



High- p_T Physics

**Fakultät für Physik
K. Rabbertz (ETP)**





Summary

- The new world average from PDG on $\alpha_s(M_z^2)$ is 0.1179 ± 0.0010
- Elastic scattering of electrons on protons:
 - “Simplest” elastic approximation: Rutherford scattering
 - Has one independent variable, e.g. θ or Q^2
 - Mott scattering includes energy transfer from electron to nucleon and spin-1/2
 - Rosenbluth formula introduces internal structure of the proton; electric and magnetic form factors
- Inelastic scattering requires two independent variables, e.g. scaling variable x and momentum transfer squared Q^2
- Can be converted to other combination with inelasticity y and hadron final state mass squared W^2
- Lorentz-invariant description involves structure functions $F_1(x, Q^2)$ and $F_2(x, Q^2)$
- Depend mostly on x , not Q^2 : Spin-1/2 partons lead to Callan-Gross relation

$$F_1(x) = \frac{F_2(x)}{2x}$$



Rosenbluth formula can be rewritten to include inelastic scattering.

Most general Lorentz-invariant and parity conserving expression:

$$\frac{d^2\sigma}{dx dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[(1-y) \frac{F_2(x, Q^2)}{x} + y^2 F_1(x, Q^2) \right] \quad Q^2 \gg M_p^2 y^2$$

$F_1(x, Q^2)$ and $F_2(x, Q^2)$ are structure functions incorporating the form factors (and kinematic ones, τ), but cannot be related to Fourier transforms any more since dependent on x .

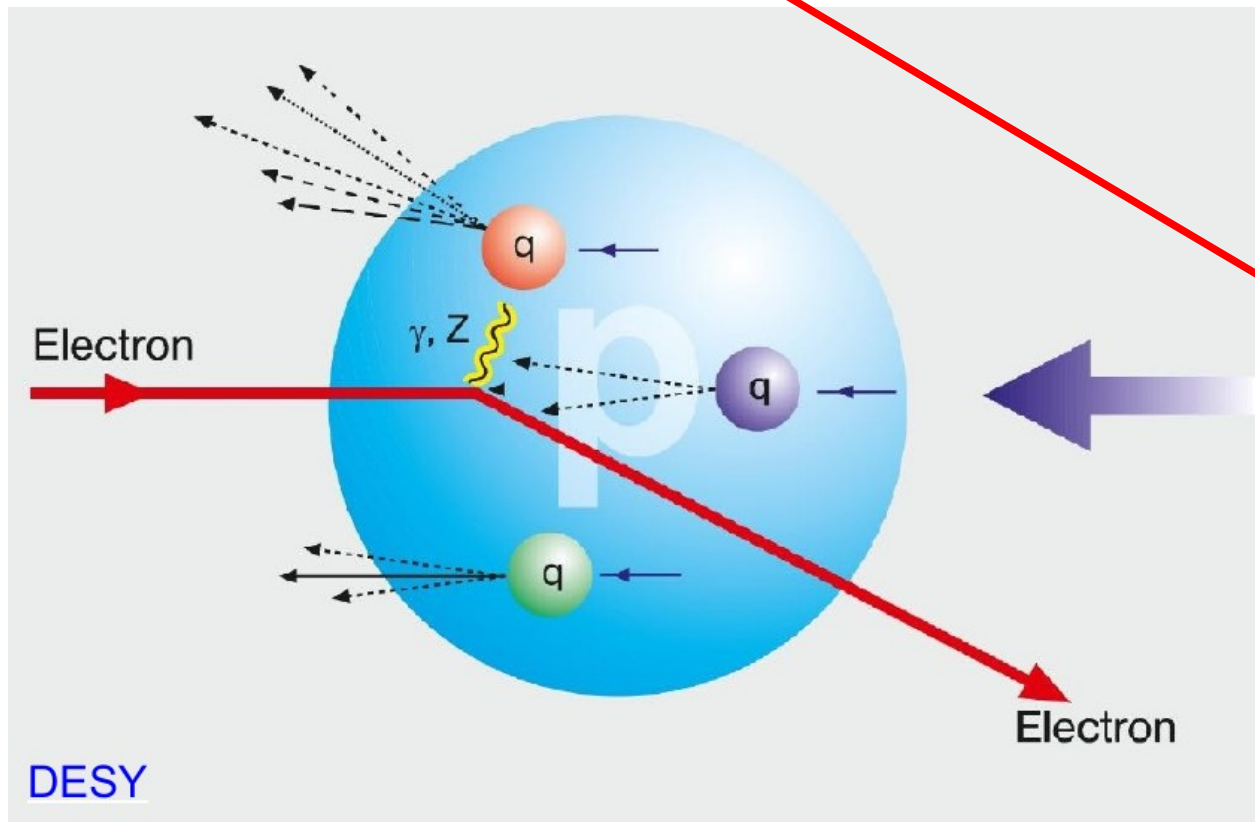
Still, $F_1(x, Q^2)$ is of purely magnetic origin, while $F_2(x, Q^2)$ originates from both, electric and magnetic effects.

What do they mean?

J.D. Bjorken, R.P. Feynman 1969:

• Infinite momentum frame

- incoherent superposition of elastic scatterings with point-like “partons”
- scale invariant, i.e. independent of resolution $\sim q^2$, no natural length scale
- partons have spin 1/2



Bjorken scaling:

$$F_1(x, Q^2) \rightarrow F_1(x)$$

$$F_2(x, Q^2) \rightarrow F_2(x)$$

Callan-Gross relation:

$$F_1(x) = \frac{F_2(x)}{2x}$$

Spin 0 would give:

$$F_1(x) = 0$$

Modern writing:

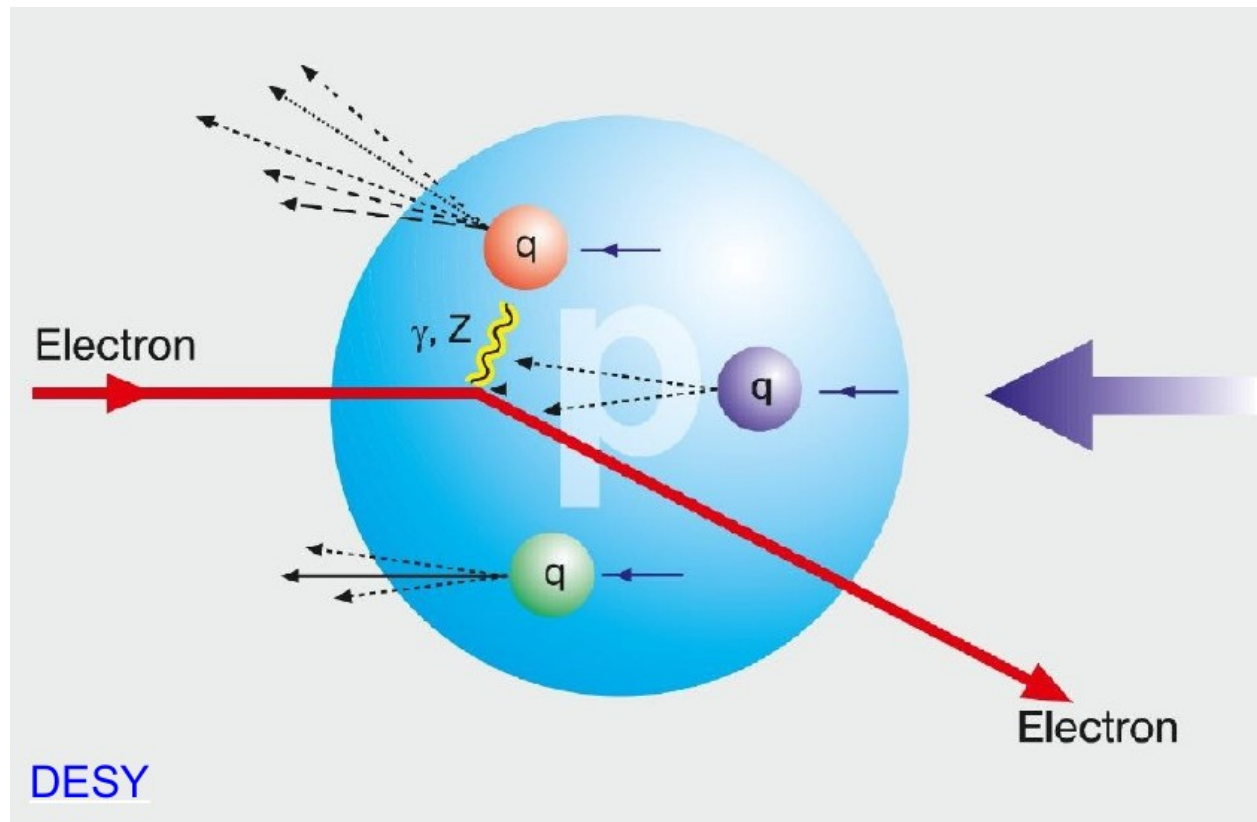
$$F_2(x) = x \sum_i e_i^2 [q_i(x) + \bar{q}_i(x)]$$

quark charges

anti-quark momentum distribution

quark momentum distribution

q_i : parton distribution functions (PDFs)



Bjorken scaling:

$$F_1(x, Q^2) \rightarrow F_1(x)$$

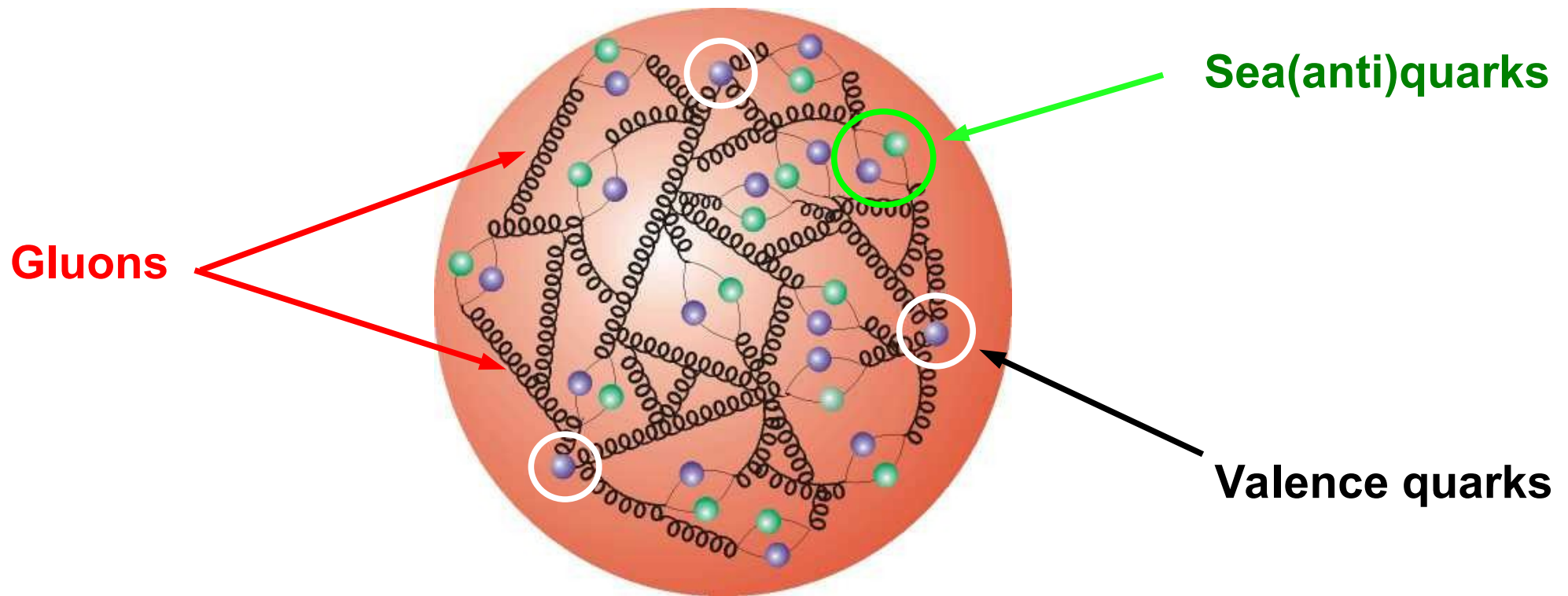
$$F_2(x, Q^2) \rightarrow F_2(x)$$

Callan-Gross relation:

$$F_1(x) = \frac{F_2(x)}{2x}$$

Spin 0 would give:

$$F_1(x) = 0$$



- **Example of generic functional form of proton structure:**

$$xf(x) = Ax^B(1-x)^C(1+Dx+Ex^2)$$

Normalisation

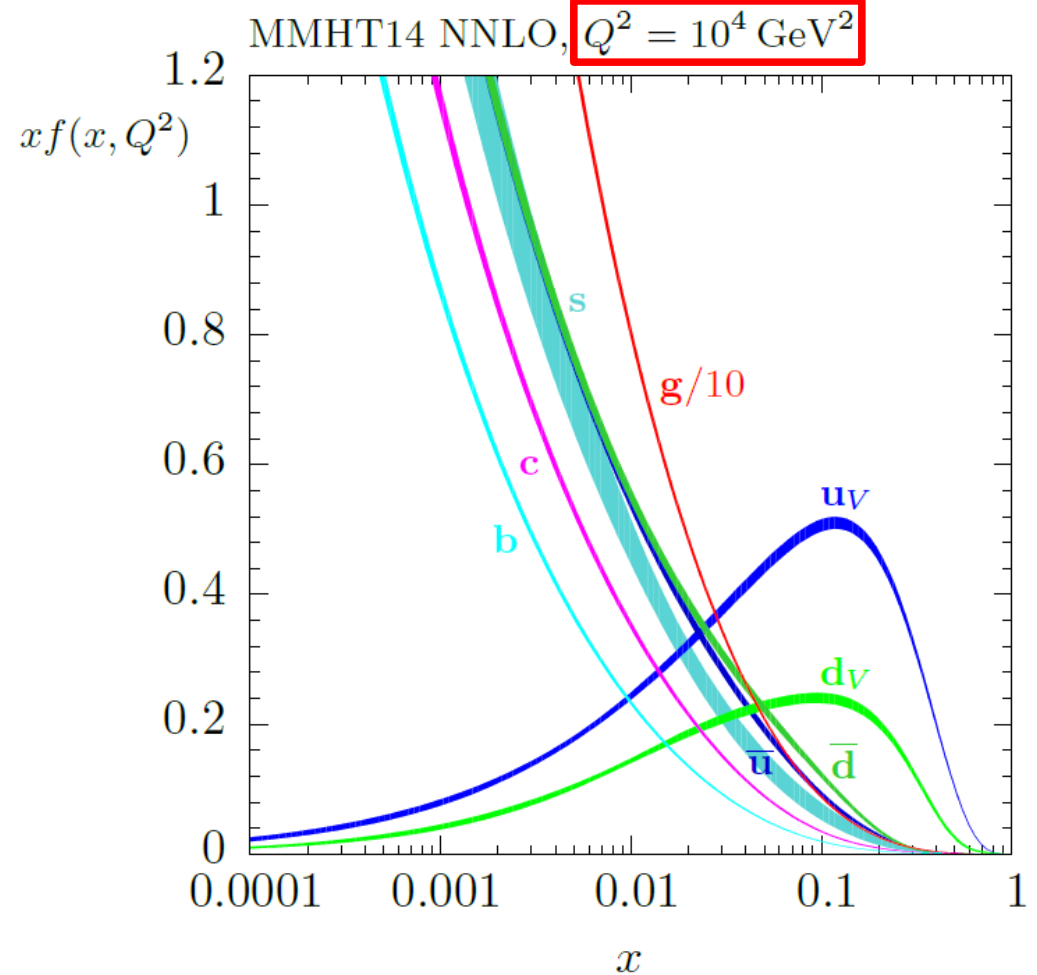
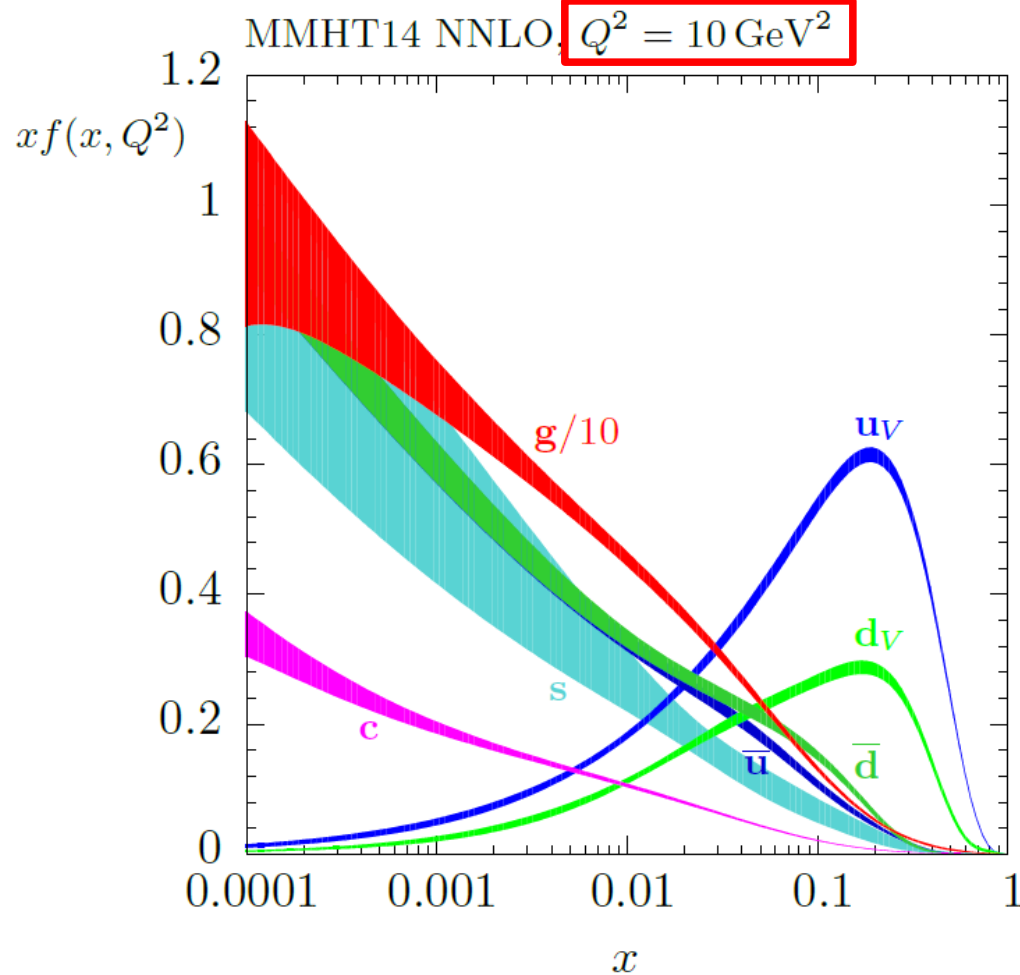
Behaviour for
 $x \rightarrow 0$

