

Theoretical Optics

Modelling material properties

Prof. Carsten Rockstuhl



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1.2 Material properties

phenomenological models for $\mathbf{P}(\mathbf{E})$ und $\mathbf{j}(\mathbf{E})$

$\mathbf{E}(\mathbf{r}, t) \rightarrow \text{medium}$ (response function) $\rightarrow \mathbf{P}(\mathbf{r}, t)$

$\bar{\mathbf{E}}(\mathbf{r}, \omega) \rightarrow \text{medium}$ (transfer function) $\rightarrow \bar{\mathbf{P}}(\mathbf{r}, \omega)$

1.2.1 Basic properties

biaxial anisotropic material

$$P_i(\mathbf{r}, t) = \epsilon_0 \sum_{j=1}^3 \int_{-\infty}^t R_{ij}^{(1)}(\mathbf{r}, t - t') E_j(\mathbf{r}, t') dt' \quad R(\mathbf{r}, \tau) \text{- response function}$$

$$\bar{P}_i(\mathbf{r}, \omega) = \epsilon_0 \sum_{j=1}^3 \chi_{ij}^{(1)}(\mathbf{r}, \omega) \bar{E}_j(\mathbf{r}, \omega) \quad \chi^{(1)}(\mathbf{r}, \omega) \text{- susceptibility}$$

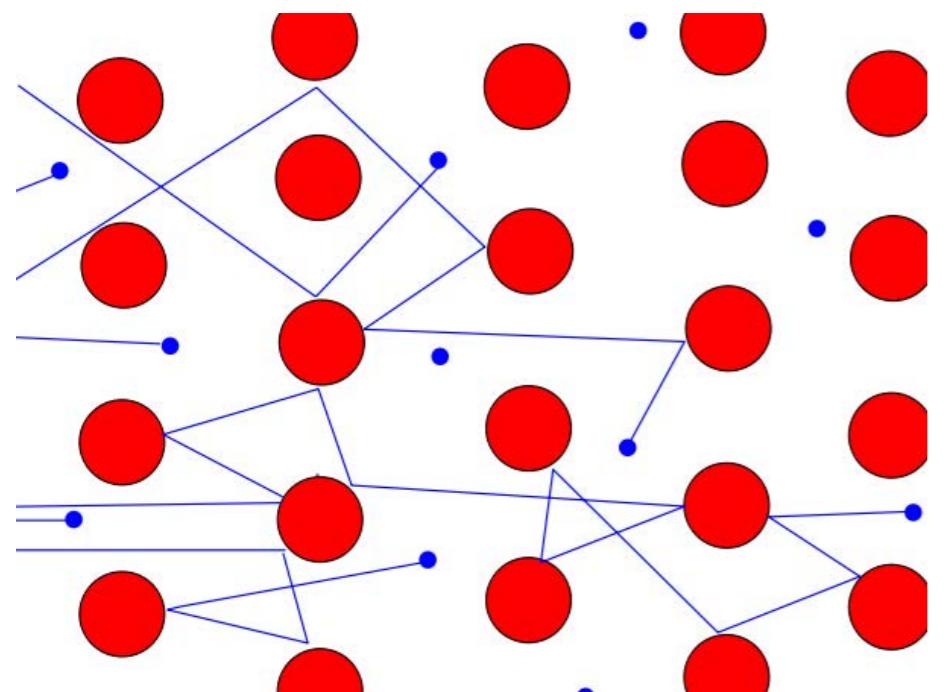
$$\boxed{\chi^{(1)}(\mathbf{r}, \omega) \xleftrightarrow{FT} R(\mathbf{r}, \tau)}$$

linear, homogenous, isotropic and dispersive media

$$R(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \chi^{(1)}(\omega) e^{-i\omega t} d\omega$$

Drude model for free electrons

$$\frac{\partial^2}{\partial t^2} \mathbf{s}(\mathbf{r}, t) + g \frac{\partial}{\partial t} \mathbf{s}(\mathbf{r}, t) = -\frac{e}{m} \mathbf{E}(\mathbf{r}, t)$$



$$\mathbf{j}(\mathbf{r}, t) = -Ne \frac{\partial}{\partial t} \mathbf{s}(\mathbf{r}, t)$$

$$\frac{\partial}{\partial t} \mathbf{j}(\mathbf{r}, t) + g \mathbf{j}(\mathbf{r}, t) = \frac{e^2 N}{m} \mathbf{E}(\mathbf{r}, t) = \epsilon_0 f \mathbf{E}(\mathbf{r}, t) = \epsilon_0 \omega_p^2 \mathbf{E}(\mathbf{r}, t)$$

