

Collective symmetry breaking & Little Higgs

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The Simplest Little Higgs

- global $SU(3) \rightarrow SU(2)$ breaking by VEV $\Sigma_0 = (0, 0, f)^T$ in fundamental representation of $SU(3)$
- broken $SU(3)$ generators correspond to $\lambda_4 \dots \lambda_8$

$$U_{\text{NGB}} = \exp \left[\frac{i}{f} \left(\begin{array}{c|c} \eta/\sqrt{6} & H \\ \hline H^\dagger & -2\eta/\sqrt{6} \end{array} \right) \right]$$

- H identified as **SM Higgs doublet**, extra $SU(2)$ singlet η
- non-linear sigma model field

$$\Sigma = U_{\text{NGB}} \Sigma_0 = \begin{pmatrix} iH \\ f - \frac{H^\dagger H}{2f} \end{pmatrix} + \dots$$

Gauging $SU(2)_L$

Gauging upper left $SU(2)$ subgroup

$$\mathcal{L}_{\text{nl}\sigma\text{m}} = (D_\mu \Sigma)^\dagger D^\mu \Sigma = |D_\mu H|^2 \left(1 + \frac{H^\dagger H}{f^2} \right) + \dots$$

with $D_\mu H = (\partial_\mu - i \frac{g}{2} W_\mu^a \tau^a) H$

- gauging $SU(2)$ **explicitly breaks $SU(3)$**
 - $|g W_\mu h|^2$ interaction breaks GB shift symmetry
- one-loop contribution to Higgs potential



The diagram shows two Feynman diagrams representing one-loop corrections to the Higgs potential. The first diagram shows a Higgs line (dashed) entering a loop of W bosons (wavy) and exiting as a Higgs line. The second diagram shows a Higgs line (dashed) entering a loop of W bosons (wavy) and exiting as a Higgs line, with an additional Higgs line (dashed) entering the loop. The diagrams are summed and equated to the expression $\simeq \frac{3g^2}{64\pi^2} \Lambda^2 H^\dagger H$.

$$H \text{---} \bullet \text{---} H + H \text{---} \bullet \text{---} \bullet \text{---} H \simeq \frac{3g^2}{64\pi^2} \Lambda^2 H^\dagger H$$

➤ reintroduces **quadratic divergence!** ☹

Understanding the origin of the quadratic divergence

Coleman-Weinberg potential from radiative gauge boson contributions

$$V_{\text{CW}}(H) = \frac{\Lambda^2}{16\pi^2} \text{Tr}[M_{ab}^2] + \dots \quad \text{with } M_{ab}^2 = \frac{g^2}{4} \underbrace{\Sigma^\dagger \left(\frac{\tau^a \tau^b}{\text{spurion } P} \right) \Sigma}_{\text{spurion } P}$$

➤ quadratically divergent because $P \neq \mathbb{1}$

Solution

extend model to global $SU(3)_1 \times SU(3)_2 \rightarrow SU(2)_1 \times SU(2)_2$ breaking by

$$\Sigma_1 = e^{i\pi_1/(\sqrt{2}f)} \begin{pmatrix} 0 \\ 0 \\ f \end{pmatrix} \quad \Sigma_2 = e^{i\pi_2/(\sqrt{2}f)} \begin{pmatrix} 0 \\ 0 \\ f \end{pmatrix}$$

Towards collective symmetry breaking

- “double” symmetry breaking pattern to $SU(3)_1 \times SU(3)_2 \rightarrow SU(2)_1 \times SU(2)_2$ with $n\ell\sigma m$ fields Σ_1, Σ_2
- gauge diagonal subgroup $SU(3)_D$

$$\mathcal{L}_{n\ell\sigma m} = |D_\mu \Sigma_1|^2 + |D_\mu \Sigma_2|^2$$

- $n\ell\sigma m$ dynamics breaks gauged $SU(3)_D \rightarrow SU(2)_D$
- 10 GBs: 5 eaten by broken $SU(3)$ generators, 5 physical scalars $H \sim \pi_1 - \pi_2, \eta$
- quadratically divergent contribution to the Coleman-Weinberg potential

$$V_{\text{CW}} = \frac{\Lambda^2}{16\pi^2} \left(\text{Tr} \left[\Sigma_1^\dagger \Sigma_1 \right] + \text{Tr} \left[\Sigma_2^\dagger \Sigma_2 \right] \right)$$

- independent of the GBs!