Behaviour for
 $x \rightarrow 1$

Middle region
largest variability

MMHT 2014 PDFs

Resolution scale

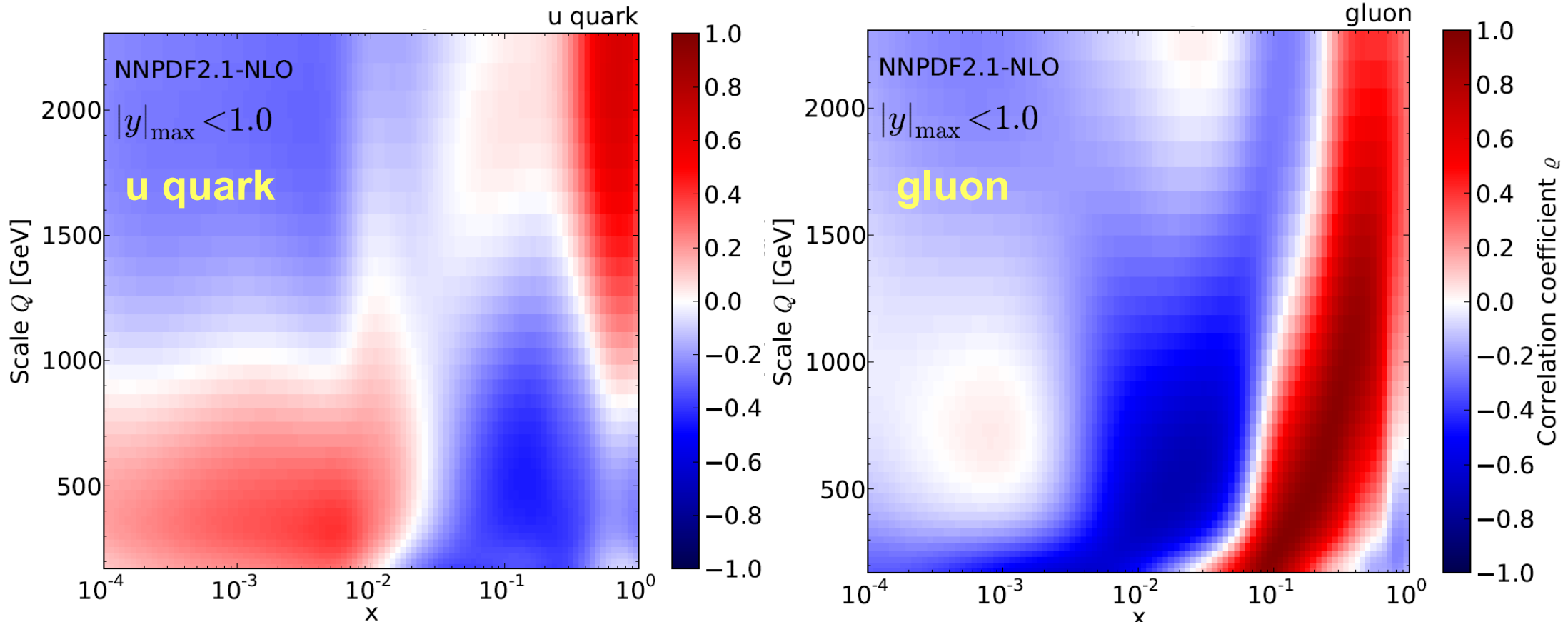


R. Thorne, 13.04.2015
PDF4LHC Meeting

Example: Measurement of high pT jets and gluon content of the proton:

- Gluon distribution at high x (> 0.1)
- Quark distribution at high x (> 0.3)

$$\rho_f(x, Q) = \frac{N}{(N-1)} \frac{\langle \sigma_{\text{jet}}(Q)_i \cdot x f(x, Q^2)_i \rangle - \langle \sigma_{\text{jet}}(Q)_i \rangle \cdot \langle x f(x, Q^2)_i \rangle}{\Delta_{\sigma_{\text{jet}}(Q)} \Delta_{x f(x, Q^2)}}.$$





Hadron Colliders

Tevatron: 1986 – 2011

Collisions of **p anti-p**

Run II: $E_{\text{cms}} = 1.96 \text{ TeV}$

Run II: Record luminosity: $4.3 \times 10^{32} \text{ cm}^{-2}\text{s}^{-1}$

LHC: 2009 – present

Collisions of **p-p, Pb-Pb, and p-Pb**

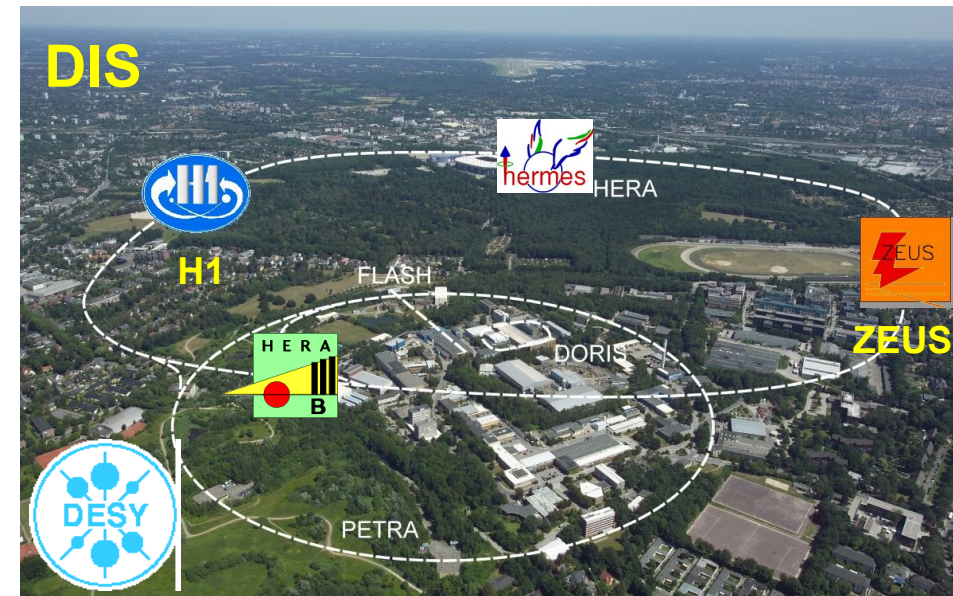
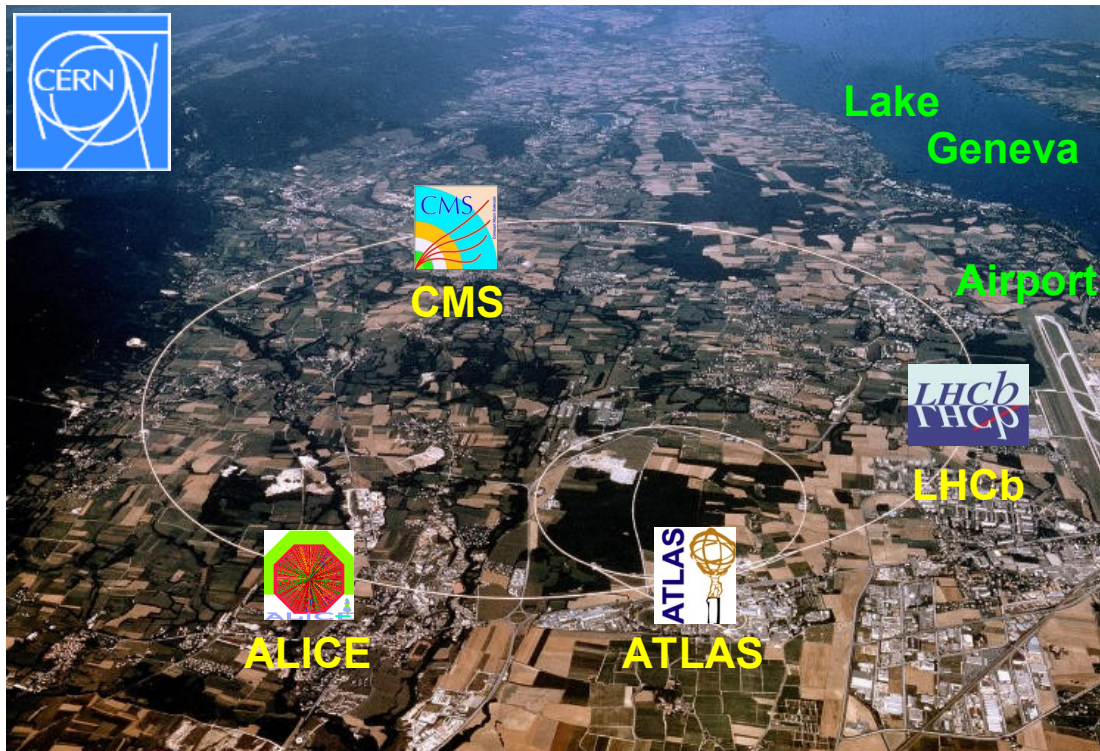
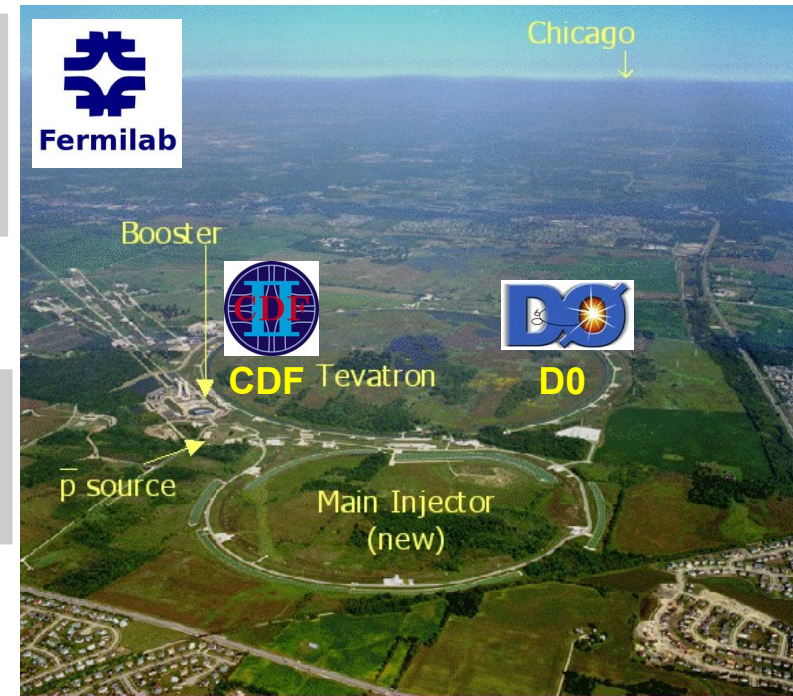
$E_{\text{cms}} = 0.9, 2.36, 2.76, 5.02, 7, 8, 13 \text{ TeV}$

Peak inst. Luminosity: $\sim 2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$

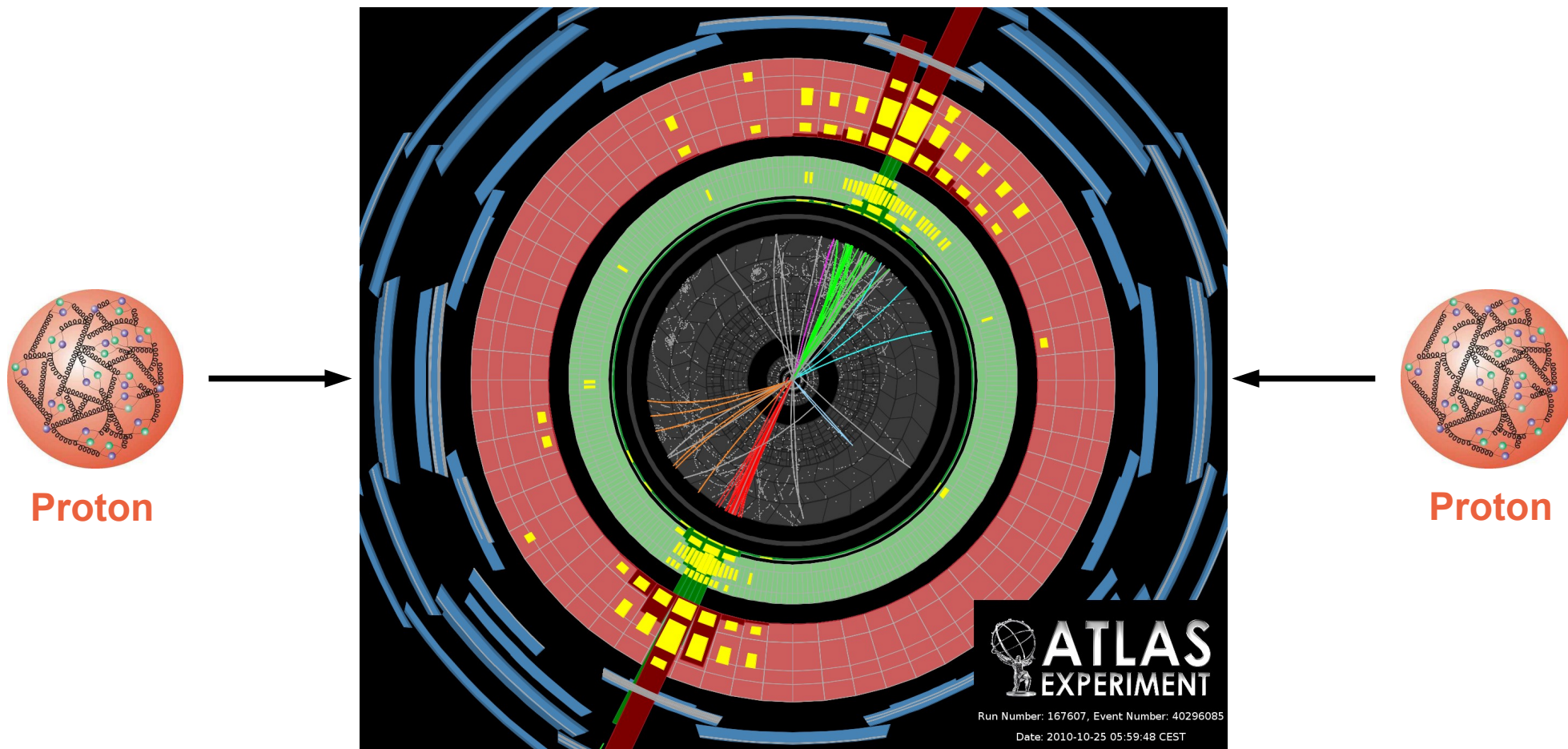
HERA: 1992 – 2007

Collisions of **e^+p , e^-p**

HERA II: $E_{\text{cms}} = 319 \text{ GeV}$



“Broadband beams” of various parton types with various energies
→ QCD parton collider!