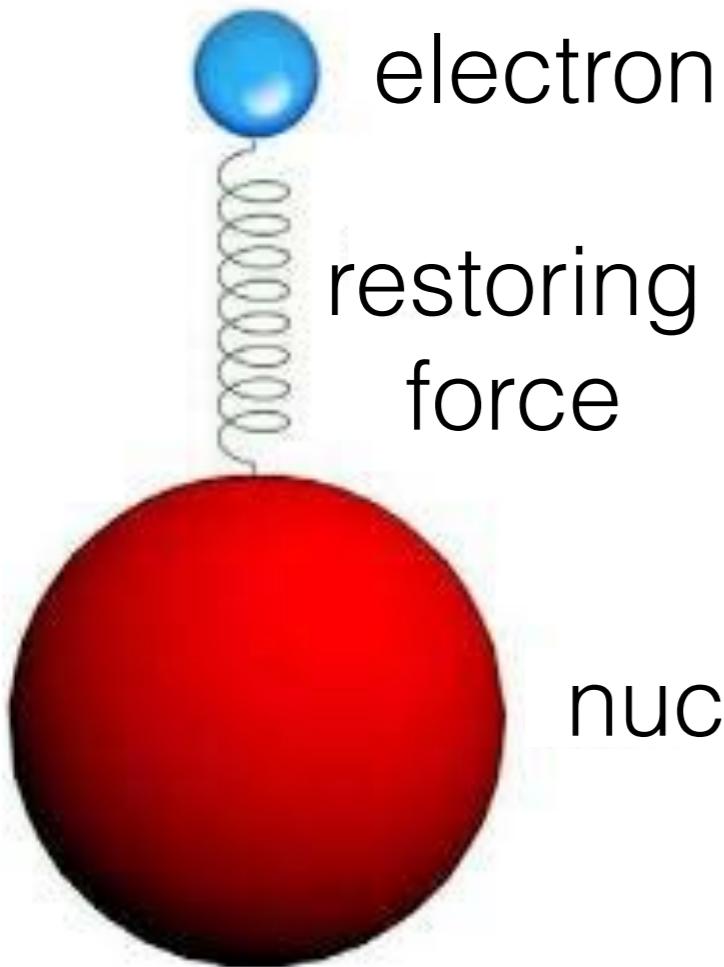
$$\omega_p^2 = f = \frac{1}{\epsilon_0} \frac{e^2 N}{m}$$

$$-i\omega \bar{\mathbf{j}}(\mathbf{r}, \omega) + g \bar{\mathbf{j}}(\mathbf{r}, \omega) = \epsilon_0 \omega_p^2 \bar{\mathbf{E}}(\mathbf{r}, \omega)$$

$$\bar{\mathbf{j}}(\mathbf{r}, \omega) = \frac{\epsilon_0 \omega_p^2}{g - i\omega} \bar{\mathbf{E}}(\mathbf{r}, \omega) = \sigma(\mathbf{r}, \omega) \bar{\mathbf{E}}(\mathbf{r}, \omega)$$

Lorentz model for bound charges

$$\frac{\partial^2}{\partial t^2} \mathbf{s}(\mathbf{r}, t) + g \frac{\partial}{\partial t} \mathbf{s}(\mathbf{r}, t) + \omega_0^2 \mathbf{s}(\mathbf{r}, t) = \frac{q}{m} \mathbf{E}(\mathbf{r}, t)$$



$$\mathbf{p}(\mathbf{r}, t) = q\mathbf{s}(\mathbf{r}, t)$$

$$\mathbf{P}(\mathbf{r}, t) = N\mathbf{p}(\mathbf{r}, t) = Nq\mathbf{s}(\mathbf{r}, t)$$

