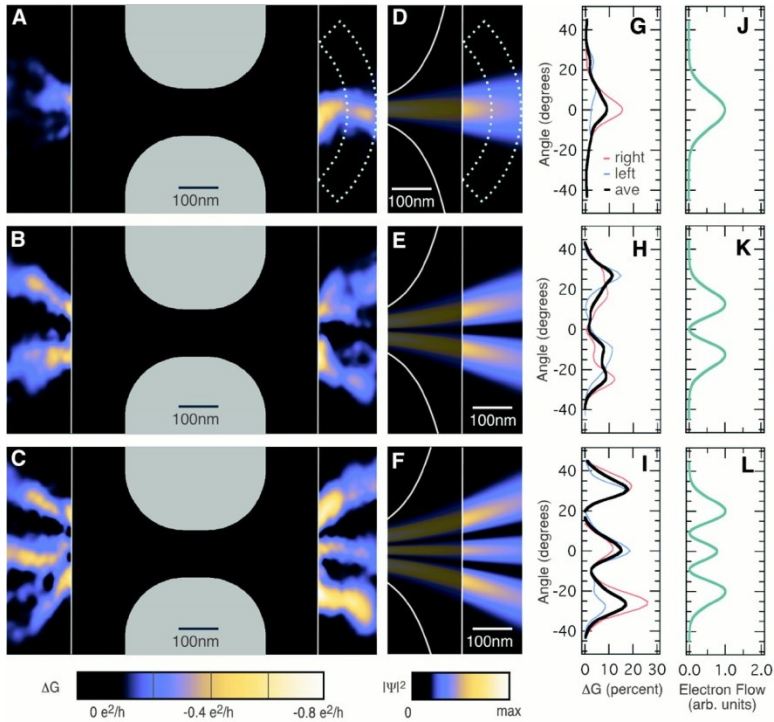


Leitwertquantisierung









Quelle: Topinka et al. Science 289, 2323 (2000)

Aharonov-Bohm-Effekt

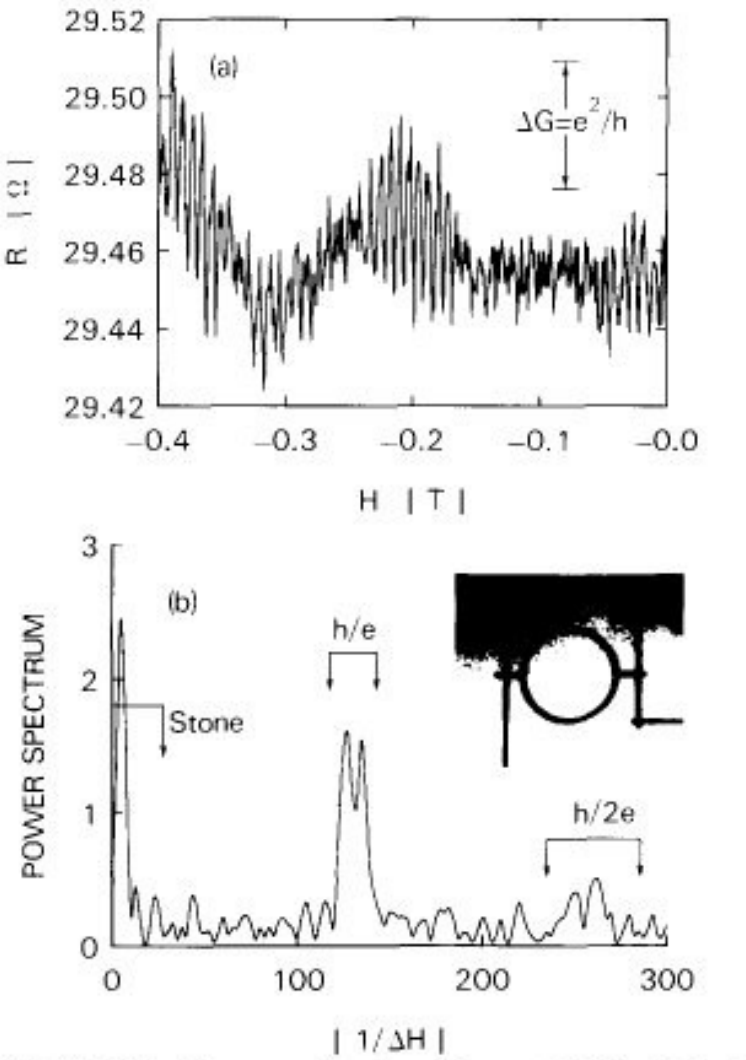
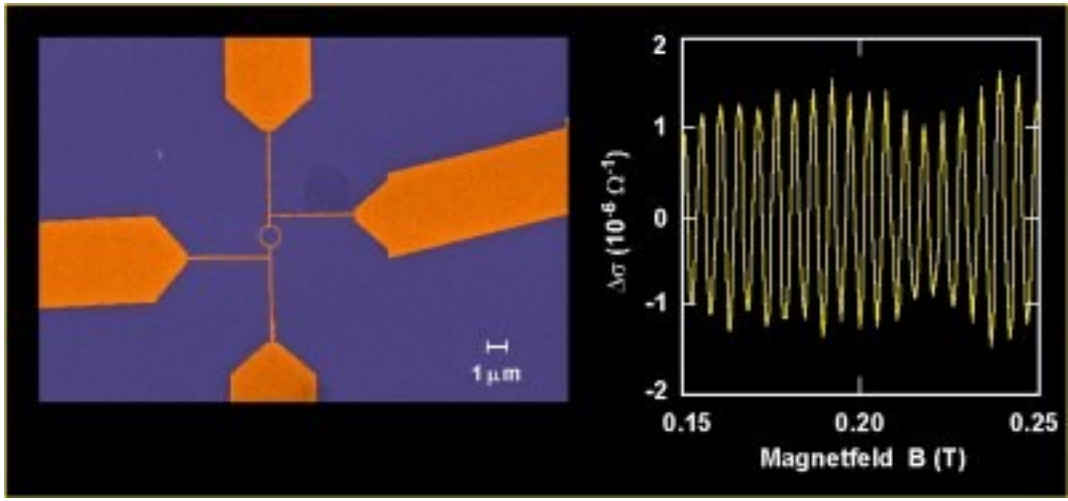


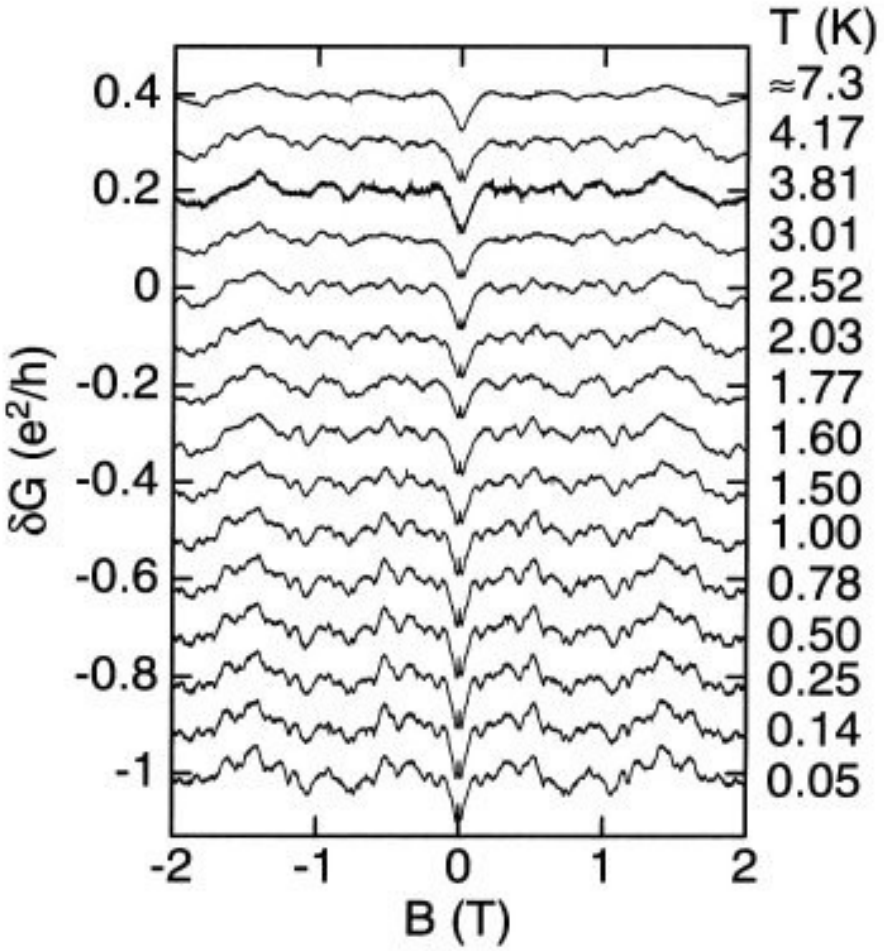
Fig. 1. (a) Resistance as a function of magnetic field obtained at 60 mK for an Au ring 0.825 μm average diameter and 0.041 μm line width. Arrows indicate the conductance scale in units of e^2/h. (b) Fourier transform of the data displayed above with arrows indicating the expected frequency range for h/e and $h/2e$ oscillations as well as the scale for the aperiodic background fluctuations predicted by Stone for this ring. The inset is an electron micrograph of this ring.

Quelle: Webb et al. JMMM, 54, 1423 (1986).



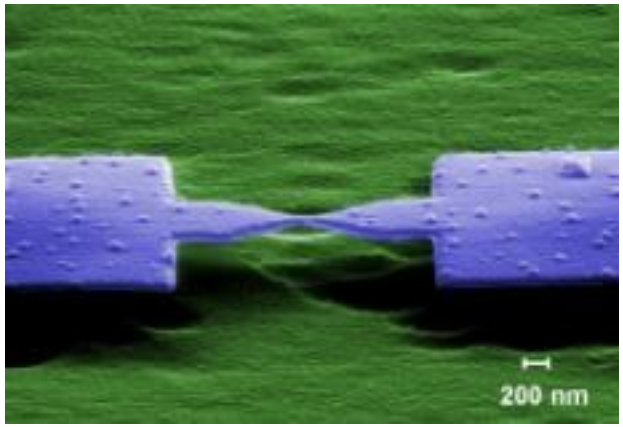
Quelle: R. Häussler, Dissertation Karlsruhe

Schwache Lokalisierung und Universelle Leitwertfluktuationen (UCF)

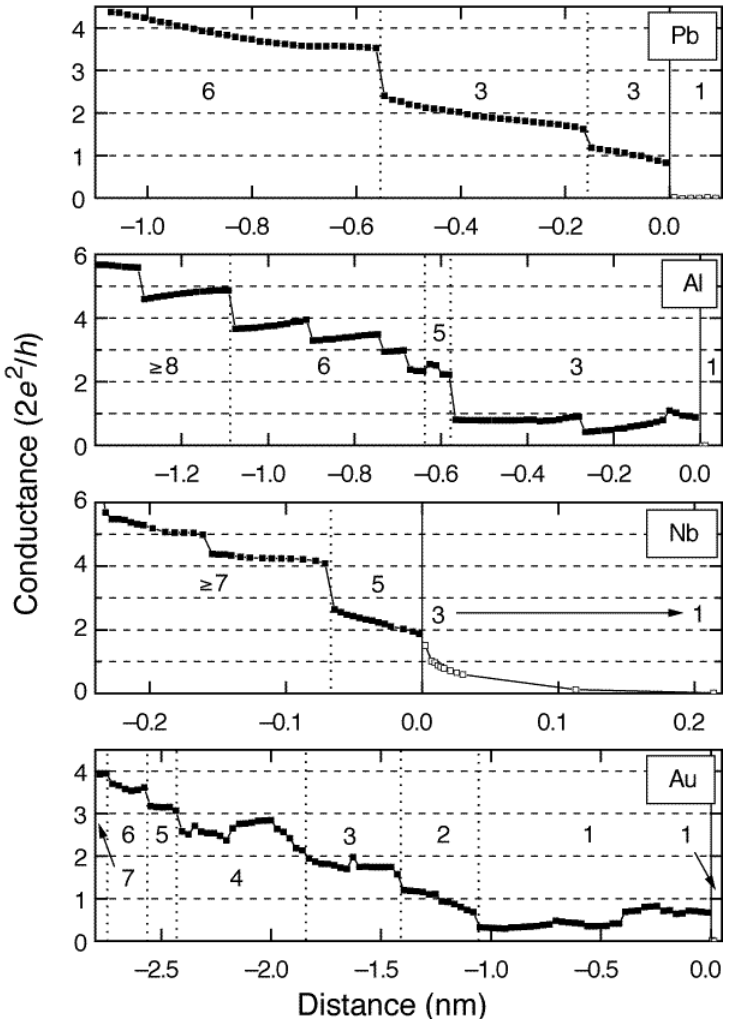


Quelle: E. Scheer, Dissertation Karlsruhe

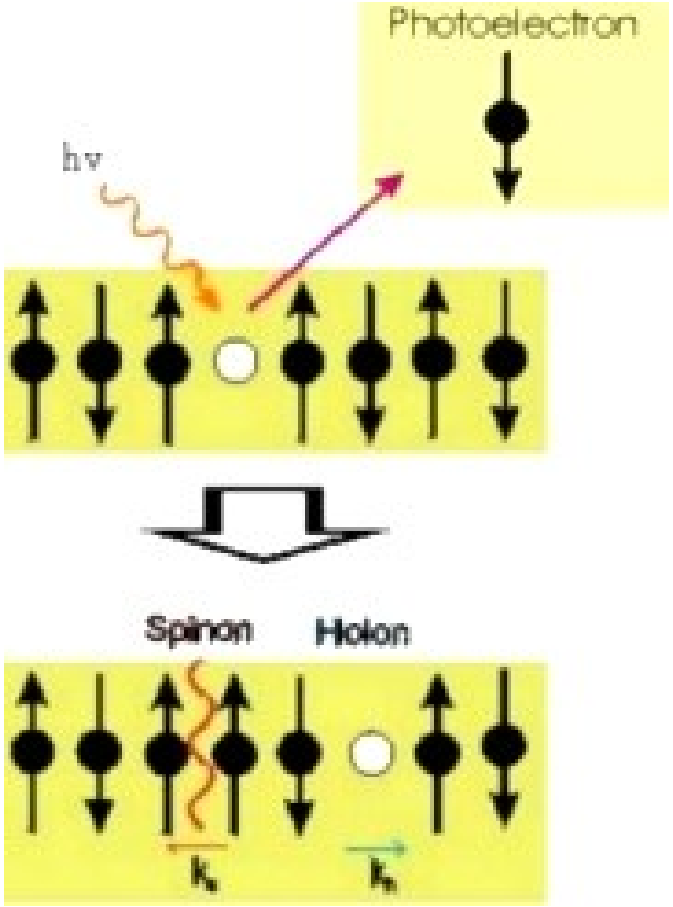
Leitwertstufen : Leitwert wird durch Transportkanäle bestimmt



Quelle: Scheer et al. Nature 394, 154(1998).

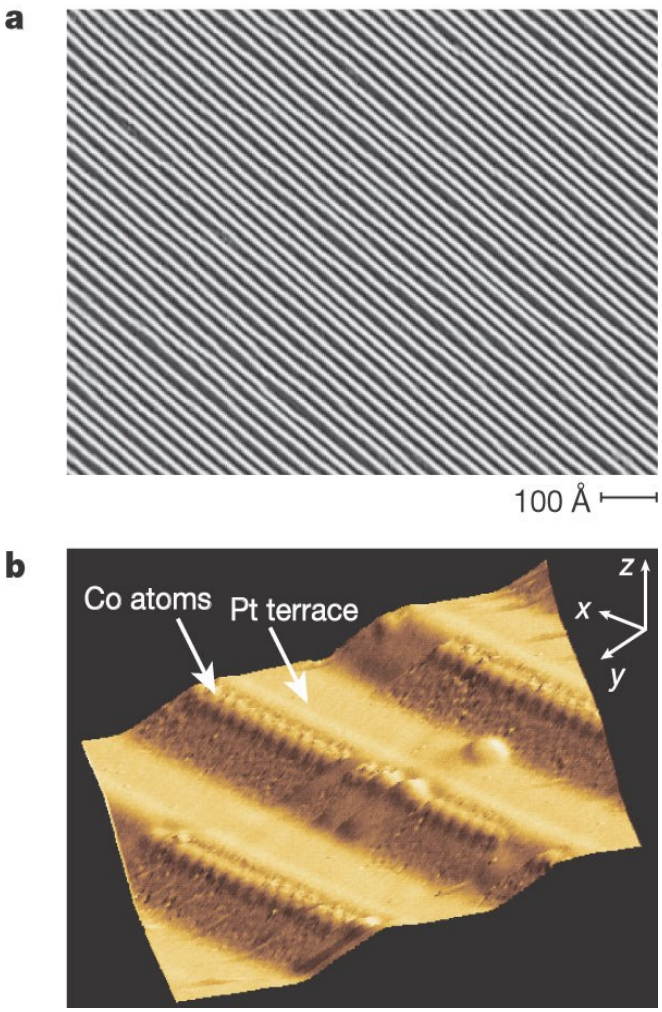


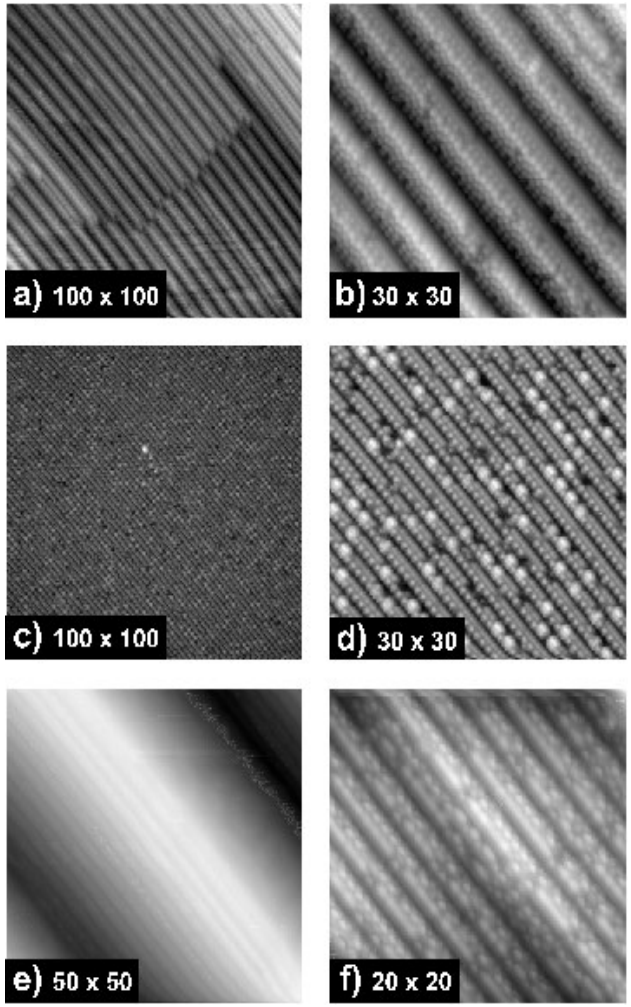
Spin-Ladungs-Entkopplung in 1D antiferromagnetischer Kette



Quelle: http://www.pi1.physik.uni-stuttgart.de/glossar/Luttinger_d.php

Eindimensionale monoatomare Metallketten



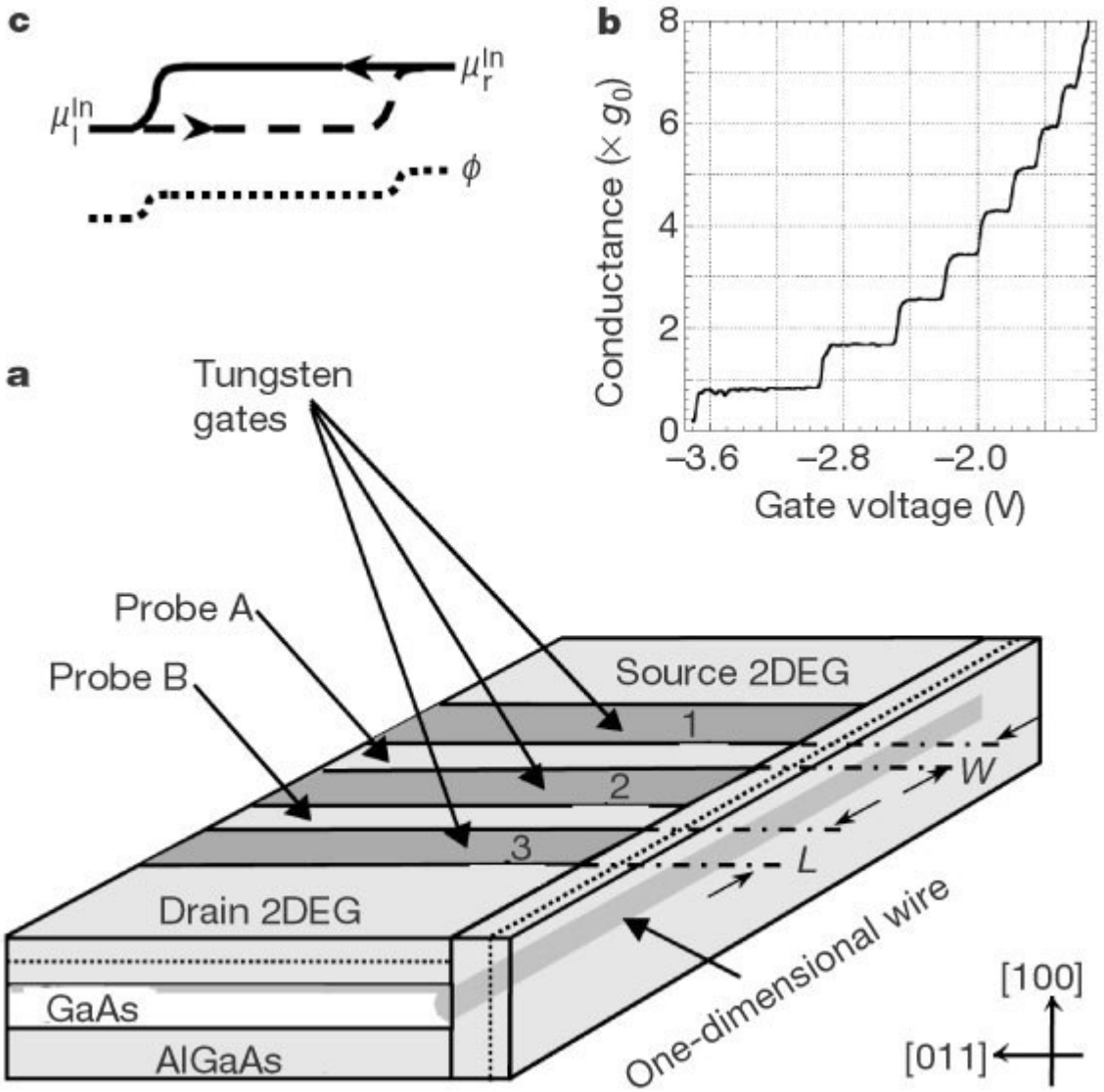


Co auf vicinalem Pt (111)

Au auf vicinalem Si (111)

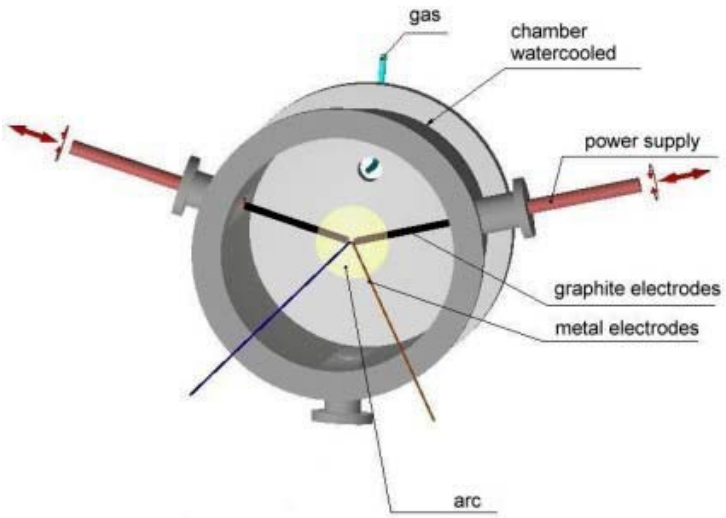
Quelle: Gambardella et al. Nature 416, 301 (2002).

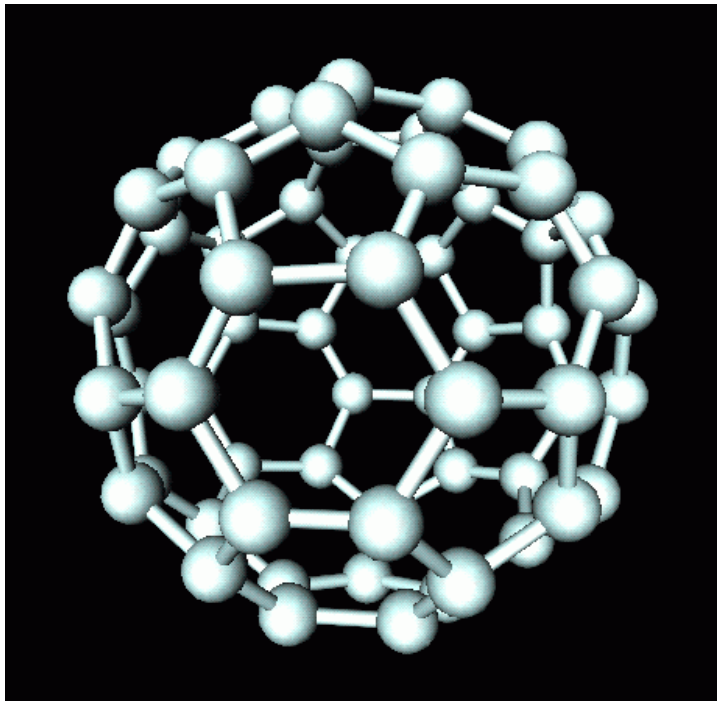
Eindimensionale Halbleiter-Strukturen hergestellt durch “Cleaved-Edge-Overgrowth”

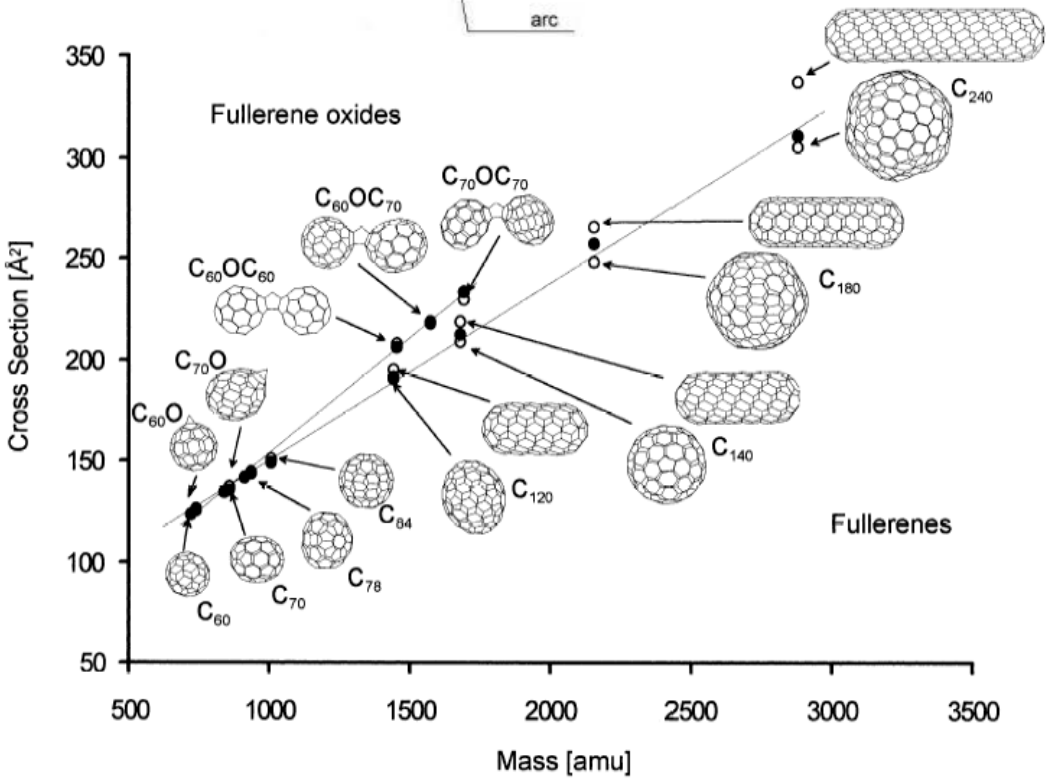


Quelle: Picciotto et al. Nature 411, 51 (2001).

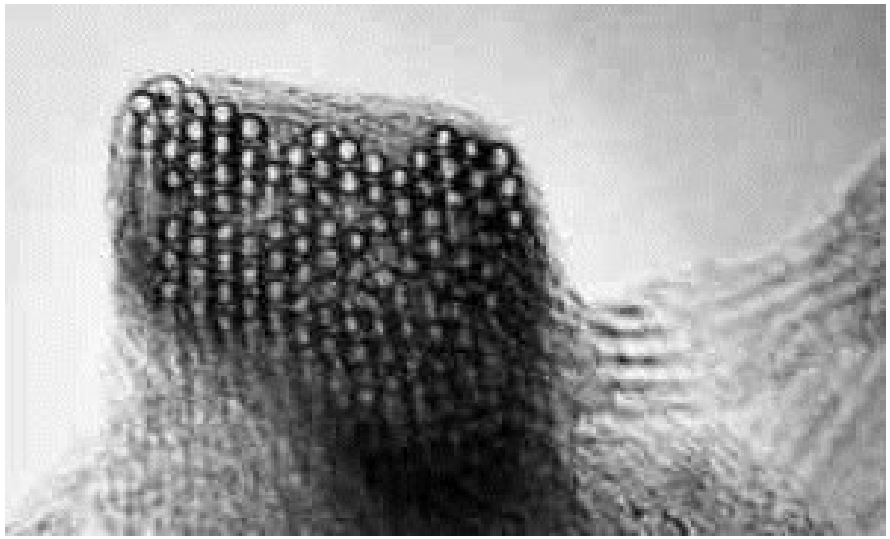
Fullerene

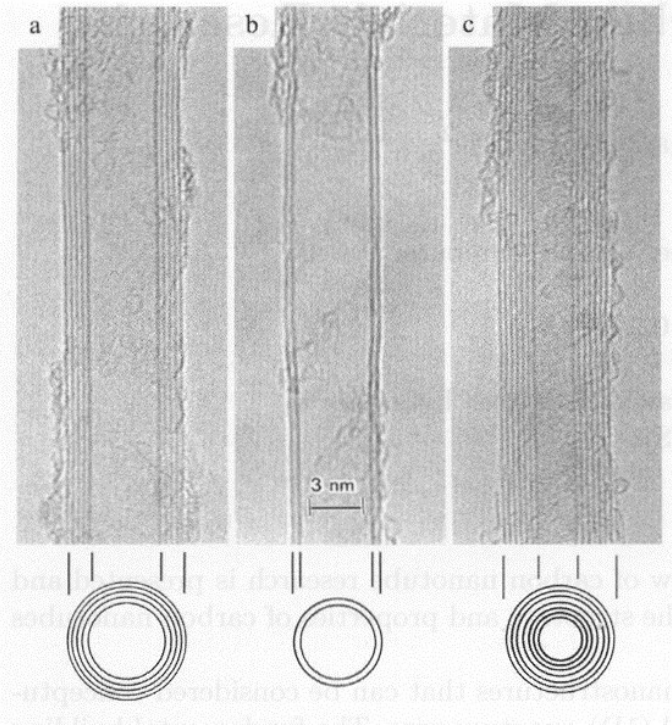

gas
chamber
watercooled
power supply
graphite electrodes
metal electrodes
arc




Fullerene oxides
Fullerenes
Cross Section [Å²]
Mass [amu]
350
300
250
200
150
100
50
500
1000
1500
2000
2500
3000
3500
C60
C70
C78
C84
C120
C140
C180
C240
C60O
C70O
C60OC60
C60OC70
C70OC70

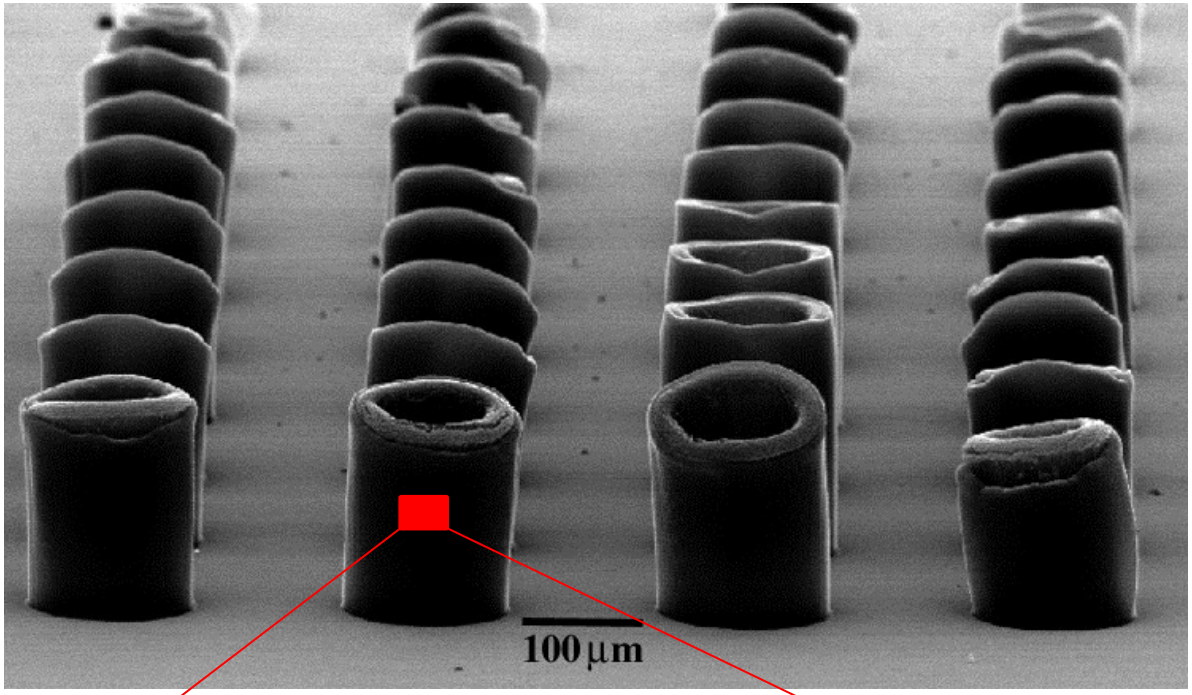
Kohlenstoff- Nanoröhren

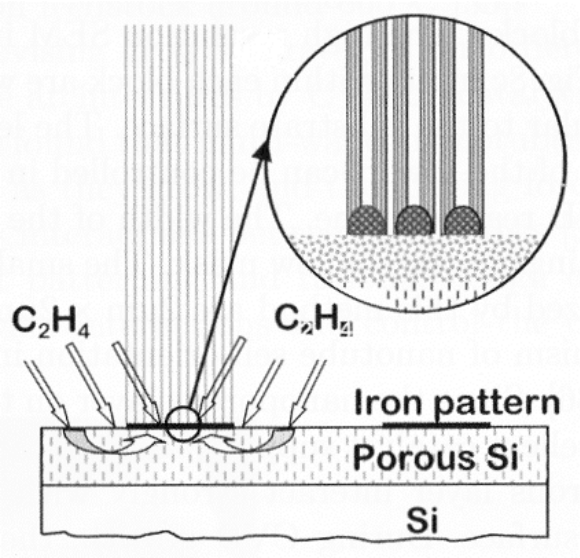


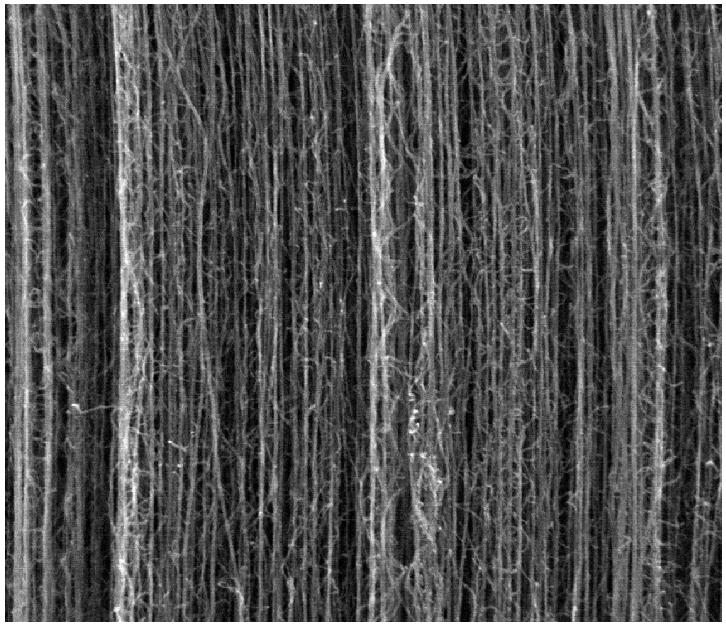


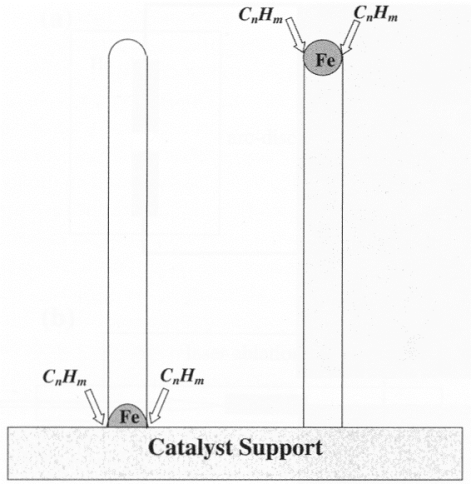
Quelle: Thess et al. Science 273, 483 (1996).