More on collective breaking

Consider quadratic term: $\mathcal{L}_{\text{nl}\sigma\text{m}} \supset |gW_\mu \Sigma_1|^2 + |gW_\mu \Sigma_2|^2$

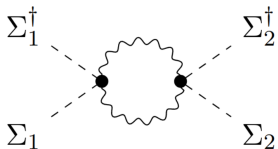
- $SU(3)_D$ gauge transformation:

$$W_\mu \rightarrow e^{i\alpha^a T^a} W_\mu e^{-i\alpha^a T^a} \Rightarrow \Sigma_1 \rightarrow e^{i\alpha^a T^a} \Sigma_1, \Sigma_2 \rightarrow e^{i\alpha^a T^a} \Sigma_2$$

➤ global $SU(3) \times SU(3)$ broken as **both** Σ_1 and Σ_2 couple to W

collective symmetry breaking

- quadratic contribution to Higgs potential must involve **both** terms



➤ logarithmically divergent

Who cancelled the divergence?

- recall *SUSY*: superpartners cancel quadratic divergence
- gauging of $SU(3)_D$ implies **extra gauge bosons**

$$\begin{pmatrix} & W^+ & X^+ \\ W^- & & Y^1 + iY^2 \\ X^- & Y^1 - iY^2 & \end{pmatrix} + W^3, A^8 \text{ on the diagonal}$$

- $\text{nl}\sigma\text{m}$ Lagrangian yields

$$\frac{g^2}{4} H^\dagger H \left[2W_\mu^+ W^{-\mu} + W_\mu^3 W^{3\mu} - X_\mu^+ X^{-\mu} - \frac{1}{2}(Y_\mu^1 Y^{1\mu} + Y_\mu^2 Y^{2\mu}) - A_\mu^8 A^{8\mu} \right]$$

- prefactors and relative signs imply cancellation of quadratic divergence by new $SU(3)/SU(2)$ gauge bosons, i. e. **particles with same spin!**

The top sector, collectively

Top quark yields dominant contribution to quadratic divergence

- introduce top partners in complete rep.s of $SU(3)_1 \times SU(3)_2$
 - formally start from $(\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{3})$
 - keep only degrees of freedom in diagonal $SU(3)_D$ (c.f. gauge bosons)

$$\Psi = \begin{pmatrix} t_L \\ b_L \\ T_L \end{pmatrix}$$

- two right-handed $SU(3)_D$ singlets: t_1, t_2
- top Yukawa

$$\mathcal{L}_{\text{top}} = \lambda_1 \bar{\Psi} \Sigma_1 t_1 + \lambda_2 \bar{\Psi} \Sigma_2 t_2$$

- **collective symmetry breaking in the top sector**

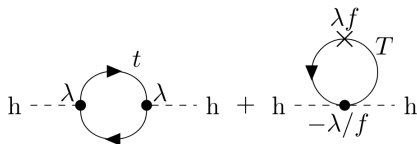
The top and its partner

- expanding $\mathcal{L}_{\text{top}}$ in H , setting $\lambda_1 = \lambda_2 = \lambda$ and rotating to the mass eigenbasis

$$\mathcal{L}_{\text{top}} = \lambda \bar{Q}_L H t_R + \lambda f \left(1 - \frac{H^\dagger H}{2f^2} \right) \bar{T}_L T_R$$

with $t_R = 1/\sqrt{2}(t_1 - t_2)$, $T_R = 1/\sqrt{2}(t_1 + t_2)$

- SM top quark t + heavy top partner T with mass λf
- cancellation of quadratic divergence



The Littlest Higgs – basics

Littlest Higgs – a simple group model

- $SU(5) \rightarrow SO(5)$ symmetry breaking triggered by VEV

$$\Sigma_0 = f \begin{pmatrix} 0 & 0 & \mathbb{1} \\ 0 & 1 & 0 \\ \mathbb{1} & 0 & 0 \end{pmatrix}$$

➤ 14 Goldstone bosons

$$U_{\text{NGB}} = \exp \left[\frac{i}{f} \begin{pmatrix} & H & \Phi \\ \hline H^\dagger & & H^T \\ \hline \Phi^* & H^* & \end{pmatrix} \right] + \text{eaten GBs}$$

H : SM Higgs doublet, Φ : scalar $SU(2)$ triplet

- $SU(5)/SO(5)$ coset parametrised by $\Sigma = U_{\text{NGB}} \Sigma_0 U_{\text{NGB}}^\dagger$