$$\frac{\partial^2}{\partial t^2} \mathbf{P}(\mathbf{r}, t) + g \frac{\partial}{\partial t} \mathbf{P}(\mathbf{r}, t) + \omega_0^2 \mathbf{P}(\mathbf{r}, t) = \frac{q^2 N}{m} \mathbf{E}(\mathbf{r}, t) = \varepsilon_0 f \mathbf{E}(\mathbf{r}, t)$$

$$f = \frac{1}{\varepsilon_0} \frac{q^2 N}{m}$$

$$-\omega^2 \bar{\mathbf{P}}(\mathbf{r}, \omega) - ig\omega \bar{\mathbf{P}}(\mathbf{r}, \omega) + \omega_0^2 \bar{\mathbf{P}}(\mathbf{r}, \omega) = \varepsilon_0 f \bar{\mathbf{E}}(\mathbf{r}, \omega)$$

$$\bar{\mathbf{P}}(\mathbf{r}, \omega) = \frac{\varepsilon_0 f}{(\omega_0^2 - \omega^2) - ig\omega} \bar{\mathbf{E}}(\mathbf{r}, \omega) = \varepsilon_0 \chi(\omega) \bar{\mathbf{E}}(\mathbf{r}, \omega)$$

1.2.2 Complex dielectric function

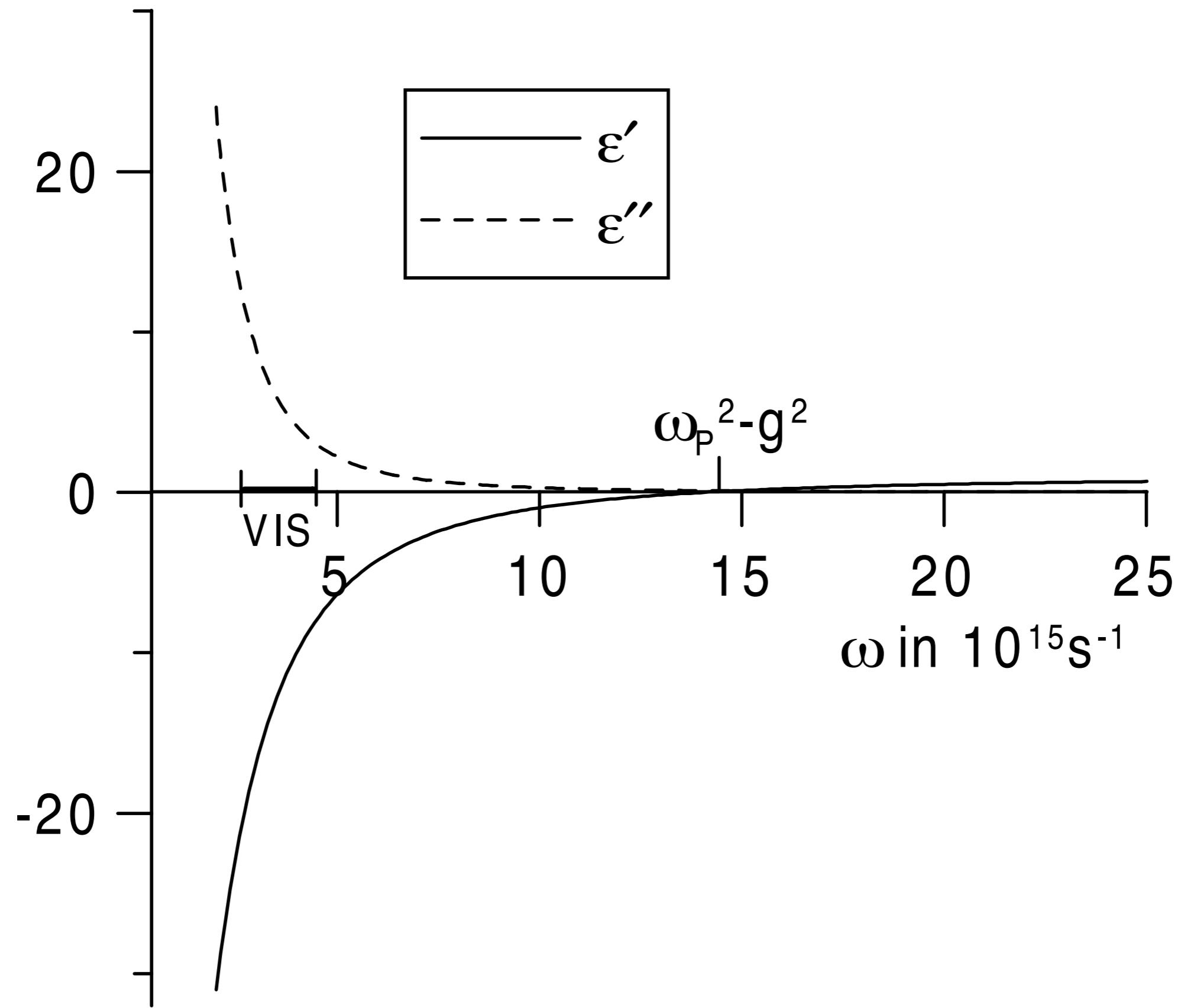
$$\begin{aligned} \mathbf{rot} \mathbf{rot} \bar{\mathbf{E}}(\mathbf{r}, \omega) - \frac{\omega^2}{c_0^2} \bar{\mathbf{E}}(\mathbf{r}, \omega) &= \omega^2 \mu_0 \bar{\mathbf{P}}(\mathbf{r}, \omega) + i\omega \mu_0 \bar{\mathbf{J}}(\mathbf{r}, \omega) \\ &= [\mu_0 \varepsilon_0 \omega^2 \chi(\omega) + i\omega \mu_0 \sigma(\mathbf{r}, \omega)] \bar{\mathbf{E}}(\mathbf{r}, \omega) \end{aligned}$$