Herstellung von Kohlenstoff- Nanoröhren

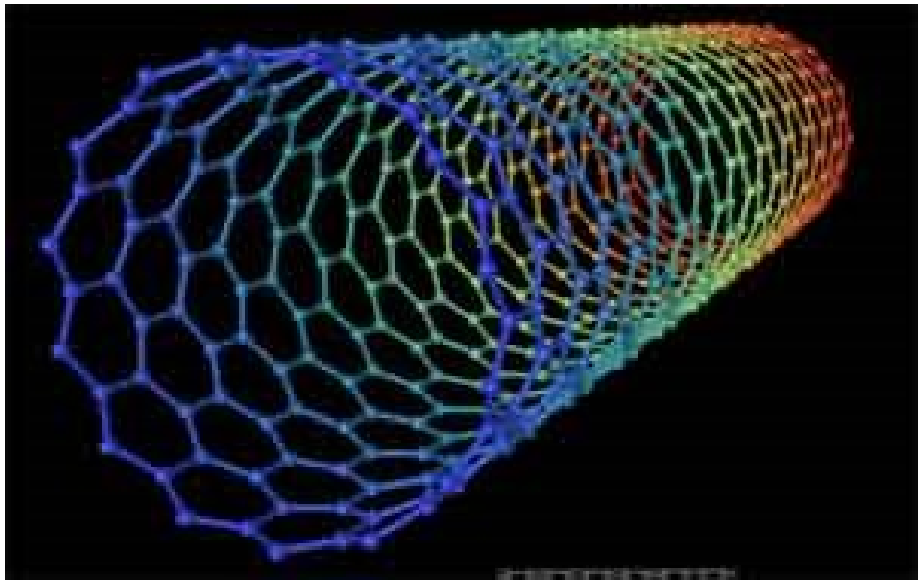


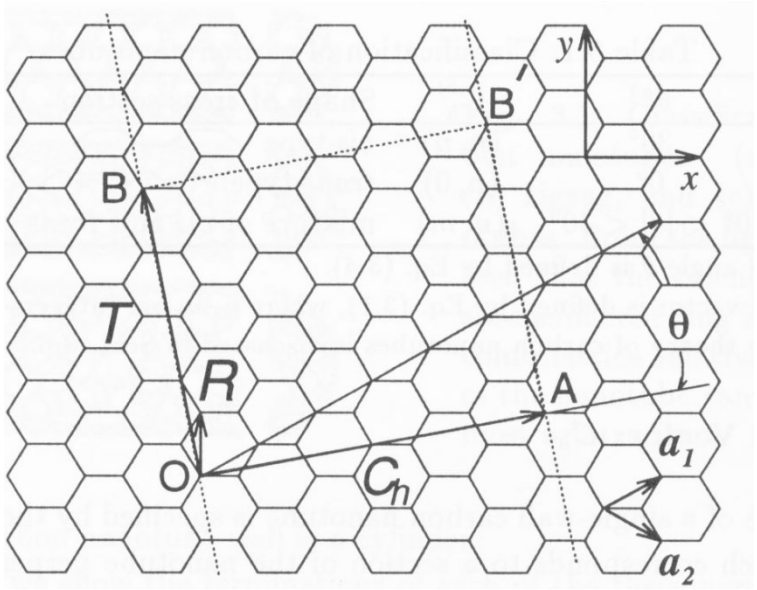






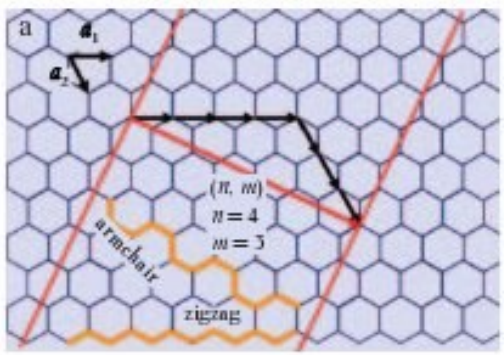
Geometrische Struktur von Kohlenstoff- Nanoröhren





Aufrollvektor (n,m): $C_h = na_1 + ma_2$

W. Hönlein, F. Kreupl, Physik Journal **3**, 39 (2004)



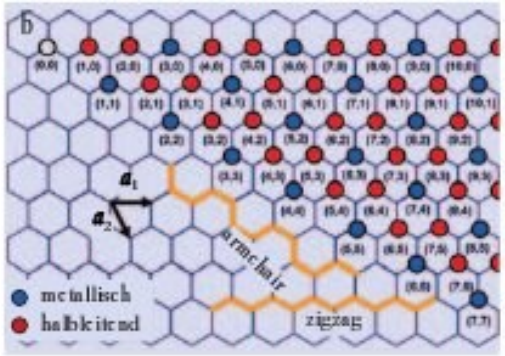
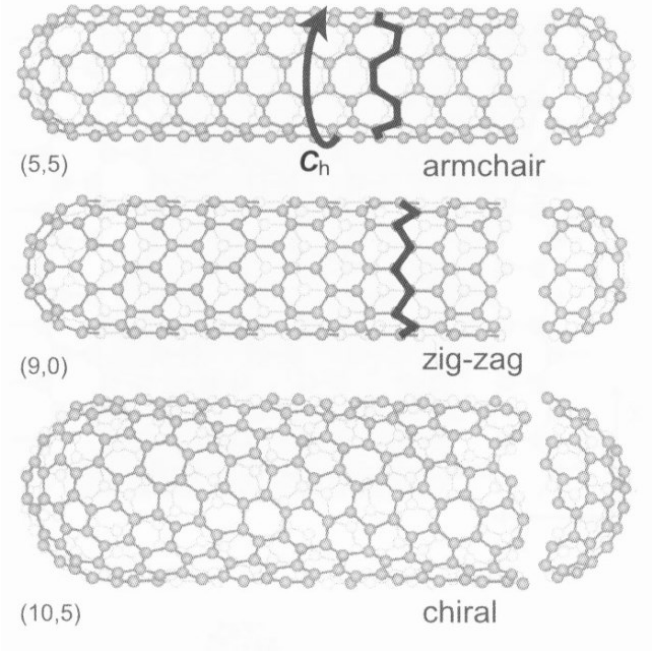
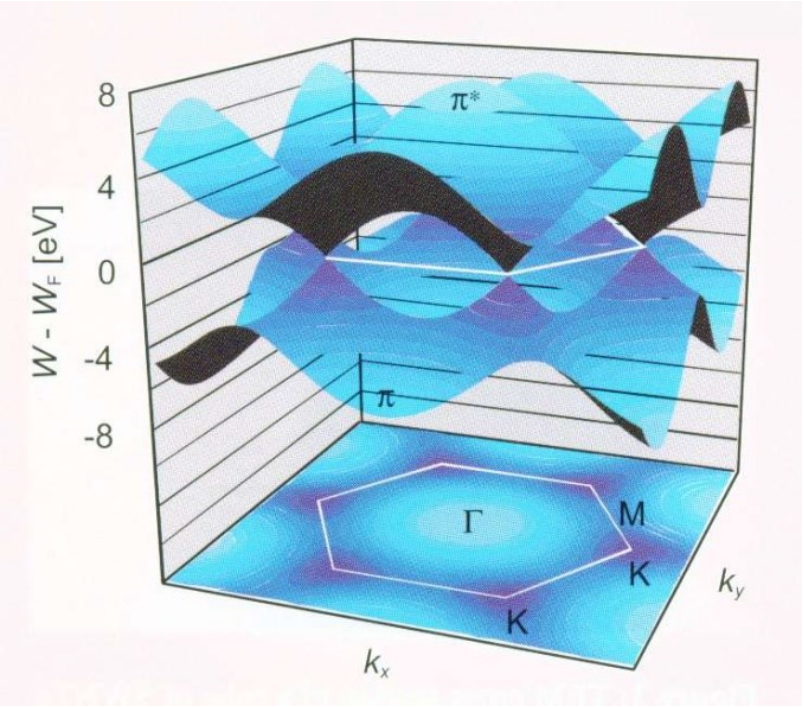
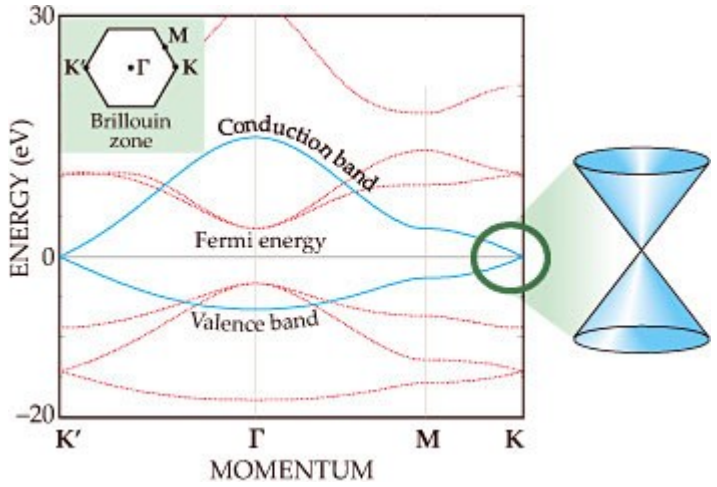


Abb. 2:
Der Aufrollvektor (n,m) der Nanoröhrchen ist über die Basisvektoren a_1 und a_2 des Grafitgitters definiert. Je nach Konfiguration können die Röhrchen metallisch oder halbleitend sein (b, vgl. Text).

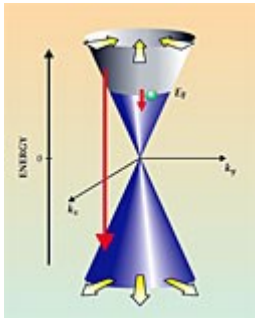


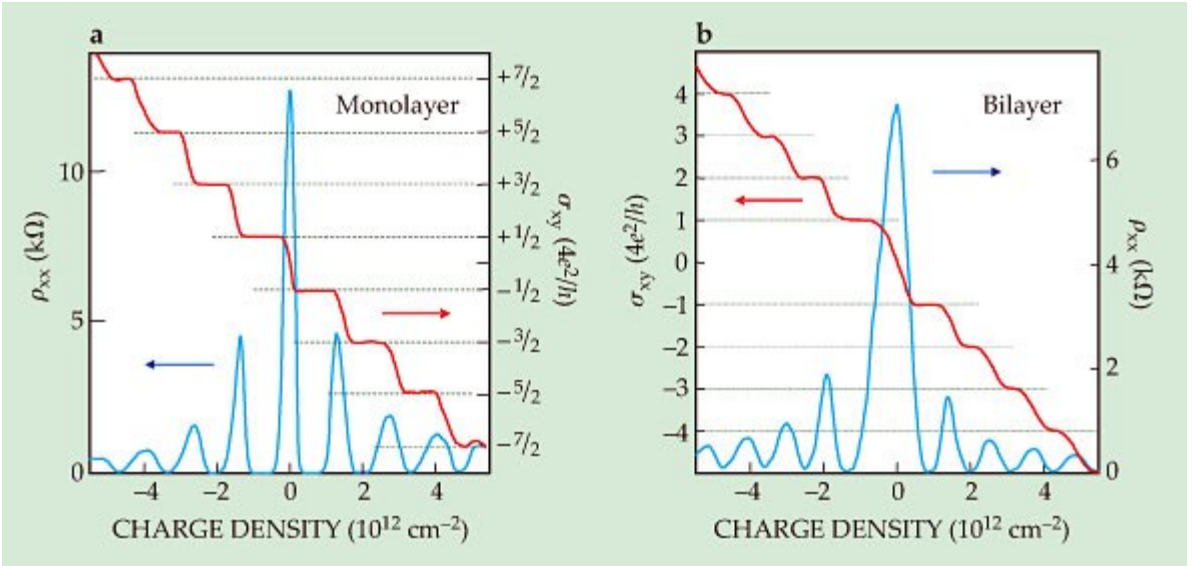
2D Dispersionsrelation von Graphen





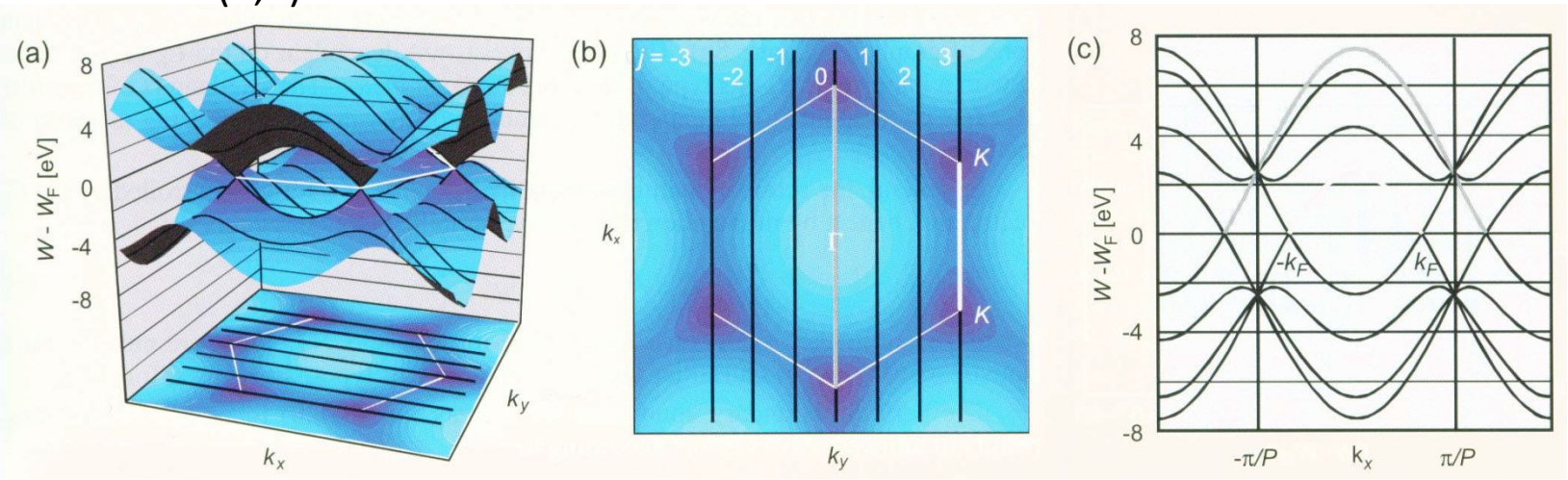
A. Geim, A.H. MacDonald
Phys. Today **60**, 35 (2007)



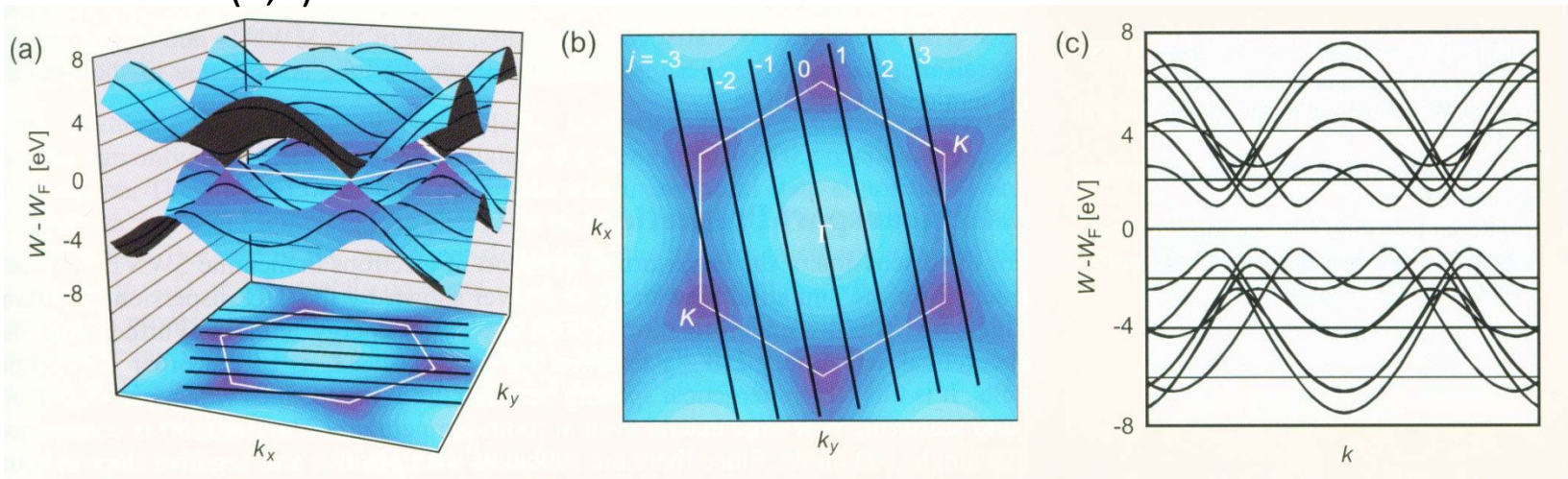


2D-Dispersionsrelation von Kohlenstoff-Nanoröhren

metallisch: (3,3) CNT

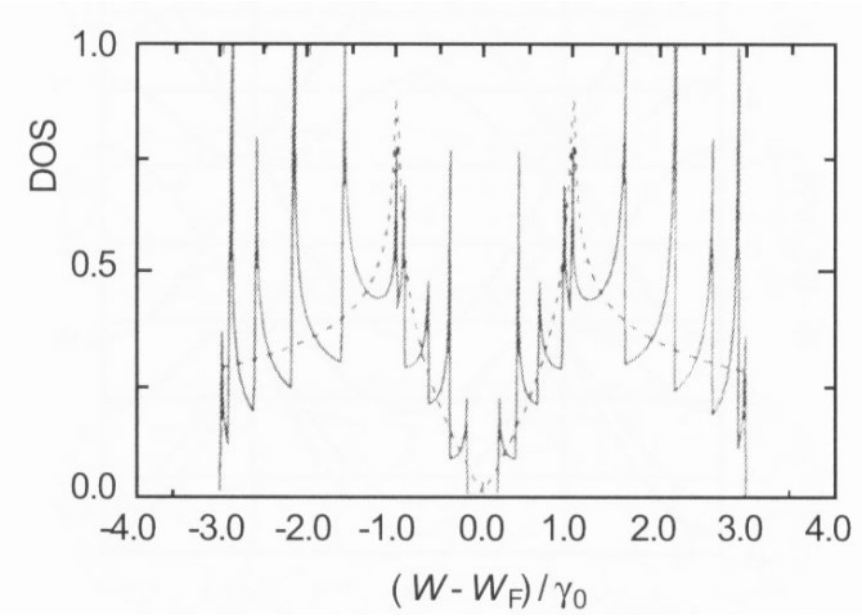


halbleitend: (4,2) CNT

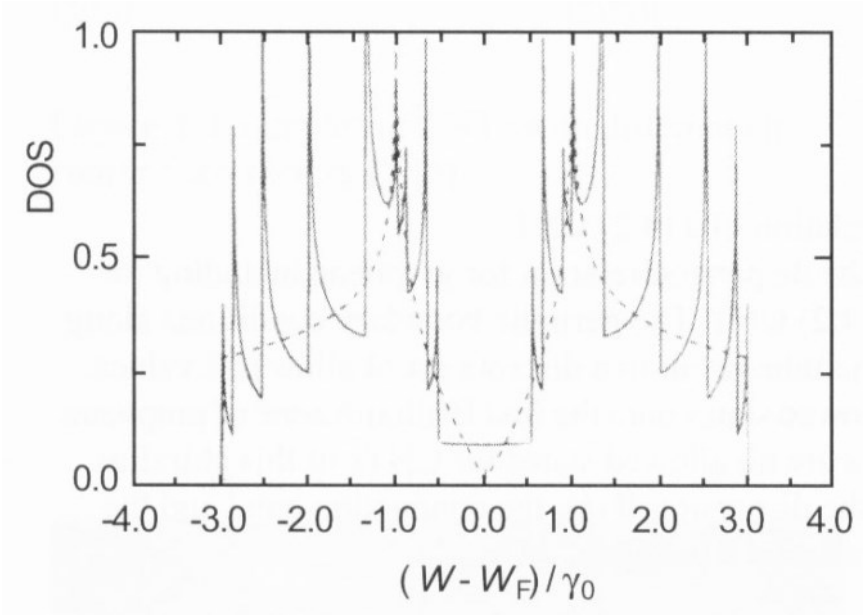


2D-Zustandsdichte von Kohlenstoff-Nanoröhren

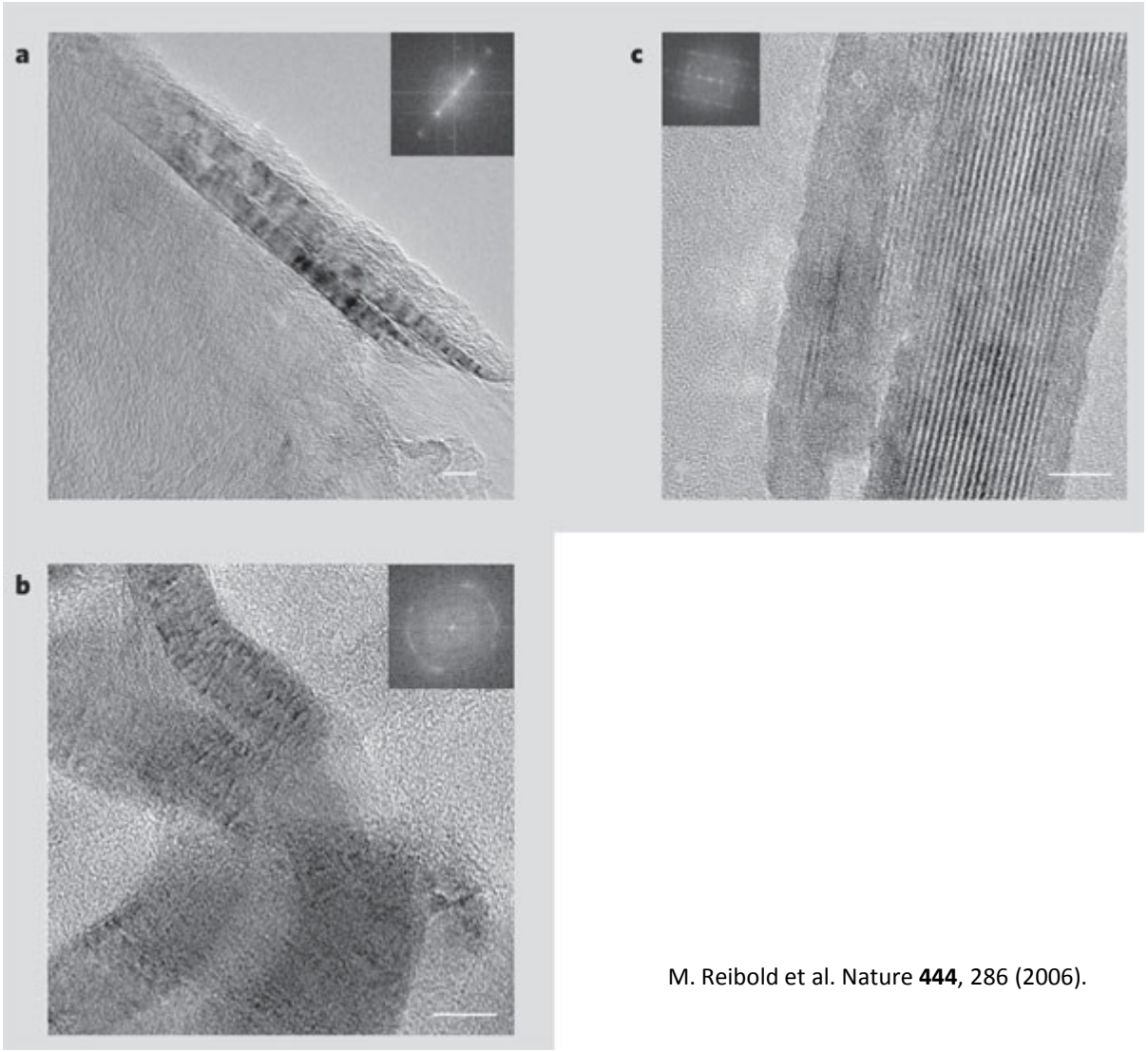
halbleitend



metallisch

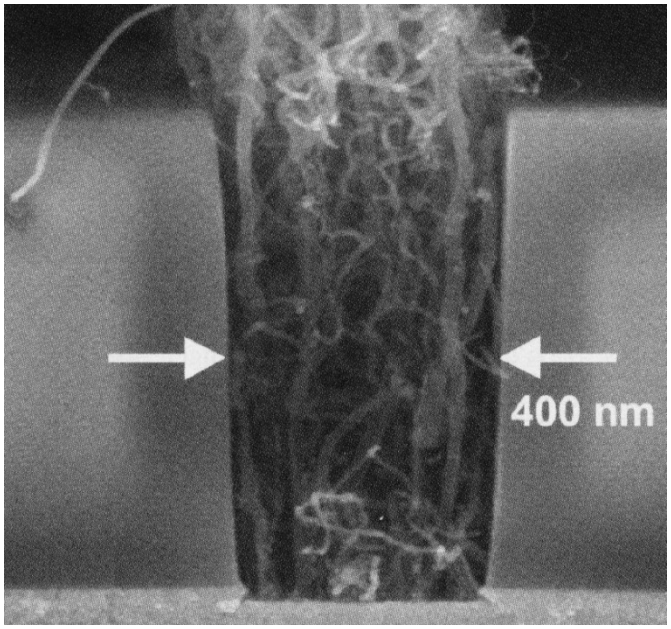


Damaszenerstahl mit Kohlenstoff-Nanoröhren und Cementite Fe_3C-Nanodrähten

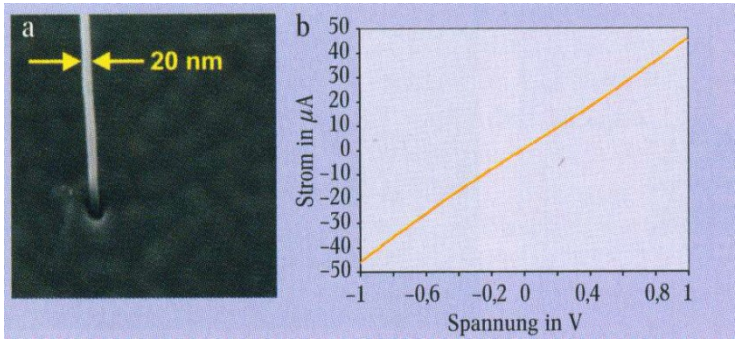


M. Reibold et al. Nature **444**, 286 (2006).

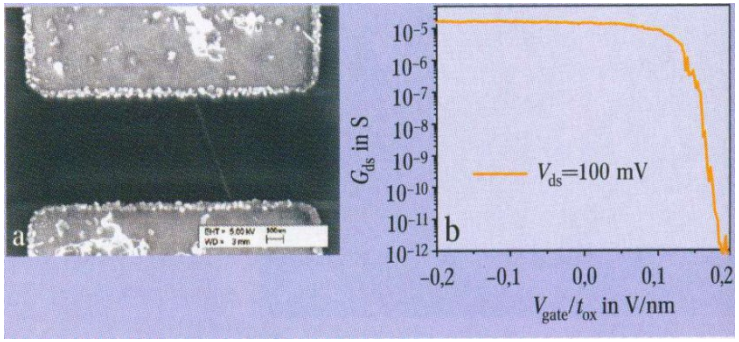
Vertikale Verbindung zwischen zwei Leiterbahnen
aus Bündeln von mehrwandigen CNT