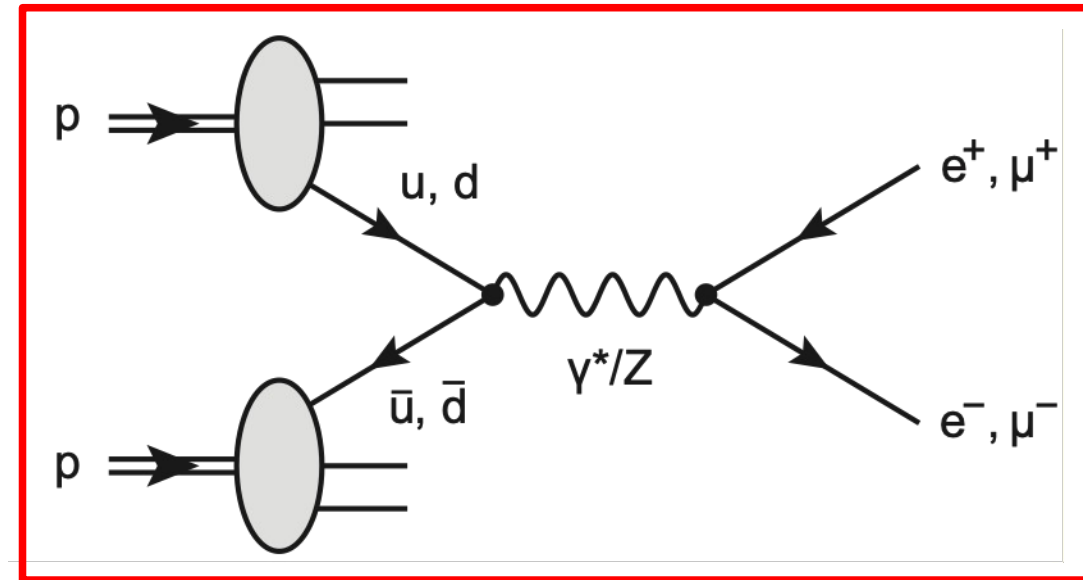


Challenge: Reliable calculations of observables

$$pp \rightarrow l^+ l^- + X$$

• Hadro-production of lepton pairs

- at large center-of-mass energies
- with large invariant mass
- color-neutral final state (except proton remnants) → no hadronisation

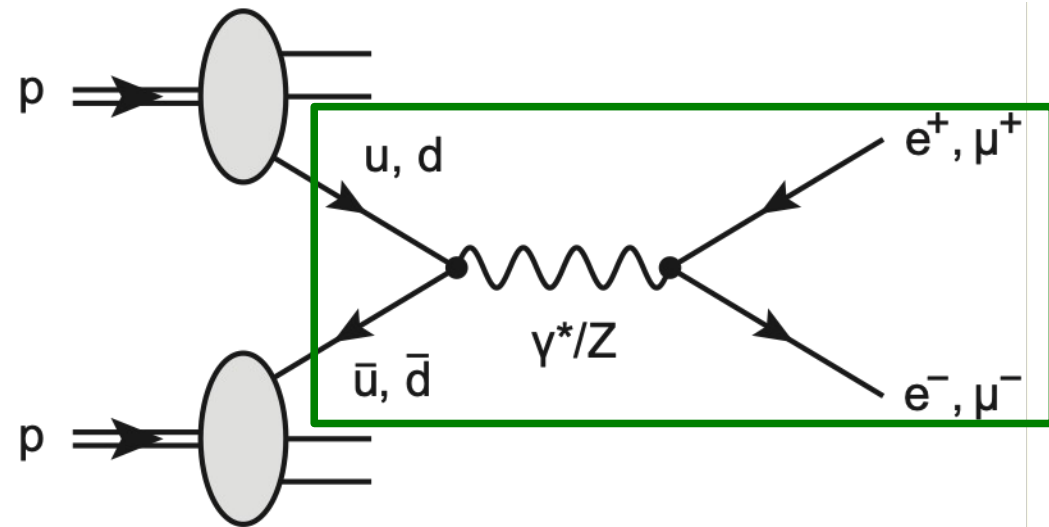


Not a Feynman diagram

$$pp \rightarrow l^+ l^- + X$$

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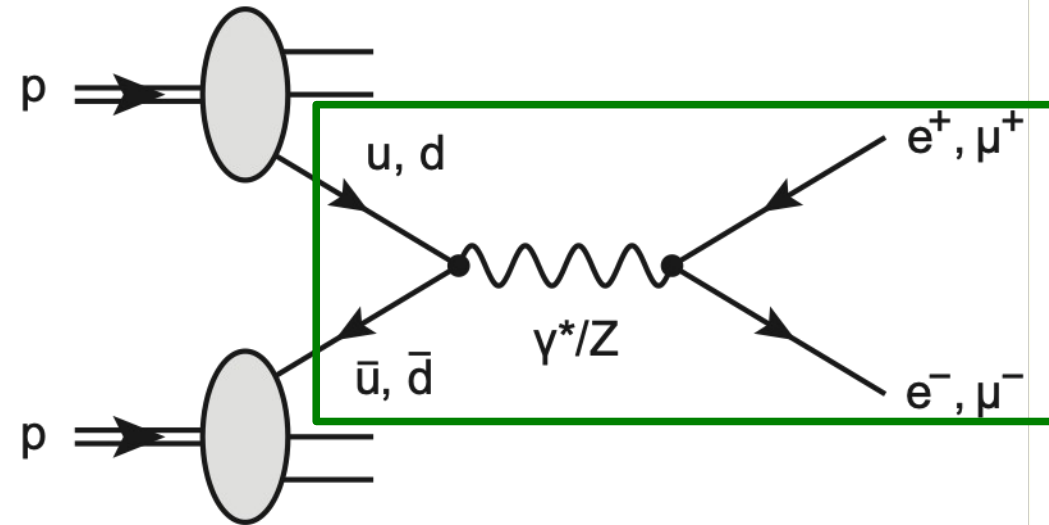


Partonic Feynman diagram
→ calculable in perturbative QCD

$$pp \rightarrow l^+ l^- + X$$

Hadro-production of lepton pairs

- at large center-of-mass energies
- with large invariant mass
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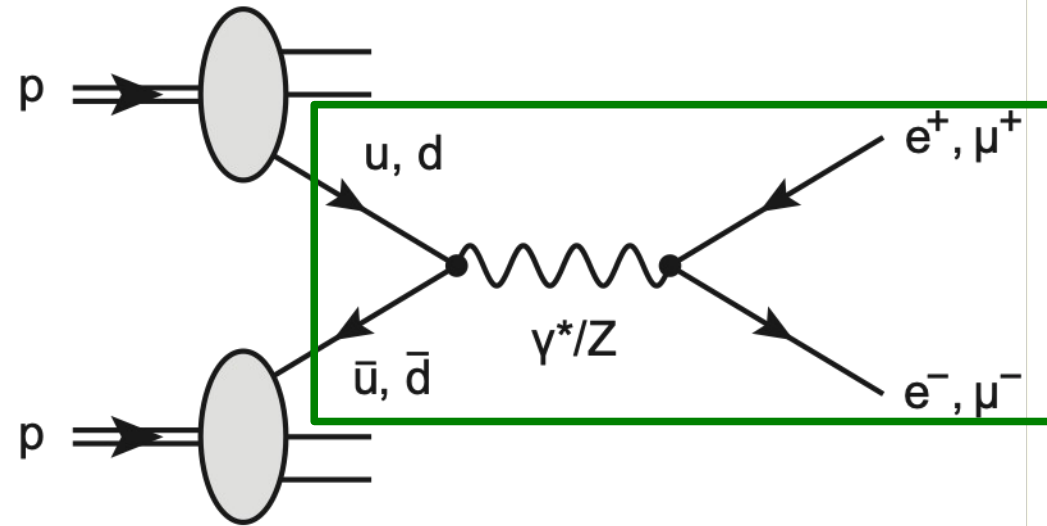
Factorisation theorem of QCD:

- Process can be calculated by factorising “hard” and “soft” components
 - ➔ Calculate hard partonic subprocess
 - ➔ Weight cross section with probability to find partons with momenta x_1, x_2 inside hadrons
 - ➔ Integrate over all possible parton momenta
 - ➔ Sum over all possible parton flavors

$$pp \rightarrow l^+ l^- + X$$

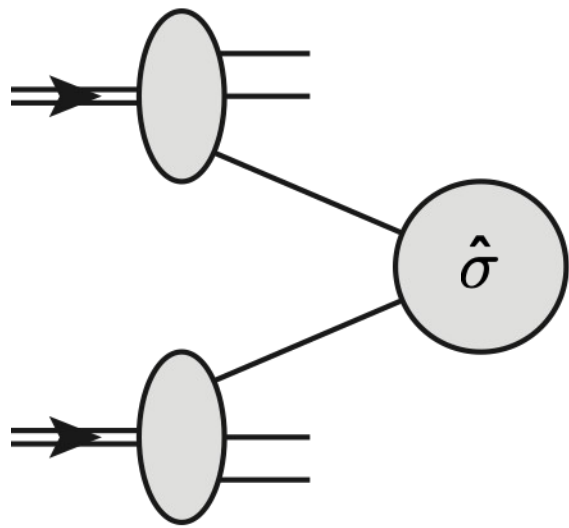
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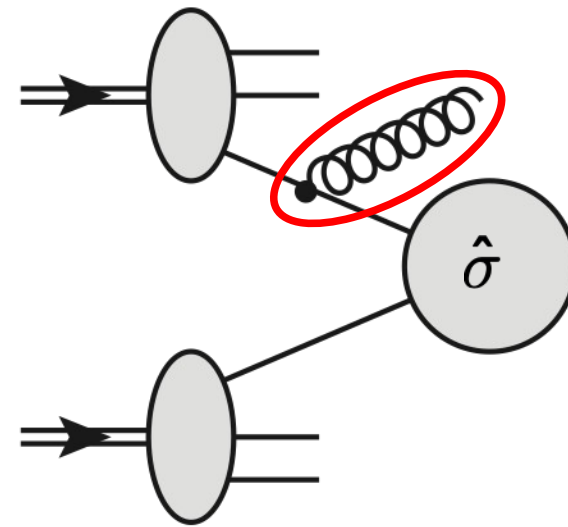


$$\sigma_{\text{DY}} = \sum_{i,j} \int dx_i dx_j f_i(x_i) f_j(x_j) \cdot \hat{\sigma}(q_i q_j \rightarrow l^+ l^-)$$

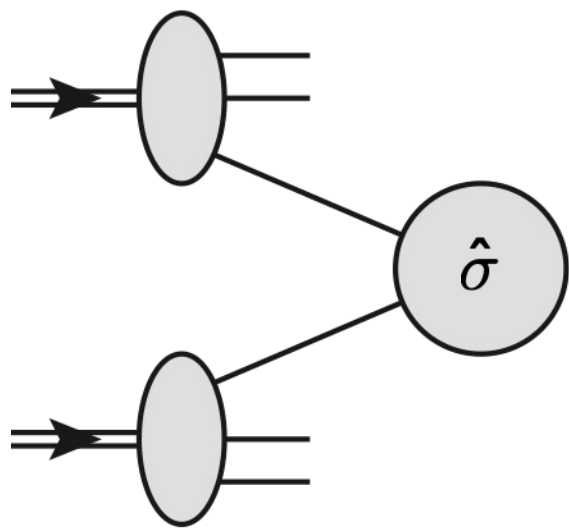
PDFs $f_i(x_i)$ are universal; can be measured independently e.g. in DIS!



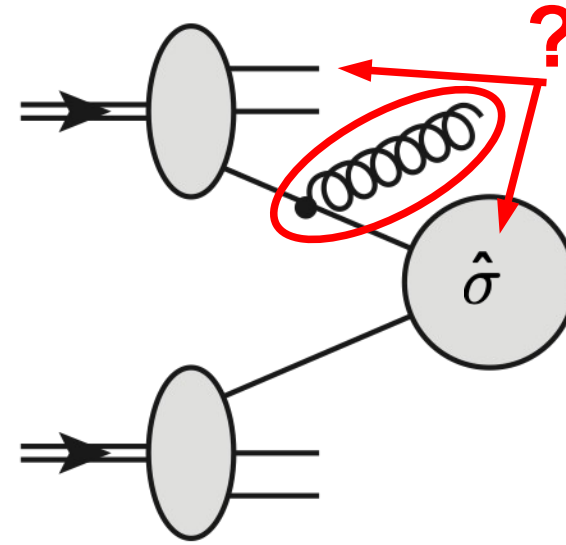
higher orders



Where does
this belong to?



higher orders



Where does
this belong to?

The PDF?

Or the parton
process?

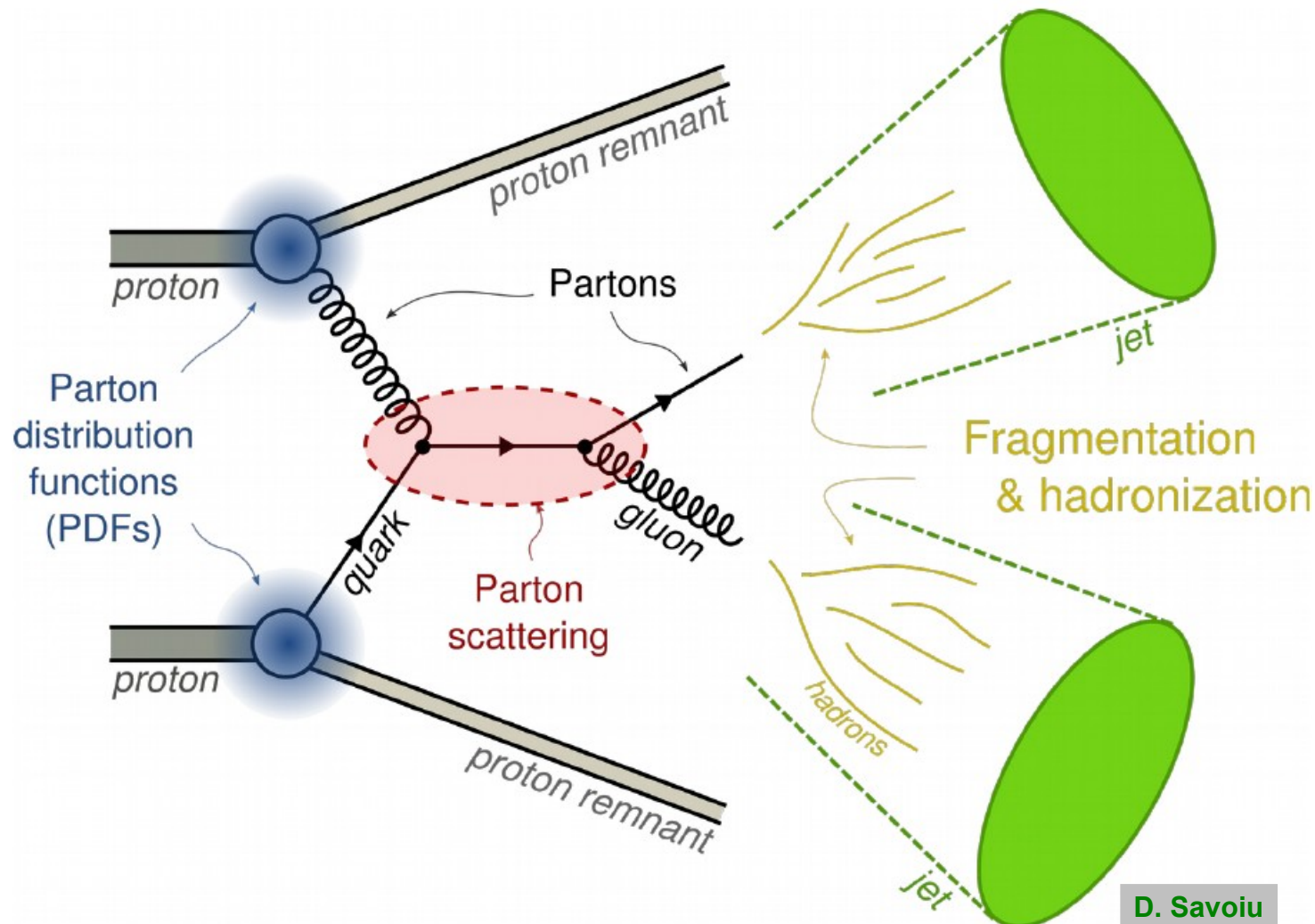
Attribution ambiguous:

- ➔ Leads to soft and/or collinear divergences (long-distance effects!)
- ➔ Solution: Introduce a new scale to separate short- and long-distance effects
 - ➔ Factorisation scale μ_f
 - ➔ All soft and collinear divergences (long-distance effects) are absorbed into the PDFs determined from experimental measurements

$$f_i(x_i) \rightarrow f_i(x_i, \mu_f^2)$$

$$\hat{\sigma}_{ij}(x_i, x_j, \mu_r^2, \alpha_s(\mu_r^2)) \rightarrow \hat{\sigma}_{ij}(x_i, x_j, \mu_f^2, \mu_r^2, \alpha_s(\mu_r^2))$$

Factorisation valid also for more general final states, e.g. jet production!



