

## Problem Set 13 Nonlinear Optics (NLO)

Due: 15. July 2015

## 1) Nonlinear Schrödinger Equation

The nonlinear Schrödinger-equation (NLSE) of a fiber was derived in the lecture:

$$\frac{\partial}{\partial z}\underline{A}(z,t) + \beta_c^{(1)}\frac{\partial}{\partial t}\underline{A}(z,t) - \frac{1}{2}j\beta_c^{(2)}\frac{\partial^2}{\partial t^2}\underline{A}(z,t) = -\frac{\alpha}{2}\underline{A}(z,t) - j\gamma |\underline{A}(z,t)|^2\underline{A}(z,t)$$
(1.1)

- 1. Explain the parameters  $oldsymbol{eta}_c^{(1)}, \, oldsymbol{eta}_c^{(2)}, \, lpha \,$  and  $\gamma$  .
- 2. For optical fibers, the parameter  $D = \mathrm{d}\beta_c^{(1)}/\mathrm{d}\lambda$  is usually specified instead of  $\beta_c^{(2)}$ . What is the connection between D and  $\beta_c^{(2)}$ ? A typical single-mode fiber (SSMF) has  $D = 18 \, \mathrm{ps/(nm \cdot km)}$  at the wavelength of  $\lambda = 1.55 \, \mu \mathrm{m}$ . Calculate  $\beta_c^{(2)}$  and explain the meaning of D.
- 3. Consider the new coordinate system t', z' generated by the following transformation as well as the new function  $\underline{A}'$ :

$$t' = t - \beta_c^{(1)} z$$
$$z' = z$$
$$\underline{A}'(z', t') = \underline{A}(z, t)$$

Imagine for the moment that  $\underline{A}(z,t)$  represents a pulse moving along z with velocity  $1/\beta_c^{(1)}$ . Sketch the functions  $\underline{A}(z,t)$  and  $\underline{A}'(z',t')$  as a function of t and t' for two different positions of z and z'. Explain why the (z', t') coordinate system is usually referred to as a retarded time frame.

- 4. Find a formulation of the NLSE for  $\underline{A}'$ . Notice that the term  $\beta_c^{(1)} \frac{\partial}{\partial t'} \underline{A}'(z',t')$  does no longer appear in the differential equation. In the following, we will omit the primes keeping in mind that the time dependence is given with respect to a retarded reference frame.
- 5. We will now assume that there are no losses  $(\alpha = 0)$  and search for solutions describing fundamental solitons, i.e., waveforms which do not change their shape as they propagate along z. We therefore require the magnitude of the complex amplitude  $\underline{A}'(z,t)$  to be independent of z, but still allow for a z-dependent phase shift. Substitute the ansatz  $\underline{A}'(z,t) = A_0(t) \exp(-jKz)$  in the NLSE. Assuming further that  $A_0(t)$  is a real-valued function, show that the following differential equation holds for  $A_0(t)$ :

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$$\frac{1}{2}\beta_c^{(2)} \frac{1}{A_0} \frac{\partial^2 A_0}{\partial t^2} - \gamma A_0^2 = -K . \tag{1.2}$$

6. Show that  $A_0(t) = A_1 \operatorname{sech}\left(\frac{t}{T}\right) = A_1/\cosh\left(\frac{t}{T}\right)$  is a solution for the differential equation (1.2). Remember that  $\cosh^2 - \sinh^2 = 1$ , and that the derivative of  $\sinh(x)$  is  $\cosh(x)$  and vice versa. Show in particular that  $A_1$  and T must fulfill the following relations:

$$K = \frac{1}{2}\gamma A_1^2 \tag{1.3}$$

$$A_{\rm l}^2 = -\frac{\beta_c^{(2)}}{\gamma T^2} \ . \tag{1.4}$$

7. Substituting the soliton ansatz and Eqs. (1.3) and (1.4) in Eq. (1.2), we can derive the following relation:

$$\frac{1}{2}\beta_c^{(2)}\frac{\partial^2}{\partial t^2}\underline{A}'(z,t) = \gamma \left|\underline{A}'(z,t)\right|^2\underline{A}'(z,t) + \text{const.}$$
(1.5)

How can this relation be interpreted taking into account the interplay of dispersion and self-phase modulation? If a pulse gets shorter, do you expect that it must have a larger or smaller peak intensity for building a soliton? Check your answer with the help of Eq. (1.4).

## **Questions and Comments:**

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