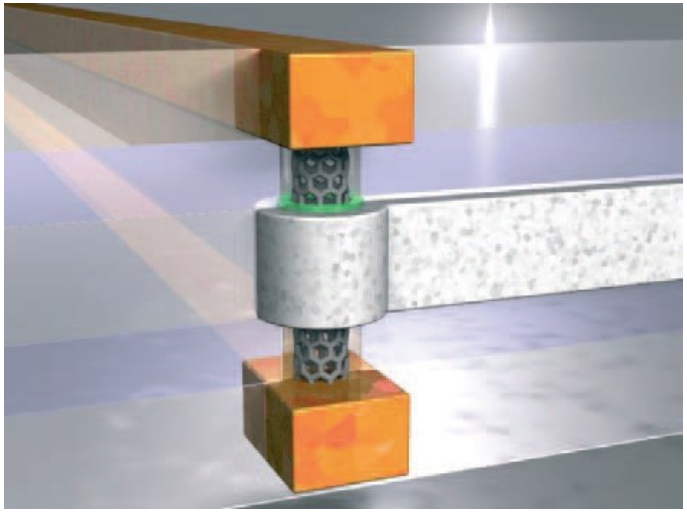
aus einzelnen mehrwandigen CNT

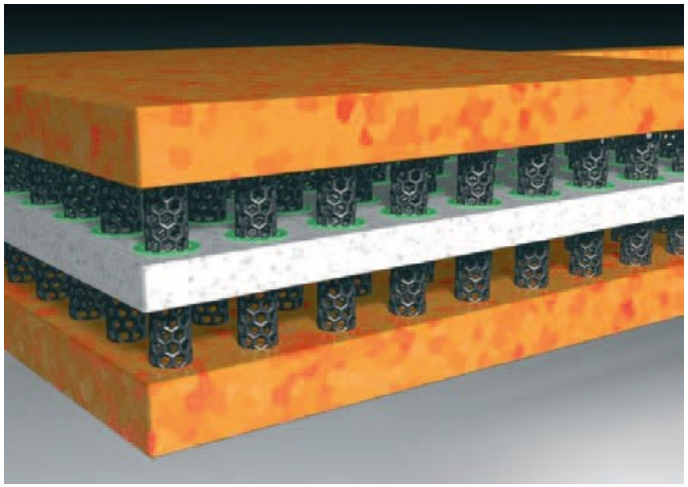


Planarer Feldeffekttransistor



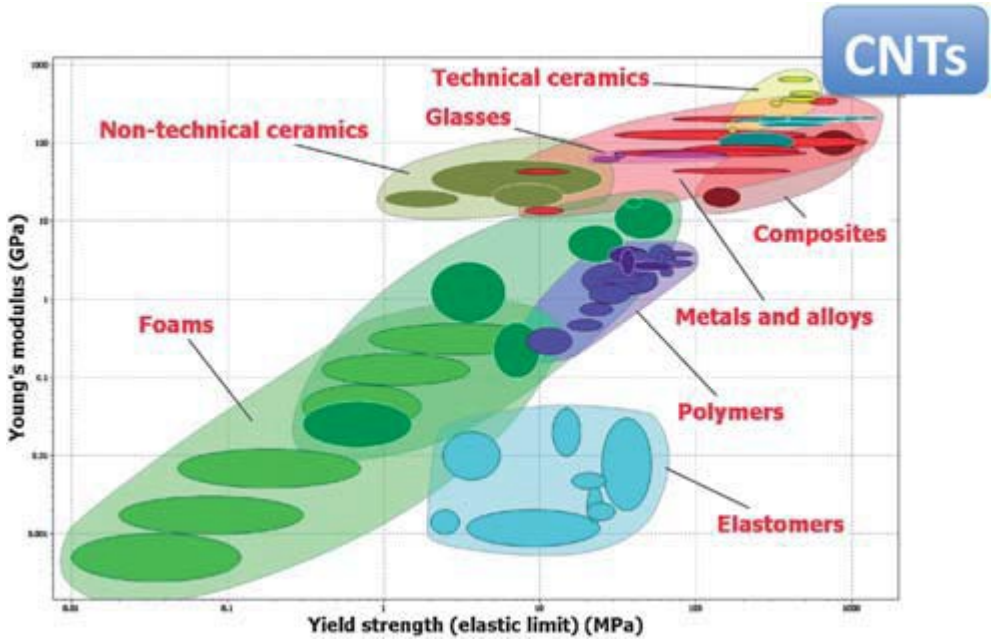
Konzept eines vertikalen Feldeffekttransistor





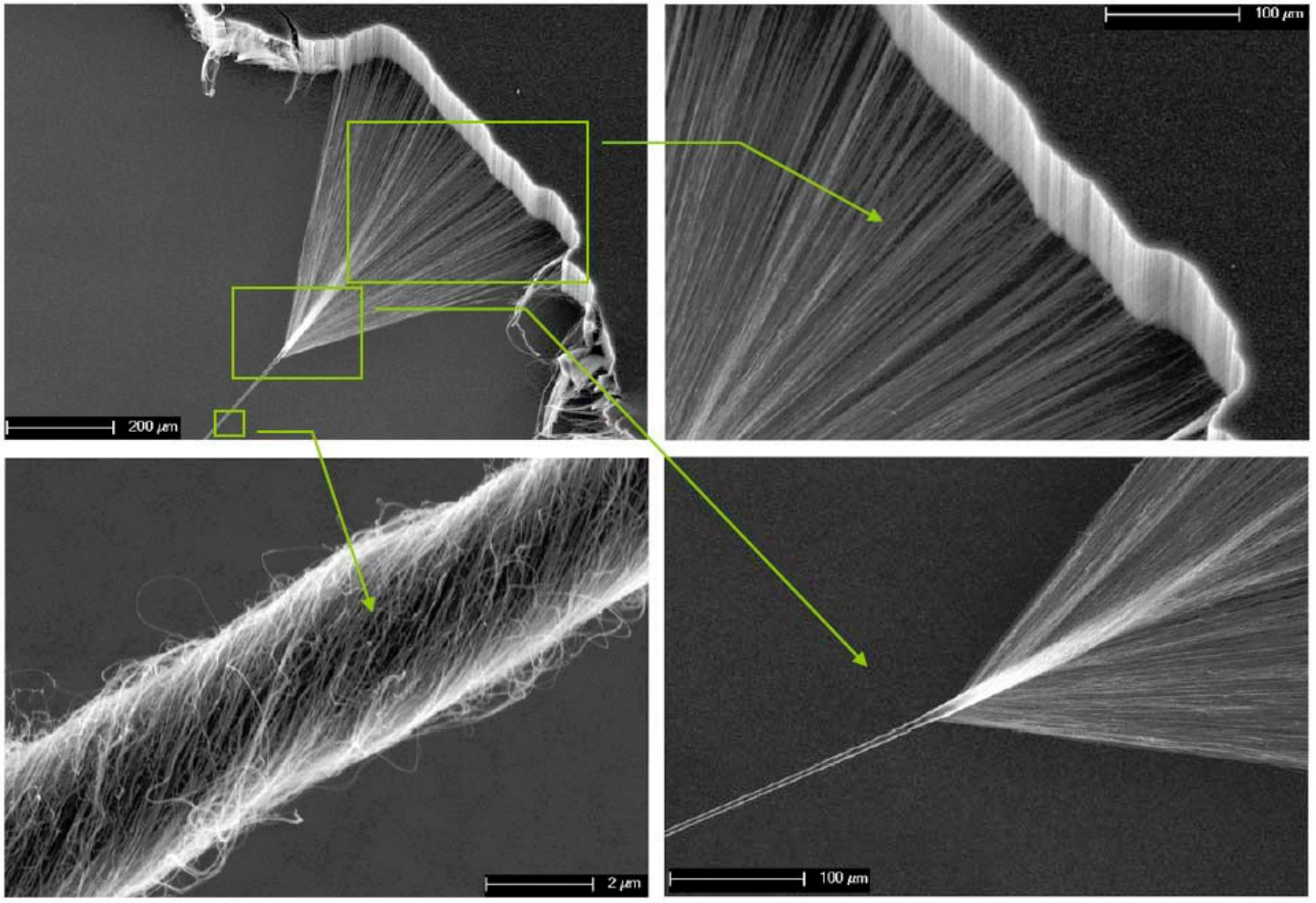
Mechanische Eigenschaften von Kohlenstoff-Nanoröhrchen

Material	Young's Modulus (GPa)	Tensile Strength (GPa)	Electrical Resistivity (Ωm)	Thermal conductivity (W/mK)	Density (g/cm3)
SWNT	1054	150	10-6	1750 - 5800	1.3
SWNT bundle	1054	75	$0.1\text{-}10^{-3}$	35	
MWNT	1200	150	10^{-6}	>3000	2.6
Diamond	1000	1.2	$1.0\text{x}10^{12}$ - $1.0\text{x}10^{14}$	1000 - 2600	
Steel	200	0.4	$12.0\text{x}10^{-8}$ - $170\text{x}10^{-8}$	43	7.8
Copper	110	0.413	$1.70\text{x}10^{-8}$ - $2.65\text{x}10^{-8}$	390	
Epoxy	3.5	0.005			1.25
Wood	16	0.008			0.6

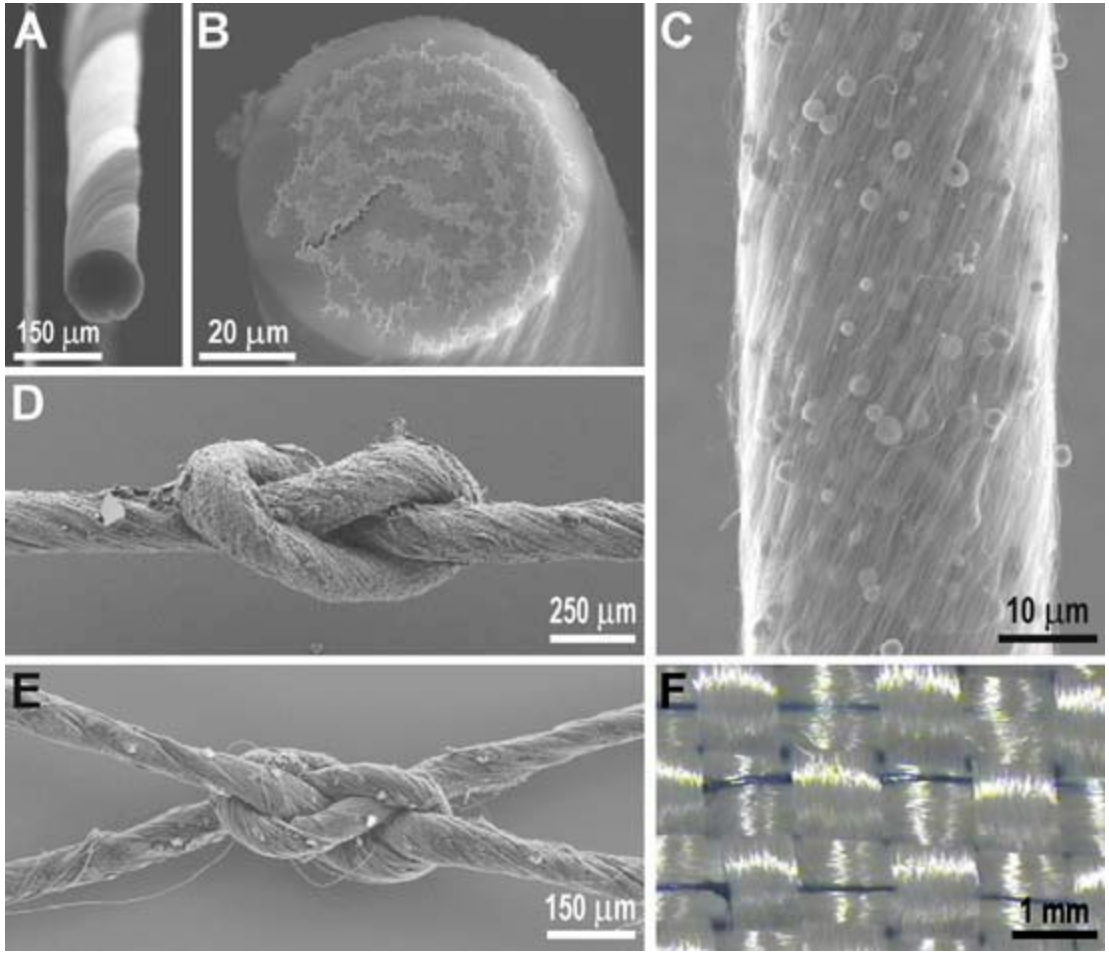


Herstellung von Garnen aus Kohlenstoff-Nanoröhrchen

Multifunctional Carbon Nanotube Yarns by Downsizing an Ancient Technology, M. Zhang, K. R. Atkinson, & R. H. Baughman, *Science*, **306**, 1358, 2004.

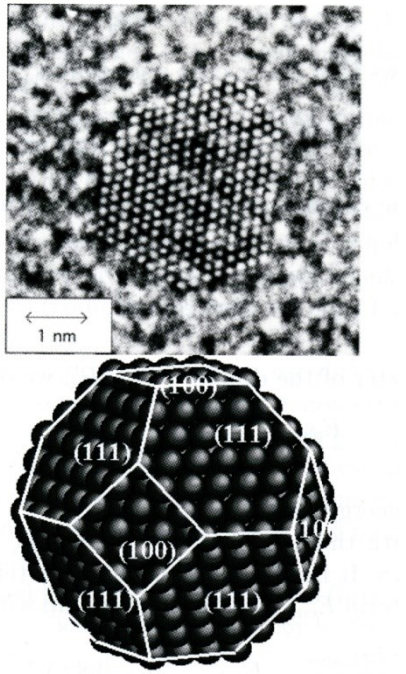


Garne und funktionale Textilien aus Kohlenstoff-Nanoröhrchen

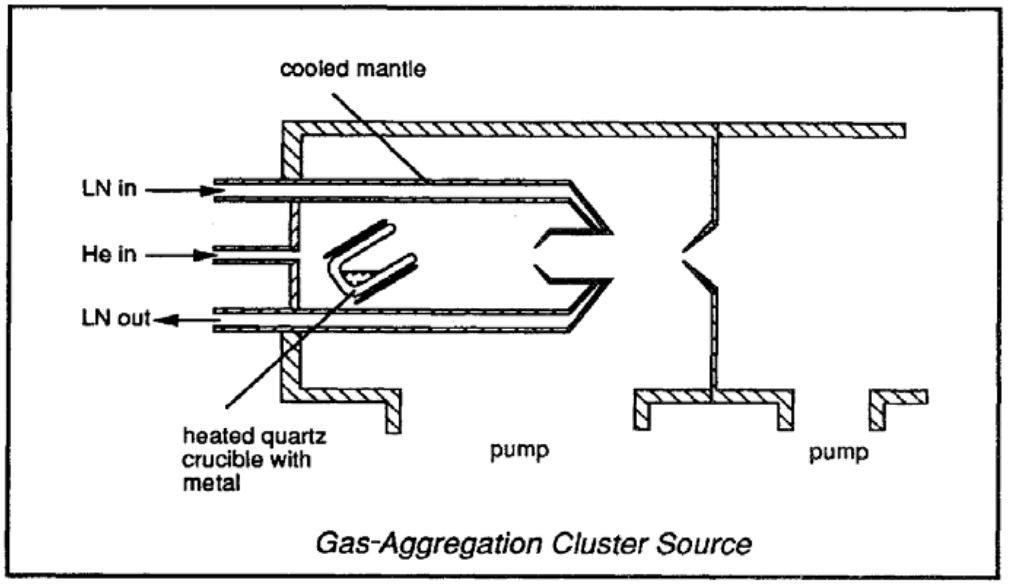


M.D. Lima et al., Science 331 (2011) 51

Herstellung von Clustern aus der Gasphase



SEM-Bild eines cubo-oktaedrischen Co-Clusters



Größenverteilung von Cluster

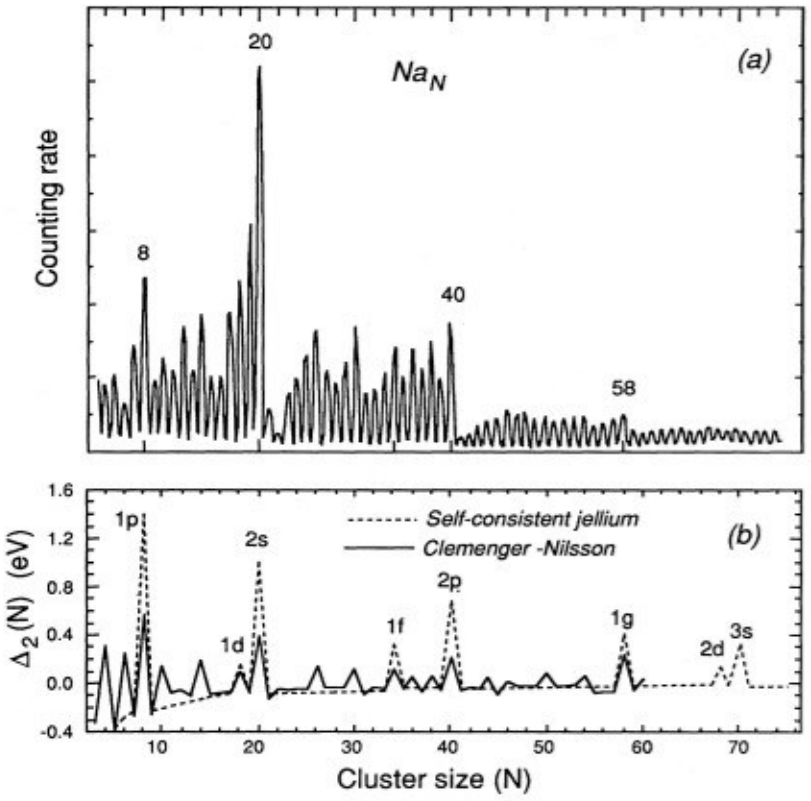
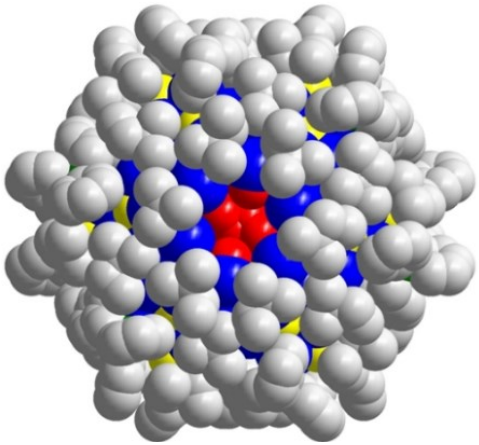


FIG. 1. Sodium cluster abundance spectrum: (a) experimental (after Knight *et al.*, 1984); (b) dashed line, using Woods-Saxon potential (after Knight *et al.*, 1984); solid line, using the ellipsoidal shell (Clemenger-Nilsson) model (after de Heer, Knight, Chou, and Cohen, 1987).

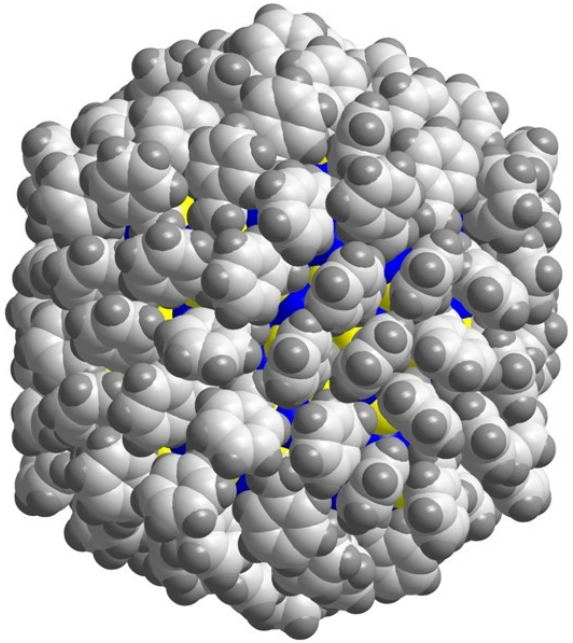
Heer et al. RMP 65, 611 (1993).

Herstellung von Clustern durch Reduktion von Metallionen

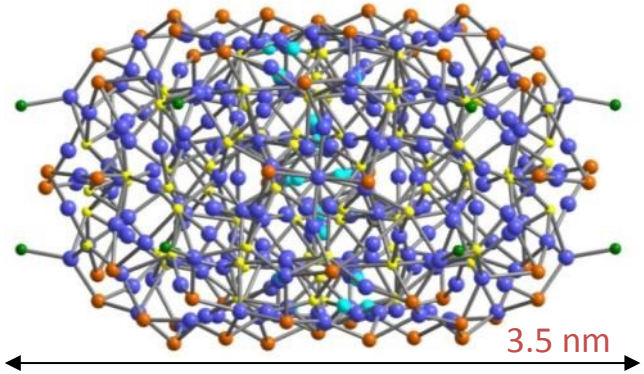
Space filling model of the molecular structure of $[Ag_{168}S_{66}(StC_5H_{11})_{36}(PP)_6]$ (PP: bidentate ligand).



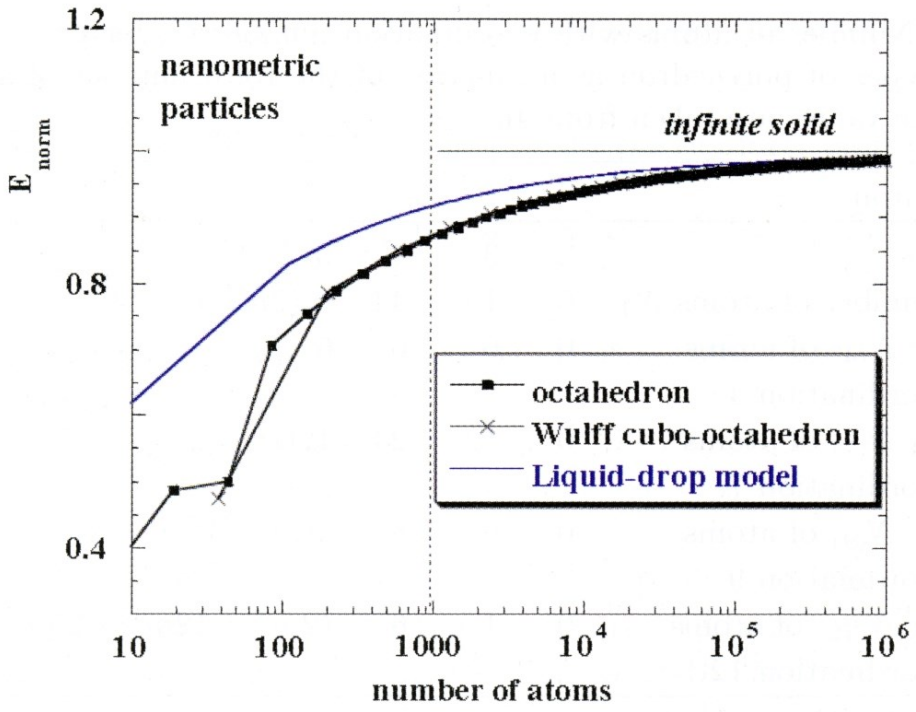
Space filling model of the molecular structure of $[Ag_{274}S_{80}(SCH_2Ph)_{114}]$, (Ag: blue; S: yellow; P: green; C: light grey; H: grey).



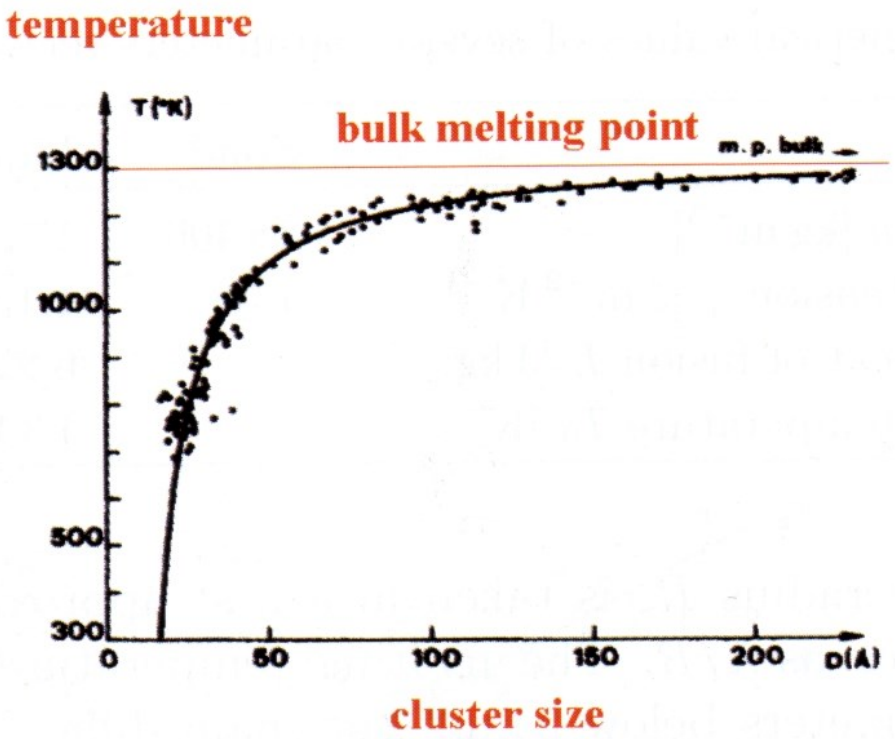
$[Ag_{262}S_{100}(tBuS)_{62}(Ph_2P(CH_2)_4PPh_2)_6]$



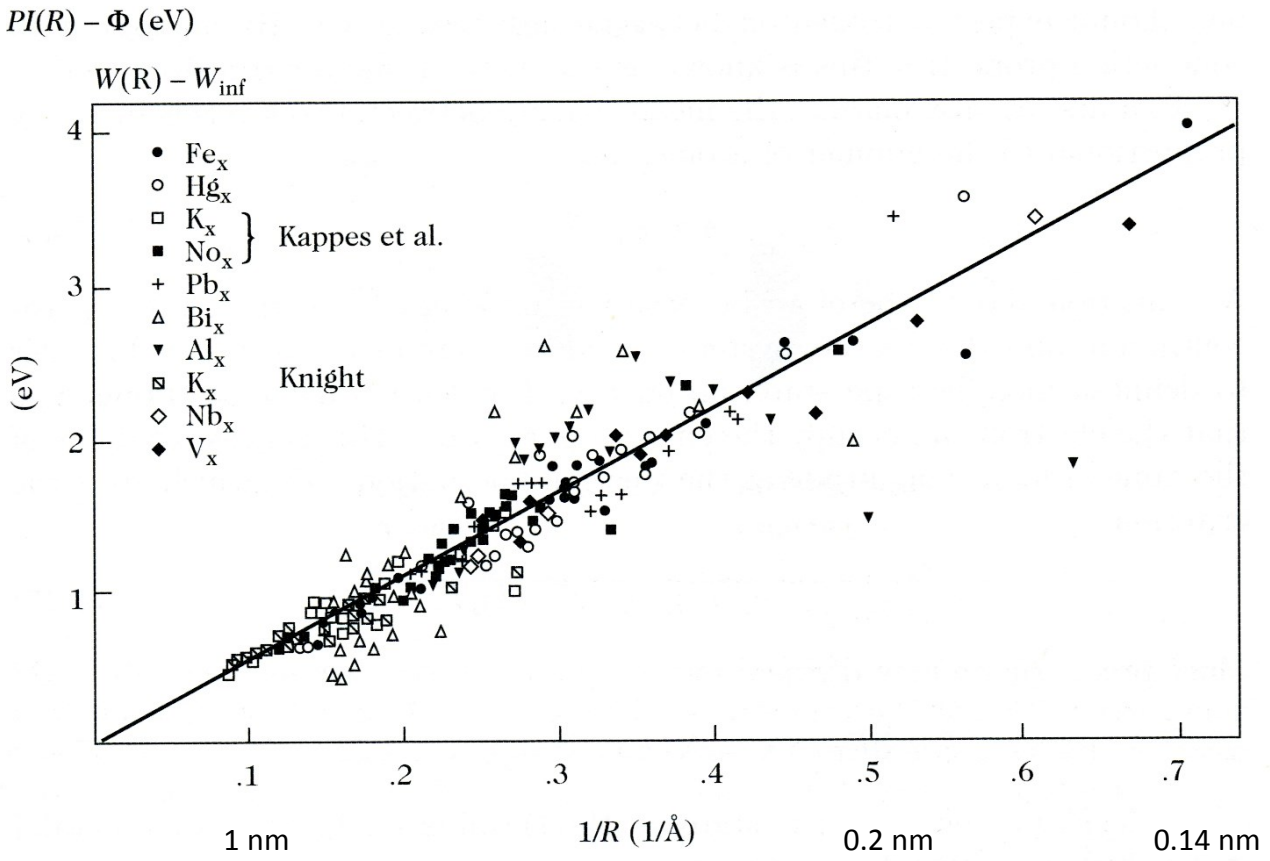
Größenabhängigkeit der Bindungsenergie



Schmelztemperatur für Au-Cluster



Größenabhängigkeit der Ionisationspotentials



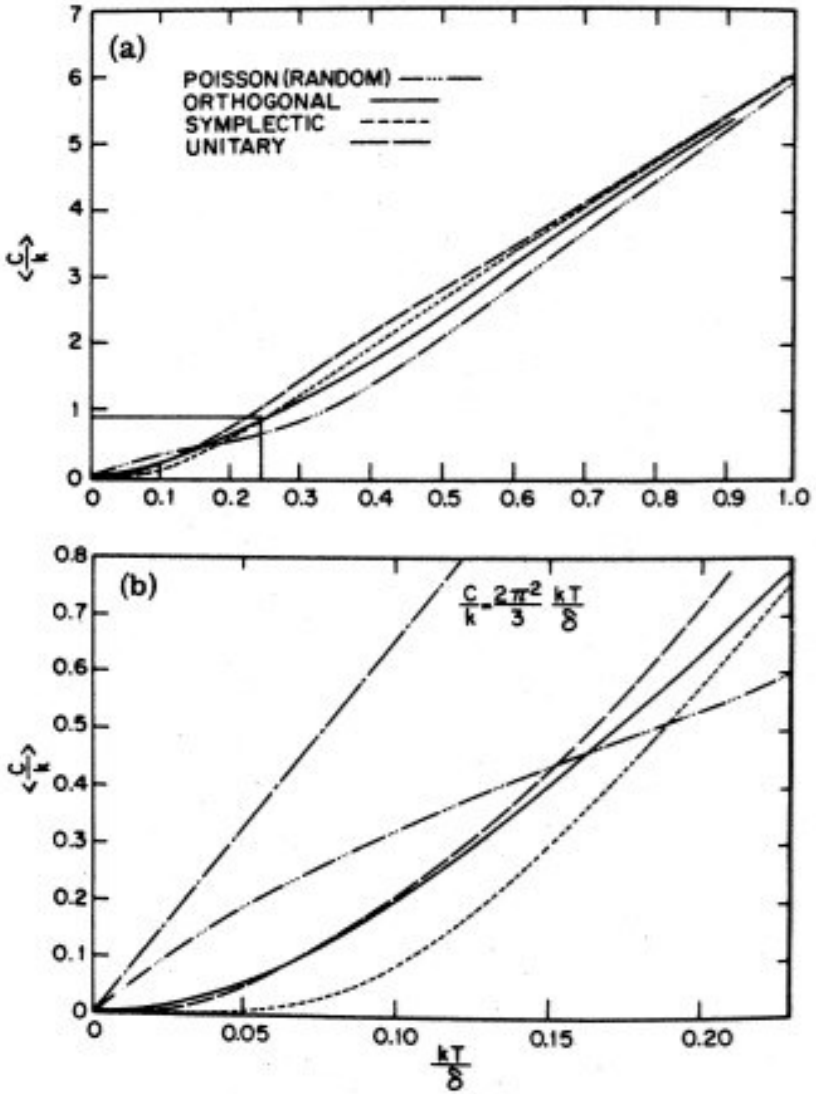


FIG. 3. The small-particle electronic heat capacity as a function of temperature, taken from Denton, Mühlschlegel, and Scalapino (1973): (a) An average of the even and odd particle calculations; (b) the linear heat capacity for the continuum case, shown for comparison (dashed-dotted line). Note the displacement of the high-temperature asymptote by $-k_B/2$ consistent with Eq. (2.29).

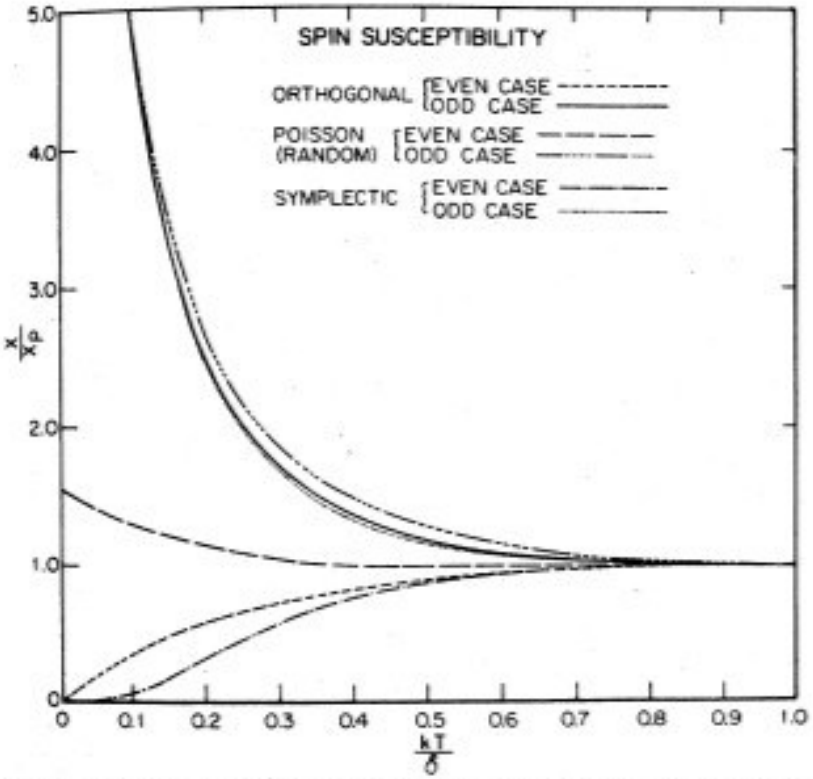
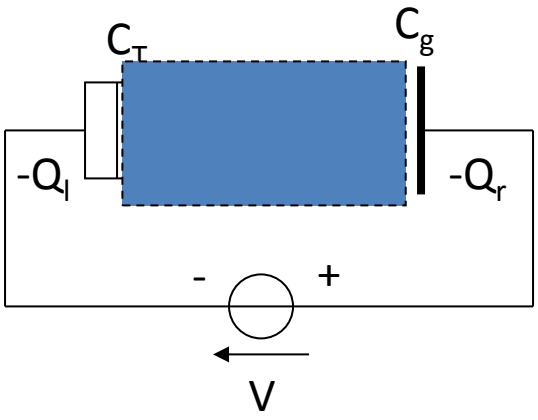


FIG. 4. The small-particle magnetic susceptibility calculated by Denton, Mühlschlegel, and Scalapino (1973). The calculated spin susceptibilities are normalized to the Pauli value taken here to be $\chi_P = 2\mu_B^2/\delta$.