The Littlest Higgs – gauge sector

- gauged $[SU(2) \times U(1)]^2$ subgroup with generators

$$Q_1^a = \begin{pmatrix} \sigma^a/2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad Q_2^a = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -\sigma^{a*}/2 \end{pmatrix}$$

$$Y_1 = \frac{1}{10} \text{diag}(3, 3, -2, -2, -2) \quad Y_2 = \frac{1}{10} \text{diag}(2, 2, 2, -3, -3)$$

- explicitly breaks global symmetry $\Rightarrow$ GBs acquire mass
- symmetry broken collectively – each $SU(2) \times U(1)$ factor preserves global $SU(3)$
- global $SU(5) \rightarrow SO(5)$ breaking induces

$$[SU(2) \times U(1)]^2 \rightarrow SU(2) \times U(1)$$

$\Rightarrow W^\pm, Z$, photon partners at the scale f

The Littlest Higgs – top sector

Analogously to Simplest Little Higgs

- left-handed third generation doublet Q_L embedded into $SU(3)$ triplet

$$\bar{\Psi} = (\bar{t}_L, \bar{b}_L, \bar{T}_L) \equiv (\bar{Q}_L, \bar{T}_L)$$

- explicit breaking of $SU(5)$ global symmetry
- two right-handed $SU(3)$ singlets t_1, t_2
- top Yukawa ($i, j, k = 1, 2, 3; m, n = 4, 5$)

$$\mathcal{L}_{\text{top}} = -\frac{\lambda_1 f}{2} \bar{\Psi}_i \epsilon_{ijk} \epsilon_{mn} \Sigma_{jm} \Sigma_{kn} t_1 - \lambda_2 f \bar{T}_L t_2$$

- each term preserves one $SU(3)$ subgroup – collective breaking

The Littlest Higgs – Higgs potential

Higgs potential generated radiatively – Coleman-Weinberg potential

$$V_{\text{CW}}^{\text{quad}} = a \frac{\Lambda^2}{16\pi^2} f^2 \sum_{j=1,2} \left(g_j^2 \text{Tr} [Q_j^a \Sigma Q_j^{a*} \Sigma^*] + g_j'^2 \text{Tr} [Y_j \Sigma Y_j^* \Sigma^*] \right)$$

g_j, g_j' : $[SU(2) \times U(1)]_j$ gauge couplings

a : $\mathcal{O}(1)$ coefficient

➤ expanding in H and Φ

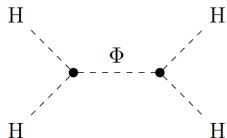
$$V_{\text{CW}}^{\text{quad}}(H, \Phi) = af^2 \left[(g_1 + g_1')^2 \left| \Phi_{ij} + \frac{i}{4f} (H_i H_j + H_j H_i) \right|^2 + (g_2 + g_2')^2 \left| \Phi_{ij} - \frac{i}{4f} (H_i H_j + H_j H_i) \right|^2 \right]$$

The Littlest Higgs – an $\mathcal{O}(1)$ quartic

Features of $V_{\text{CW}}^{\text{quad}}(H, \Phi)$

- quadratic divergence for H cancels as expected
- quadratically divergent mass term for Φ and $H\Phi\Phi$ coupling remain
- integrating out Φ below $M_\Phi \sim af$

$$\lambda = a \frac{(g_1^2 + g_1'^2)(g_2^2 + g_2'^2)}{g_1^2 + g_1'^2 + g_2^2 + g_2'^2}$$



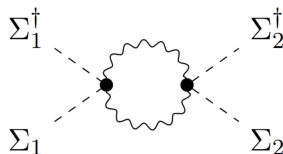
➤ $\mathcal{O}(1)$ Higgs quartic generated at tree level

- similar Higgs quartic from fermion sector
- interesting prediction: $m_h = \sqrt{2\lambda}v \sim \sqrt{2}g_{\text{SM}}v \sim (200 - 300) \text{ GeV}$
 - some tuning required to reconcile model with measured Higgs mass

Summary

Study goal: Little Higgs models

- Simplest Little Higgs
- collective symmetry breaking
- top and gauge partners
- Littlest Higgs model



Reading assignment

- chapter 3–3.3.1 of C. Csaki, S. Lombardo, O. Telem, *TASI Lectures on Non-Supersymmetric BSM Models*, arXiv:1811.04279
- chapter 1–3 of M. Perelstein, *Little Higgs Models and Their Phenomenology*, arXiv:hep-ph/0512128