

Theoretical Optics

Coherence Optics

Statistical properties of light

Prof. Carsten Rockstuhl



Definitions

$u_R(\mathbf{r}, t)$ physical component of a non-monochromatic em-field at a certain position $\mathbf{r}$ and time t
real valued

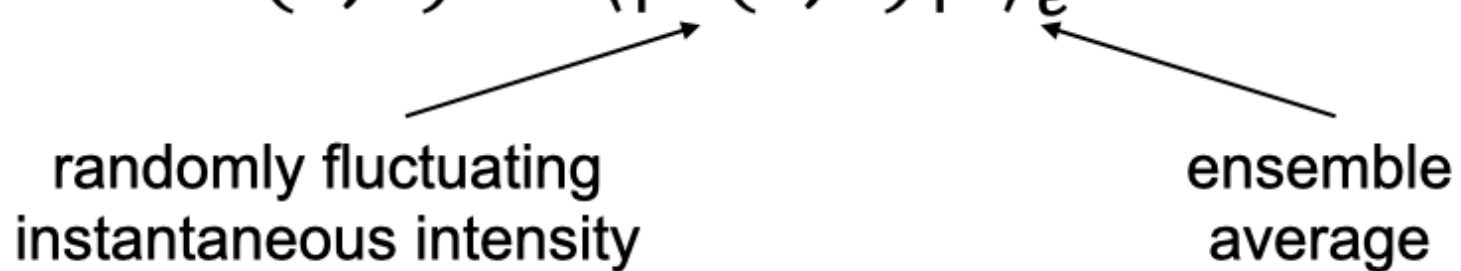
$u(\mathbf{r}, t)$ associated complex analytical signal

$$u(\mathbf{r}, t) = \int_0^\infty U(\mathbf{r}, \omega) e^{-i\omega t} d\omega \quad u_R(\mathbf{r}, t) = \Re \left[\int_0^\infty U(\mathbf{r}, \omega) e^{-i\omega t} d\omega \right]$$

Intensity

$$I(\mathbf{r}, t) = \langle |u(\mathbf{r}, t)|^2 \rangle_e$$

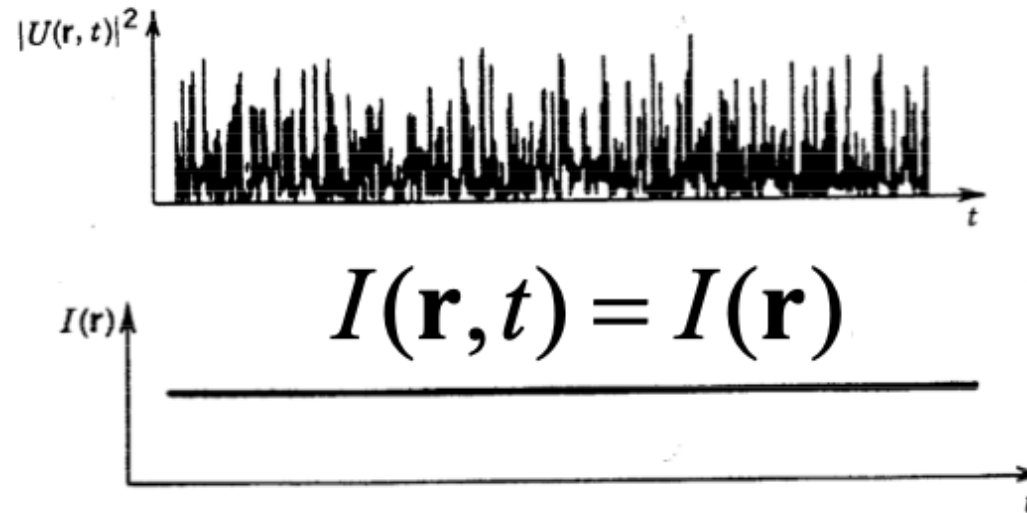
randomly fluctuating
instantaneous intensity



ensemble
average

many measurements on nominally identically prepared systems

example 1: stationary statistical light source
(light valve driven by constant current)