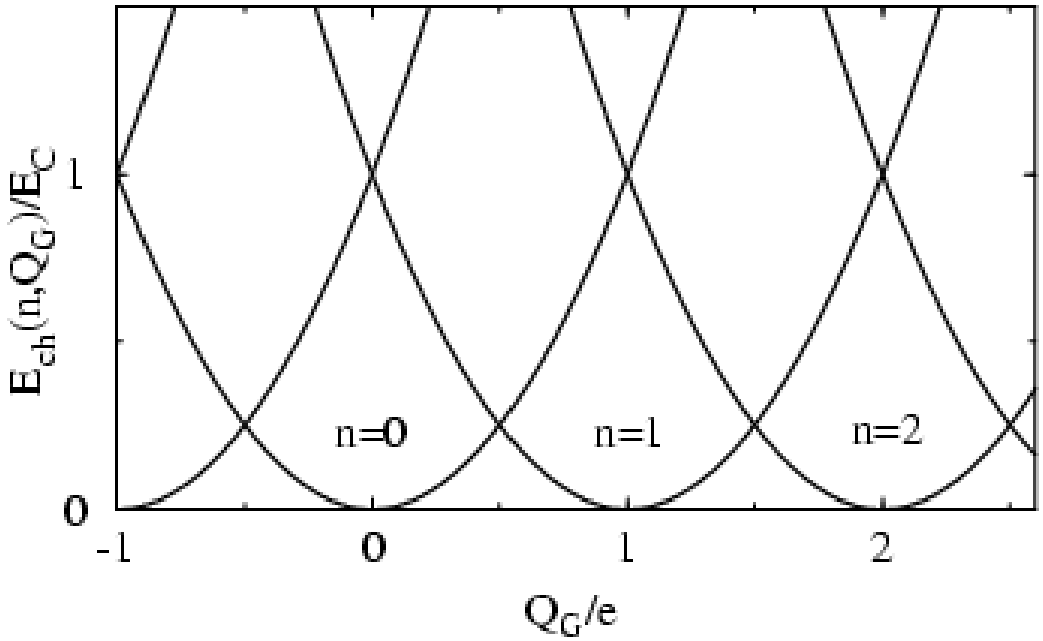
size effect in Metallclustern

W. P. Halperin, Rev. Mod. Phys. , Yol. 58, No. 3, July 1986

Schema für die Einzelelektronenbox (SEB)



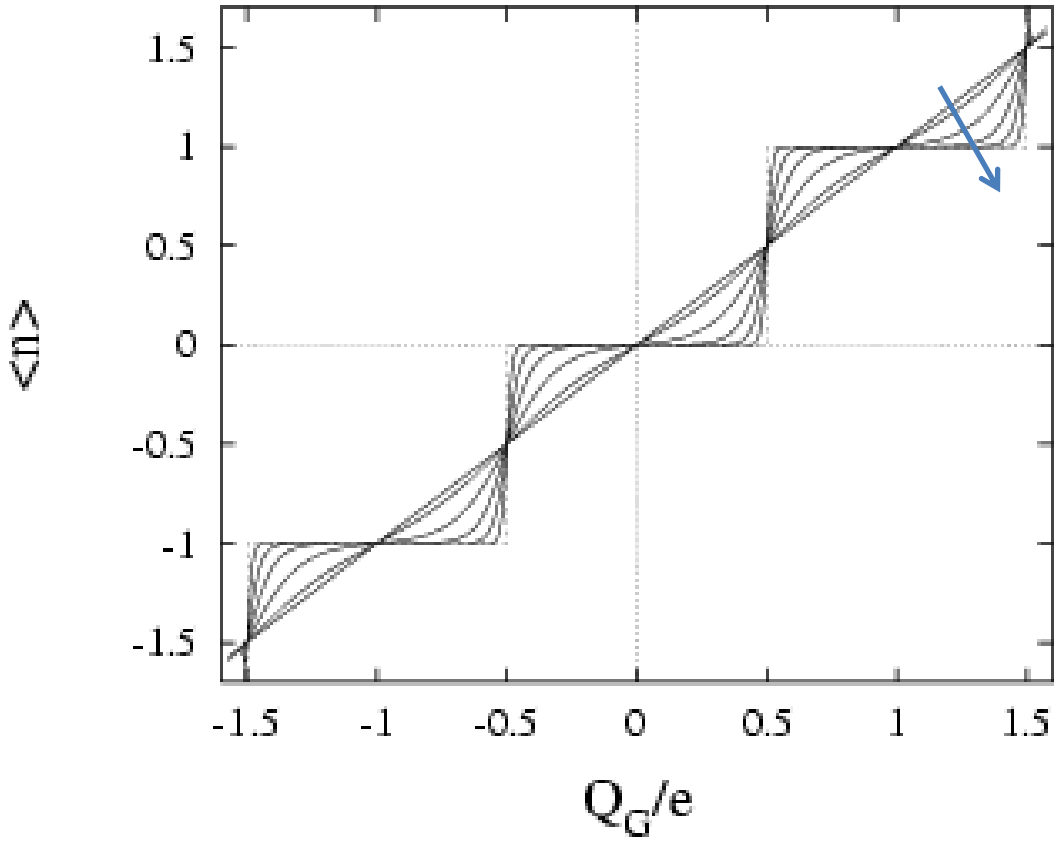
$$E(n,Q_g)=\frac{\left(ne-Q_g\right)^2}{2C_\Sigma}$$



Stabiler Zustand für n Elektronen in der Box: $e(n-1/2)<C_GU<e(n+1/2)$

Quelle: G. Schön: Single-Electron Tunneling. In: Quantum transport and dissipation. Weinheim 1998. S. 149-212.

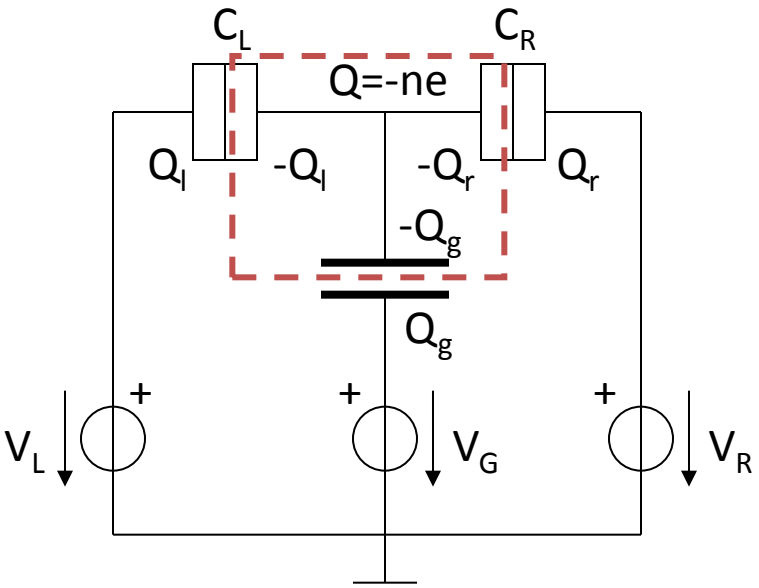
Coulomb-Treppe

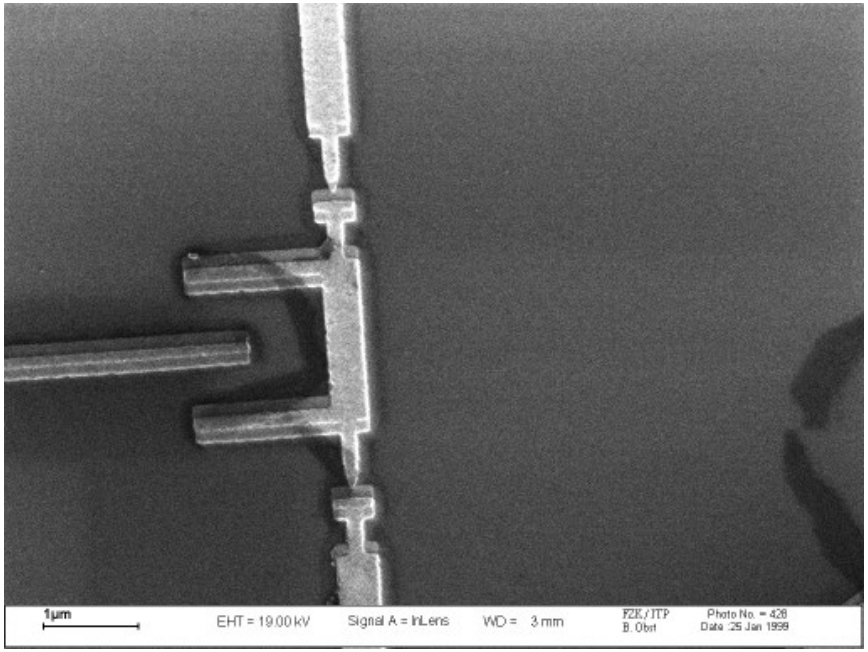


abnehmende Temp.

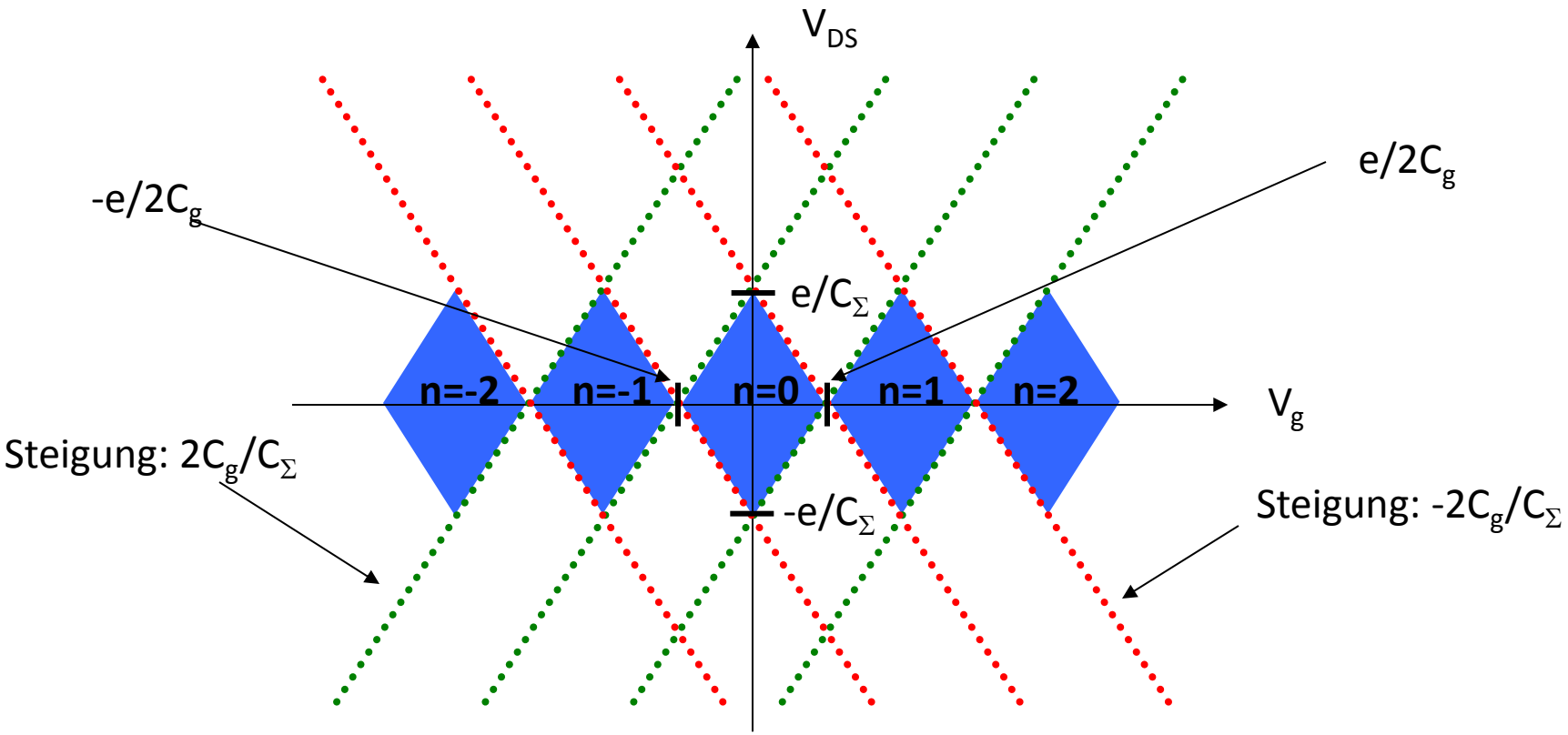
Quelle: G. Schön: Single-Electron Tunneling. In: Quantum transport and dissipation. Weinheim 1998. S. 149-212.

Schematischer Aufbau und realer SET



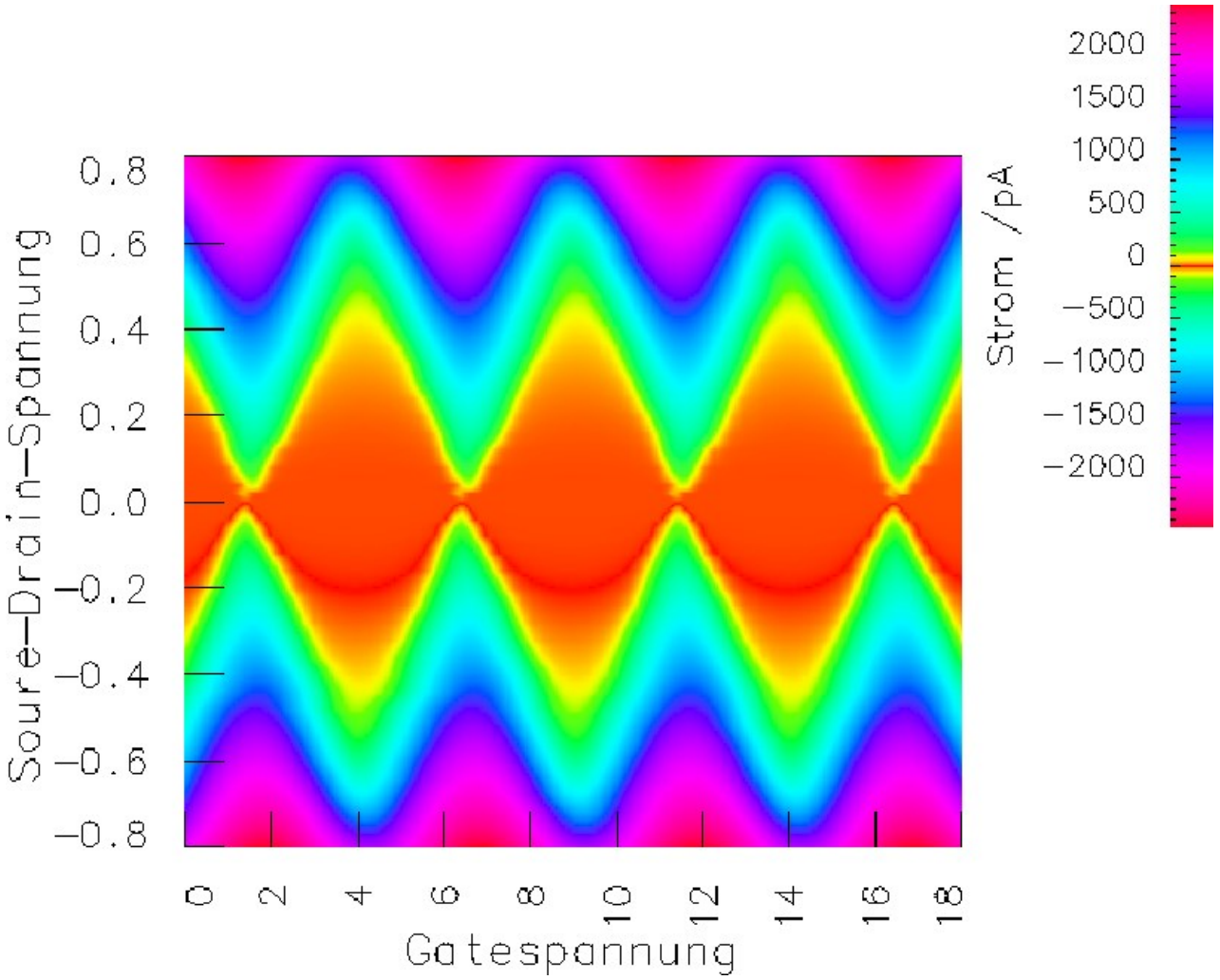


Schwellenspannung für einen symmetrischen SET ($C_L=C_R=C$) in Abhängigkeit von V_g:

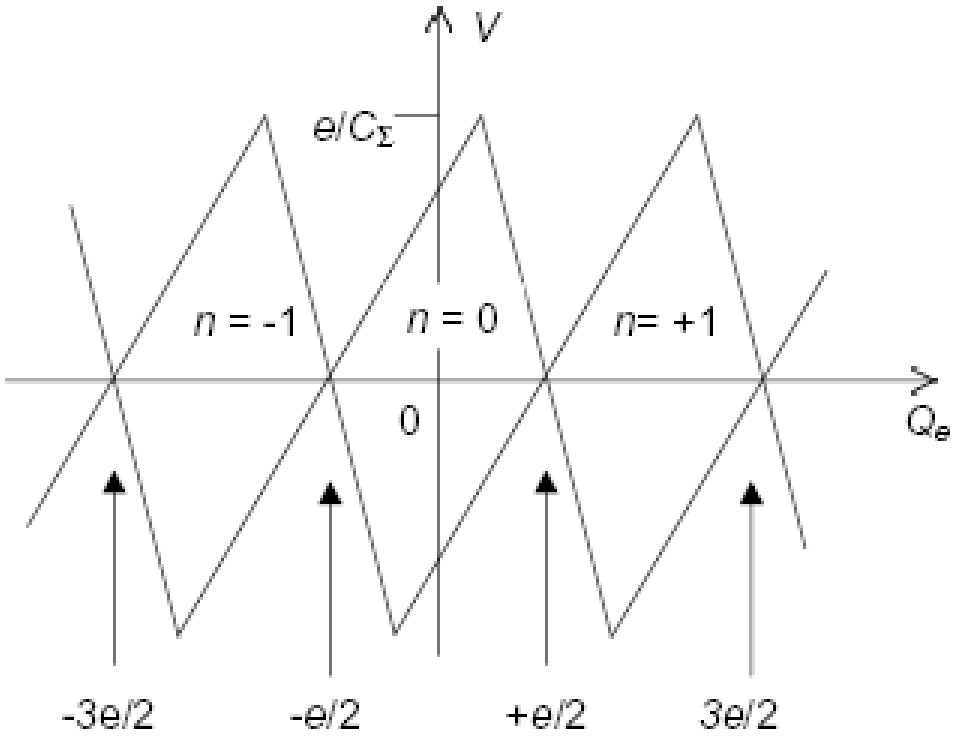


Innerhalb der Rauten („diamond shaped regions“) ist der Stromfluss durch die Coulomb-Blockade unterdrückt, die Ladung auf der Insel ist stabil.

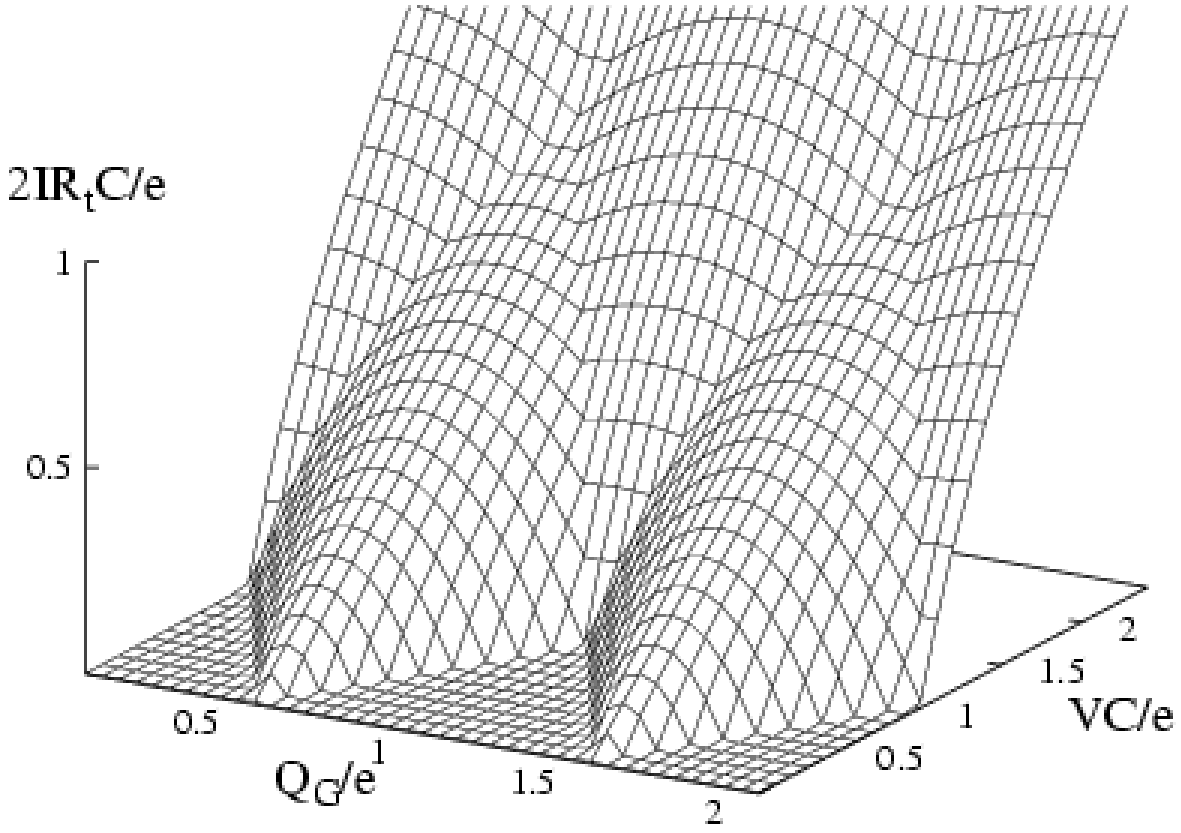
Colourplot der Coulomb-Blockade



Asymmetrischer SET

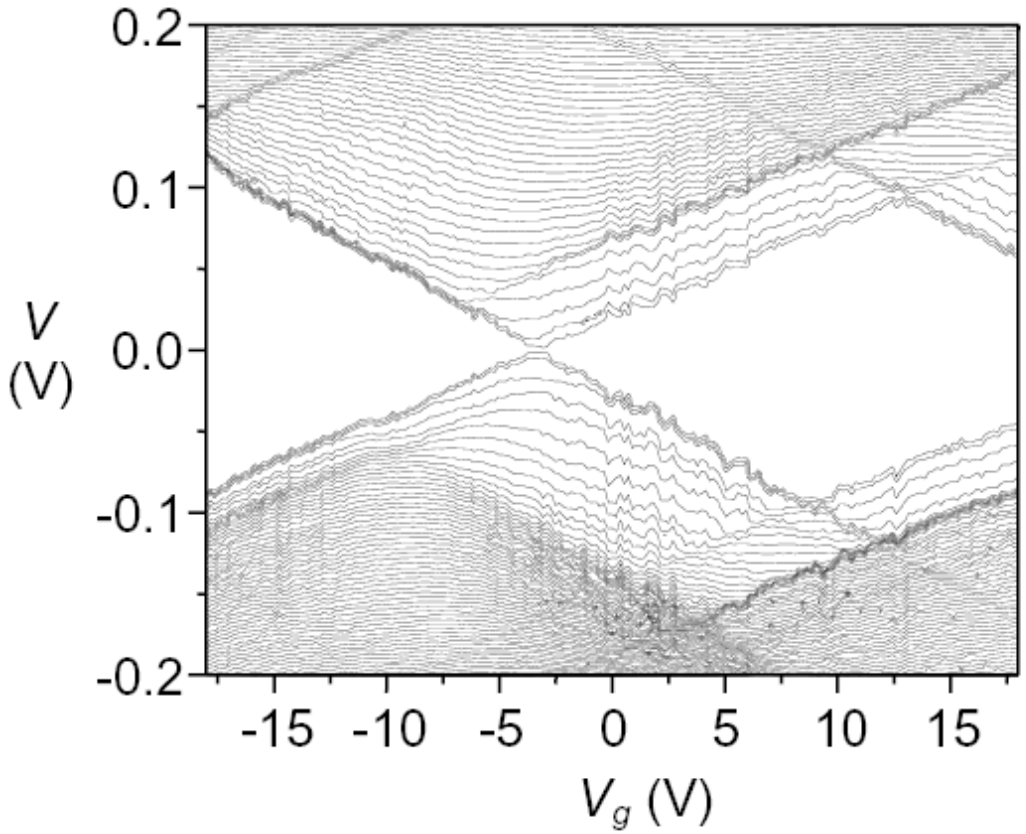


Quelle: K.K. Likharev: Single-Electron Devices and Their Applications.
In: Proceedings of the IEEE, Vol 87, No.4, April 1999.


2IR$_t$C/e
1
0.5
0.5
1
1.5
2
Q$_G$/e
0.5
1
1.5
2
VC/e

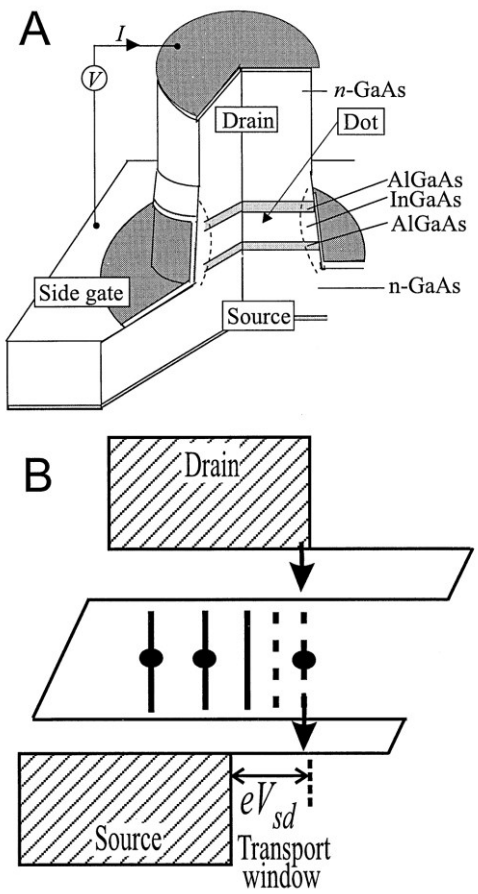
Beispiel für realen SET

Strom-Kontur-Plot für Aluminium-SET mit e/C_Σ=0.1V bei T=4.2K

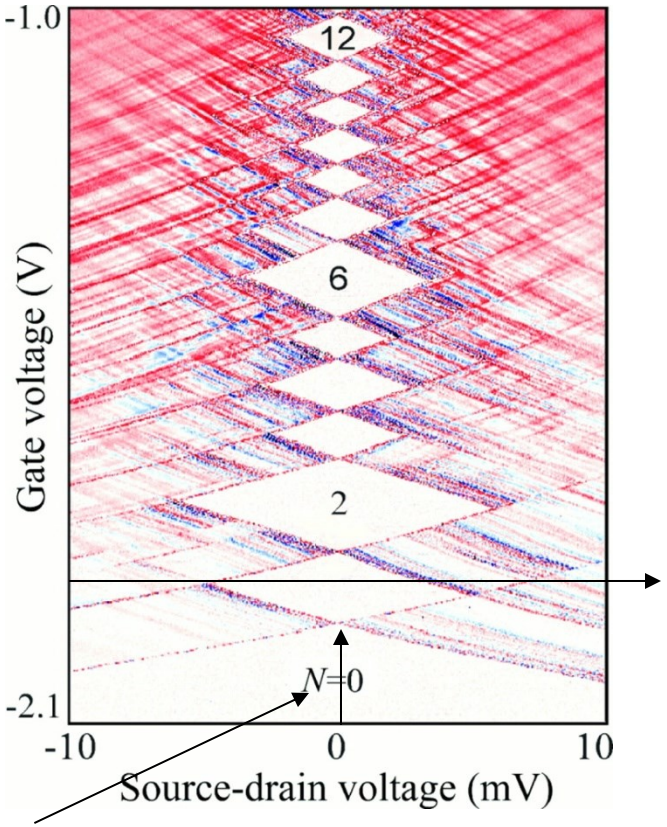


Quelle: K.K. Likharev: Sub 20-nm Electron Devices.
In: Advanced Semiconductor and Organic Nano-Techniques, Pt. 1. Academic Press, 2002.

Single electron transistor



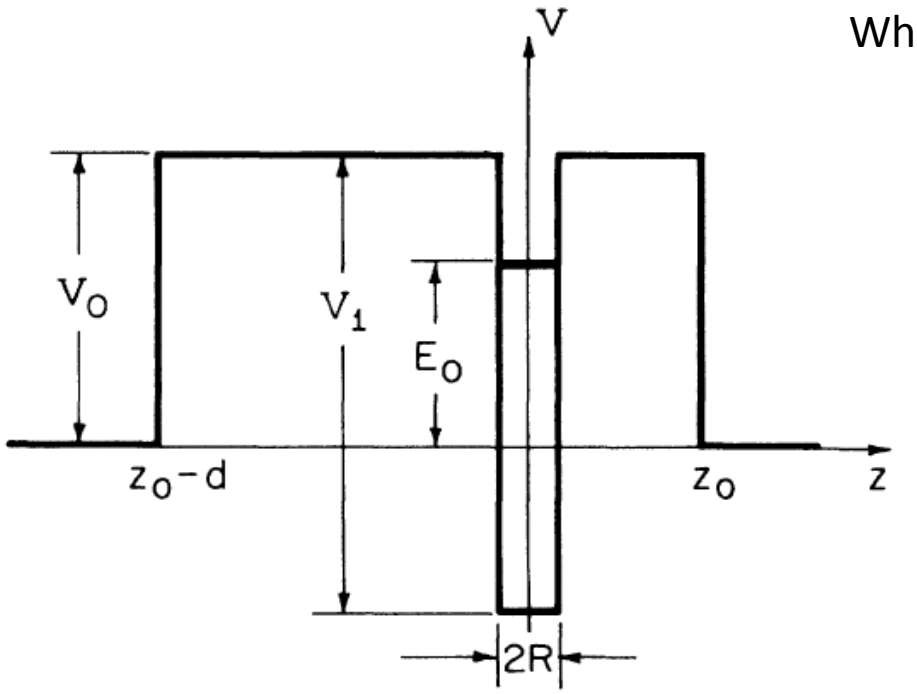
White, di/dv = 0



Coulomb Step = max di/dv

Gate bias for level at Vds = 0

Resonantes Tunneln



What is T(E)?

Josephson charge qubits

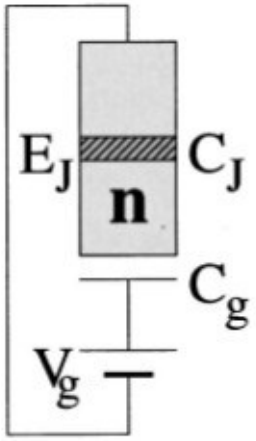


FIG. 1. A Josephson charge qubit in its simplest design formed by a superconducting single-charge box.

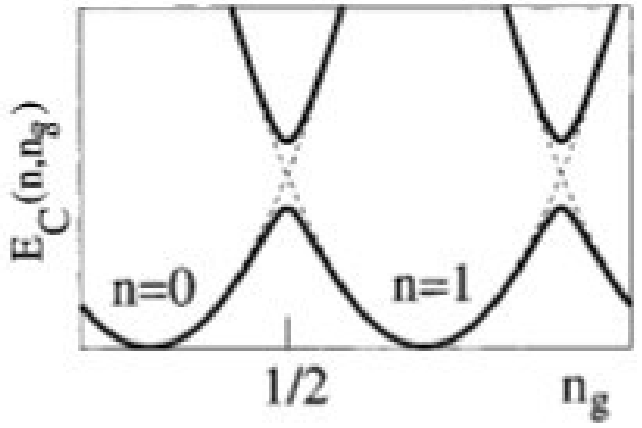


FIG. 2. The charging energy of a superconducting electron box is shown as a function of the gate charge n_g for different numbers of extra Cooper pairs n on the island (dashed parabolas). Near degeneracy points the weaker Josephson coupling mixes the charge states and modifies the energy of the eigenstates (solid lines). In the vicinity of these points the system effectively reduces to a two-state quantum system.

Quelle: Y. Makhlin et al., Rev.Mod.Phys. 73, 357 (2001).

Josephson charge qubits

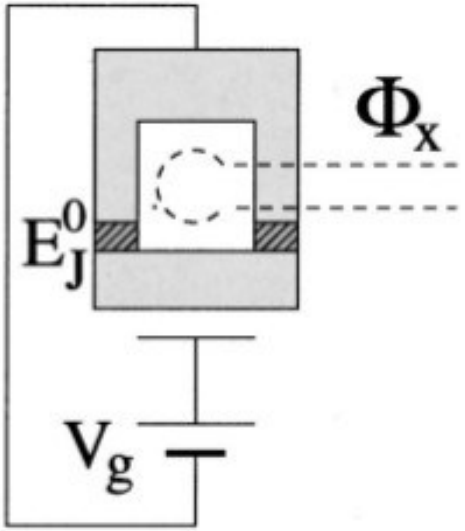


FIG. 3. A charge qubit with tunable effective Josephson coupling. The single Josephson junction is replaced by a flux-threaded SQUID. The flux in turn can be controlled by a current-carrying loop placed on top of the structure.