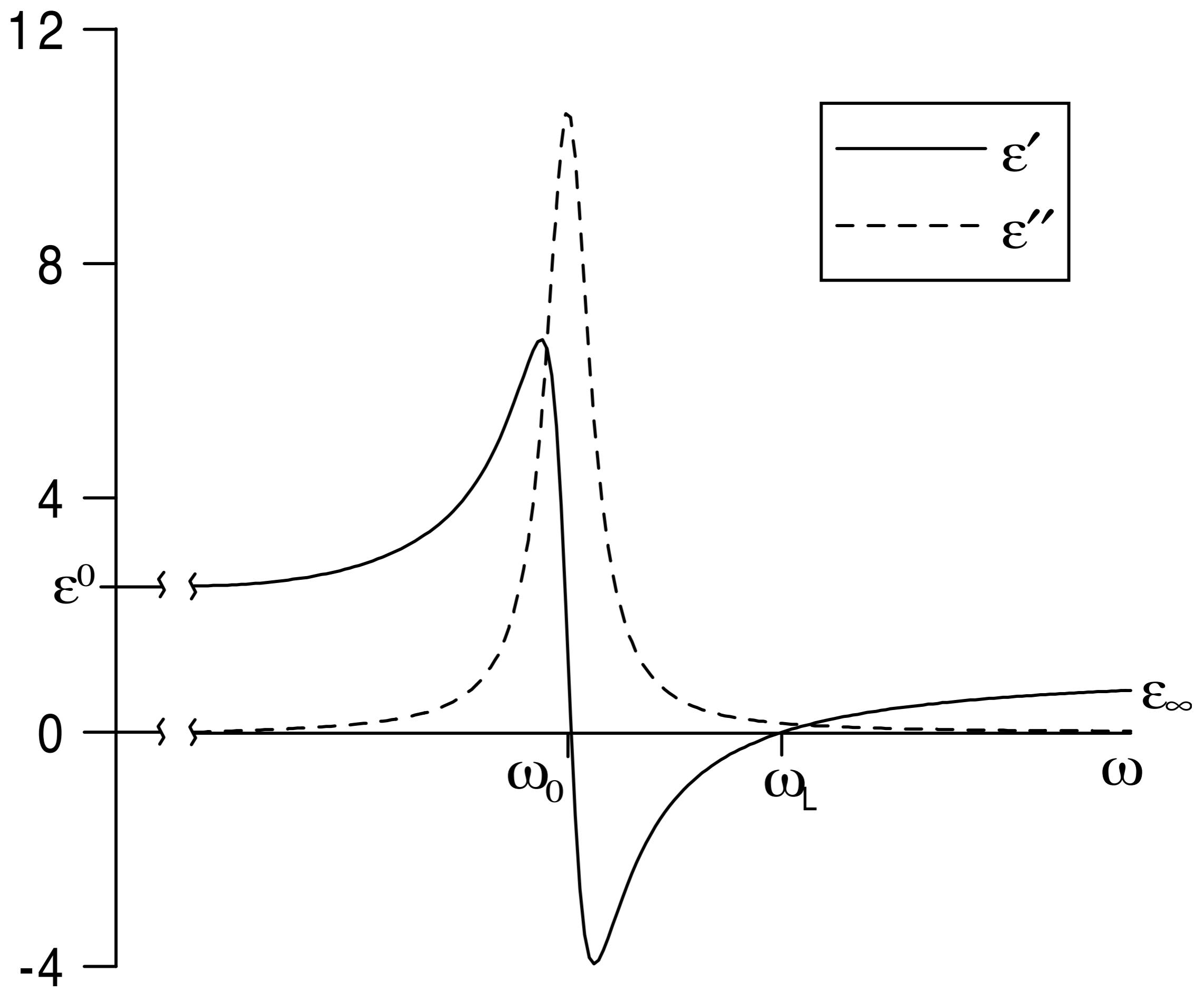
$$\begin{aligned} \mathbf{rot} \mathbf{rot} \bar{\mathbf{E}}(\mathbf{r}, \omega) &= \frac{\omega^2}{c_0^2} \left\{ \mathbf{1} + \chi(\omega) + \frac{i}{\omega \varepsilon_0} \sigma(\mathbf{r}, \omega) \right\} \bar{\mathbf{E}}(\mathbf{r}, \omega) \\ &= \frac{\omega^2}{c_0^2} \varepsilon(\omega) \bar{\mathbf{E}}(\mathbf{r}, \omega) \end{aligned}$$