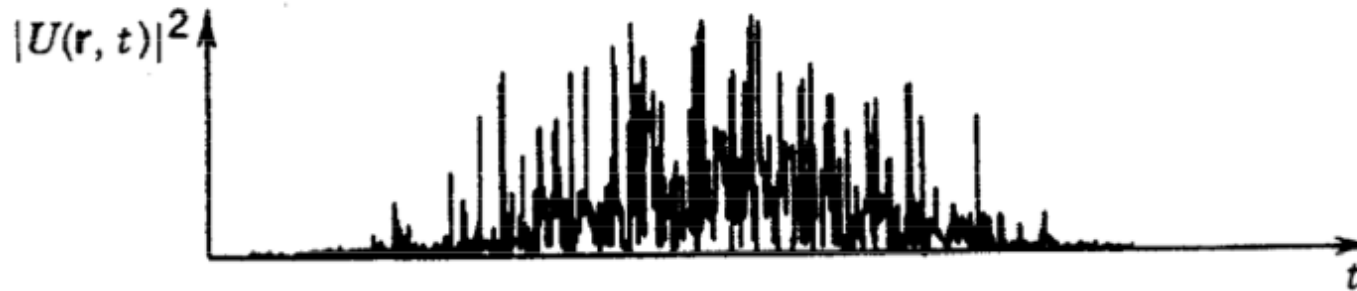
!!ergodic hypothesis applies!!

averaged over time = averaged over ensembles

$$\langle \cdots \rangle_e = \langle \cdots \rangle_t = \langle \cdots \rangle$$

$$I(\mathbf{r}) = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T |u(\mathbf{r}, t)|^2 dt$$

example 2: non-stationary statistical light source
(light valve driven by time dependent current)



Temporal coherence and spectrum

consider stationary statistical process: $\longrightarrow I(\mathbf{r})$

consider fix point in space: $\longrightarrow I(\mathbf{r}) = I$

measure for correlation of the field at $u(t)$ and $u(t + \tau)$

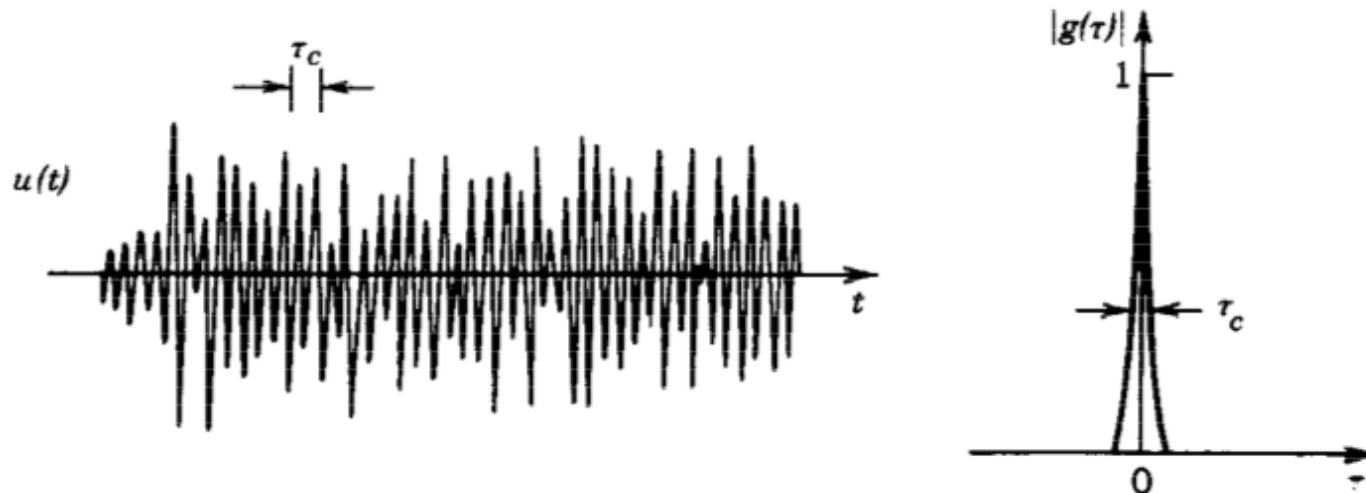
auto-correlation function

$$G(\tau) = \langle u^*(t)u(t + \tau) \rangle = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T u^*(t)u(t + \tau) dt$$