D. Savoiu



Event rates at the LHC

Total cross section

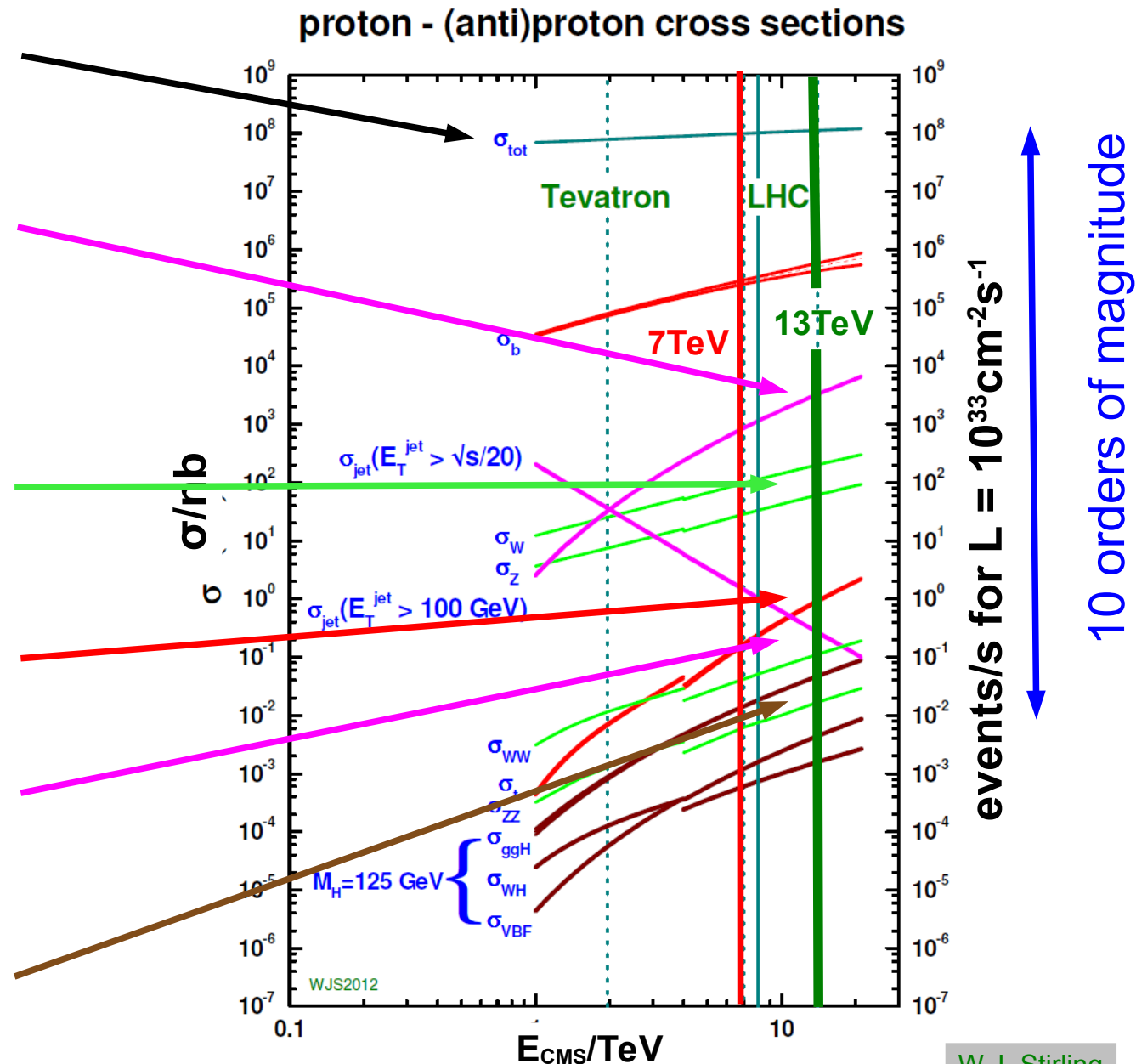
Jets: $\sigma_{\text{jet}}(E_T^{\text{jet}} > 100\text{GeV})$
 $\sim 2000 / \text{s}$

W & Z bosons: σ_W, σ_Z
 $\sim 200 / \text{s}, 50 / \text{s}$

Top quarks (σ_{tt})
 $\sim 1 / \text{s}$

Jets: $\sigma_{\text{jet}}(E_T^{\text{jet}} > 650\text{GeV})$
 $\sim 18 / \text{min}$

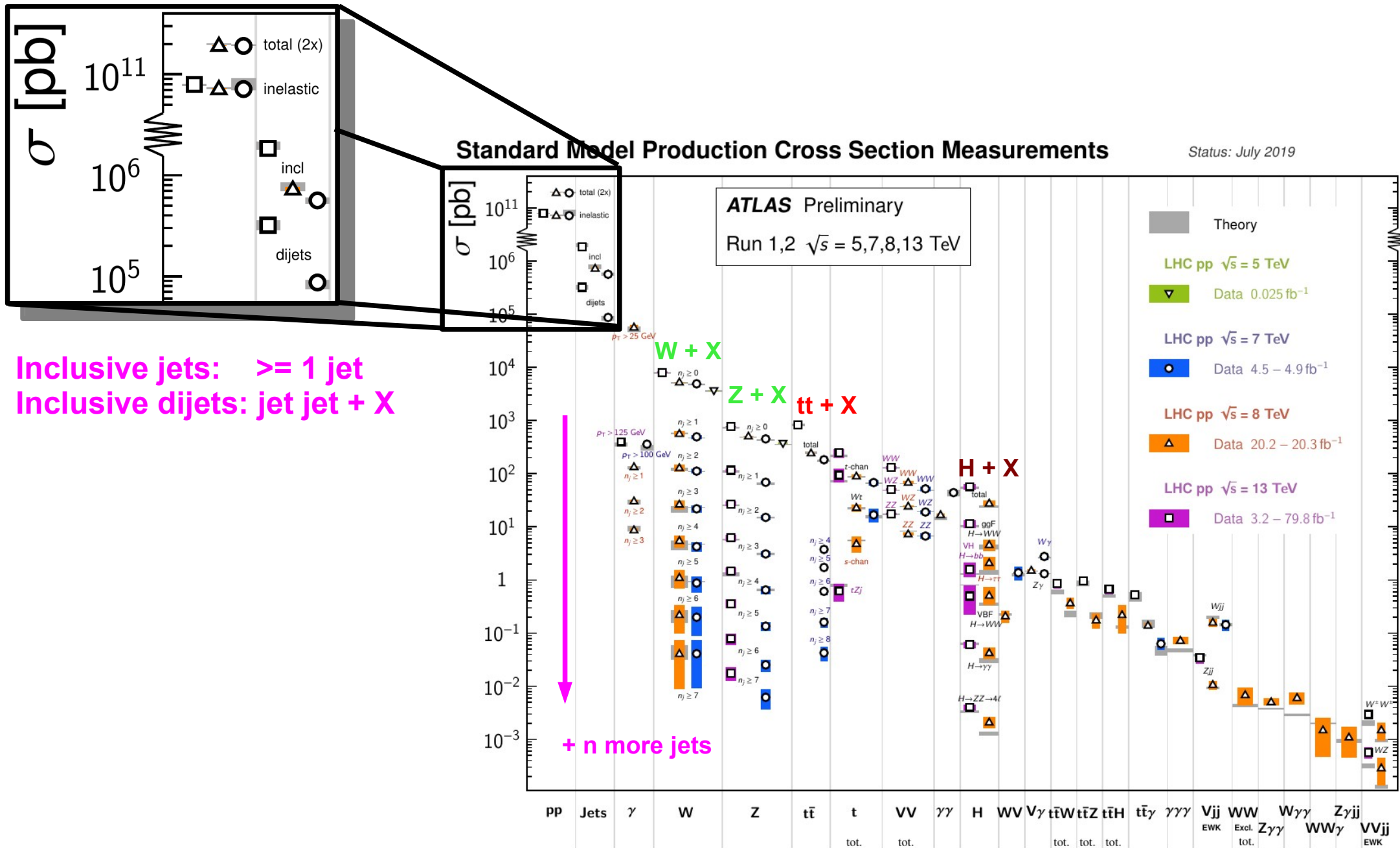
Higgs Bosonen ($\sigma_{ggH}, \sigma_{WH}, \sigma_{VBF}$)
 $\sim 150 / \text{h}$



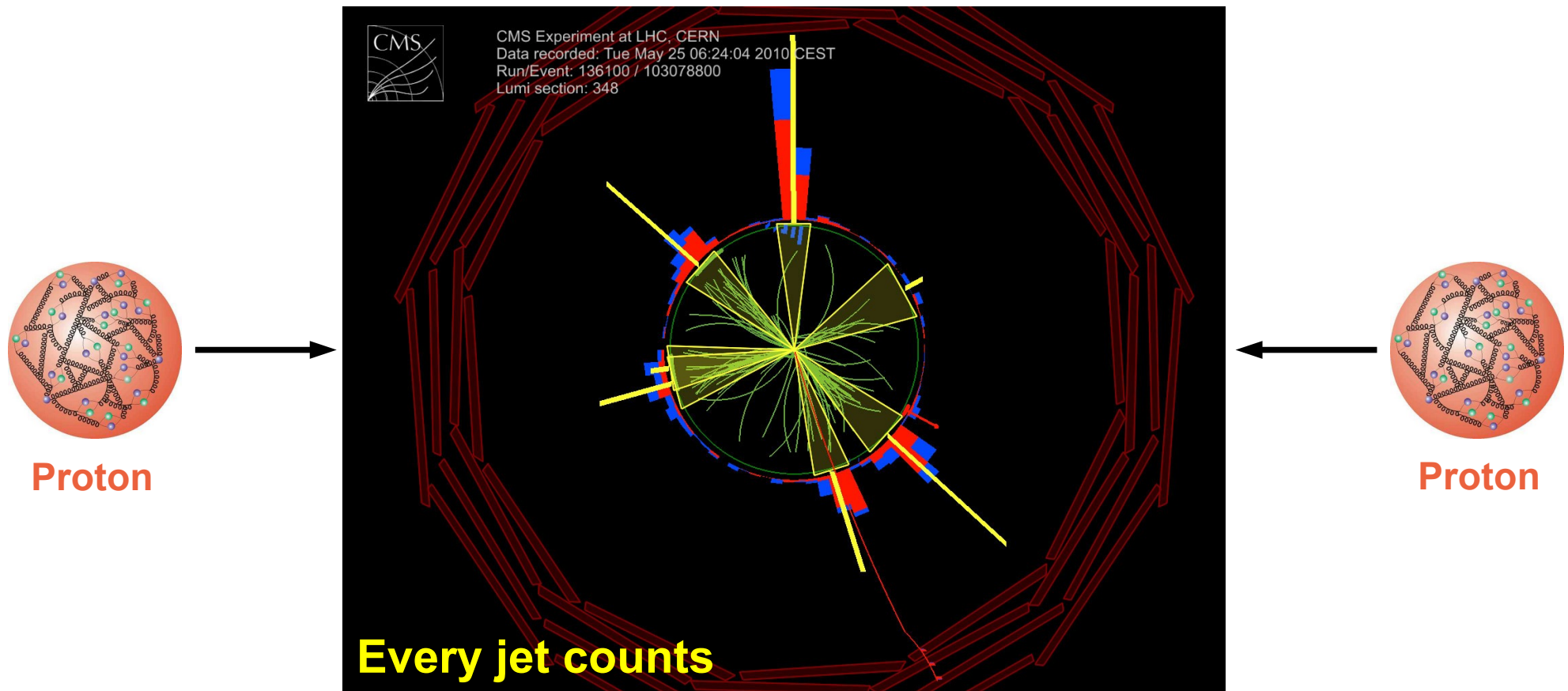
W.J. Stirling



Jets at the LHC



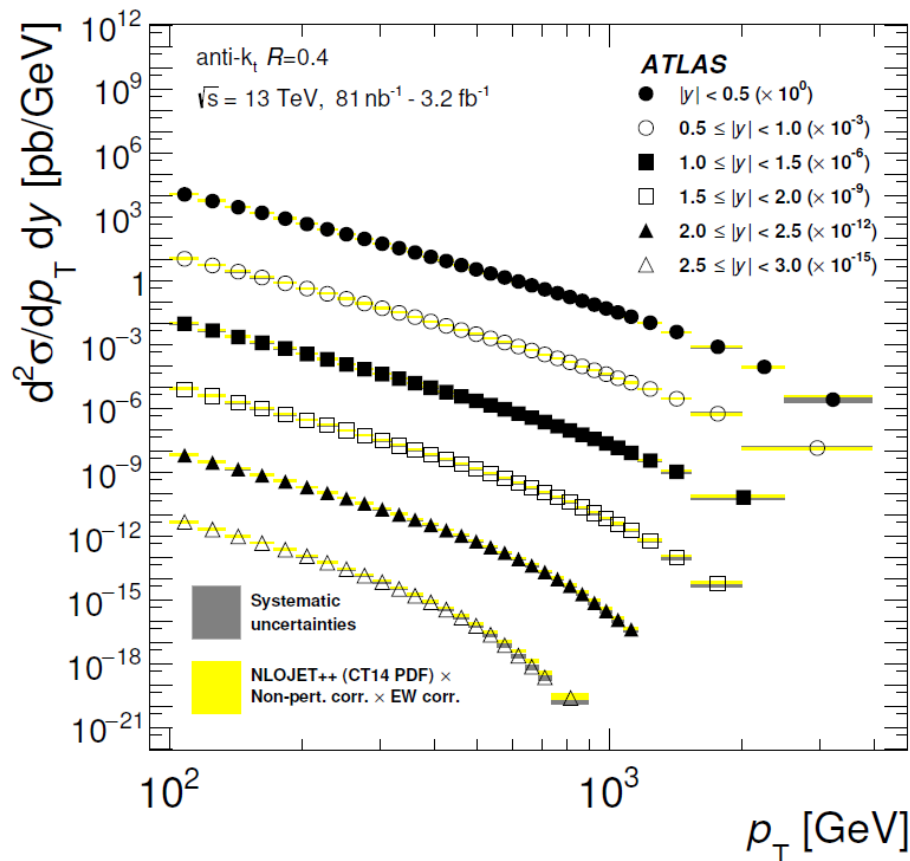
Large transverse momenta



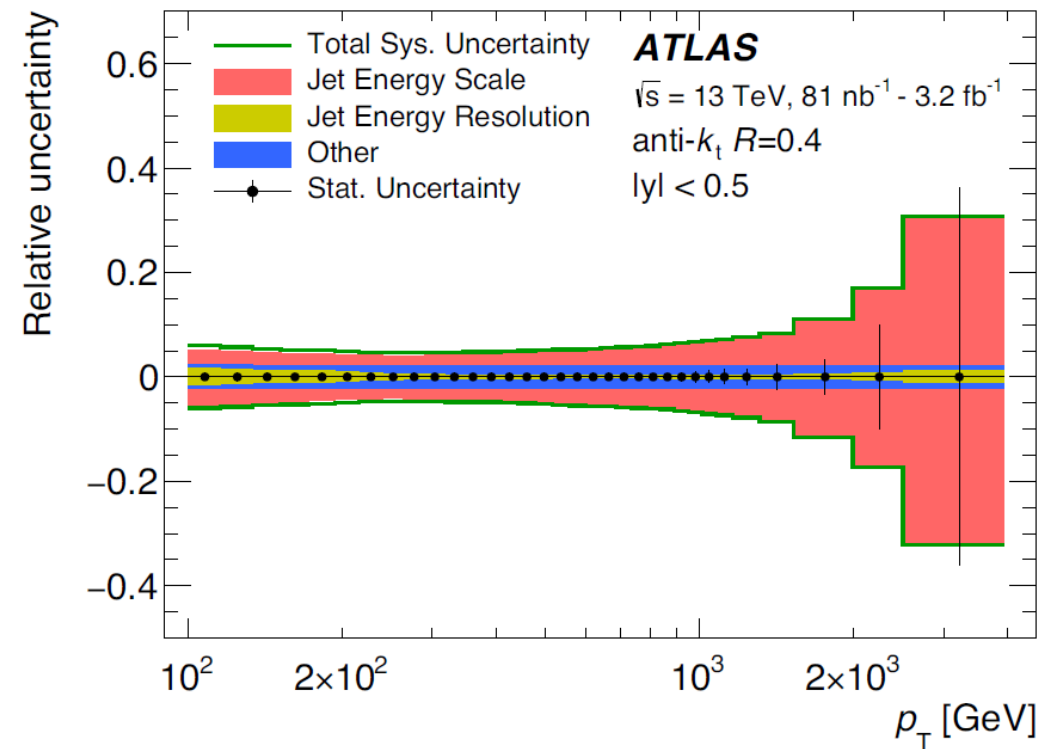
Overall agreement with predictions of **QCD** at NLO over many orders of magnitude in cross section and even beyond 2 TeV in jet p_T and for rapidities $|y|$ up to 3 ~ 5 at $\sqrt{s} = 2.76, 7, 8$, and 13 TeV.