Generalized complex dielectric function

$$\varepsilon(\omega) = \mathbf{1} + \chi(\omega) + \frac{i}{\omega \varepsilon_0} \sigma(\mathbf{r}, \omega) = \varepsilon'(\omega) + i\varepsilon''(\omega)$$

$$\varepsilon(\omega) = 1 + \sum_j \frac{f_j}{(\omega_{0j}^2 - \omega^2) - ig_j \omega} + \frac{\omega_p^2}{-\omega^2 - ig\omega}$$





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Kramers-Kronig Relation

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1.2.3 Kramers-Kronig dispersion relation

real and imaginary part of transfer function are linked by an integral

Applies when the response function is:

- time invariant
- real valued
- when causality applies.

Time invariance

$$\mathbf{P}(r, t) = \varepsilon_0 \int_{-\infty}^t R(t - t') \mathbf{E}(r, t') dt' \leftrightarrow \mathbf{P}(r, t) = \varepsilon_0 \int_0^\infty R(\tau) \mathbf{E}(r, t - \tau) d\tau$$

Real valued

$$R(\tau) = \int_{-\infty}^\infty \chi(\omega) e^{-i\omega\tau} d\omega = \int_{-\infty}^\infty \chi^*(\omega) e^{i\omega\tau} d\omega$$

$$\rightarrow \chi(\omega) = \chi^*(-\omega)$$

Causality

$$R(\tau) = \theta(\tau)y(\tau) \text{ with } \theta(\tau) = \begin{cases} 1 & \text{for } \tau > 0 \\ 1/2 & \text{for } \tau = 0 \\ 0 & \text{for } \tau < 0 \end{cases} \text{ Heaviside distribution}$$

$$\begin{aligned} \chi(\omega) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} R(\tau) e^{i\omega\tau} d\tau = \frac{1}{2\pi} \int_{-\infty}^{\infty} \theta(\tau) y(\tau) e^{i\omega\tau} d\tau \\ &= \int_{-\infty}^{\infty} \bar{\theta}(\omega - \bar{\omega}) \bar{y}(\bar{\omega}) d\bar{\omega} \end{aligned}$$

$$2\pi \bar{\theta}(\omega) = \int_{-\infty}^{\infty} \theta(t) e^{i\omega t} dt = \lim_{\varepsilon \rightarrow 0} \frac{i}{\omega + i\varepsilon} = P \frac{i}{\omega} + \pi \delta(\omega)$$

proof in tutorials

Delta distribution: $\int_{-\infty}^{\infty} \delta(\omega - \omega_0) f(\omega) d\omega = f(\omega_0)$