temporal coherence (self-coherence) function

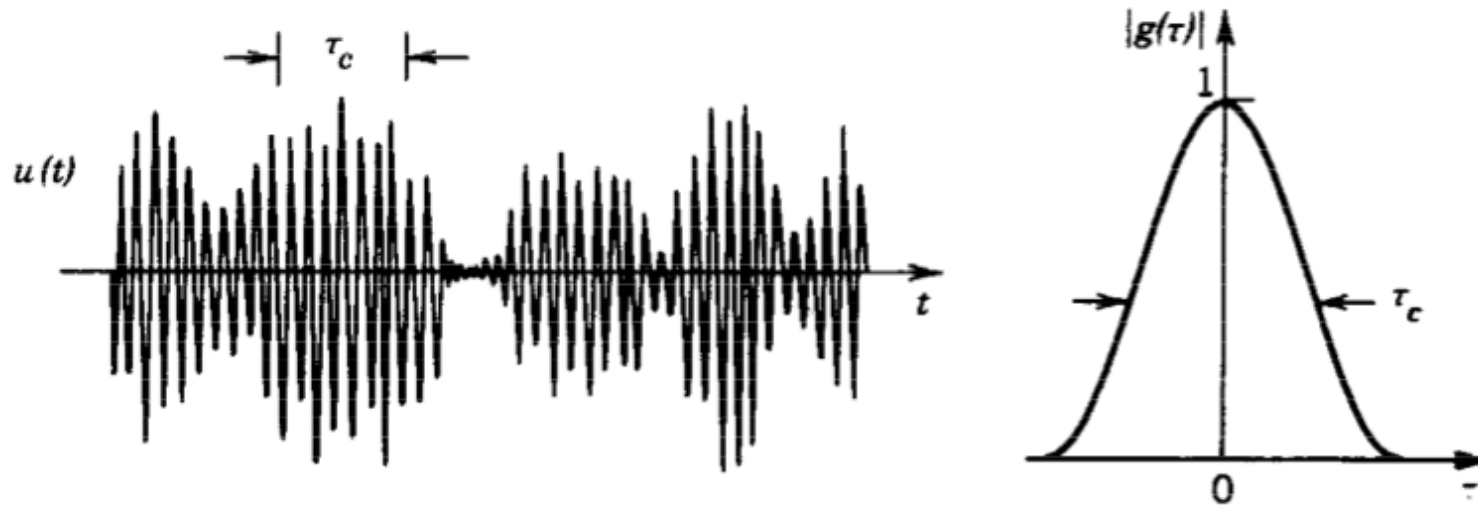
temporal width: time scale of change of envelope

coherence time τ_c



beyond this time, coherence function uncorrelated $G(\tau) = 0$

it holds: $G(0) = I$



function contains 2 information (intensity and coherence)

normalisation $\rightarrow$ complex degree of temporal coherence

$$g(\tau) = \frac{G(\tau)}{G(0)} = \frac{\langle u^*(t)u(t+\tau) \rangle}{\langle u^*(t)u(t) \rangle} \quad 0 \leq g(\tau) \leq 1$$