$$\frac{d^2\sigma}{dp_T dy} \propto \alpha_s^2$$

Data vs. NLO pQCD x non-pert. x EW corrections

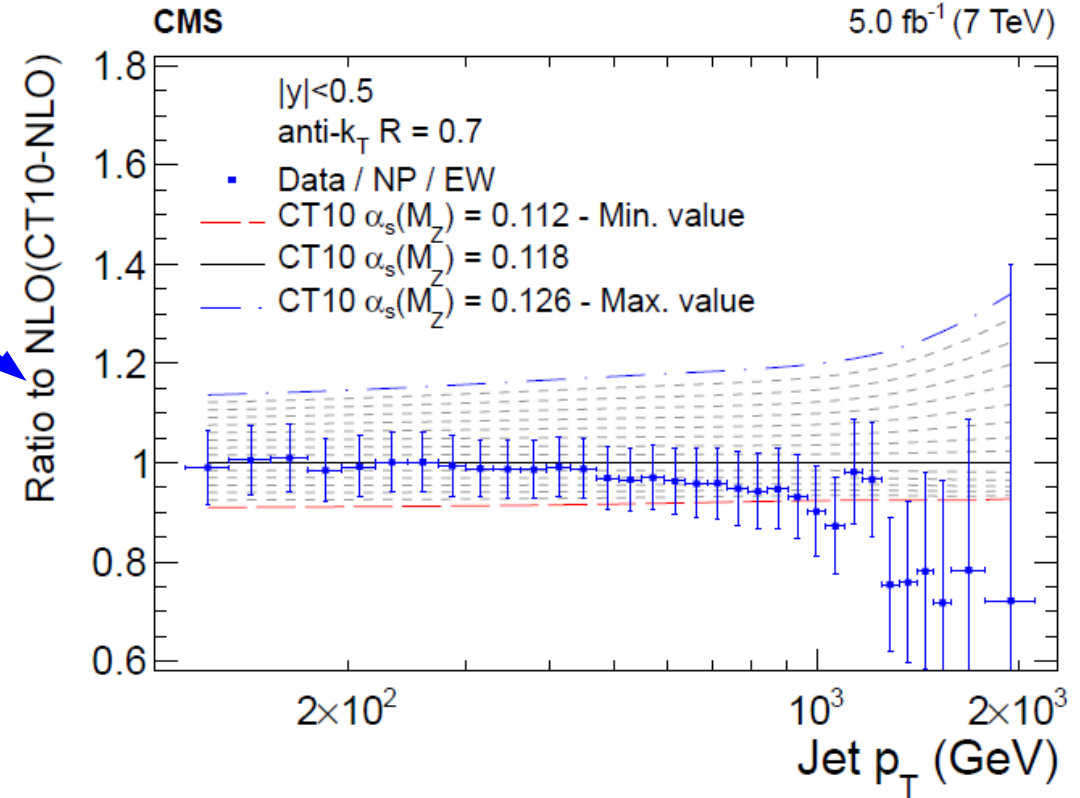
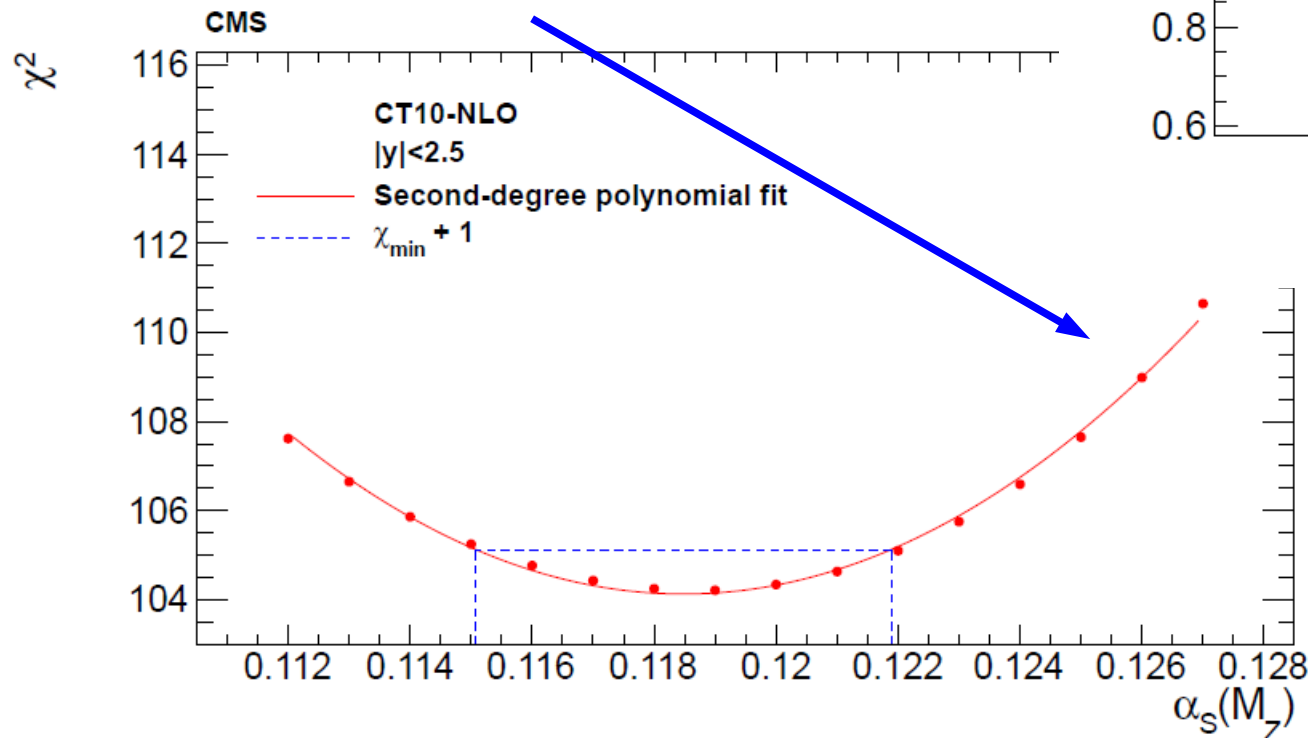


Exp. uncertainties for $|y| < 0.5$



Sensitivity of inclusive jet cross section to α_s

Example of χ^2 -fit



Much more on jet analyses in:

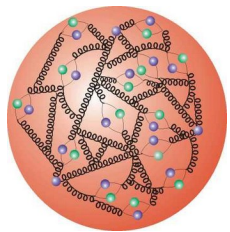
Particle physics II – Top quark and jet physics at the LHC

KR + A. Meyer

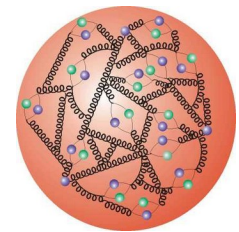
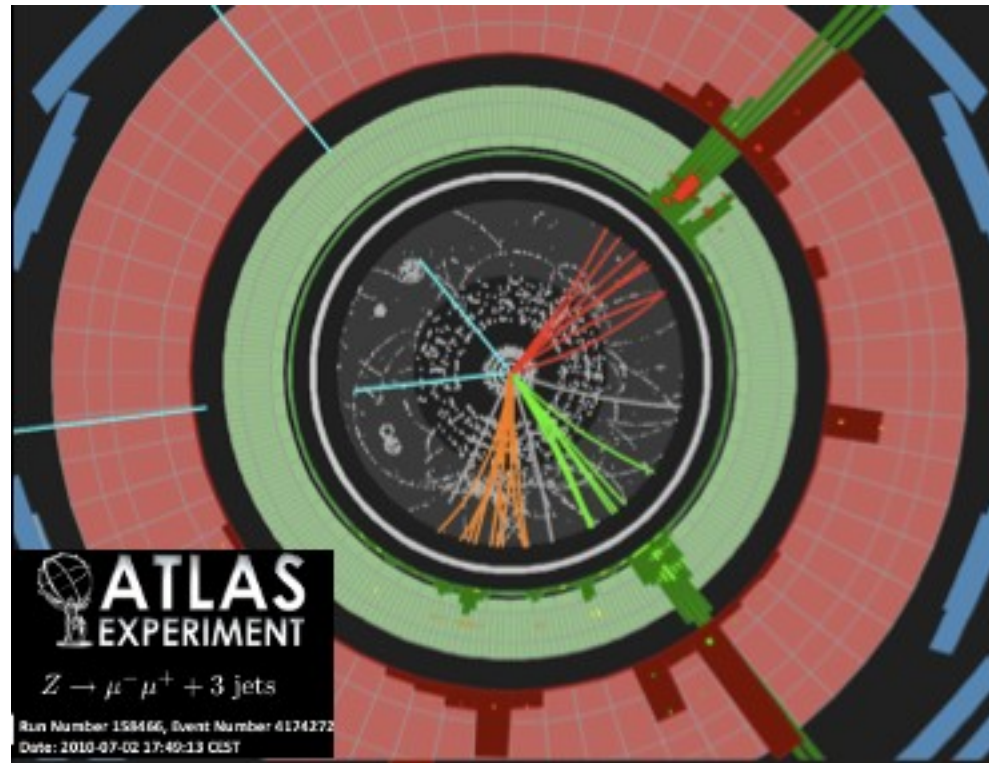


W and Z bosons

Standard Candles



Proton



Proton



- **Hadro-production of leptons pairs from W, Z decays**
 - ➔ large invariant mass
 - ➔ high p_T of leptons, **but not necessarily of W, Z**
 - ➔ high p_T of W, Z requires **balancing object** → most frequently **jets**
 - ➔ **Very interesting to study: W, Z p_T distribution**
 - ➔ **Also: di- and triple-vector boson production**
- **Further interest:**
 - ➔ **Masses of W and Z bosons: Important SM parameters**
 - ➔ **W and Z boson couplings in production and decay**
 - ➔ **E.g. asymmetries in production**

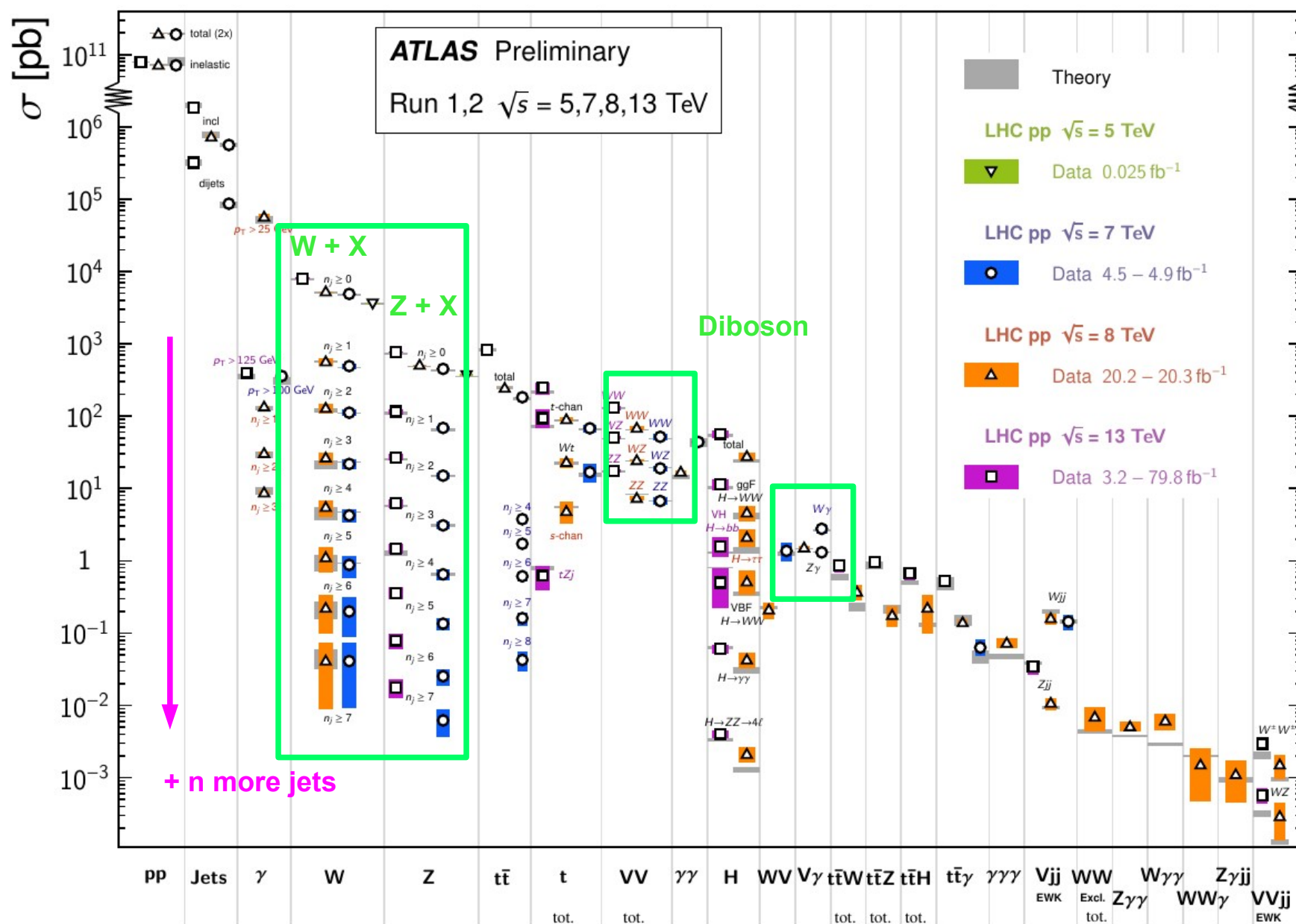
Much more on W, Z analyses in:

Particle physics II – W, Z and Higgs physics at the LHC

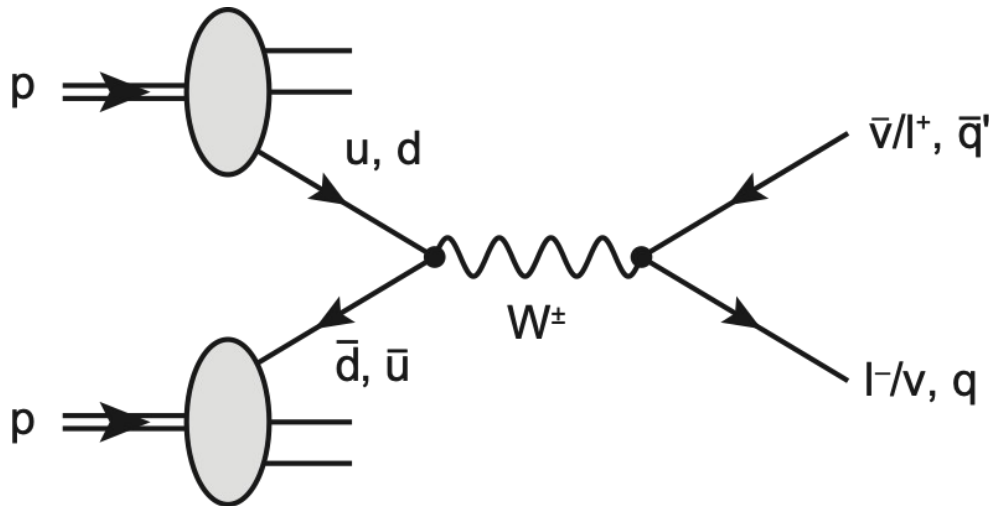
M. Schröder, R. Wolf



Status: July 2019



Proton-proton collisions, all p_T :