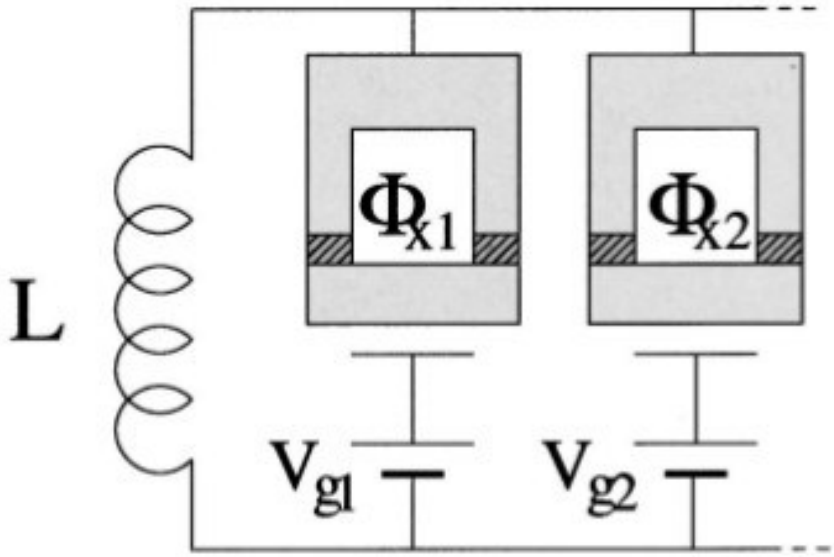
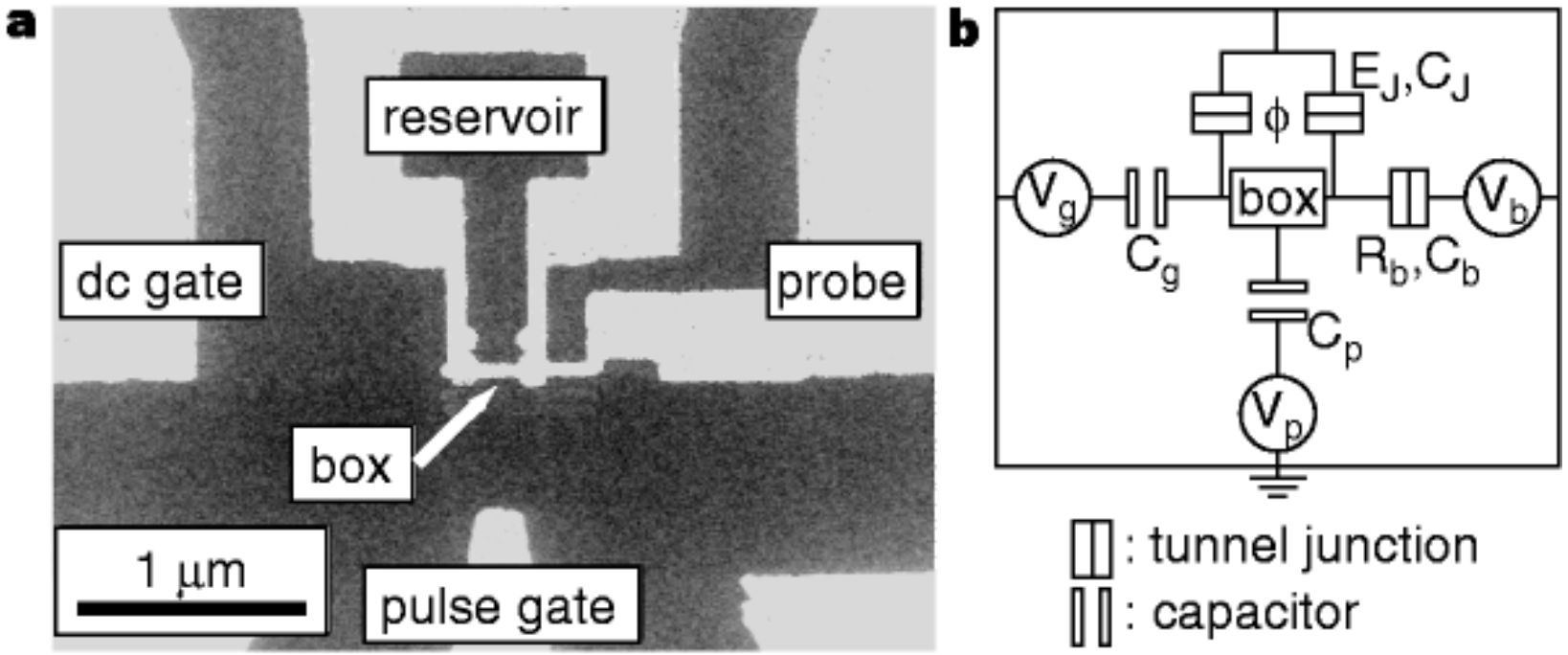


FIG. 4. A register of many charge qubits coupled by oscillator modes in the LC circuit formed by the inductor and the qubit capacitors.

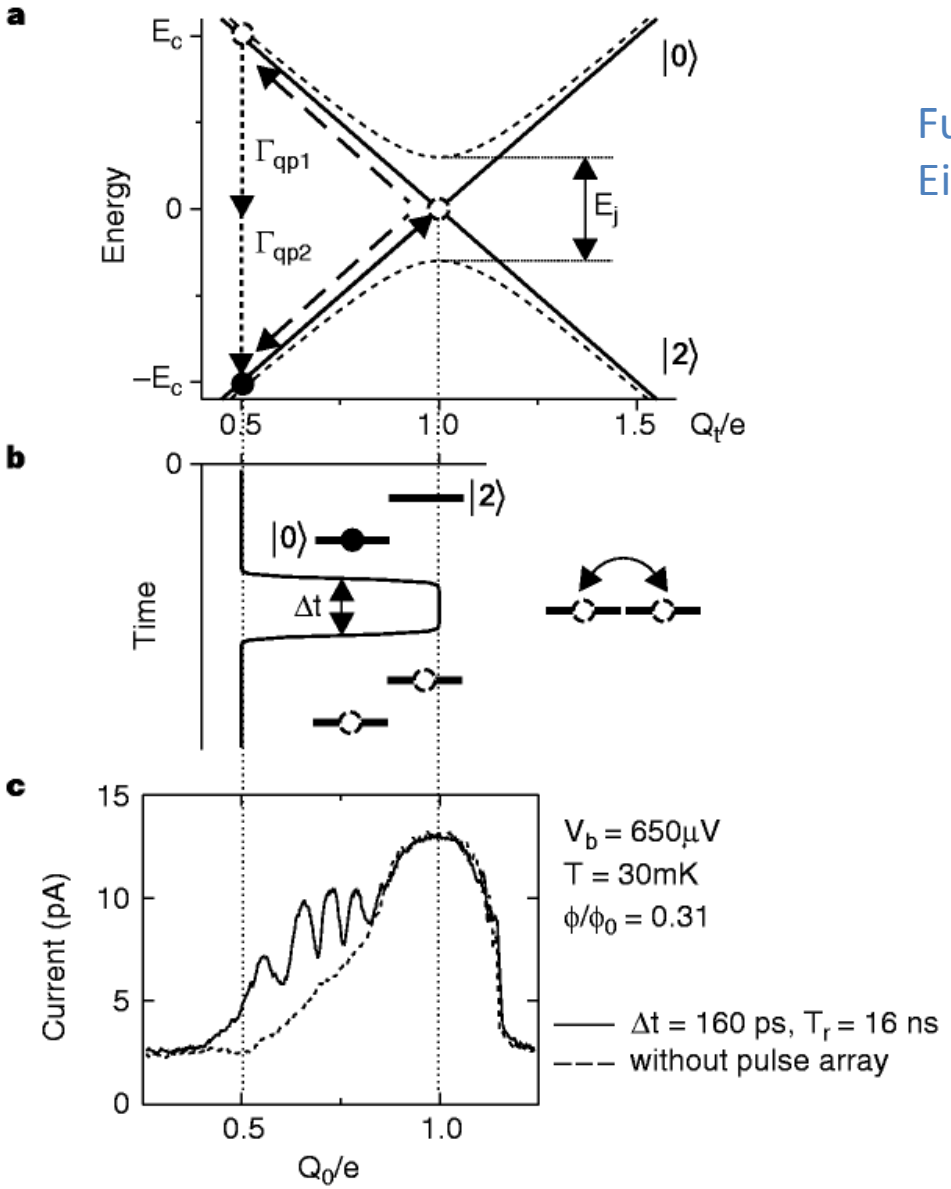
Quelle: Y. Makhlin et al., Rev.Mod.Phys. 73, 357 (2001).

Beispiel eines realen supraleitenden Qubits: Die Einzel-Cooper-Paar box



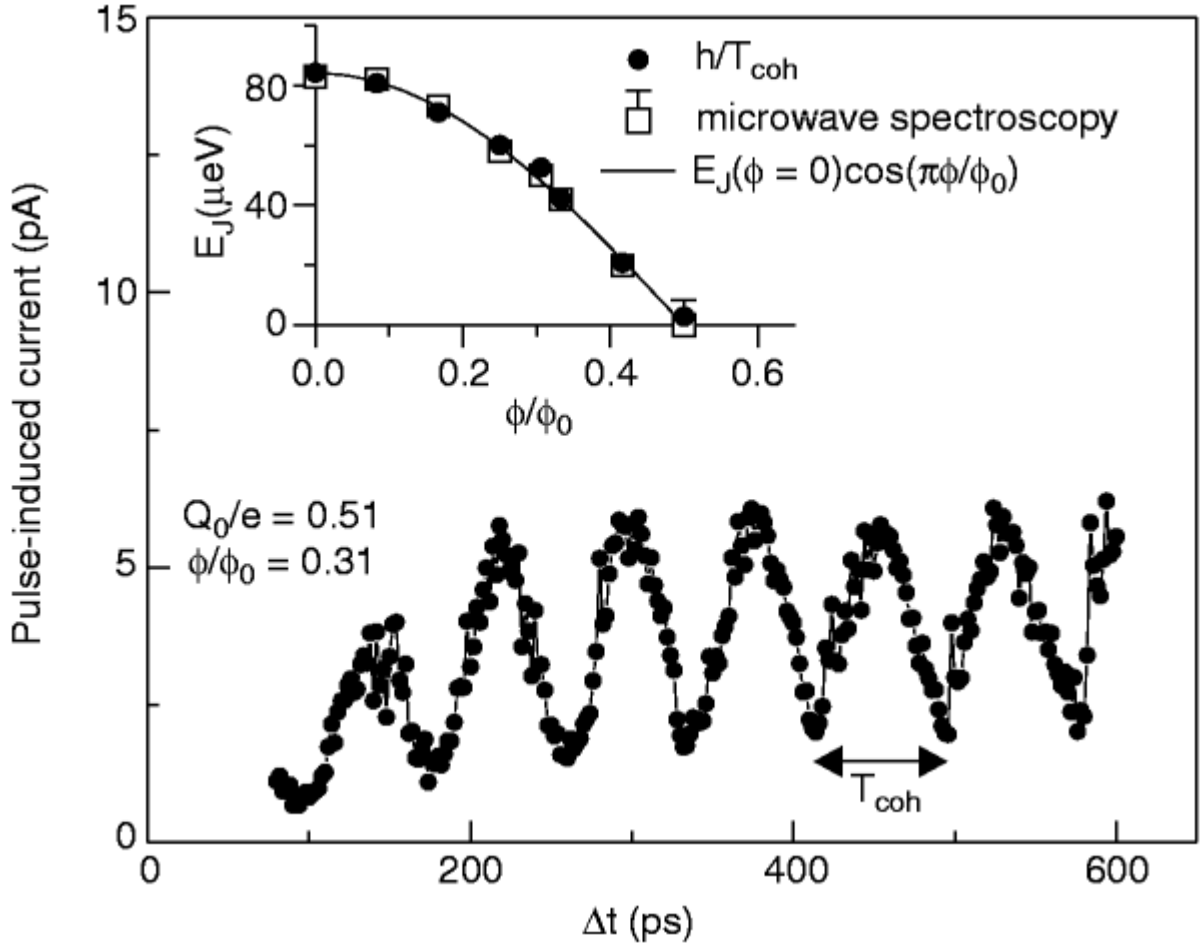
Quelle: Y. Nakamura, Yu. A. Pashkin and J. S. Tsai, Nature 398, 786-788(29 April 1999)

Funktionsweise der Einzel-Cooper-Paar box



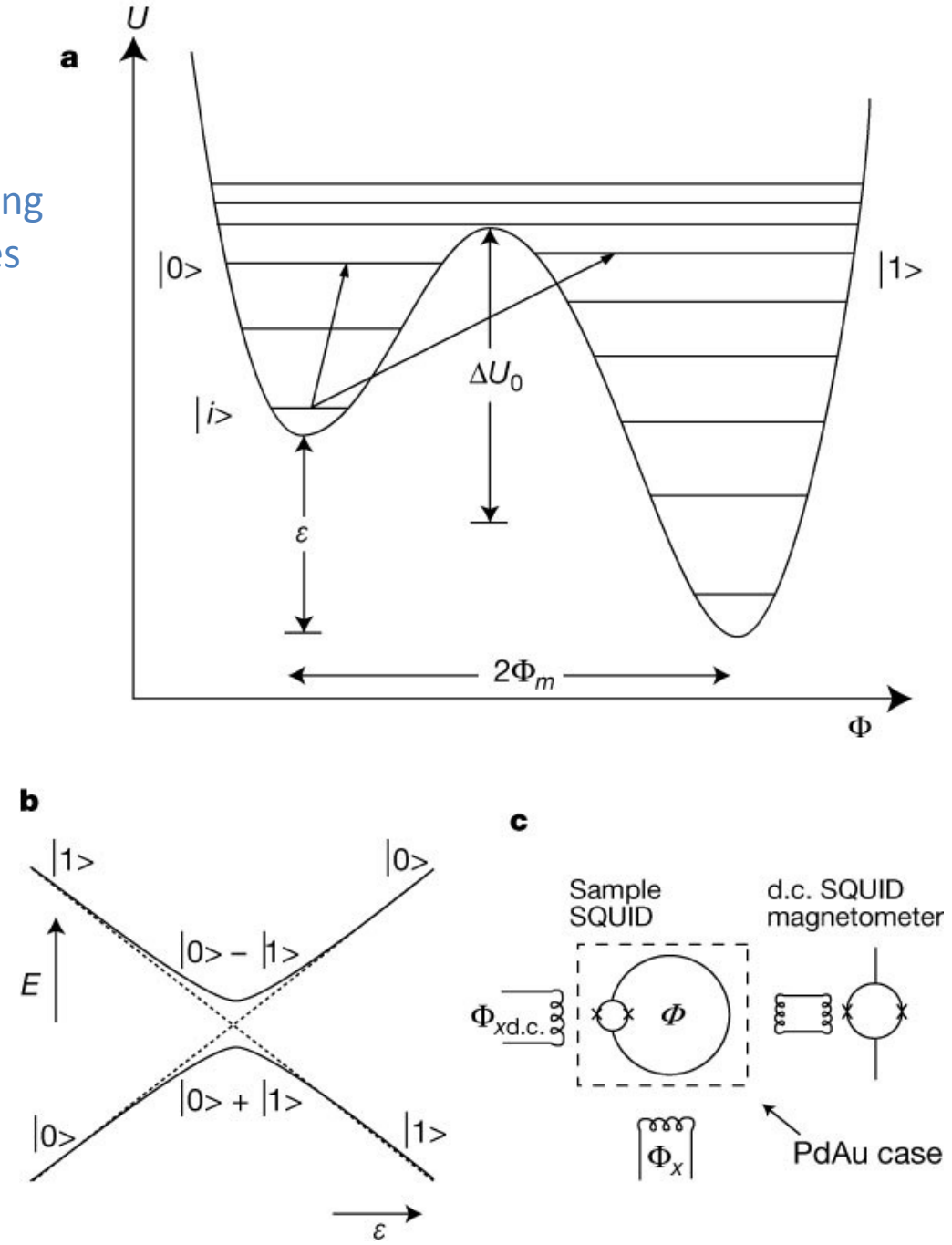
Quelle: Y. Nakamura, Yu. A. Pashkin and J. S. Tsai, Nature 398, 786-788(29 April 1999)

Einzel-Cooper-Paar box: Kohärente Oszillationen



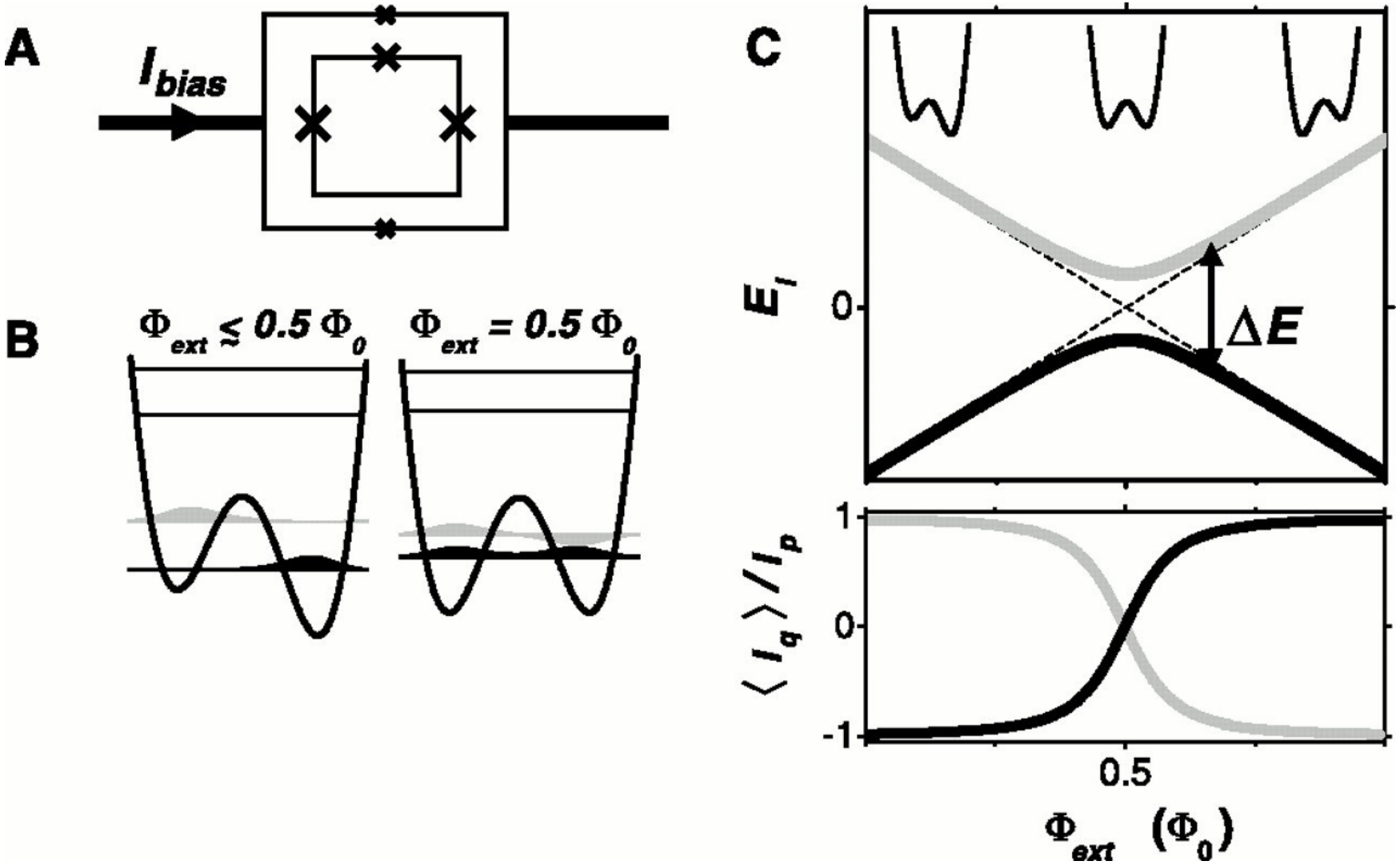
Quelle: Y. Nakamura, Yu. A. Pashkin and J. S. Tsai, Nature 398, 786-788(29 April 1999)

SQUID-Potential, level anti-crossing und experimenteller Aufbau eines flux qubits



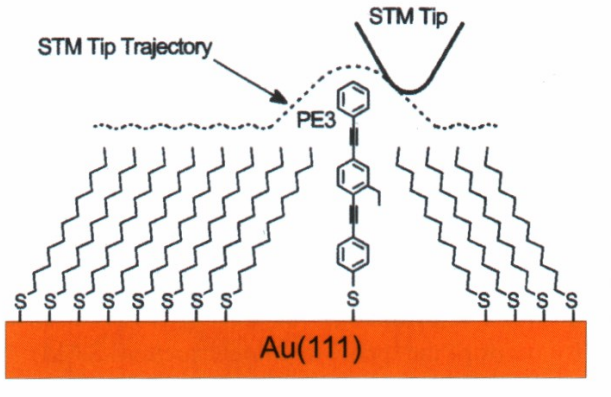
Quelle: Friedman et al., Nature 406, 43(2000)

Beispiel eines flux qubits mit 3 Josephson-Brücken



Quelle: van der Wal et al., Science 290, 773(2000)

Konzepte molekularer Elektronik



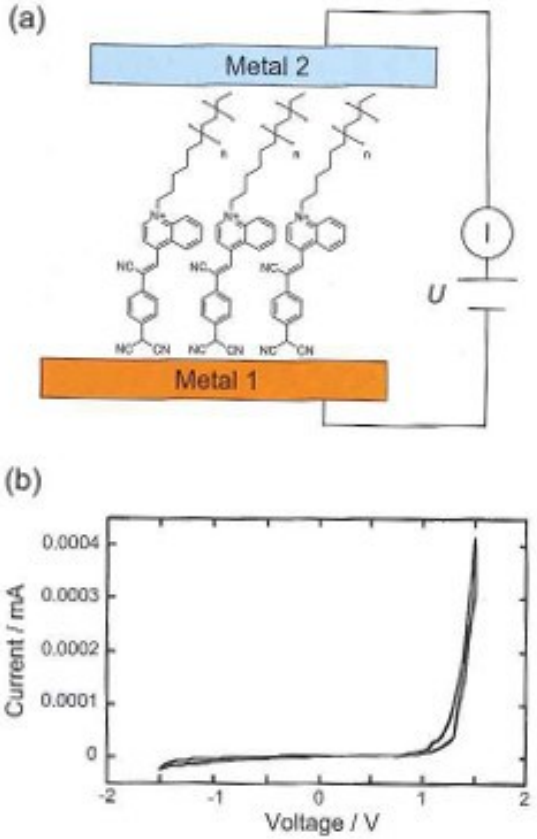


Figure 14:
(a) Schematic presentation of a rectifying device based on an LB-film of the donor-acceptor molecule **12**.
(b) The *I-V* curve of the sandwiched LB-monolayer displaying rectifying character.

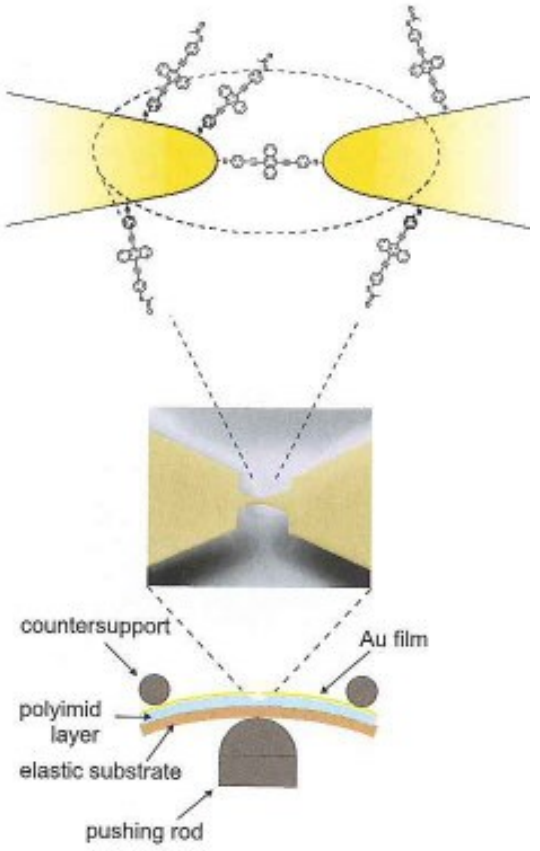
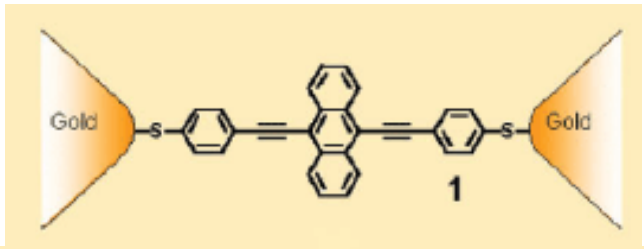
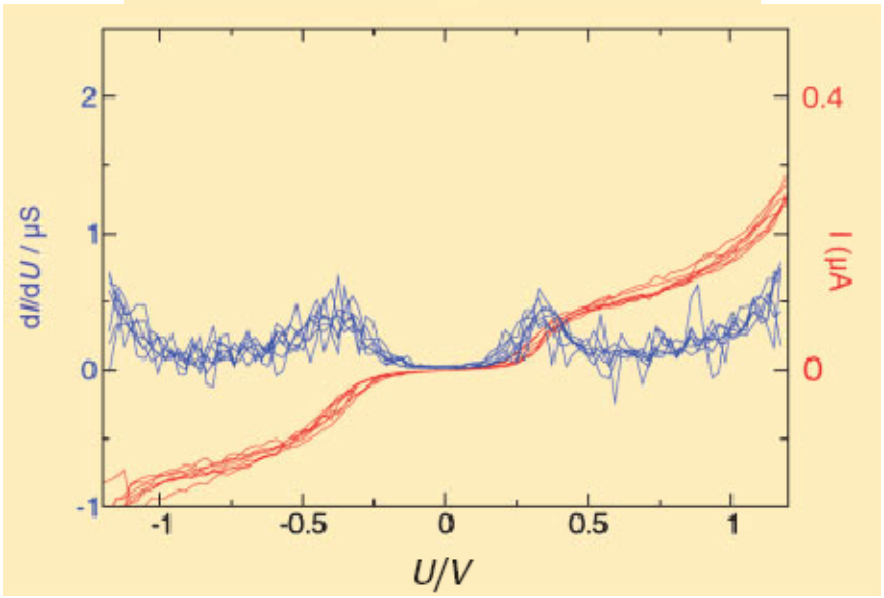
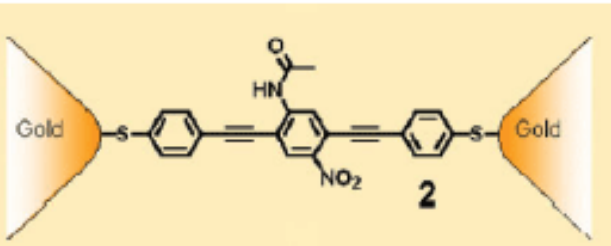


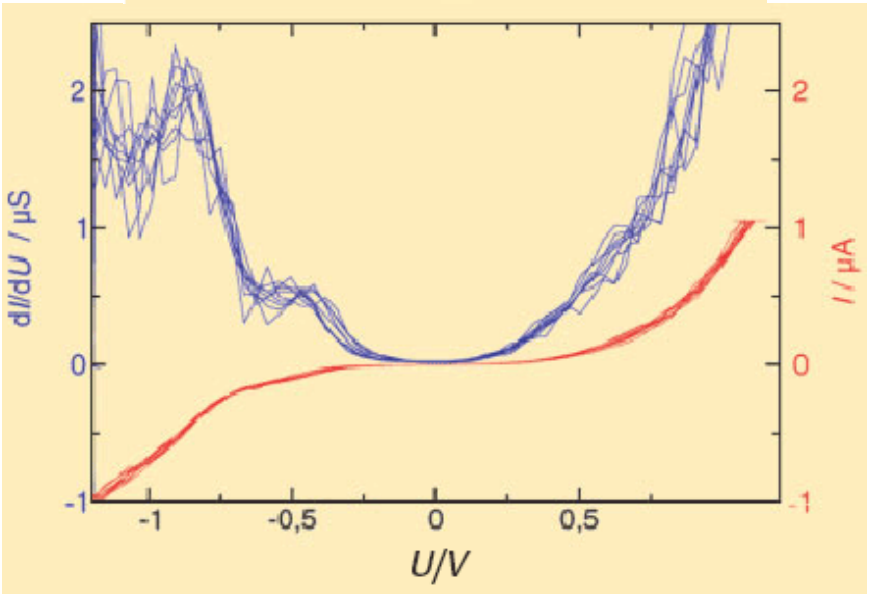
Figure 19: Scheme of a mechanically controlled break junction. On top of a bendable substrate a Au film is patterned with a freely suspended bridge in the center (see SEM micrograph in the middle panel). By bending the substrate in a three point support (lower panel: the pushing rod is driven by a motor), the Au bridge can be broken into two electrodes. By bending the substrate back and forth, the gap between the electrodes can be tuned with a distance resolution much better than Angstroms. This setup can be used to match the electrode gap precisely to the length of a molecule and finally to contact a molecule from two sides via well defined chemical S-Au bonds.

Molekulare Elektronik mit Bruchkontakten



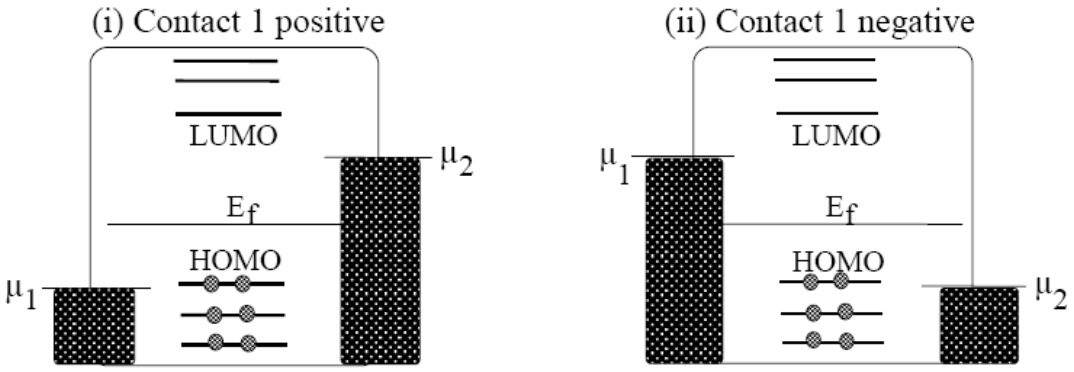






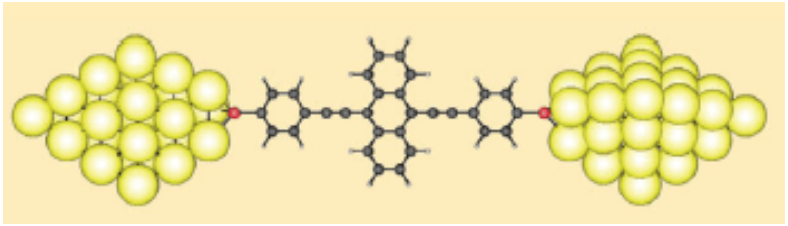
Quelle: H.B. Weber et al.

