

## Problem Set 13

### Nonlinear Optics (NLO)

Due: July 18, 2018, 09:45 AM

#### 1) Nonlinear Schrödinger Equation

The nonlinear Schrödinger-equation (NLSE) of an optical fiber was derived in the lecture:

$$\frac{\partial}{\partial z} \underline{A}(z, t) + \beta_c^{(1)} \frac{\partial}{\partial t} \underline{A}(z, t) - \frac{1}{2} j \beta_c^{(2)} \frac{\partial^2}{\partial t^2} \underline{A}(z, t) = -\frac{\alpha}{2} \underline{A}(z, t) - j \gamma |\underline{A}(z, t)|^2 \underline{A}(z, t) \quad (1.1)$$

1. Explain the parameters  $\beta_c^{(1)}$ ,  $\beta_c^{(2)}$ ,  $\alpha$  and  $\gamma$ .
2. For optical fibers, the parameter  $D = d\beta_c^{(1)} / d\lambda$  is usually specified instead of  $\beta_c^{(2)}$ . What is the connection between  $D$  and  $\beta_c^{(2)}$ ? A standard single-mode fiber (SSMF) has  $D = 18 \text{ ps}/(\text{nm} \cdot \text{km})$  at the vacuum wavelength of  $\lambda = 1.55 \text{ } \mu\text{m}$ . Calculate  $\beta_c^{(2)}$  and explain the meaning of  $D$ .
3. Consider the new coordinate system  $t'$ ,  $z'$  generated by the following transformation as well as the new function  $\underline{A}'$ :

$$\begin{aligned} t' &= t - \beta_c^{(1)} z \\ z' &= z \\ \underline{A}'(z', t') &= \underline{A}(z, t) \end{aligned}$$

Imagine for the moment that  $\underline{A}(z, t)$  represents a pulse moving along  $z$  with velocity  $1/\beta_c^{(1)}$ . Sketch the functions  $\underline{A}(z, t)$  and  $\underline{A}'(z', t')$  as a function of  $t$  and  $t'$  for two different positions of  $z$  and  $z'$ . Explain why the  $(z', t')$  coordinate system is usually referred to as a retarded time frame.

4. Find a formulation of the NLSE for  $\underline{A}'$ . Notice that the term  $\beta_c^{(1)} \frac{\partial}{\partial t'} \underline{A}'(z', t')$  does no longer appear in the differential equation. In the following, we will omit the primes keeping in mind that the time dependence is given with respect to a retarded reference time frame.
5. We will now assume that there are no losses ( $\alpha = 0$ ) and search for solutions describing fundamental solitons, i.e., waveforms which do not change their shape as they propagate along  $z$ . We therefore require the magnitude of the complex amplitude  $\underline{A}'(z, t)$  to be independent of  $z$ , but still allow for a  $z$ -dependent phase shift. Substitute the ansatz  $\underline{A}'(z, t) = A_0(t) \exp(-jKz)$  in the NLSE. Assuming further that  $A_0(t)$  is a real function, show that the following differential equation holds for  $A_0(t)$ :

$$\frac{1}{2} \beta_c^{(2)} \frac{1}{A_0(t)} \frac{\partial^2 A_0(t)}{\partial t^2} - \gamma A_0^2(t) = -K. \quad (1.2)$$

6. Show that  $A_0(t) = A_1 \operatorname{sech}(t/T) = \frac{A_1}{\cosh(t/T)}$  is a valid solution ansatz for the differential equation (1.2). Remember that  $\cosh^2 - \sinh^2 = 1$ , and that the derivative of  $\sinh(x)$  is  $\cosh(x)$  and vice versa. Show in particular that the pulse amplitude  $A_1$  and the pulse duration  $T$  must fulfill the following relations:

$$K = \frac{1}{2} \gamma A_1^2 \quad (1.3)$$

$$A_1^2 = -\frac{\beta_c^{(2)}}{\gamma T^2}. \quad (1.4)$$

7. Eq. 1.2 can be reformulated as:

$$\frac{1}{2} \beta_c^{(2)} \frac{\partial^2 A_0(t)}{\partial t^2} = (\gamma A_0^2(t) - K) A_0(t). \quad (1.5)$$

How can this relation be interpreted taking into account the interplay of dispersion and self-phase modulation? If a pulse gets shorter, do you expect that it must have a larger or smaller peak intensity for building a soliton? Check your answer with the help of Eq. (1.4).

### Questions and Comments:

Pablo Marin-Palomo

Building: 30.10, Room: 2.33

Phone: 0721/608-42487

Philipp Trocha

Building: 30.10, Room: 2.32-2

Phone: 0721/608-42480

[nlo@ipq.kit.edu](mailto:nlo@ipq.kit.edu)