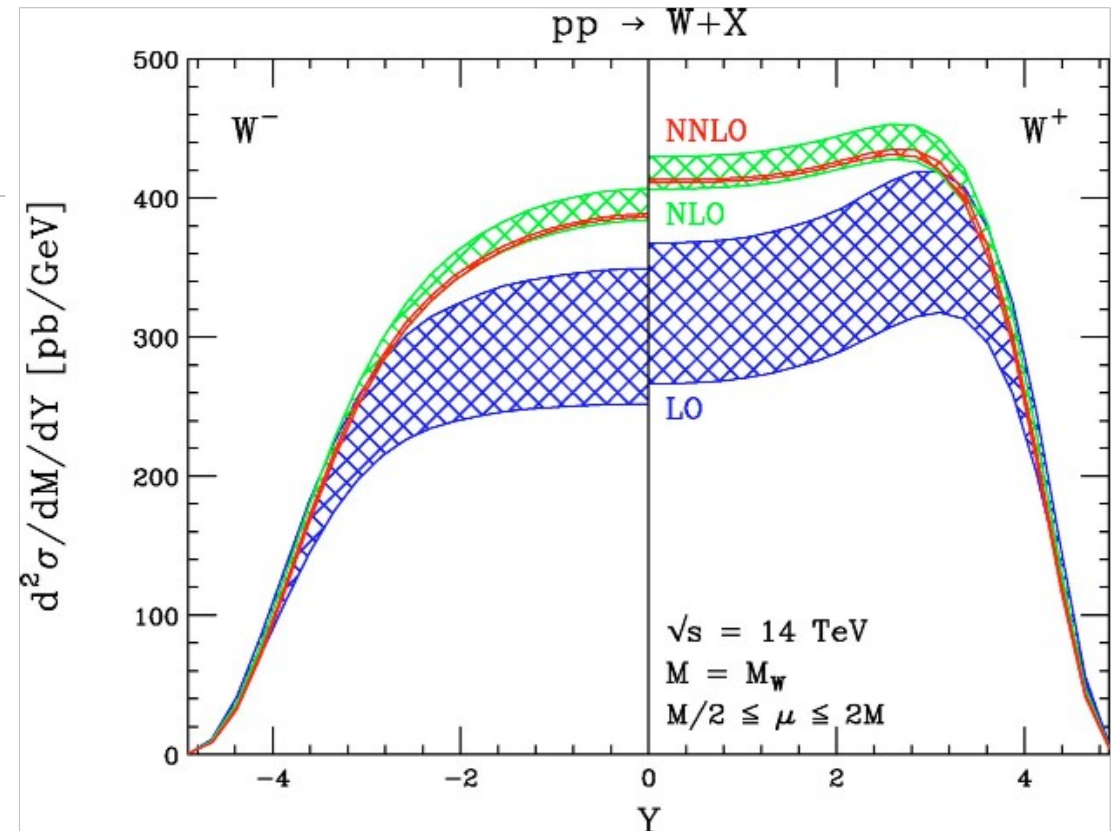
Theory known up to NNLO,
1st order independent of α_s

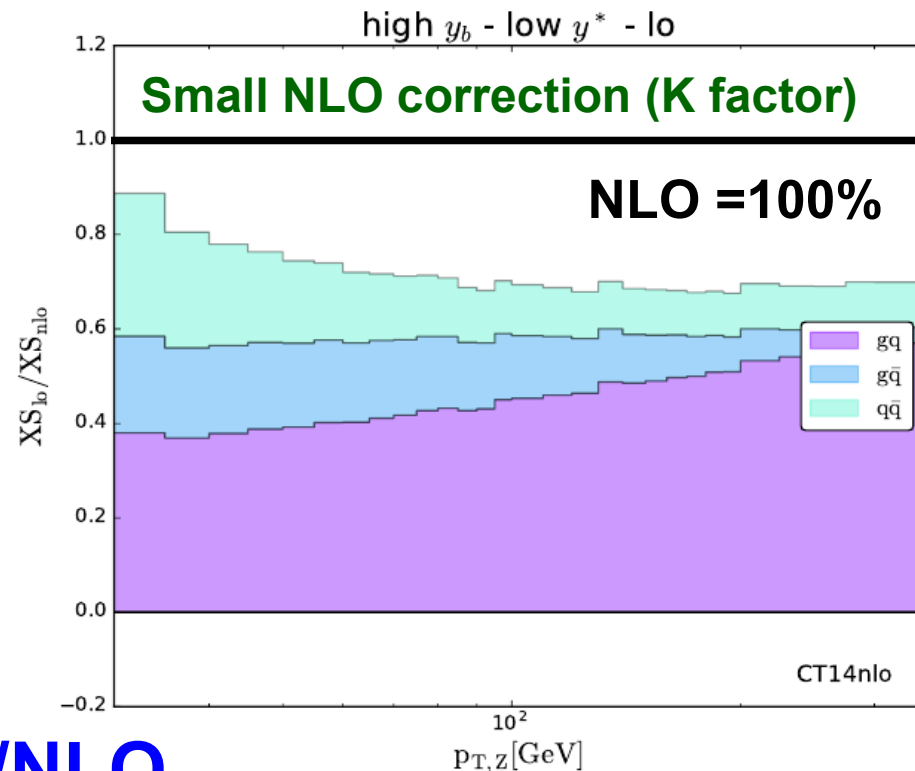
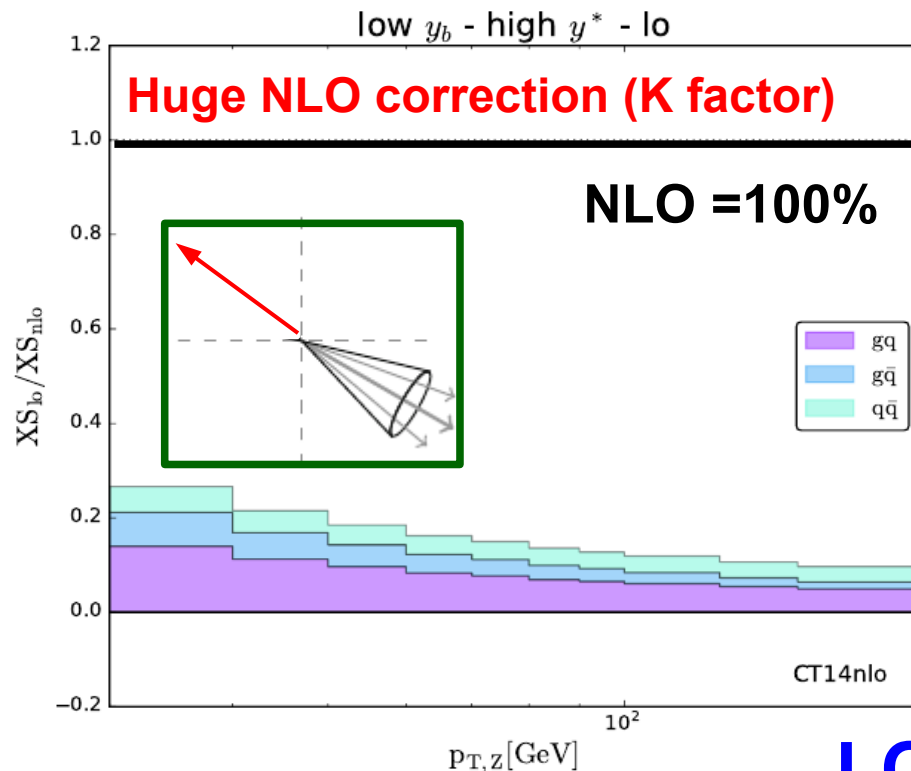
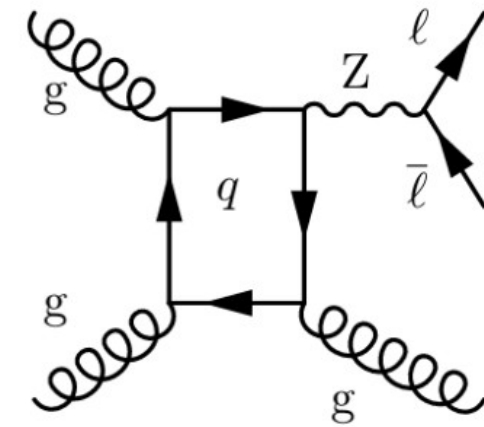
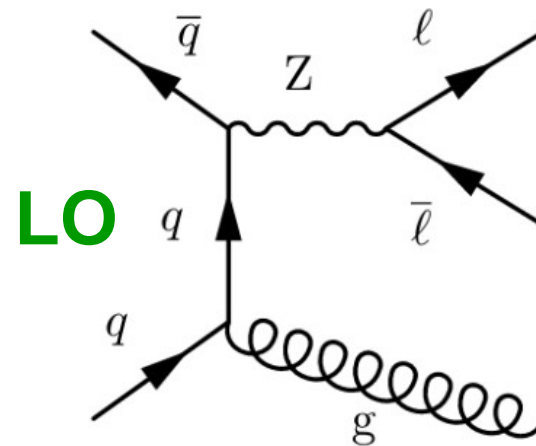
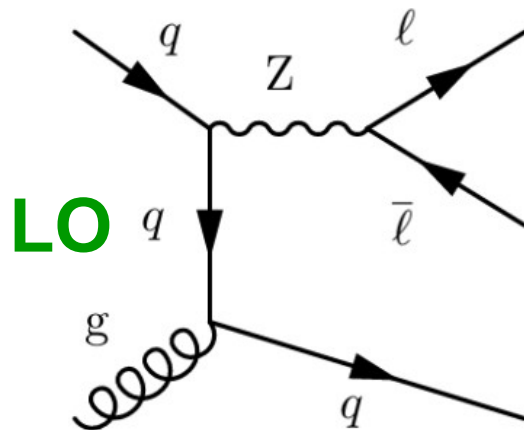
Only W^- shown

Only W^+ shown

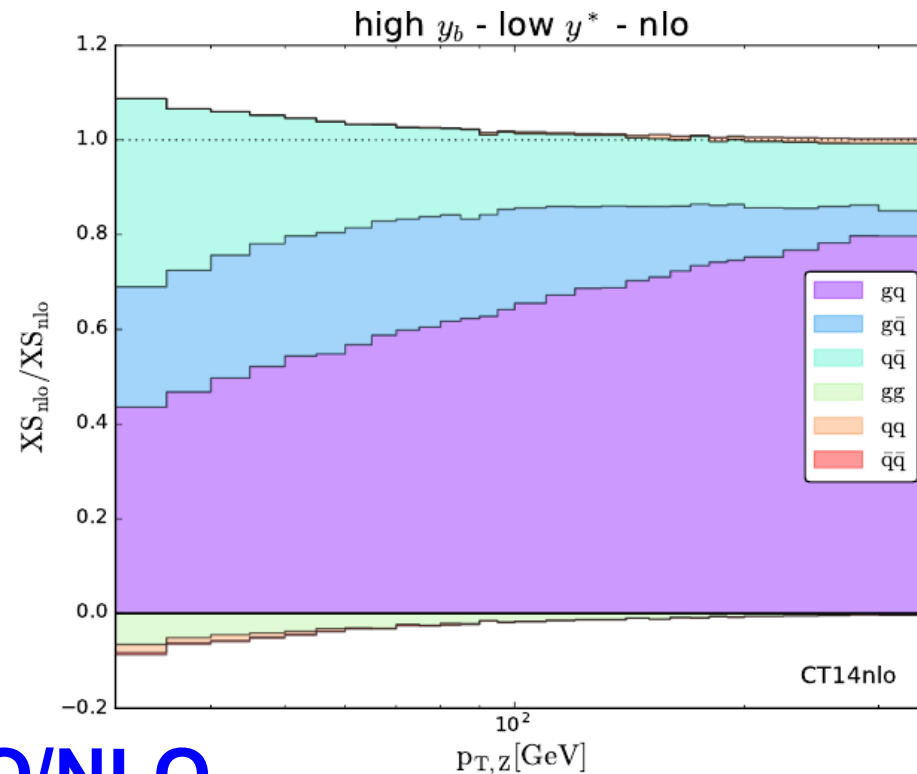
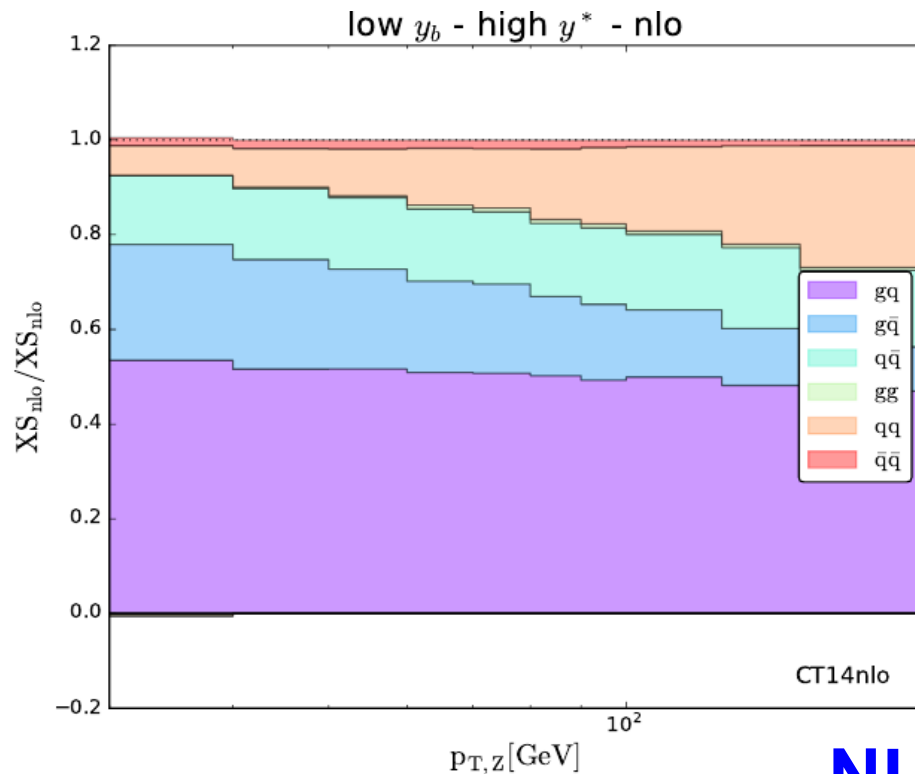
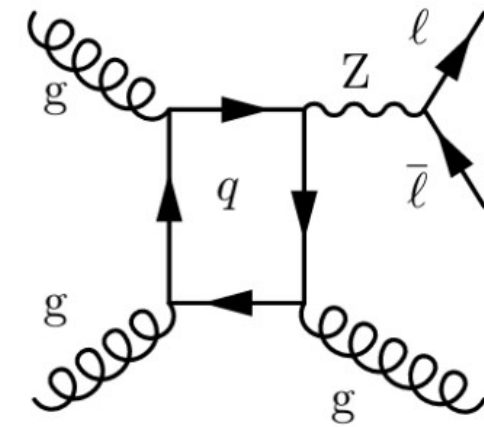
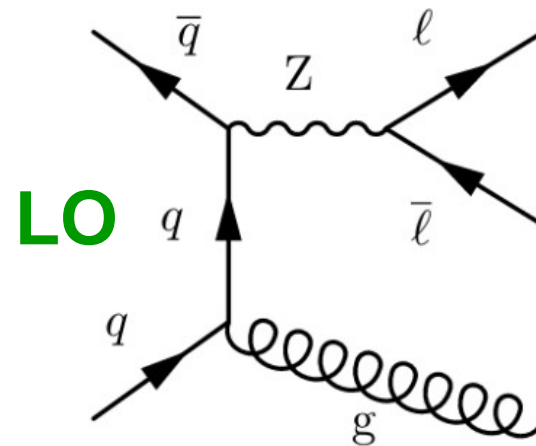
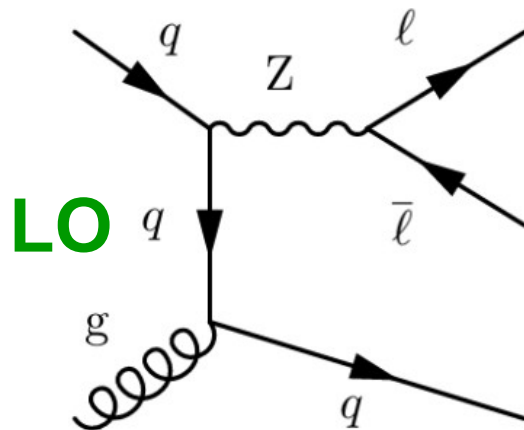
Production of W bosons vs. rapidity y

- symmetric around zero for W^+ , W^-
- **BUT: W^+ different from W^-**
- Proton content different for u and d!