- Wichtige Elemente zum Verständnis der I(V)-Kurven sind
  - Energieniveau-Diagramm mit der Lage der Fermi-Energie relativ zu LUMO und HOMO



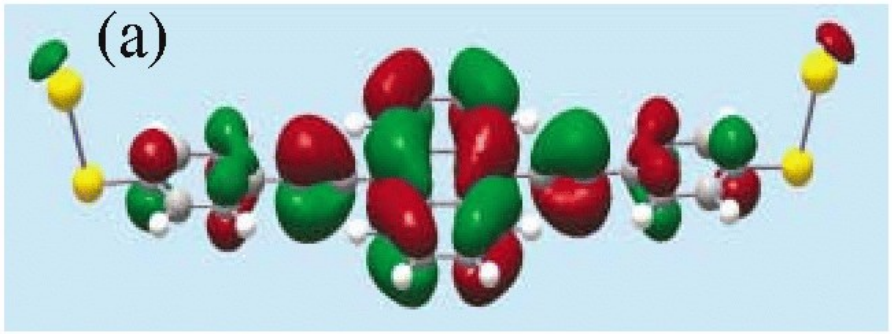
  - Verbreiterung der molekularen Niveaus aufgrund der Elektrodenkontakte

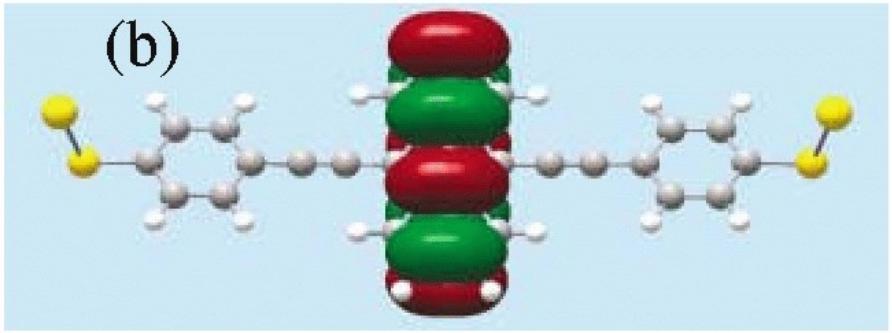
Molekulare Elektronik mit Bruchkontakten



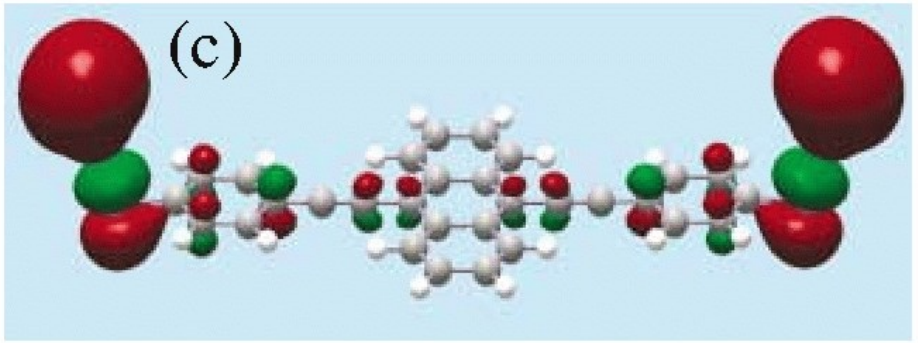
„Super-Molekül“ aus Atomen der Elektroden und dem eigentlichen Molekül

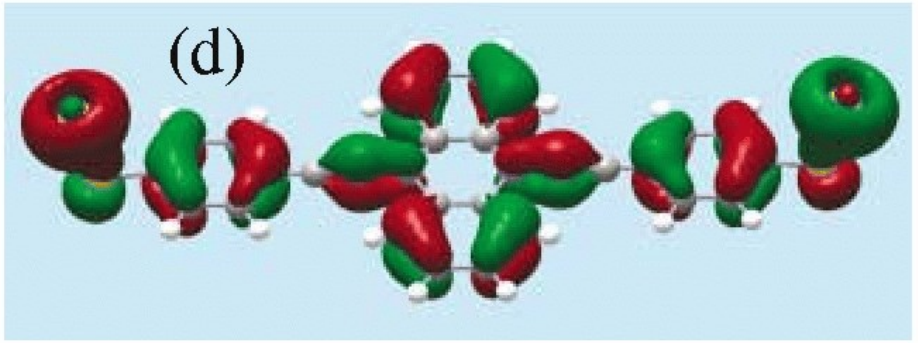
HOMO



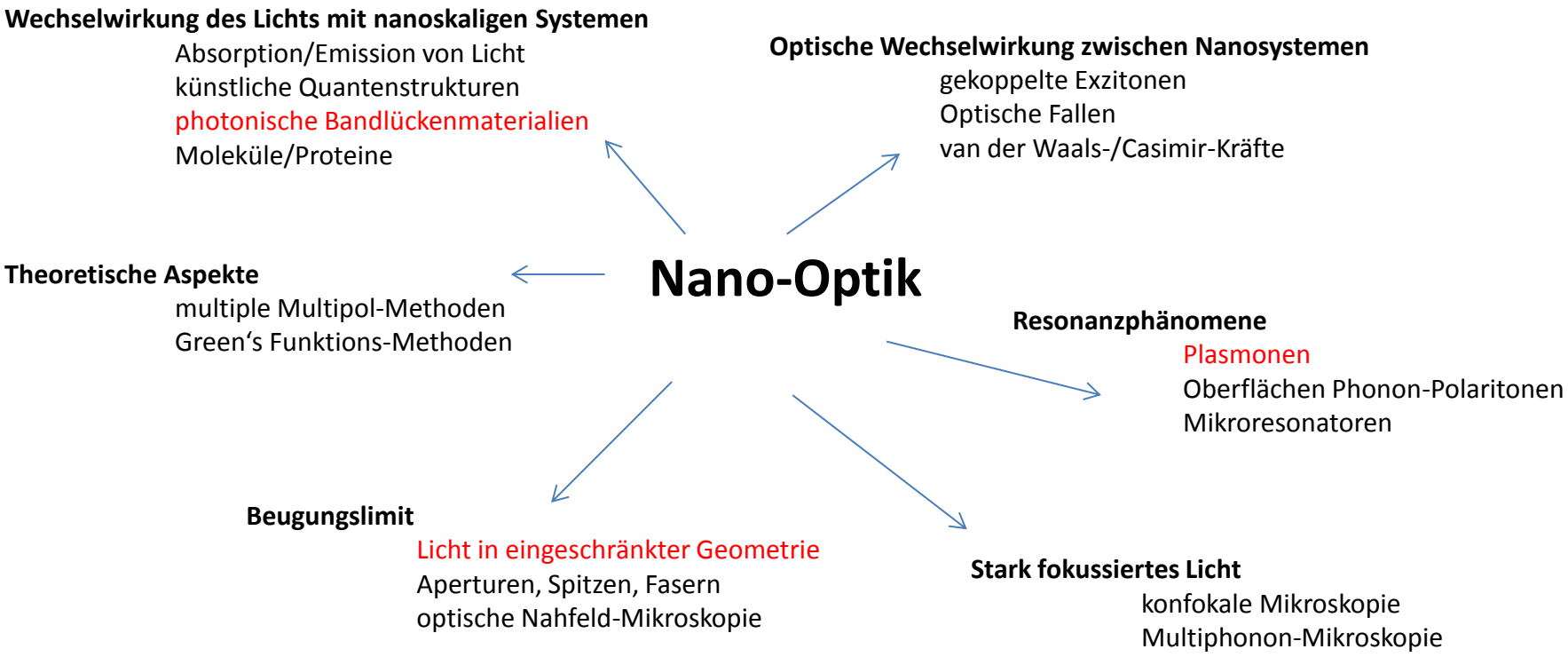


LUMO

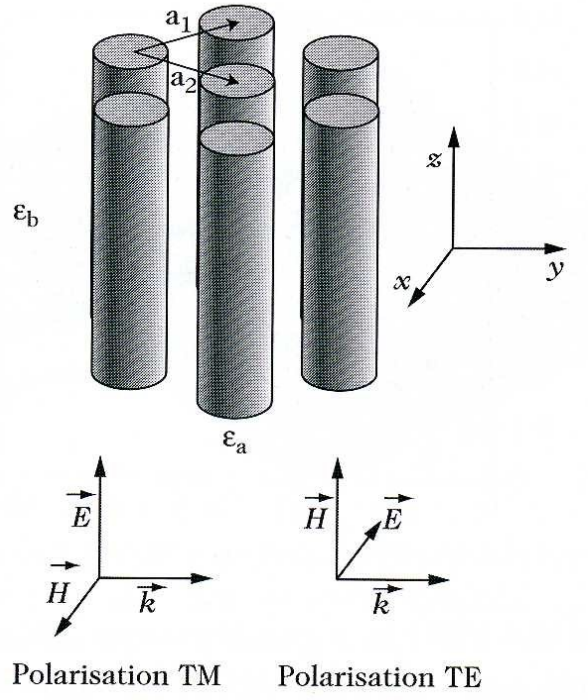




Aspekte der Nano-Optik

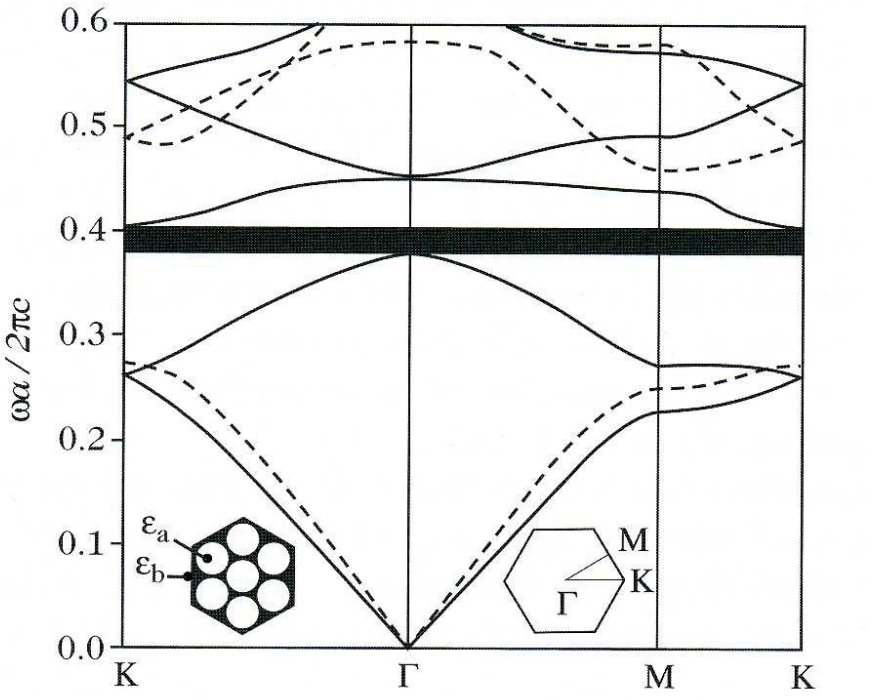


2D photonischer Kristall

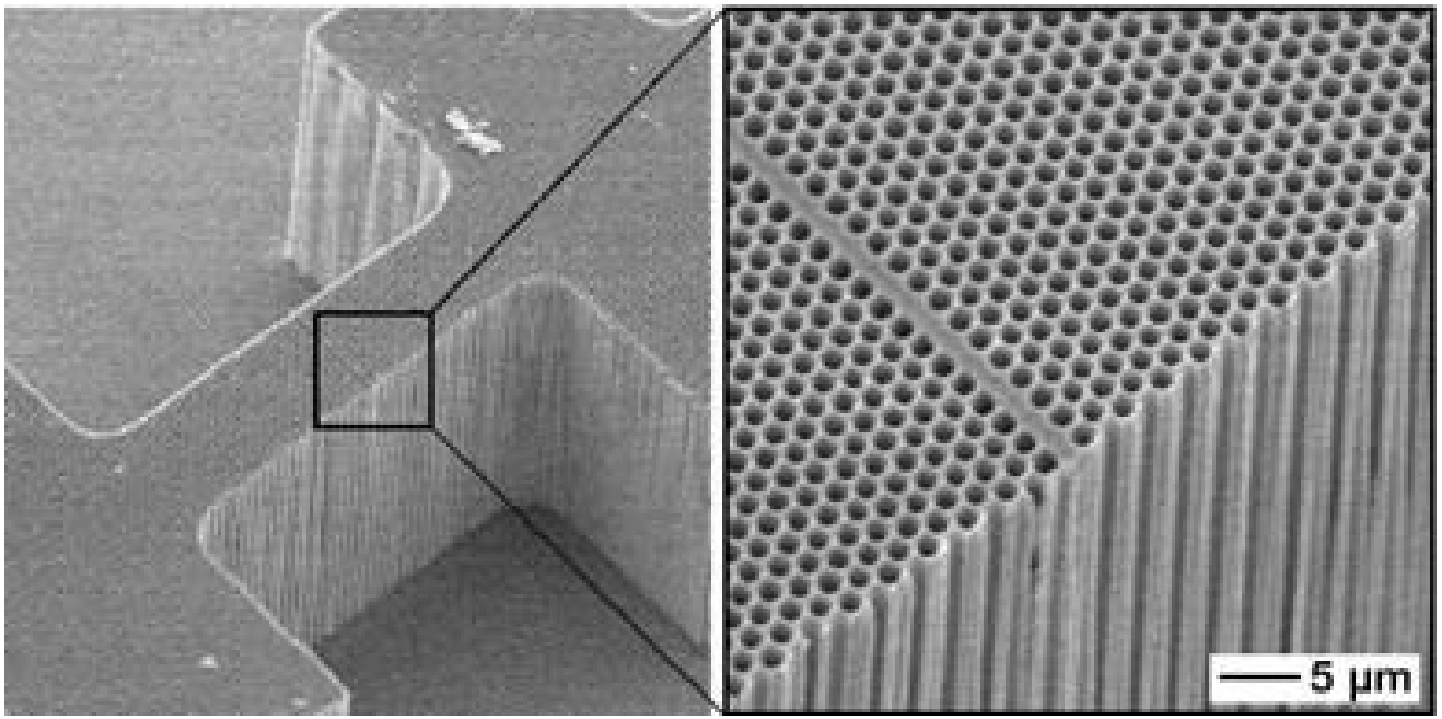


$\varepsilon_a = 1$, $\varepsilon_b = 12{,}25$

Bandstruktur eines 2D photonischen Kristalls



Beispiel eines 2D photonischen Kristalls



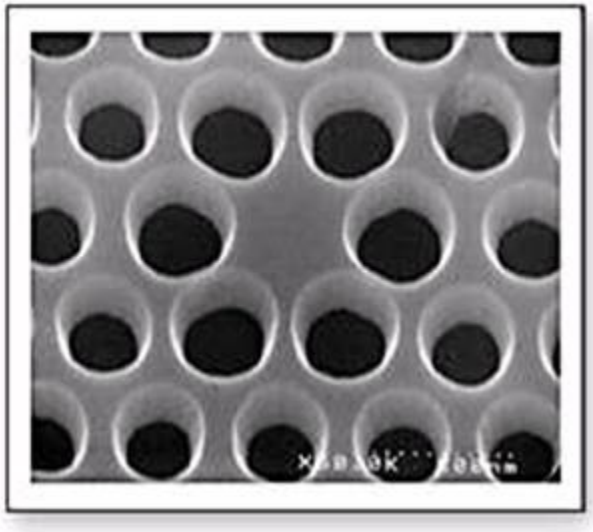
Taken from: R. Wehrsporn, U. Gösele et al., MPI Halle

Größenordnung der Gitterkonstante photonischer Kristalle

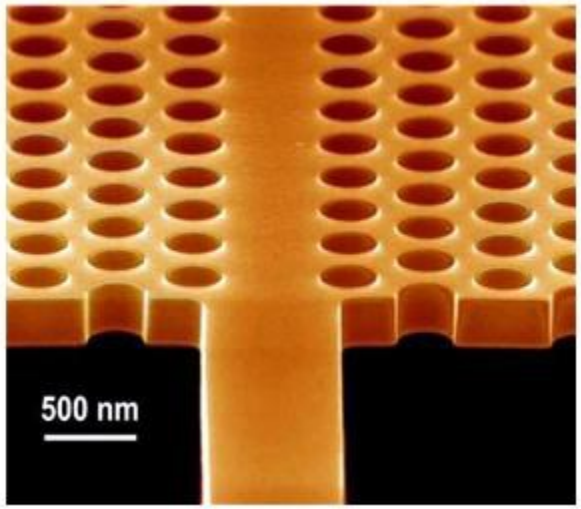
Spektralbereich	Frequenz [Hz]	Wellenlänge λ_0 [m]
Mikrowellen	$10^{9} - 10^{11}$	$10^{-1} - 10^{-3}$
Infrarot	$10^{13} - 10^{14}$	$10^{-5} - 10^{-6}$
sichtbares Licht	10^{14}	10^{-6}
UV-Strahlung	$10^{15} - 10^{16}$	10^{-7}
Röntgenstrahlung	$10^{17} - 10^{19}$	$10^{-8} - 10^{-11}$

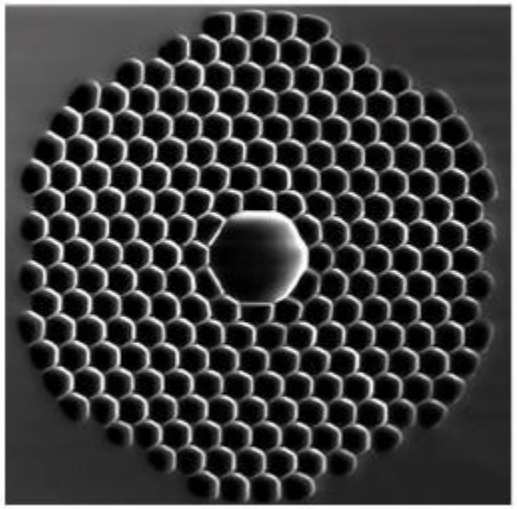
→ Sub-µm-Skala für Bandlücken im Sichtbaren

Beispiele für photonische Defektstrukturen



Mikrokavität

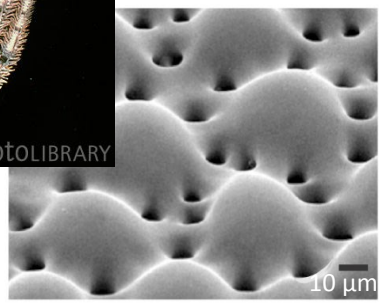


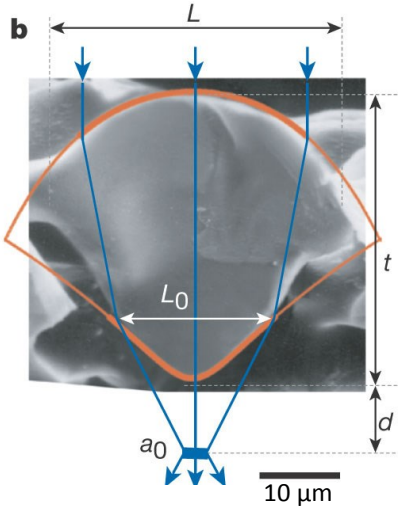


Wellenleiter

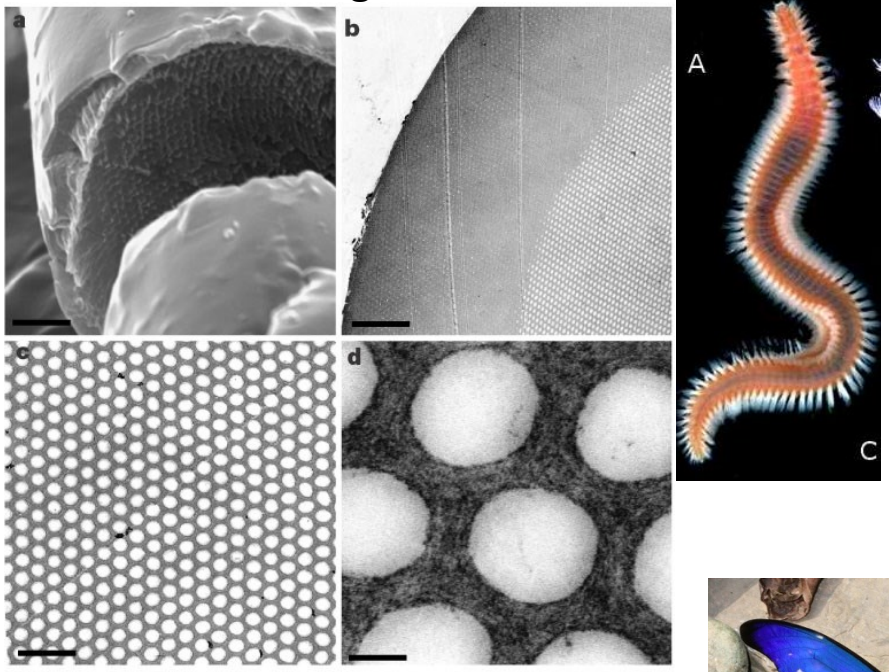
Beispiele für 3D photonische Kristalle in der Natur

Schlangenstern

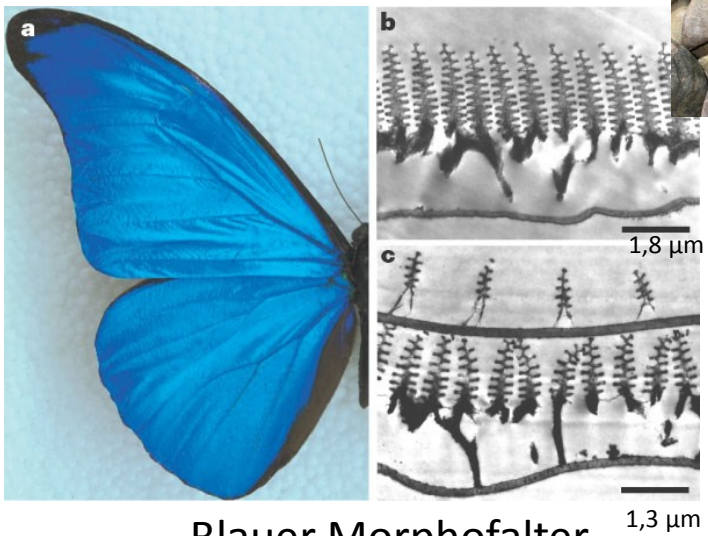




Ringelwürmer

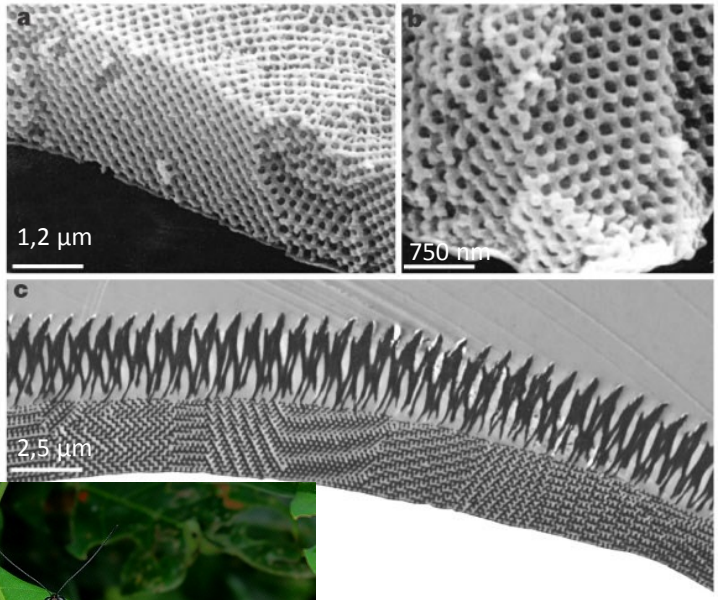


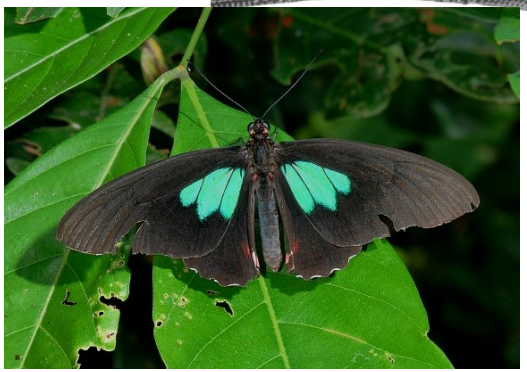
Blauer Morphofalter



Quelle: Vucosic et al., Nature 424, 852(2003)

Beispiele für 3D photonische Kristalle in der Natur

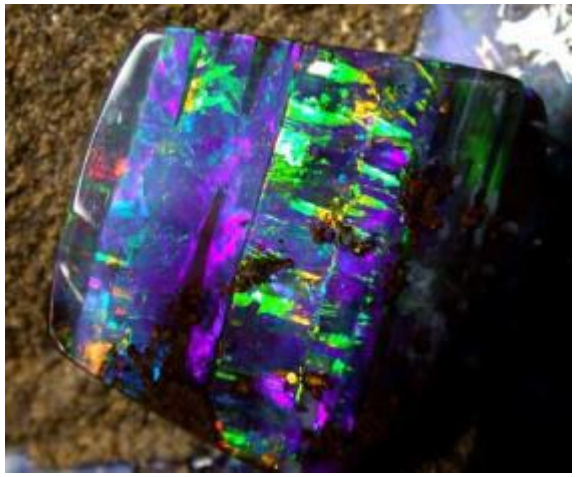


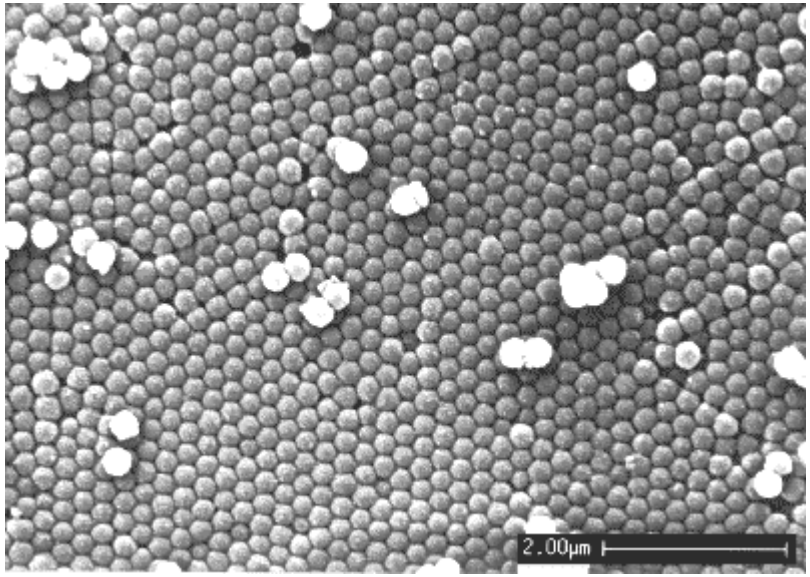


Parides Schmetterling

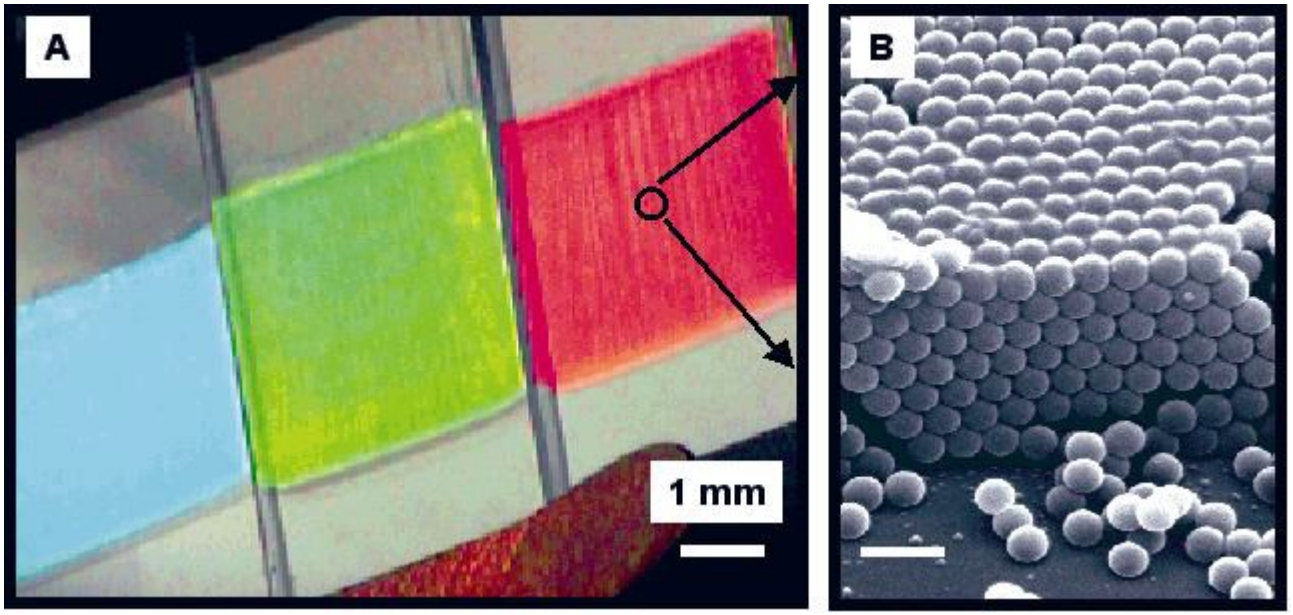
Quelle: Vucosic et al., Nature 424, 852(2003)

Beispiel für 3D photonische Kristalle in der Natur: Opale



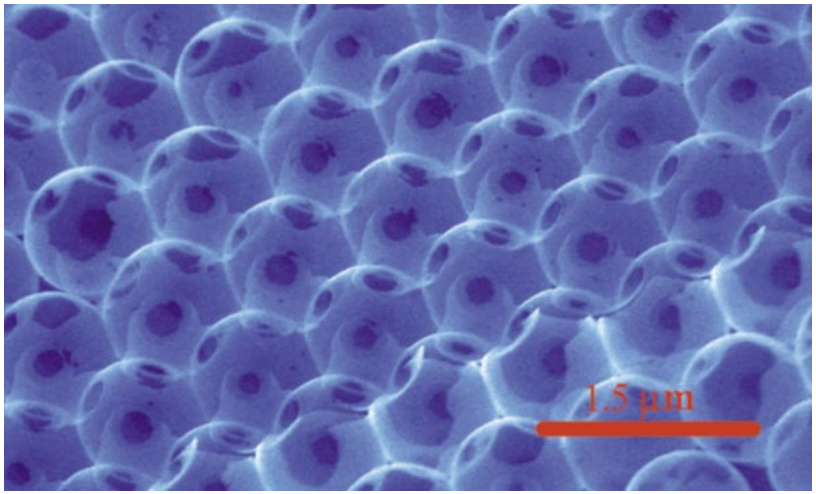


Künstliche Opale

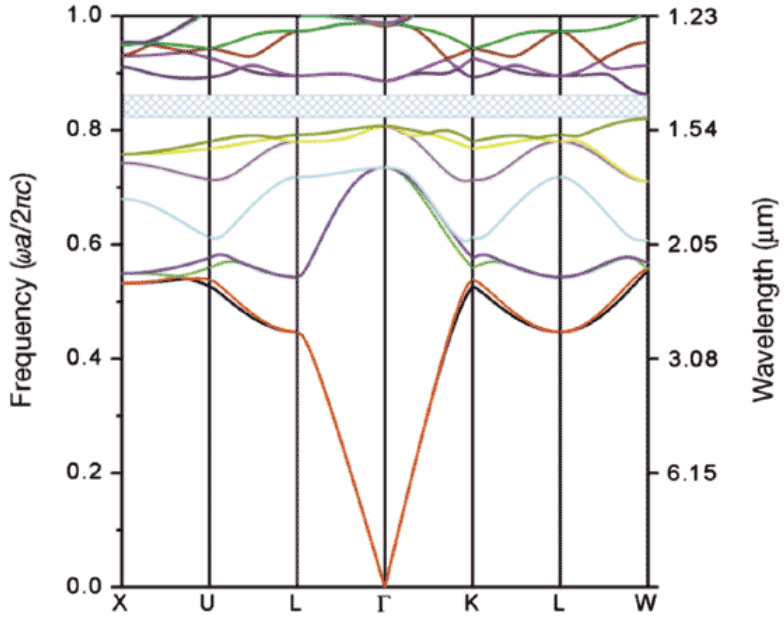


Colvin, MRS Bulletin **26(8),** 637 (2001)

Künstlicher Opal: invertierter Si-Opal mit Bandstruktur

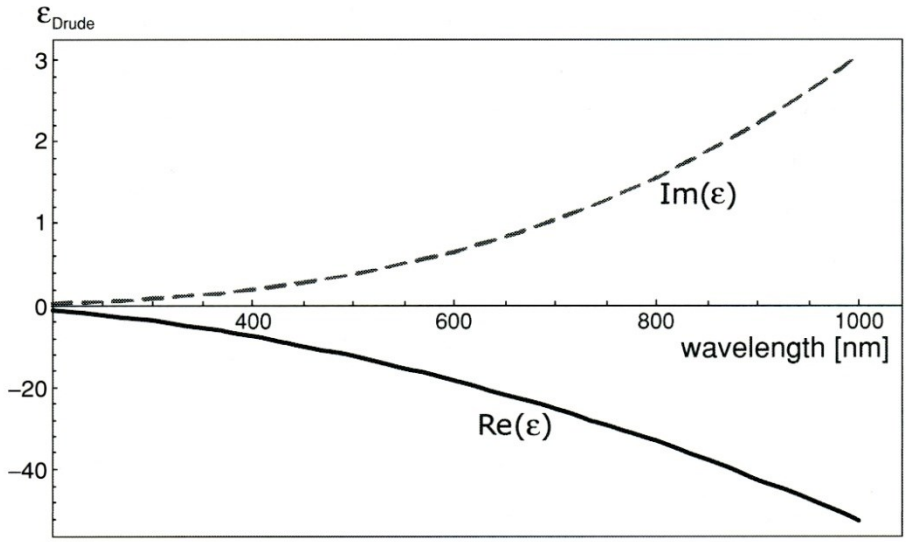


(111) Oberfläche des invertierten Si Opals



Vollständige Bandlücke bei λ_0 = 1,5 µm

Quelle: Blanco et al., Nature 405, 437 (2000)



Plasmonen:

(Drude-)Dielektrizitätskonstante

Figure 12.1 Real and imaginary part of the dielectric constant for gold according to the Drude–Sommerfeld free-electron model ($\hbar\omega_p = 8.95\,\text{eV}$, $\hbar\Gamma = 65.8\,\text{meV}$). The solid line is the real part, the dashed line is the imaginary part. Note the different scales for real and imaginary parts.

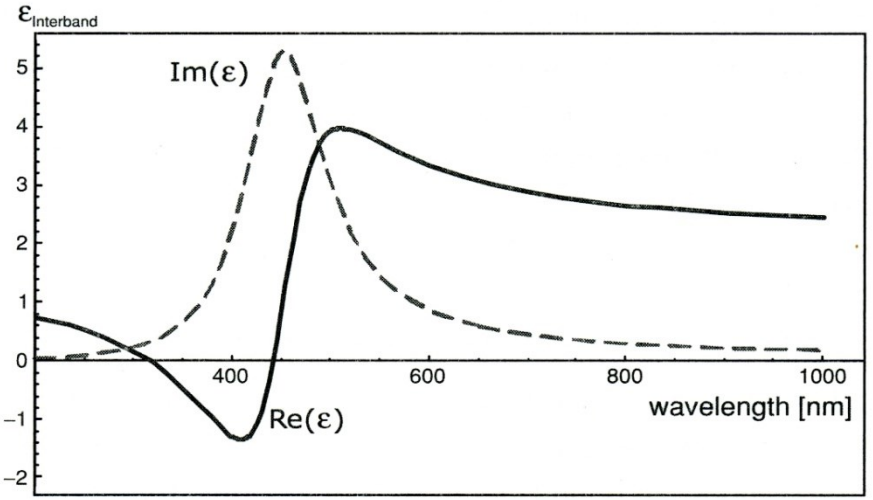


Figure 12.2 Contribution of bound electrons to the dielectric function of gold. The parameters used are $\hbar\tilde{\omega}_p = 2.96\,\text{eV}$, $\hbar\gamma = 0.59\,\text{eV}$, and $\omega_0 = 2\pi c/\lambda$, with $\lambda = 450\,\text{nm}$. The solid line is the real part, the dashed curve is the imaginary part of the dielectric function associated with bound electrons.