Cauchy principal value:

$$P \int_{-\infty}^{\infty} d\omega \frac{i}{\omega} f(\omega) = \lim_{\alpha \rightarrow 0} \left[\int_{-\infty}^{-\alpha} d\omega \frac{i}{\omega} f(\omega) + \int_{\alpha}^{\infty} d\omega \frac{i}{\omega} f(\omega) \right]$$

$$\chi(\omega) = \frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} + \frac{\bar{y}(\omega)}{2} \quad (*)$$

(1) $y(-\tau) = y(\tau)$ even function

(2) $y(-\tau) = -y(\tau)$ odd function

(1) Even function

$$\bar{y}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\tau y(\tau) e^{i\omega\tau} = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\tau y(\tau) e^{-i\omega\tau} = \bar{y}^*(\omega)$$

$\bar{y}(\omega)$ real valued \rightarrow
$$\chi^*(\omega) = -\frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} + \frac{\bar{y}(\omega)}{2}$$

$$\chi(\omega) + \chi^*(\omega) = \frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} + \frac{\bar{y}(\omega)}{2} - \frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} + \frac{\bar{y}(\omega)}{2} = \bar{y}(\omega) \quad (\text{a})$$

$$\chi(\omega) - \chi^*(\omega) = \dots = \frac{1}{\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} \quad (\text{b})$$

inserting (a) into the right hand side of (b)

$$\Im[\chi(\omega)] = -\frac{1}{\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{\Re[\chi(\bar{\omega})]}{\bar{\omega} - \omega}$$

(2) Odd function

$$\bar{y}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\tau y(\tau) e^{i\omega\tau} = -\frac{1}{2\pi} \int_{-\infty}^{\infty} d\tau y(\tau) e^{-i\omega\tau} = -\bar{y}^*(\omega)$$

$$\chi^*(\omega) = \frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} - \frac{\bar{y}(\omega)}{2}$$

$$\chi(\omega) - \chi^*(\omega) = \frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} + \frac{\bar{y}(\omega)}{2} - \frac{1}{2\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} + \frac{\bar{y}(\omega)}{2} = \bar{y}(\omega) \quad (\text{a})$$

$$\chi(\omega) + \chi^*(\omega) = \dots = \frac{1}{\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{i \bar{y}(\bar{\omega})}{\omega - \bar{\omega}} \quad (\text{b})$$

inserting (a) into the right hand side of (b)

$$\Re[\chi(\omega)] = \frac{1}{\pi} P \int_{-\infty}^{\infty} d\bar{\omega} \frac{\Im[\chi(\bar{\omega})]}{\bar{\omega} - \omega}$$

Using

$$\chi(\omega) = \chi^*(-\omega) \quad \rightarrow \quad \chi'(\omega) = \chi'(-\omega) \quad \text{and} \quad \chi''(\omega) = -\chi''(-\omega)$$

and

$$\chi(\omega) = \varepsilon(\omega) - 1 = [\varepsilon'(\omega) - 1] + i\varepsilon''(\omega)$$

it follows for the two Kramers-Kronig relations for the permittivity that

$$\varepsilon'(\omega) - 1 = \frac{2}{\pi} P \int_0^\infty d\bar{\omega} \frac{\bar{\omega} \varepsilon''(\bar{\omega})}{\bar{\omega}^2 - \omega^2}$$

$$\varepsilon''(\omega) = -\frac{2}{\pi} \omega P \int_0^\infty d\bar{\omega} \frac{[\varepsilon'(\bar{\omega}) - 1]}{\bar{\omega}^2 - \omega^2}$$

Example: absorption line at a discrete frequency.

$$\varepsilon''(\omega) \sim \delta(\omega - \omega_0).$$

Contribution to the integral $\bar{\omega} = \omega_0$.

$$\text{Real part } \varepsilon'(\omega) - 1 \sim \frac{\omega_0}{\omega_0^2 - \omega^2}.$$

Lorentzian line

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