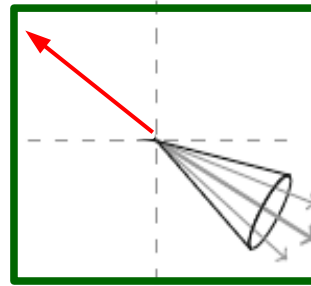


LO/NLO



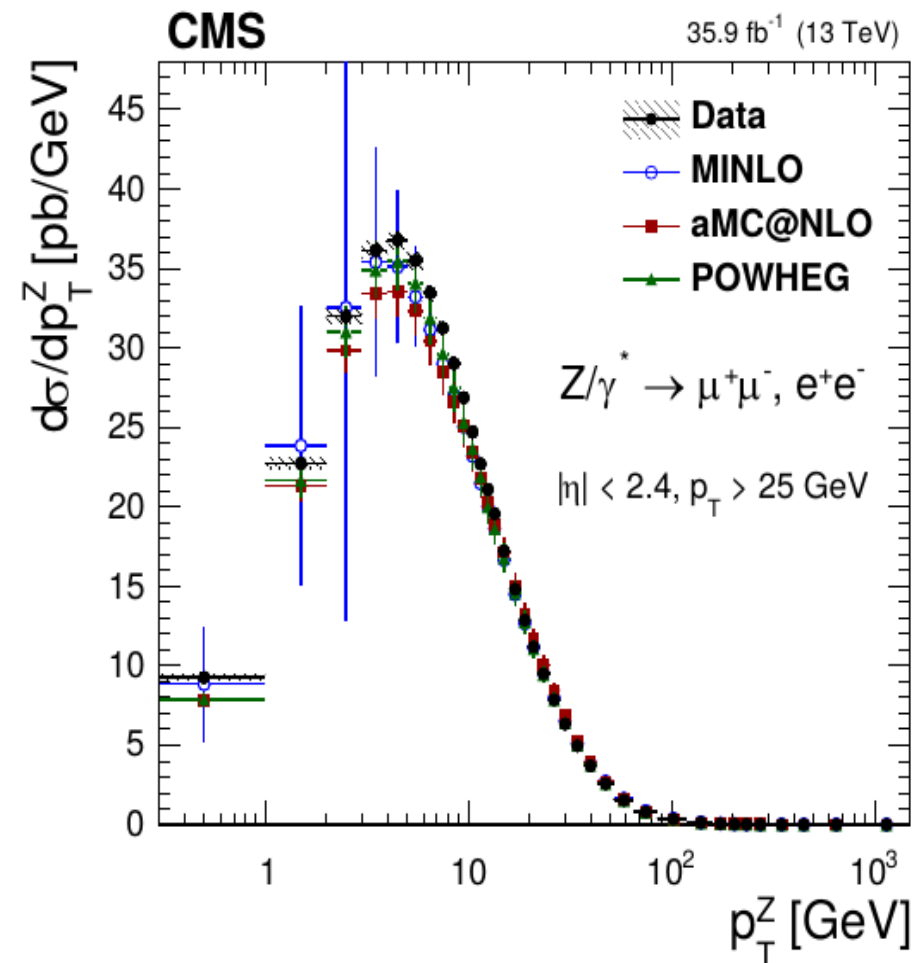
NLO/NLO

- Theory for forward-backward topology
- ➔ starts being precise only at NNLO

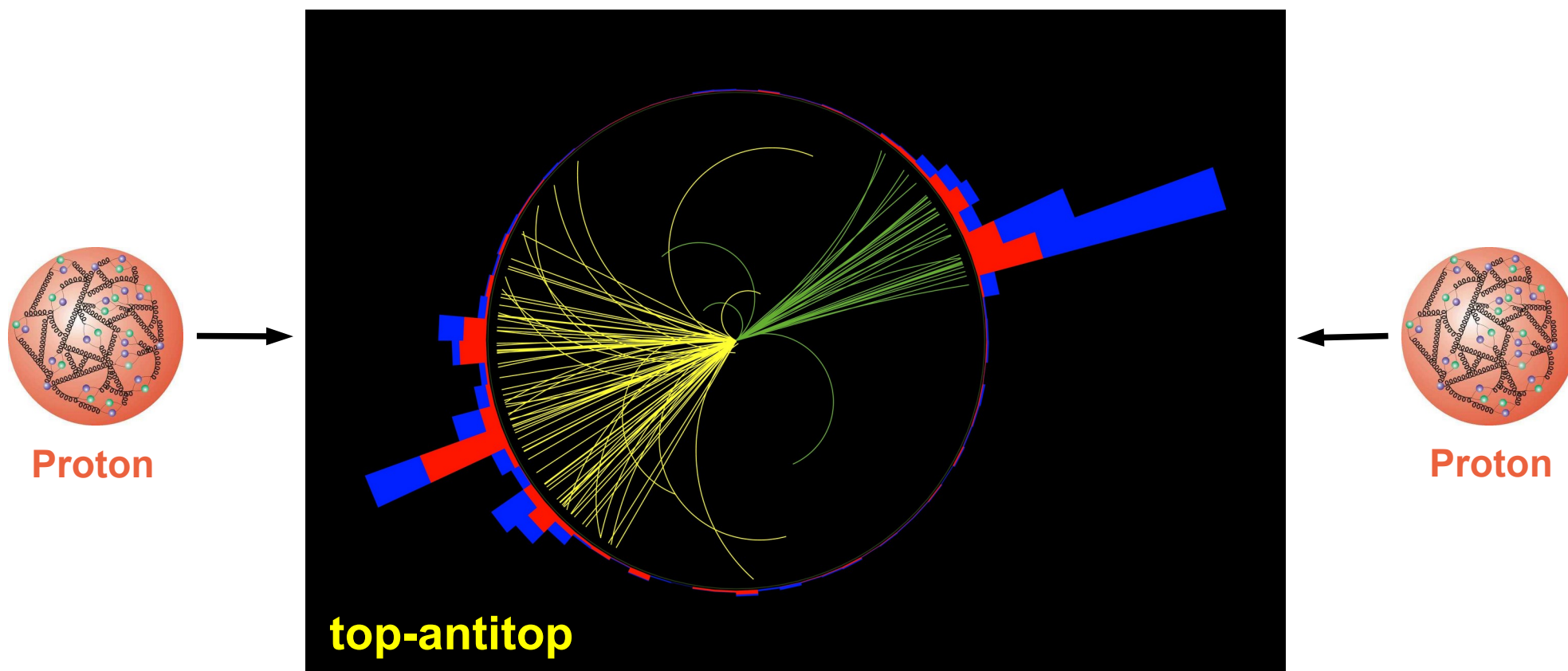


- Z p_T distribution:

- ➔ Z p_T balanced by jets at high p_T
- ➔ At small p_T multiple soft gluon radiation must be considered, fixed-order pQCD insufficient
- ➔ Resummation of leading terms to all orders
- ➔ Parton showers
- ➔ Description by MC generators with showers and/or resummation ok



Heavy quarks

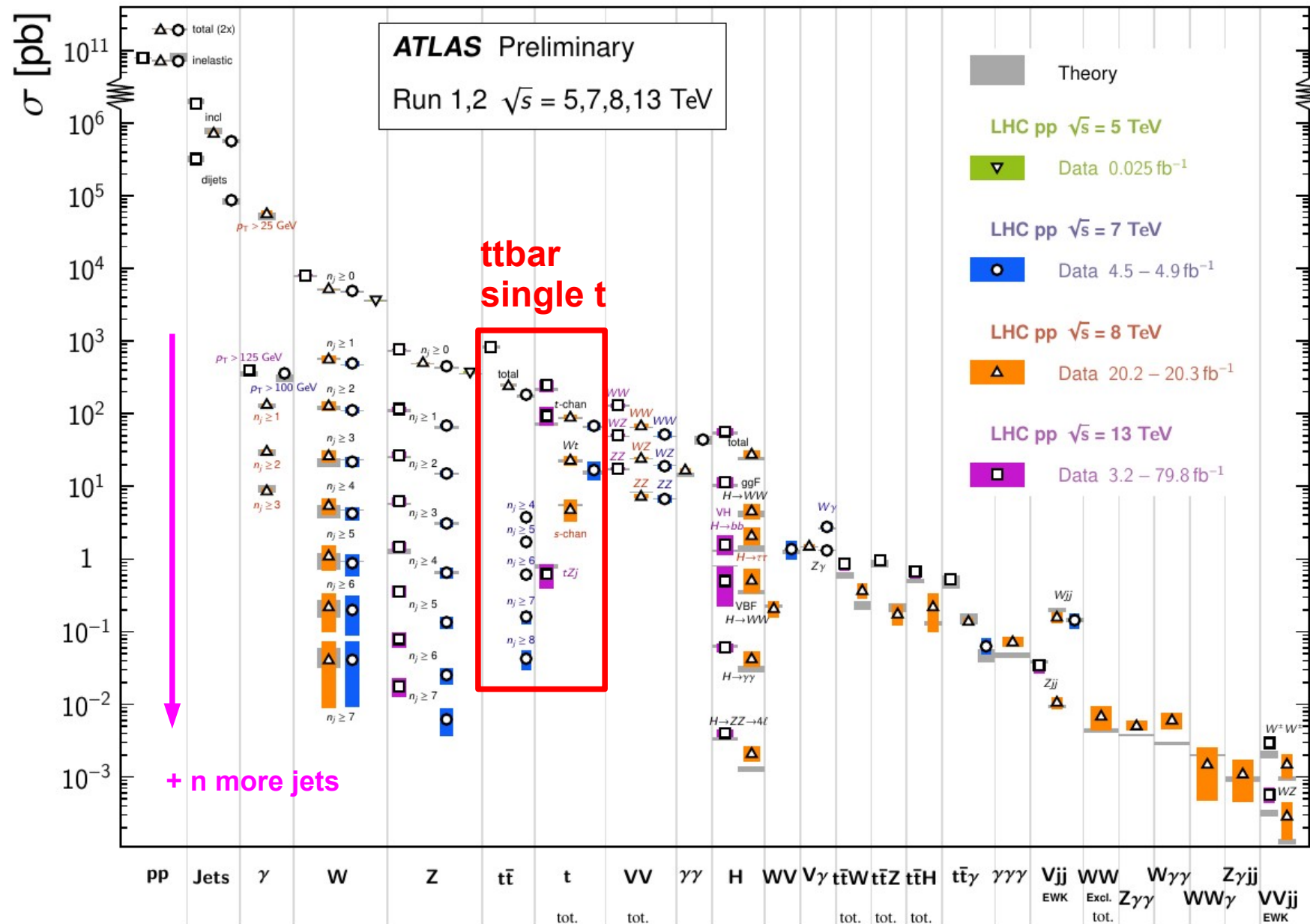


Relevant CMS measurements:

PLB 728, 496 (2013), JHEP 11, 067 (2012)
[Erratum: PLB 738, 526 (2014)],
CMS-TOP-17-001, arXiv:1812.10505
CMS-PAS-TOP-18-004.

Standard Model Production Cross Section Measurements

Status: July 2019



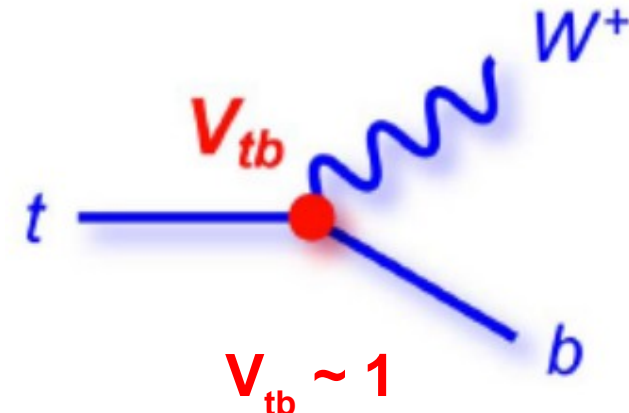
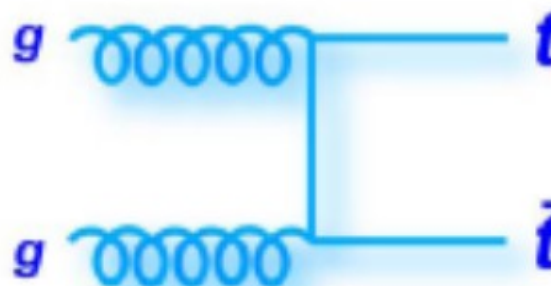


Top quark in the Standard Model

Institut für Experimentelle Teilchenphysik

• top like a normal quark

- ➔ Spin 1/2
- ➔ Couples to color
- ➔ Produced “strongly”, predominantly via gluon-gluon process at the LHC
- ➔ Decays “weakly” to ~100% via $t \rightarrow Wb$
- ➔ No flavor-changing neutral currents



$$V_{tb} \sim 1$$

	d	s	b
u	0.97434	0.22506	0.00357
c	0.22492	0.97351	0.0411
t	0.00875	0.0403	0.99915

CKM matrix

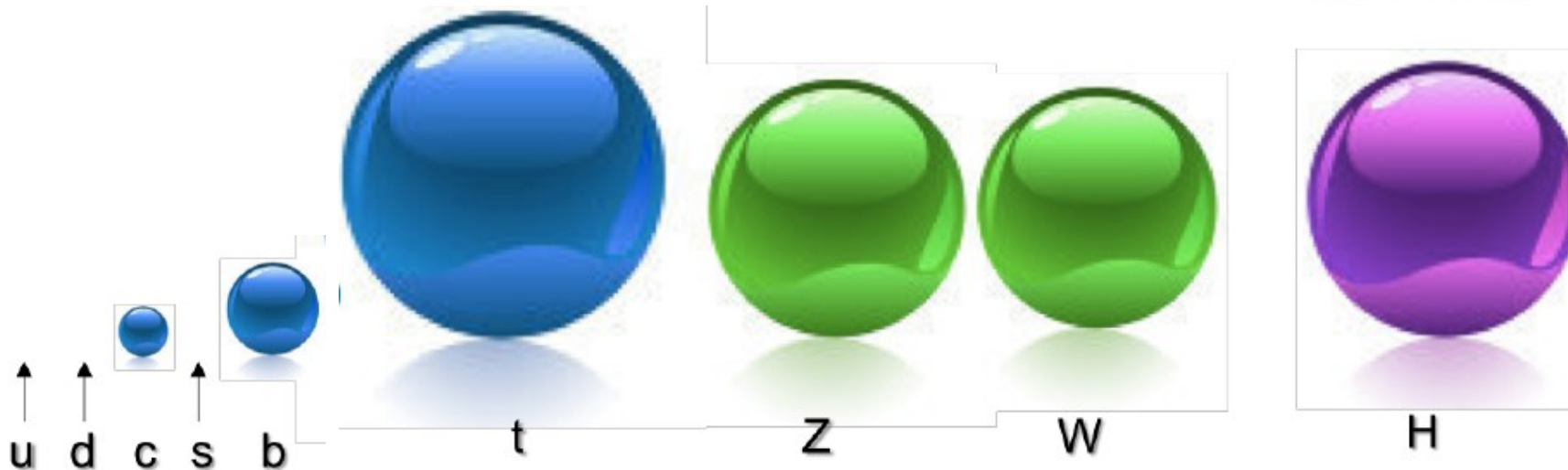
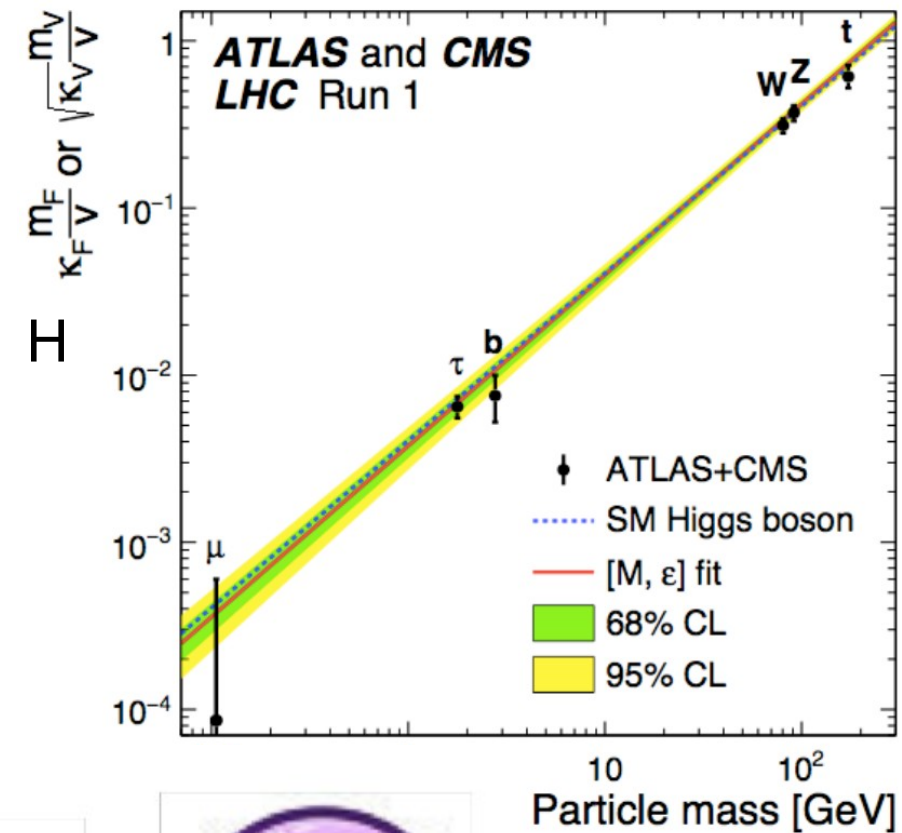