Die Fresnel-Gleichungen

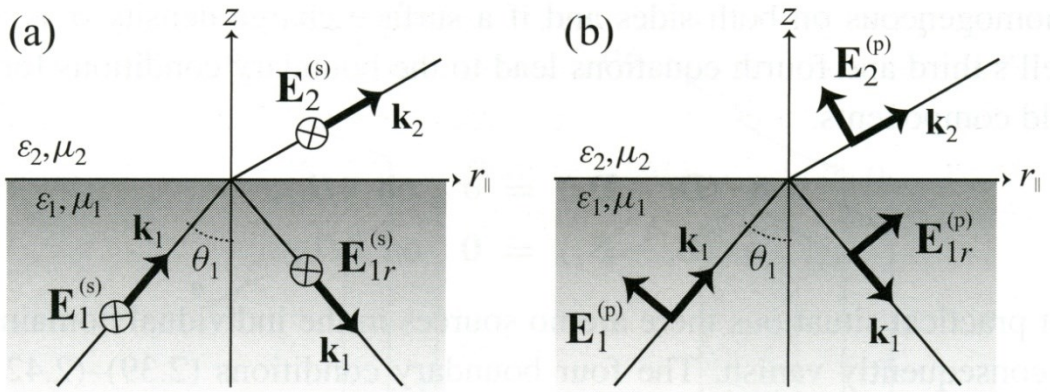


Figure 2.2 Reflection and refraction of a plane wave at a plane interface. (a) s-polarization, and (b) p-polarization.

Amplituden

$$\mathrm{E}_{1r}^{(\mathrm{s})} = \mathrm{E}_1^{(\mathrm{s})}\, r^{\mathrm{s}}(k_x, k_y)\,, \qquad \mathrm{E}_{1r}^{(\mathrm{p})} = \mathrm{E}_1^{(\mathrm{p})}\, r^{\mathrm{p}}(k_x, k_y),$$
$$\mathrm{E}_2^{(\mathrm{s})} = \mathrm{E}_1^{(\mathrm{s})}\, t^{\mathrm{s}}(k_x, k_y)\,, \qquad \mathrm{E}_2^{(\mathrm{p})} = \mathrm{E}_1^{(\mathrm{p})}\, t^{\mathrm{p}}(k_x, k_y)\,,$$

Reflexionskoeffizienten

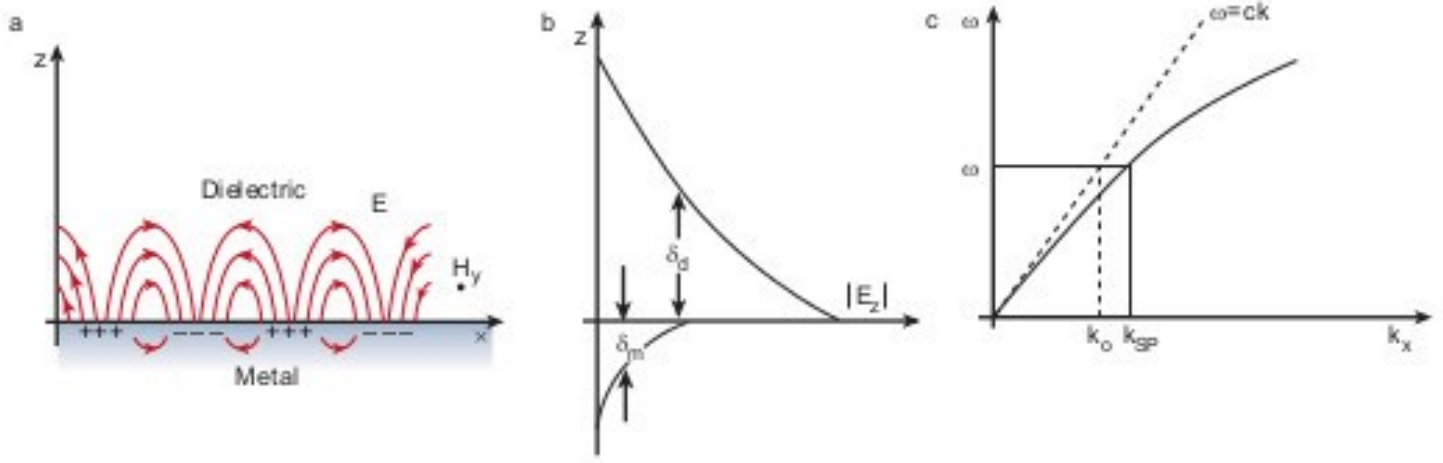
$$r^{\mathrm{s}}(k_x, k_y) = \frac{\mu_2 k_{z_1} - \mu_1 k_{z_2}}{\mu_2 k_{z_1} + \mu_1 k_{z_2}}\,, \qquad r^{\mathrm{p}}(k_x, k_y) = \frac{\varepsilon_2 k_{z_1} - \varepsilon_1 k_{z_2}}{\varepsilon_2 k_{z_1} + \varepsilon_1 k_{z_2}},$$

Transmissionskoeffizienten

$$t^{\mathrm{s}}(k_x, k_y) = \frac{2\mu_2 k_{z_1}}{\mu_2 k_{z_1} + \mu_1 k_{z_2}}\,, \qquad t^{\mathrm{p}}(k_x, k_y) = \frac{2\varepsilon_2 k_{z_1}}{\varepsilon_2 k_{z_1} + \varepsilon_1 k_{z_2}} \sqrt{\frac{\mu_2 \varepsilon_1}{\mu_1 \varepsilon_2}}\,.$$

Oberflächenplasmonen

Box 1
Surface plasmon basics

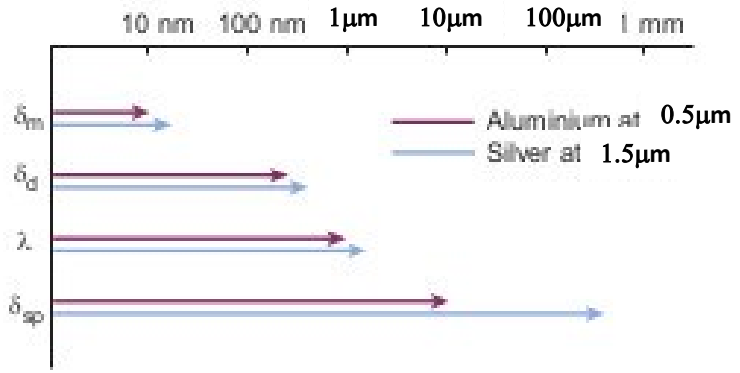


SPs at the interface between a metal and a dielectric material have a combined electromagnetic wave and surface charge character as shown in a. They are transverse magnetic in character (H is in the y direction), and the generation of surface charge requires an electric field normal to the surface. This combined character also leads to the field component perpendicular to the surface being enhanced near the surface and decaying exponentially with distance away from it (b). The field in this perpendicular direction is said to be evanescent, reflecting the bound, non-radiative nature of SPs, and prevents power from propagating away from the surface. In the dielectric medium above the metal, typically air or glass, the decay length of the field, δ_d, is of the order of half the wavelength of light involved, whereas the decay length into the metal, δ_m, is determined by the skin depth. c, The dispersion curve for a SP mode shows the momentum mismatch problem that must be overcome in order to couple light and SP modes together, with the SP mode always lying beyond the light line, that is, it has greater momentum ($\hbar k_{SP}$) than a free space photon ($\hbar k_0$) of the same frequency ω.

Quelle: Barnes et al., Nature 424, 824 (2003)

Oberflächenplasmonen: relevante Längenskalen

Box 2
Surface plasmon length scales

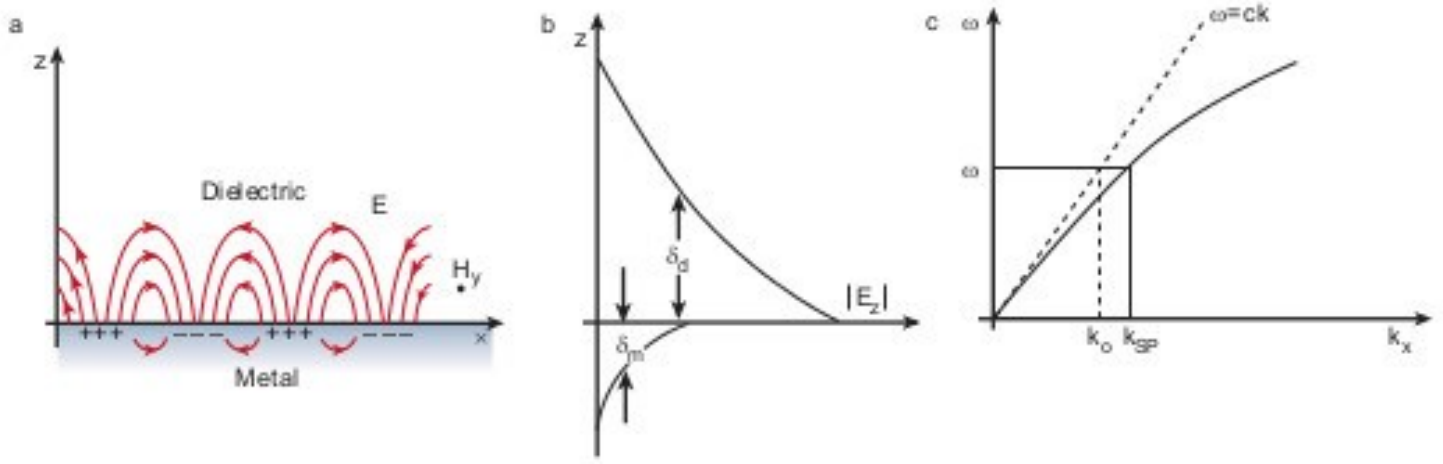


There are three characteristic length scales that are important for SP-based photonics in addition to that of the associated light. The propagation length of the SP mode, δ_{sp}, is usually dictated by loss in the metal. For a relatively absorbing metal such as aluminium the propagation length 2 µm at a wavelength of 500 nm. For a low loss metal, for example, silver, at the same wavelength it is increased to 20 µm. By moving to a slightly longer wavelength, such as 1.55 µm, the propagation length is further increases towards 1 mm. The propagation length sets the upper size limit for any photonic circuit based on SPs. The decay length in the dielectric material, δ_d, is typically of the order of half the wavelength of light involved and dictates the maximum height of any individual features, and thus components, that might be used to control SPs. The ratio of δ_{sp}:δ_d thus gives one measure of the number of SP-based components that may be integrated together. The decay length in the metal, δ_m, determines the minimum feature size that can be used; as shown in the diagram, this is between one and two orders of magnitude smaller than the wavelength involved, thus highlighting the need for good control of fabrication at the nanometre scale. The combinations chosen give an indication of range from poor (Al at 0.5 µm) to good (Ag at 1.5 µm) SP performance.

Quelle: Barnes et al., Nature 424, 824 (2003)

Oberflächenplasmonen

Box 1
Surface plasmon basics



SPs at the interface between a metal and a dielectric material have a combined electromagnetic wave and surface charge character as shown in a. They are transverse magnetic in character (H is in the y direction), and the generation of surface charge requires an electric field normal to the surface. This combined character also leads to the field component perpendicular to the surface being enhanced near the surface and decaying exponentially with distance away from it (b). The field in this perpendicular direction is said to be evanescent, reflecting the bound, non-radiative nature of SPs, and prevents power from propagating away from the surface. In the dielectric medium above the metal, typically air or glass, the decay length of the field, δ_d, is of the order of half the wavelength of light involved, whereas the decay length into the metal, δ_m, is determined by the skin depth. c, The dispersion curve for a SP mode shows the momentum mismatch problem that must be overcome in order to couple light and SP modes together, with the SP mode always lying beyond the light line, that is, it has greater momentum ($\hbar k_{SP}$) than a free space photon ($\hbar k_0$) of the same frequency ω.

Quelle: Barnes et al., Nature 424, 824 (2003)

Dispersionsrelation von Oberflächenplasmonen

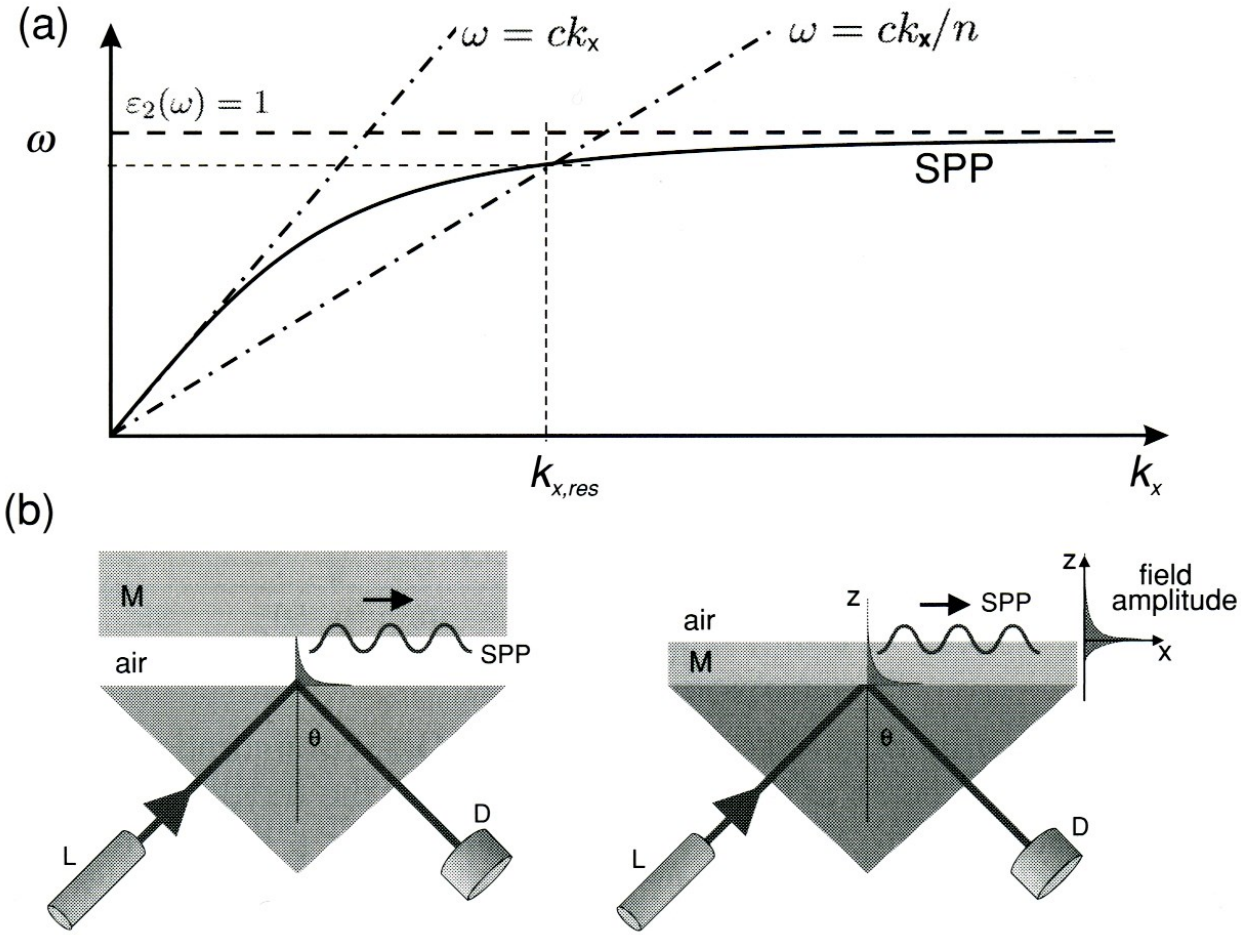
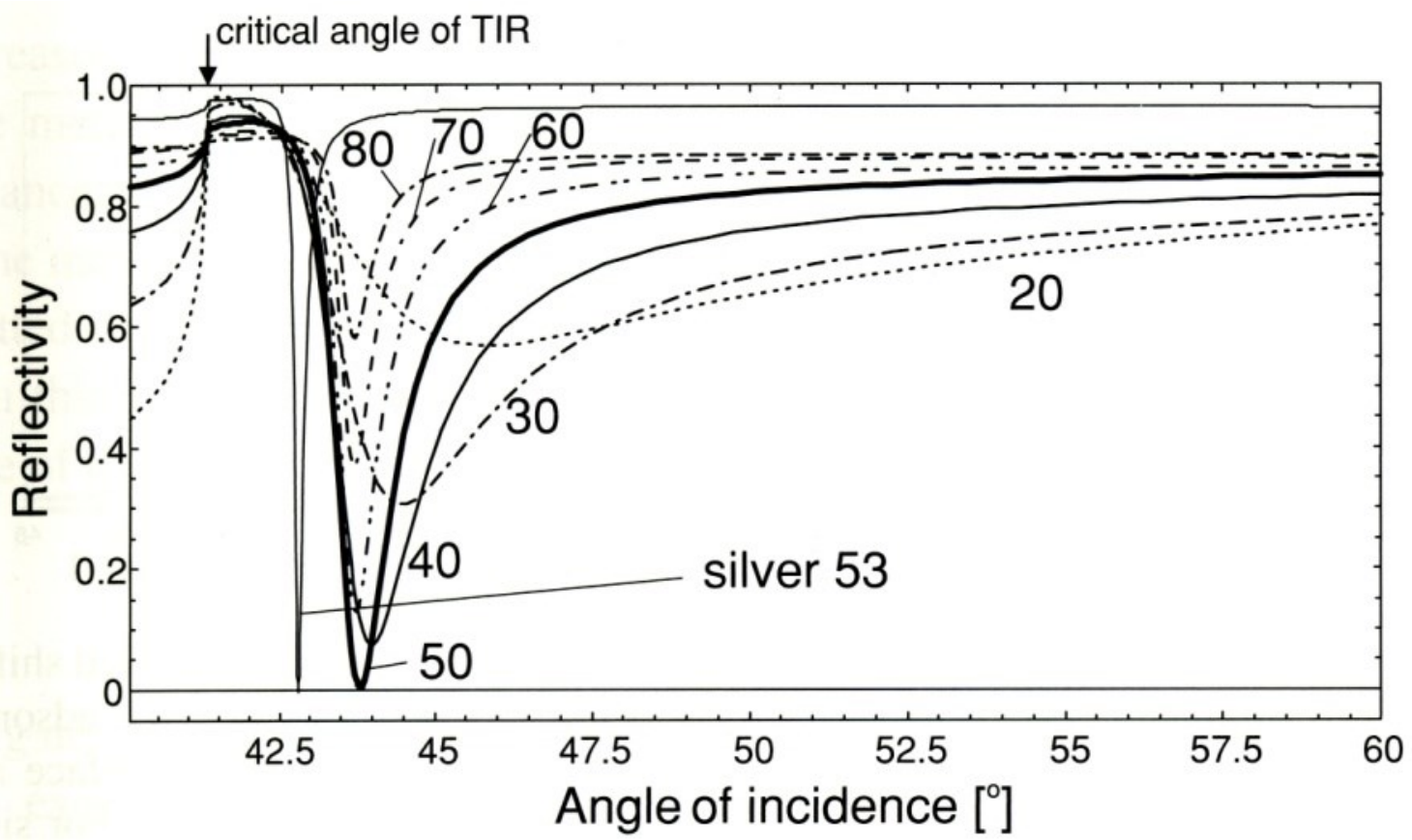
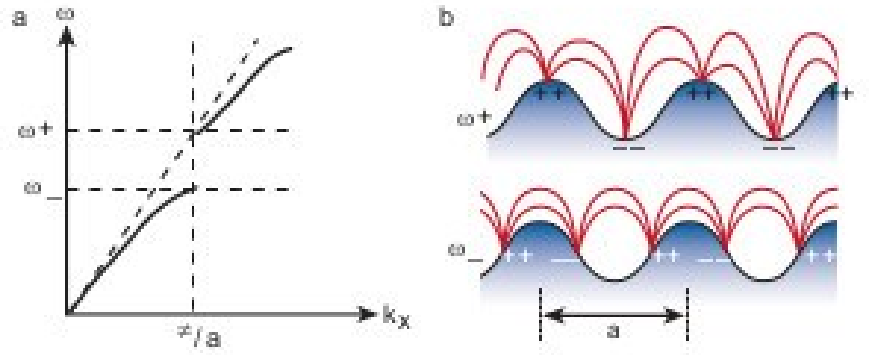


Figure 12.6 Excitation of surface plasmons. (a) Close-up of the dispersion relation with the free-space light line and the tilted light line in glass. (b) Experimental arrangements to realize the condition sketched in (a). Left: Otto configuration. Right: Kretschmann configuration. L: laser, D: detector, M: metal layer.

Anregung von Oberflächenwellen in der Kretschmann-Konfiguration für Goldfilme verschiedener Dicke (in nm)



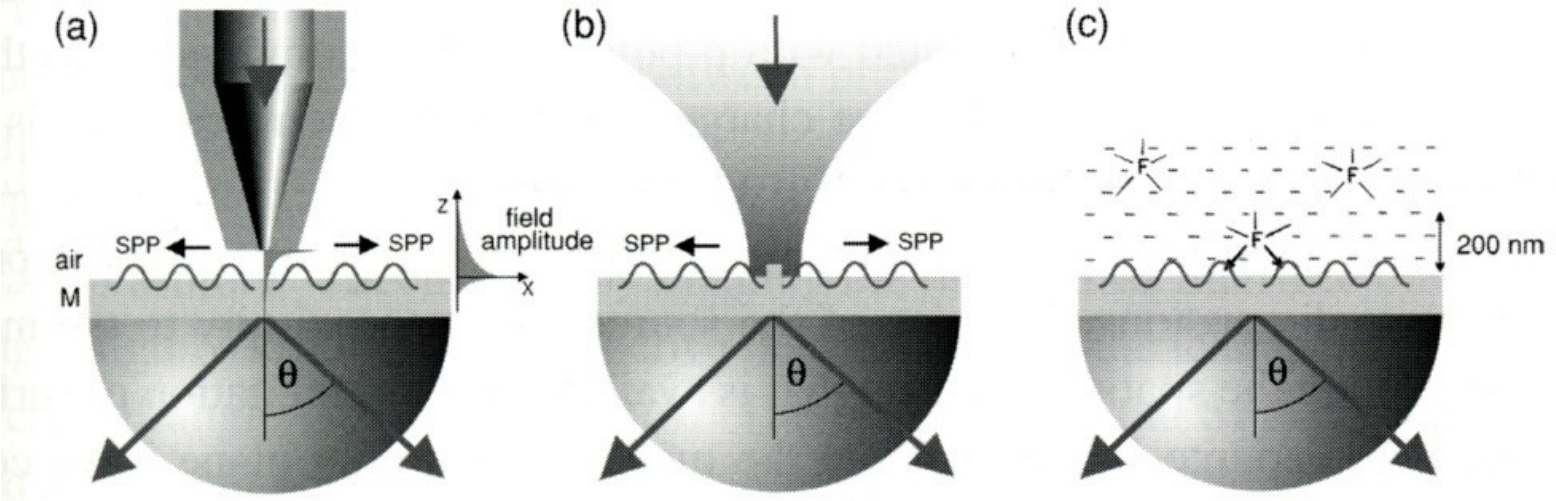
Box 3
Surface plasmon bandgaps



Periodic texturing of the metal surface can lead to the formation of an SP photonic bandgap when the period, a, is equal to half the wavelength of the SP, as shown in the dispersion diagram (a). Just as for electron waves in crystalline solids, there are two SP standing wave solutions, each with the same wavelength but, owing to their different field and surface charge distributions, they are of different frequencies. The upper frequency solution, ω_+, is of higher energy because of the greater distance between the surface charges and the greater distortion of the field, as shown schematically in b. SP modes with frequencies between the two band edges, ω_+ and ω_-, cannot propagate, and so this frequency interval is known as a stop gap. By providing periodic texture in two dimensions, SP propagation in all in-plane directions can be blocked, leading to the full bandgap for SPs. At the band edges the density of SP states is high, and there is a significant increase in the associated field enhancement.

Anregung von Oberflächenplasmonen durch periodischer Gitterkoppler

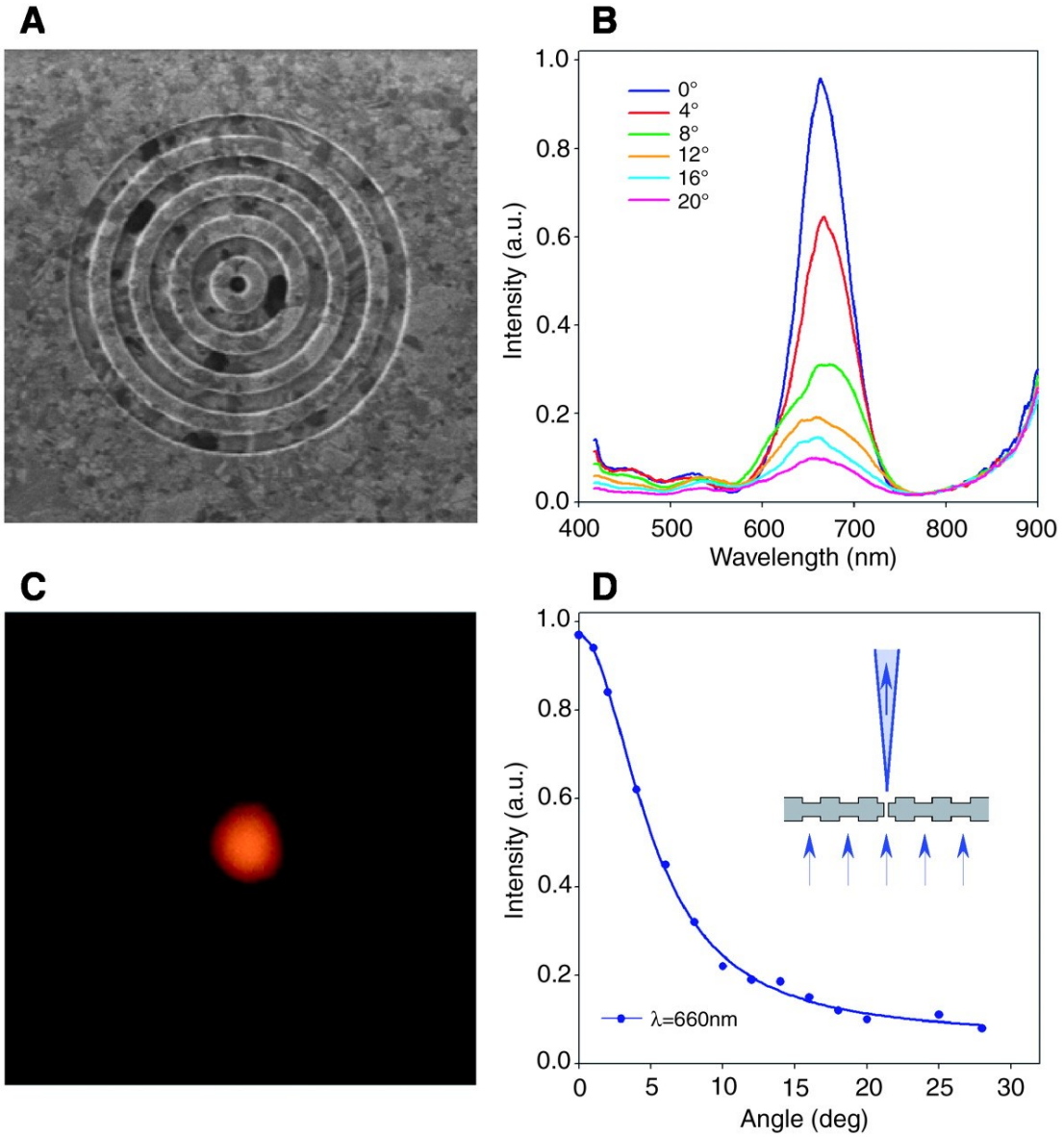
Quelle: Barnes et al., Nature 424, 824 (2003)



Lokale Anregung von Oberflächenplasmonen durch verschiedene eingeschränkte Lichtfelder: (a) Sub-Wellenlänge-Öffnung, (b) bestrahltes Nanoteilchen, (c) fluoreszierende Moleküle

Anregung von Oberflächenwellen durch Sub-Wellenlängen-Öffnungen

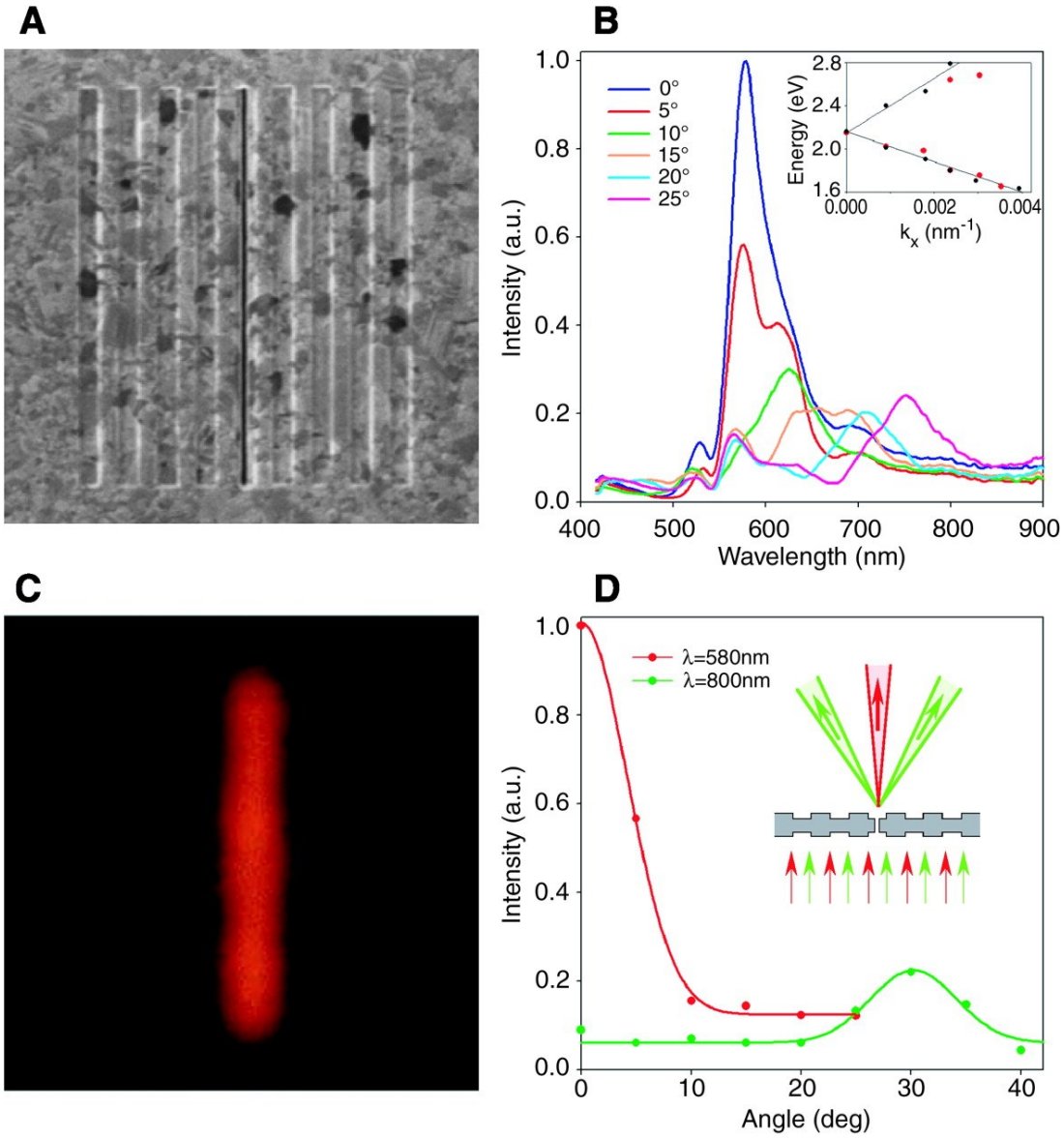
cylindrical hole in
a suspended Ag film
(groove periodicity, 500 nm;
groove depth, 60 nm;
hole diameter, 250 nm;
film thickness, 300 nm).



Quelle: Lezec et al., Science 297, 820 (2002)

Anregung von Oberflächenwellen durch Sub-Wellenlängen-Öffnungen

parallel
grooves on both sides
of a suspended Ag film
(slit width, 40 nm;
slit length, 4400 nm;
groove periodicity, 500 nm;
groove depth, 60 nm;
film thickness, 300 nm)



Quelle: Lezec et al., Science 297, 820 (2002)

Streuung und Absorption von Licht durch Plasmonenmoden in kleinen Goldteilchen

Abb. 2 Der „Lycurgus-Kelch" aus dem 4. Jahrhundert erscheint im reflektierten Licht grün. Wird er jedoch von innen beleuchtet, dann erstrahlt er in einem satten Rot.

Der Lotus-Effekt

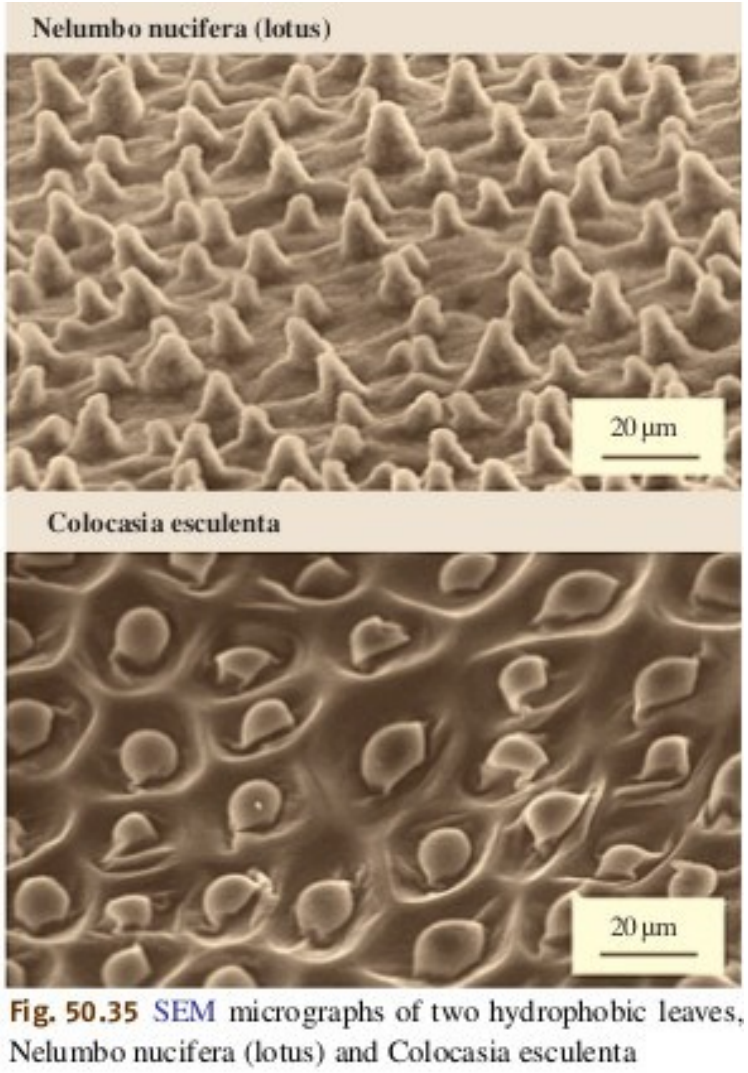
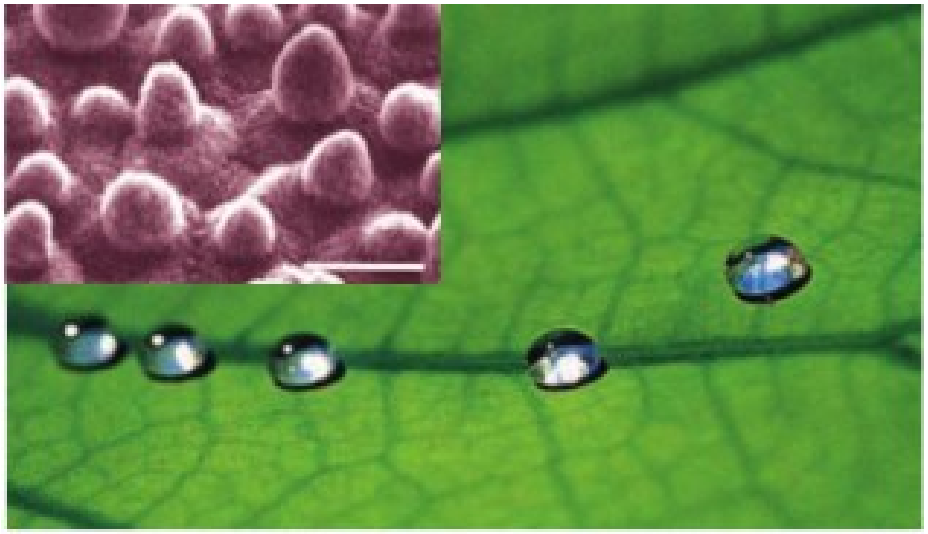


Fig. 50.35 SEM micrographs of two hydrophobic leaves, Nelumbo nucifera (lotus) and Colocasia esculenta

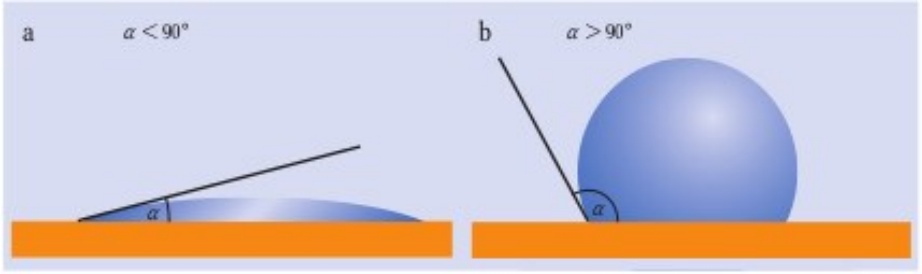
Quelle: Bushan, Nanotechnology, Springer 2003

Quelle: Wikipedia

Der Lotus-Effekt



Regentropfen perlen vom Blatt der Lotuspflanze *(Nelumbo nucifera)* einfach ab. Das liegt an der mikrofeinen Noppenstruktur auf der Blattoberfläche, die erst unter dem Rasterelektronenmikroskop (Inset) zu erkennen ist. (Quelle: W. Barthlott, Uni Bonn)



Bei Randwinkeln unterhalb von 90° (a) gilt eine Oberfläche als hydrophil und das Wasser breitet sich aus. Ist der Randwinkel dagegen größer als 90° (b), so spricht man von hydrophoben Oberflächen, auf denen sich das Wasser zu Tropfen zusammenzieht und leichter abperlen kann.

Quelle: K. Bammel, Physik Journal 4, 48 (2005)

Der Lotus-Effekt

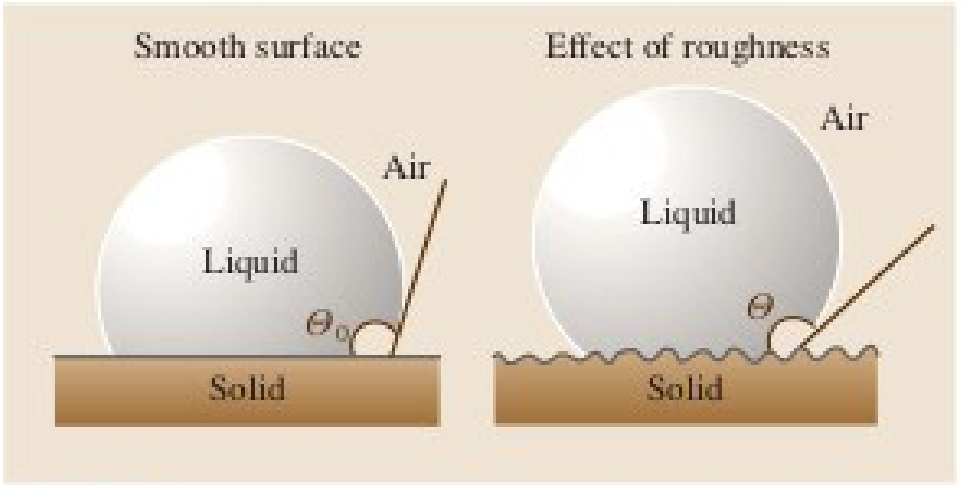


Fig. 50.36 Droplet of liquid in contact with a smooth solid surface (contact angle θ_0) and rough solid surface (contact angle θ) [50.173]

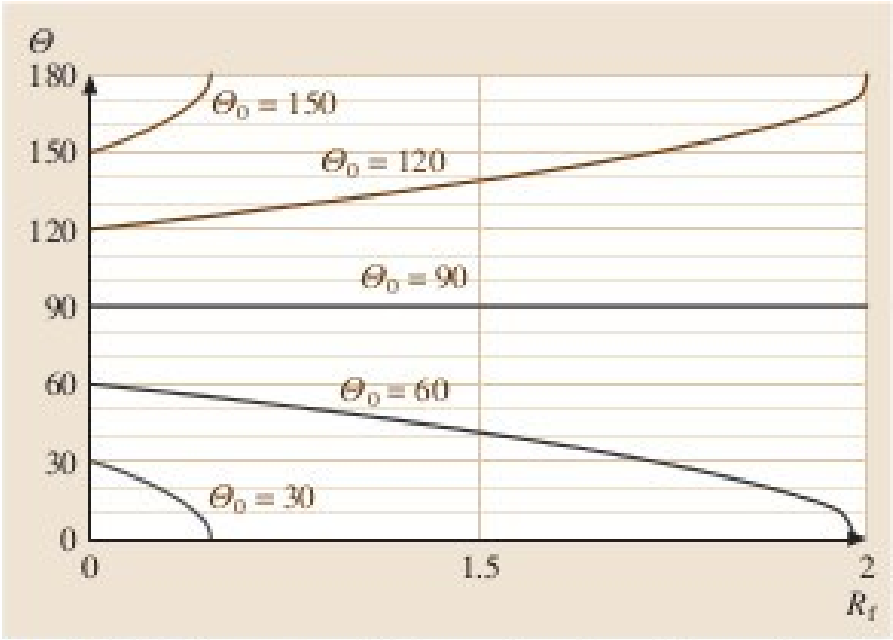


Fig. 50.37 Contact angle for rough surface (θ) as a function of the roughness factor (R_f) for various contact angles for a smooth surface (θ_0) [50.173]

Quelle: Bushan, Nanotechnology, Springer 2003

Superhydrophobizität („Lotus-Effekt“)

Der Lotus-Effekt beschreibt die Selbstreinigungsfähigkeit von Blättern.

Namensgeber ist die Lotuspflanze: Wasser perlt von ihrem Blättern ab, ohne sie zu benetzen.

Diese Eigenschaft nennt man hydrophob.

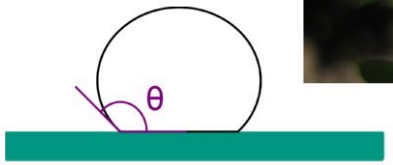
Auch einheimische Pflanzen wie die Kapuzinerkresse, die Akelei und der Kohlrabi zeigen diesen Effekt.

Blatt einer Akelei, Verschmutzung (durch Curry-Pulver simuliert) perlt mit Wasser einfach ab..

Superhydrophobizität („Lotus-Effekt“)

Geranie (hydrophil)

Akelei (hydrophob)

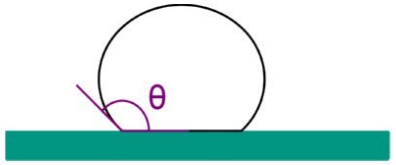


Superhydrophobizität („Lotus-Effekt“)

In Abhängigkeit vom Kontaktwinkel θ werden Oberflächen eingeteilt:

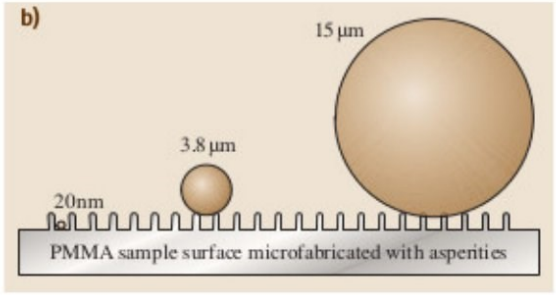
θ kleiner als 90°: 

→ hydrophil („Wasser liebend“)

θ größer als 90°: 

→ hydrophob („Wasser fürchtend“)

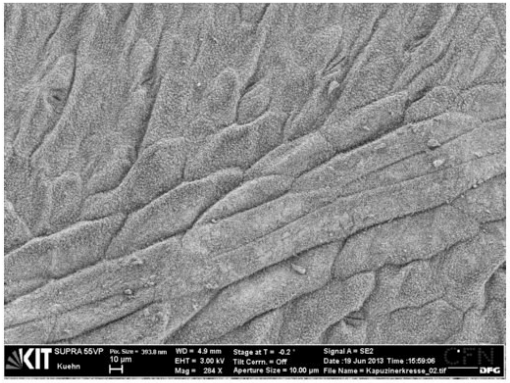
Durch Strukturierung können die hydrophoben Eigenschaften noch weiter verstärkt werden (Superhydrophobizität):



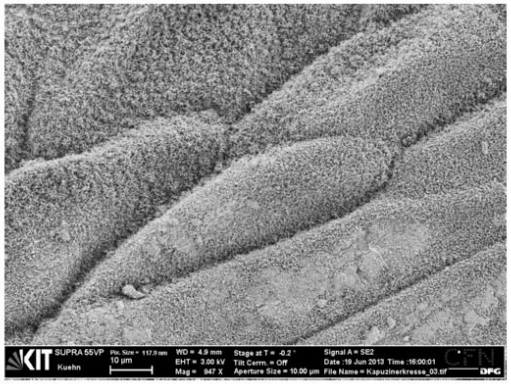
Quelle:
B. Bushan: *„Springer Handbook of Nanotechnology“*, Springer, 2007.

Blattoberfläche von Kapuzinerkresse

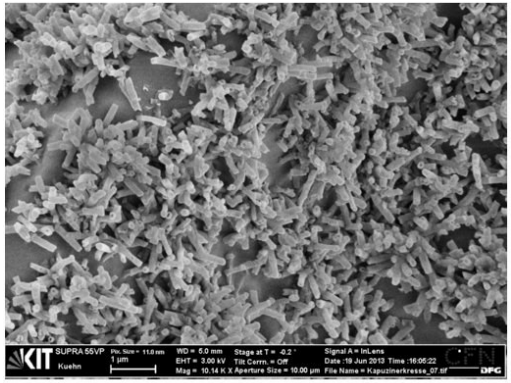
280 ×



950 ×



10 000 × vergrößert

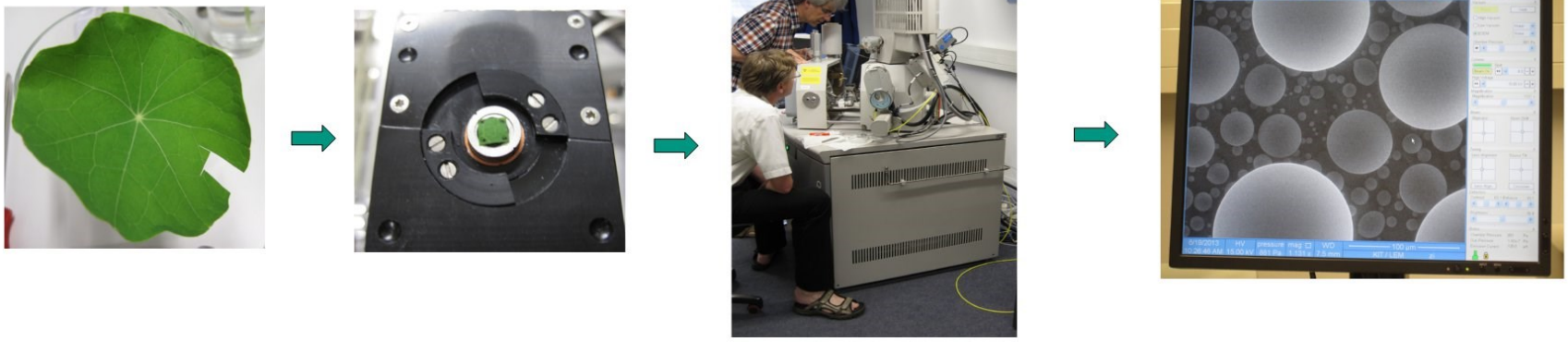


Nanoskalige Wachsstruktur verstärkt hydrophobe Eigenschaften der Blattoberfläche → Lotus-Effekt

Blattoberfläche einer getrockneten Kapuzinerkresse im Rasterelektronenmikroskop bei etwa 280- (links), 950- (Mitte) und schließlich 10 000-facher Vergrößerung (rechts). Zu erkennen ist die nanoskalige Wachsstruktur, der die superhydrophoben Eigenschaften der Kresse zugeschrieben werden.

Environmental Scanning Electron Microscope (ESEM) im Laboratorium für Elektronenmikroskopie (LEM)

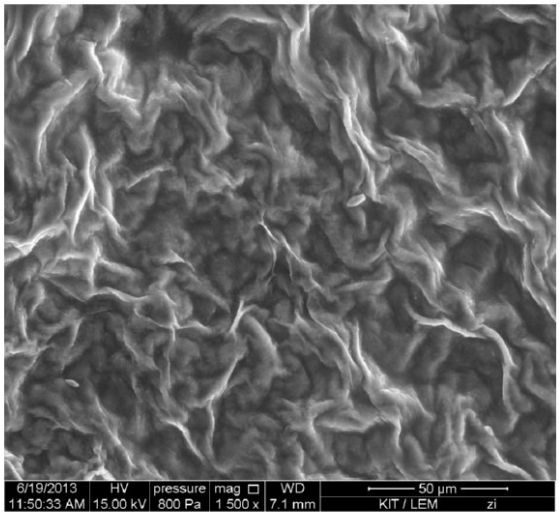
Im Gegensatz zu herkömmlichen Rasterelektronenmikroskopen können im ESEM Proben untersucht werden, die vorher nicht getrocknet und metallisiert werden müssen, was vor allem biologische Proben stark verändern würde. Es ist sogar möglich, auf der gekühlten Probe Wasserdampf zu kondensieren.

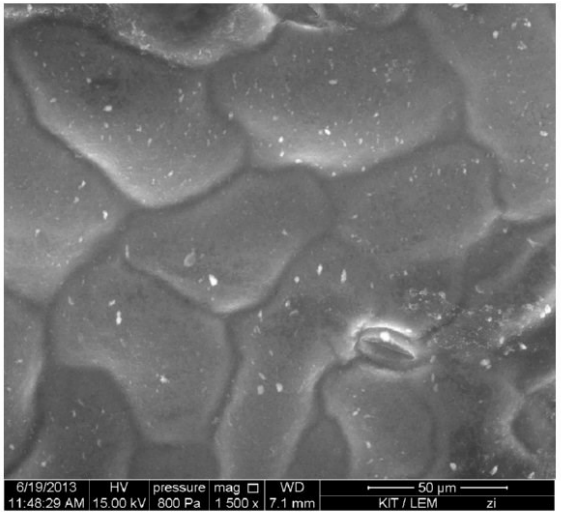


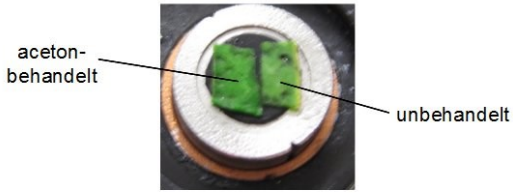
Untersuchung von Superhydrophobizität (Lotus-Effekt) bei der Kapuzinerkresse. Im ESEM wurde auf einem kleinen Blattstück Wasserdampf kondensiert, um die Tröpfchenbildung bei über 1000-facher Vergrößerung auf der Blattoberfläche zu beobachten.

Zerstörung der Superhydrophobizität von Kapuzinerkresse

Durch Behandlung mit organischen Lösemitteln oder Seife kann die Wachsschicht auf den Blättern beschädigt werden. Die superhydrophoben Eigenschaften werden dadurch zerstört.



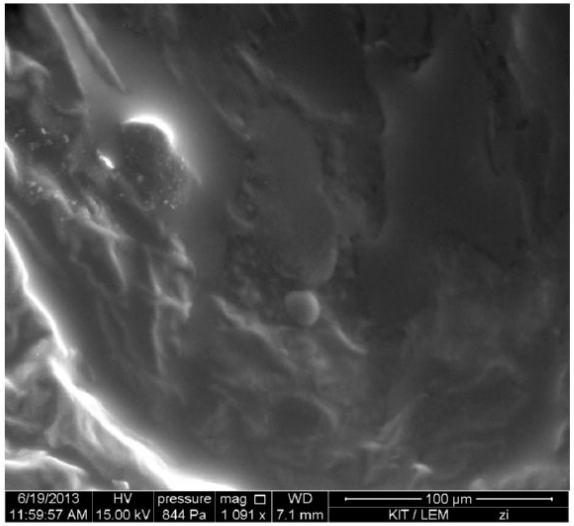


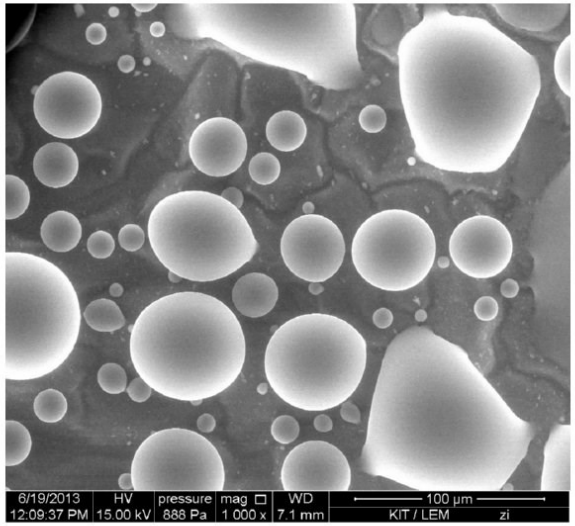


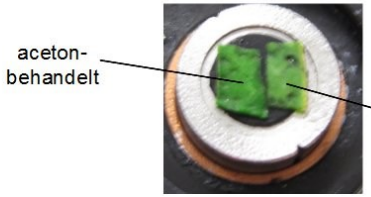
ESEM-Bilder von Kapuzinerkresse. Das linke Bild zeigt eine mit Aceton behandelte Oberfläche, das rechte eine unbehandelte (Vergrößerung 1500-fach).

Zerstörung der Superhydrophobizität von Kapuzinerkresse

Durch Behandlung mit organischen Lösemitteln oder Seife kann die nanostrukturierte Wachsschicht auf den Blättern beschädigt werden. Die superhydrophoben Eigenschaften werden dadurch zerstört.







Bei der Kondensation von Wasserdampf zeigt sich der Unterschied in den ESEM-Bildern. Auf der unbehandelten Blattoberfläche perlt das Wasser ab (rechts). Auf der mit Aceton behandelten Oberfläche (links) dagegen breitet sich das Wasser aus (Vergrößerung jeweils ca. 1000-fach).

Der Lotus-Effekt

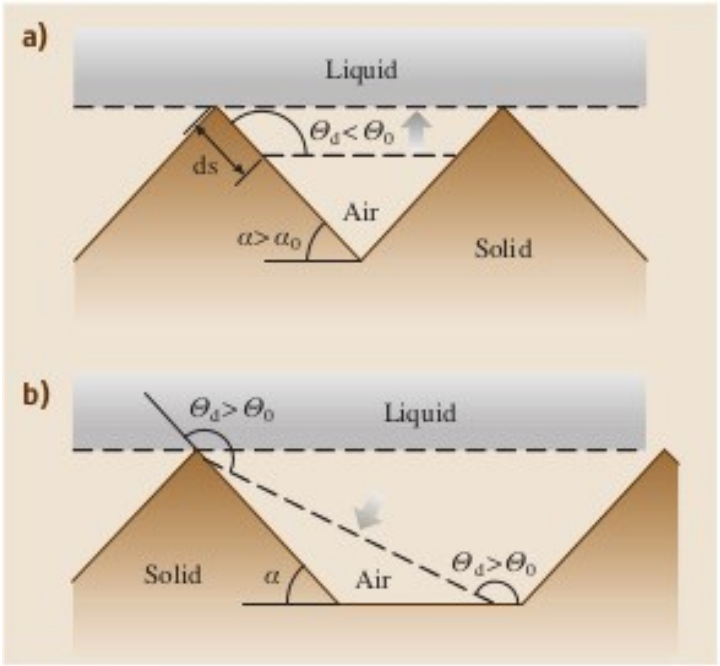


Fig. 50.38 (a) Formation of a composite solid–liquid–air interface for sawtooth and smooth profiles, and (b) destabilization of the composite interface for the sawtooth and smooth profiles due to dynamic effects. The dynamic contact angle $\theta_d > \theta_0$ corresponds to an advancing liquid–air interface, whereas $\theta_d < \theta_0$ corresponds to a receding interface [50.176]

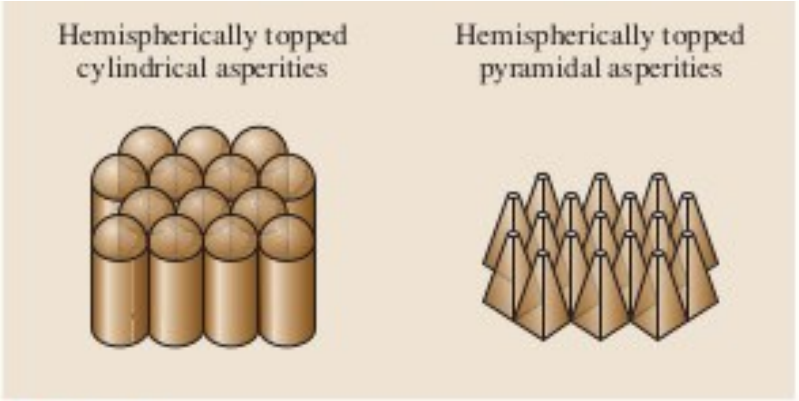


Fig. 50.39 Optimized roughness distribution – hemispherically topped cylindrical asperities and pyramidal asperities with square foundation and rounded tops. The square base gives a higher packing density but introduces undesirable sharp edges [50.173]

Quelle: Bushan, Nanotechnology, Springer 2003

Der Lotus-Effekt

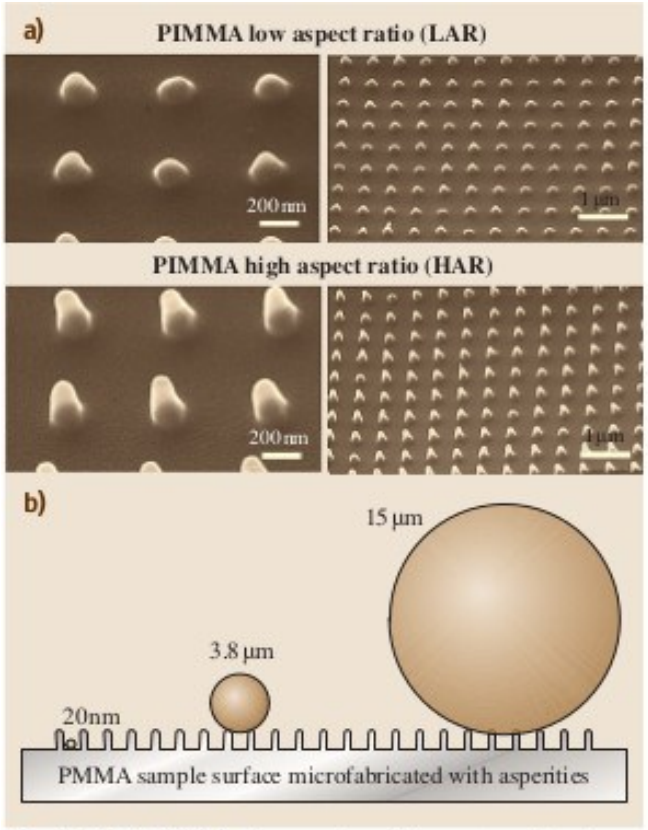


Fig. 50.40 (a) SEM micrographs of the patterned polymer surfaces. Both LAR and HAR are shown at two magnifications to see both the asperity shape and the asperity pattern on the surface. (b) Cartoon showing the effect of different radii on the patterned surface. Small radii can fit between asperities, while large radii rest on top of the asperities

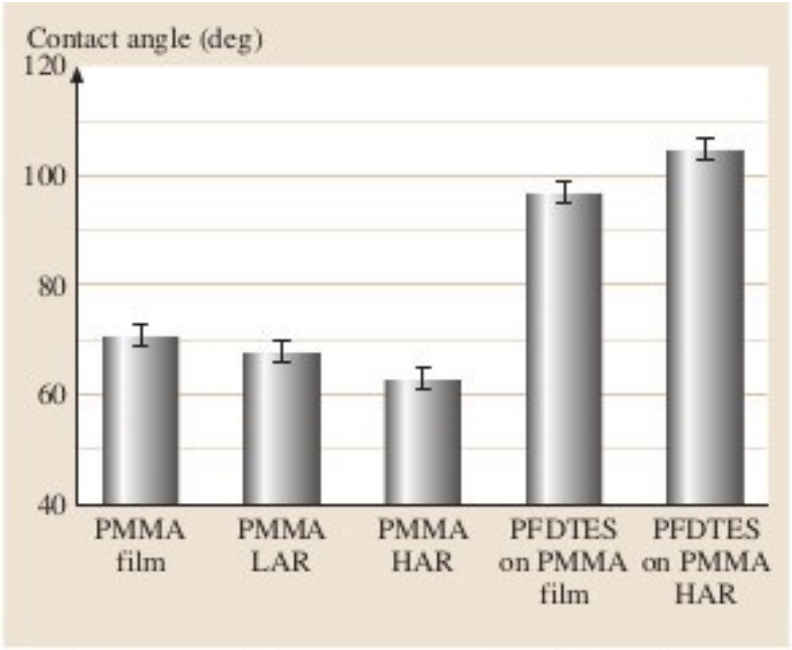


Fig. 50.41 Bar chart showing the contact angles for different materials and for different roughnesses [50.177]

Quelle: Bushan, Nanotechnology, Springer 2003

Der Lotus-Effekt

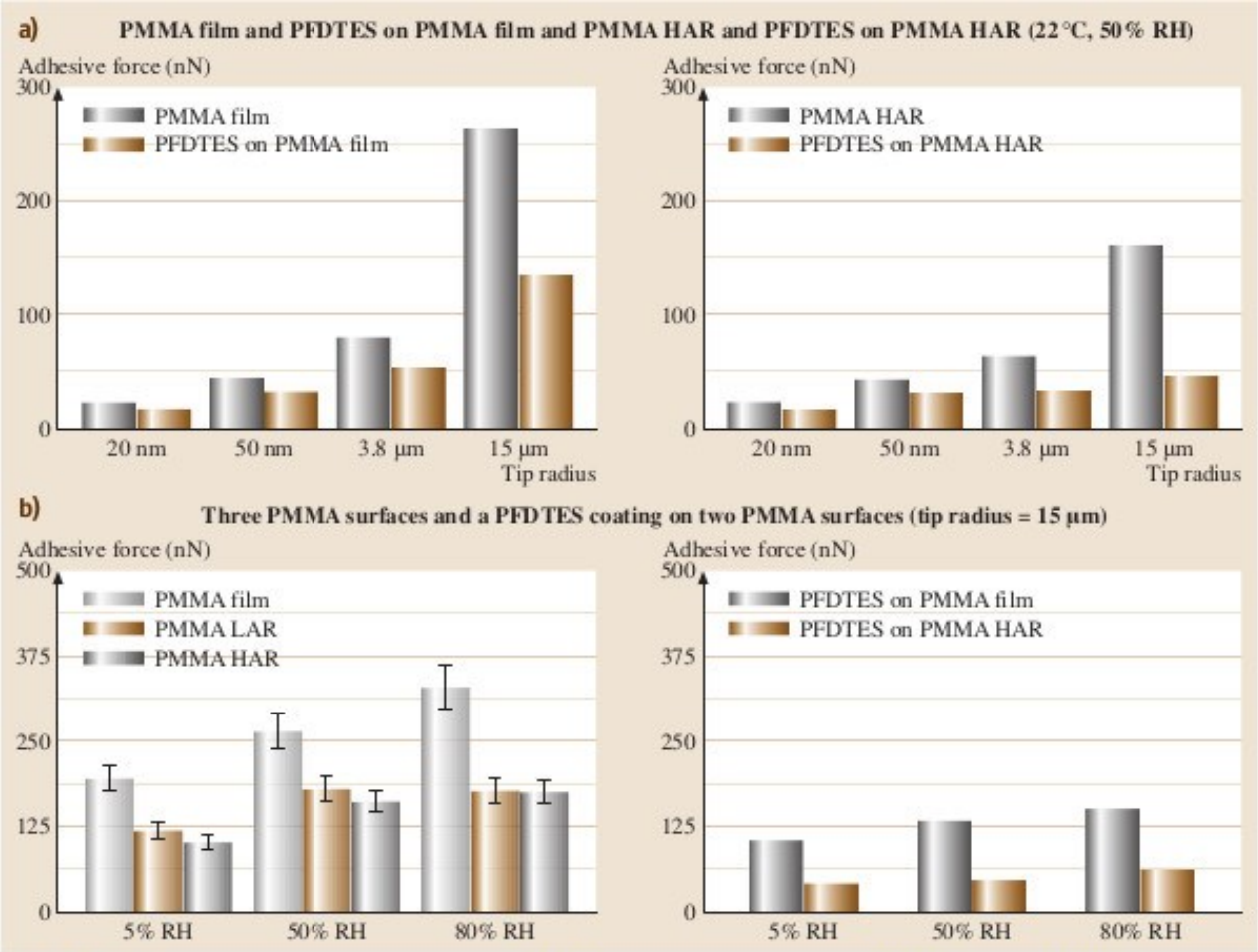


Fig. 50.42 (a) Scale-dependent adhesive force for PMMA film versus PFDTES on PMMA film and PMMA HAR versus PFDTES on PMMA HAR (*top*), and (b) the effect of relative humidity on the adhesive force for PMMA film, LAR and HAR and for PFDTES on PMMA film and HAR [50.177]

Quelle: Bushan, Nanotechnology, Springer 2003