Top quark in the Standard Model

Institut für Experimentelle Teilchenphysik

top not a normal quark

- Heaviest known elementary particle
 $m_{\text{top}} = \sim 173 \text{ GeV}$



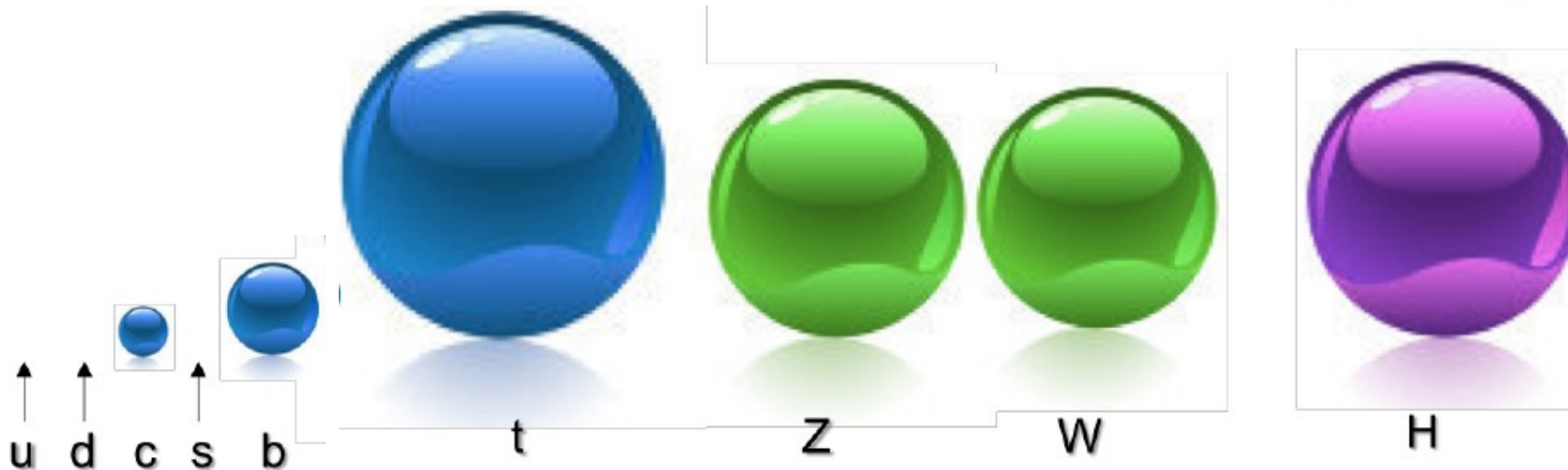
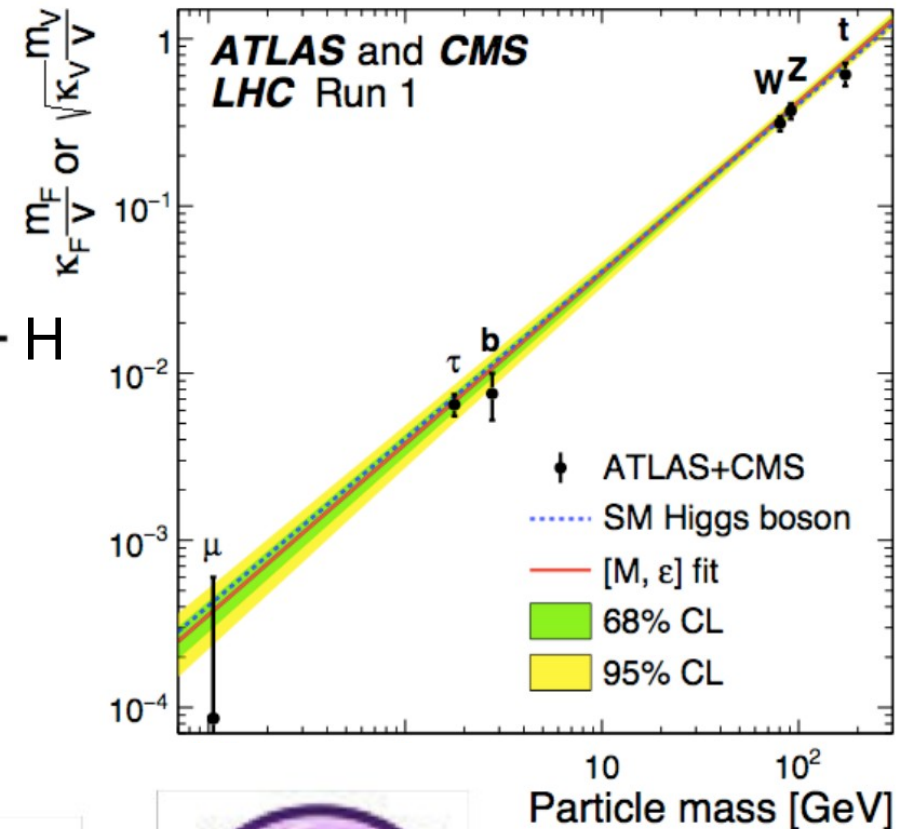
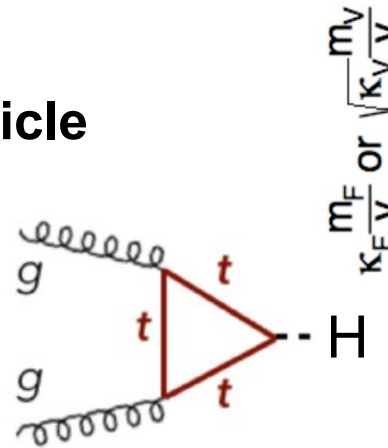


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- Heaviest known elementary particle
 $m_{\text{top}} = \sim 173 \text{ GeV}$
- Dominant loop graph for Higgs production



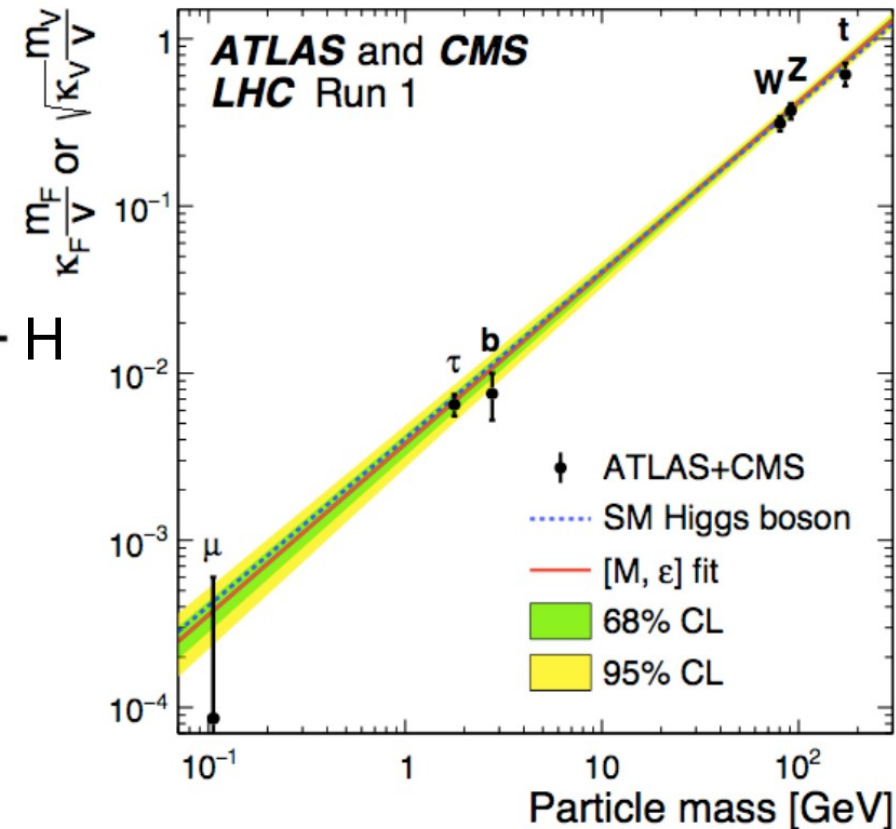
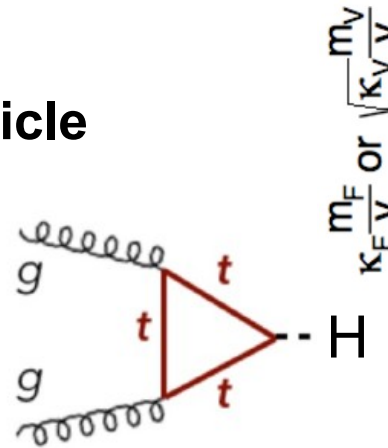


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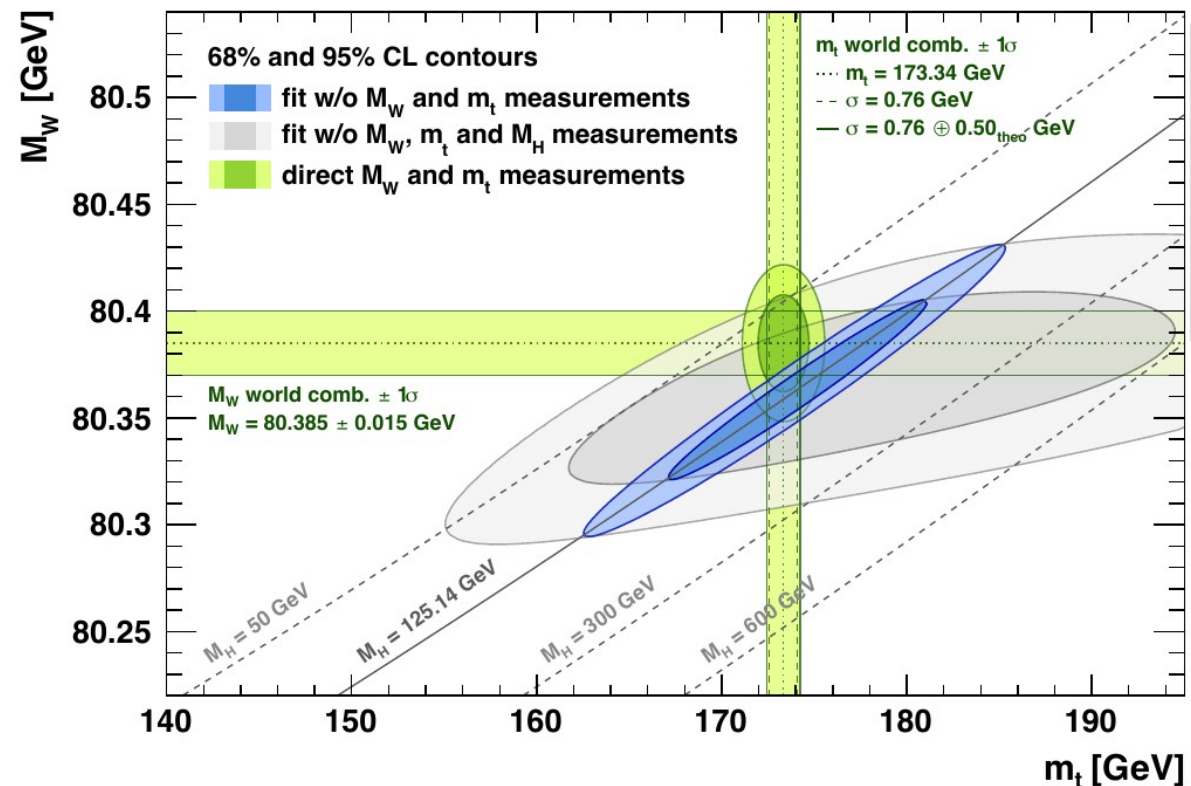
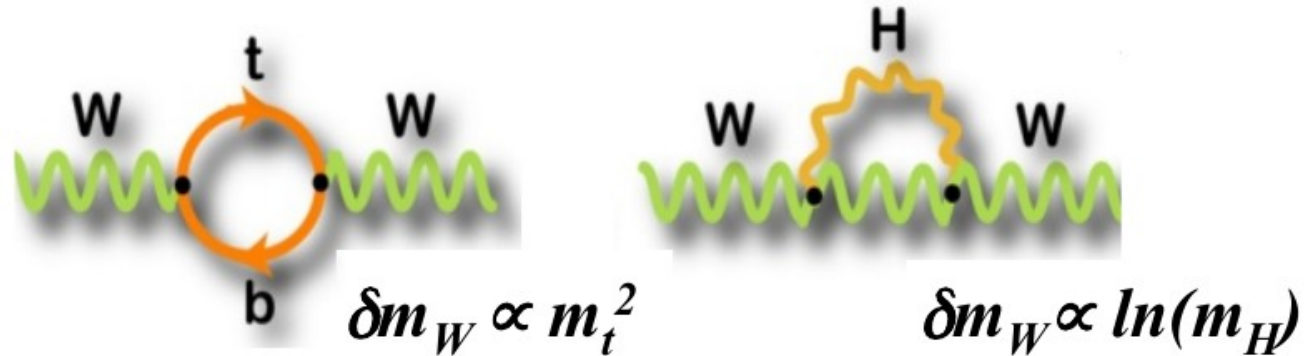
top not a normal quark

- Heaviest known elementary particle
 $m_{\text{top}} = \sim 173 \text{ GeV}$
- Dominant loop graph for Higgs production
- Very short lifetime $\tau \sim 5 \cdot 10^{-25} \text{ s}$
- Decays before forming hadrons
- Allows to measure spin correlations



$$\underbrace{\frac{1}{m_t}}_{\text{production } 10^{-27} \text{ s}} < \underbrace{\frac{1}{\Gamma_t}}_{\text{lifetime } 10^{-25} \text{ s}} < \underbrace{\frac{1}{\Lambda_{\text{QCD}}}}_{\text{hadronization } 10^{-24} \text{ s}}$$

- top mass connected to W and Higgs mass via loop corrections
- Production as single-top can provide input to PDFs
- Background to many searches for new physics



arXiv:1407.3792

Much more on top analyses in:

Particle physics II – Top quark and jet physics at the LHC

KR + A. Meyer

- Collisions including hadrons such as protons in the initial state require knowledge of the hadron's internal structure
- Protons are described by parton distribution functions (PDFs) describing the probability to find a parton i with momentum fraction x at a scale μ_f
 - The x dependence cannot be derived in perturbative QCD, but is universally usable in scattering processes once determined from data
 - The scale dependence is calculable in perturbative QCD, TP II Top & Jets
- Through the factorisation theorem of QCD cross sections involving hadrons can be derived separating soft/collinear effects into PDFs and high- p_T parton scatters in partonic matrix elements
- Examples for high- p_T processes at the LHC are
 - Inclusive jet production $\rightarrow$ determination of the strong coupling constant
 - W production $\rightarrow$ asymmetry in $W^+ W^-$ rapidity distribution
 - $Z+X$ production $\rightarrow$ process decomposition vs. rapidity; Z p_T from low to high
 - $t\bar{t}$ production $\rightarrow$ top quark mass, lifetime, ...



Backup