example 1: coherent light $u(t) = Ae^{-i\omega_0 t}$

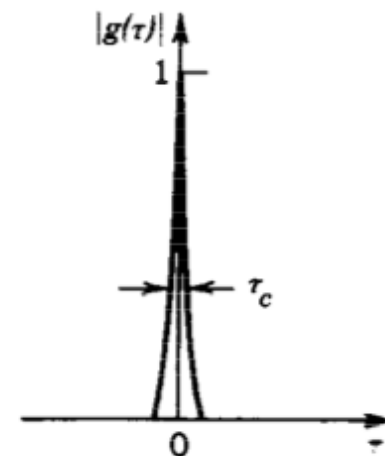
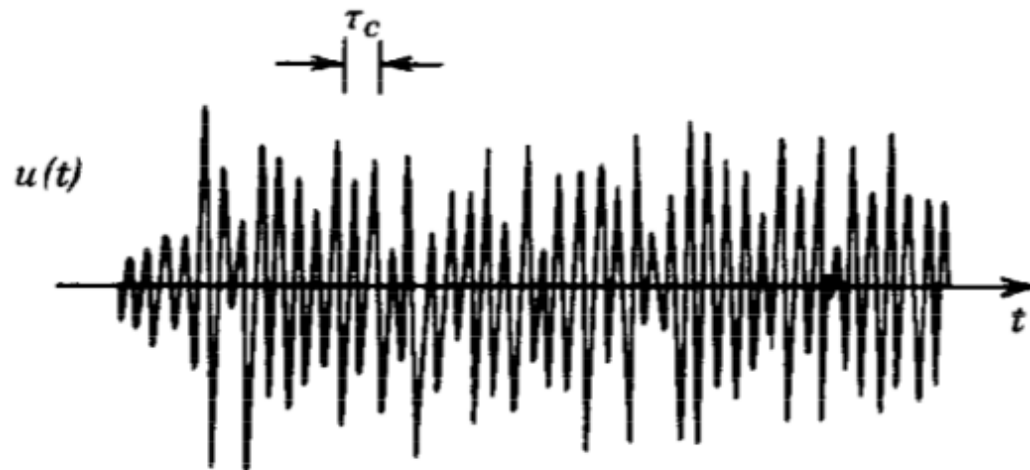
$$g(\tau) = \frac{\frac{1}{2T} \int_{-T}^T |A|^2 e^{-i\omega_0 \tau} d\tau}{\frac{1}{2T} \int_{-T}^T |A|^2 d\tau} = e^{-i\omega_0 \tau} \rightarrow \boxed{|g(\tau)| = 1}$$

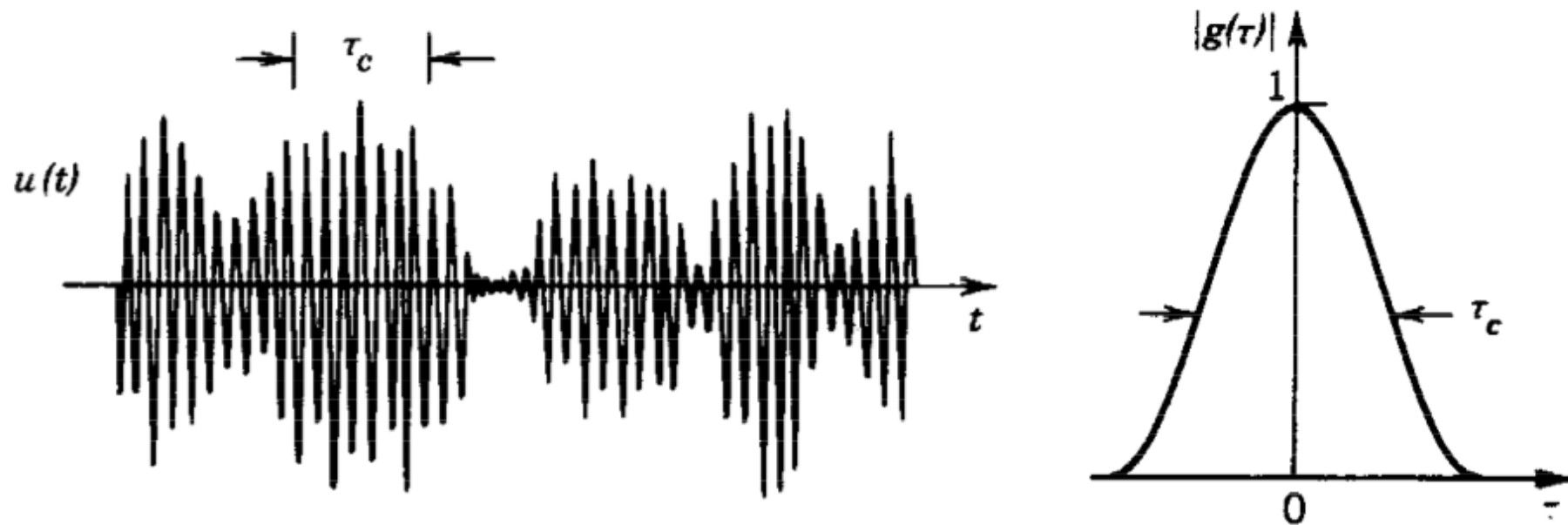
example 2: general case

$|g(\tau)|$ is a monotonously decaying function

coherence time: $|g(\tau_c)| = 1/e$

$$\boxed{\tau_c = \int_{-\infty}^{\infty} |g(\tau)|^2 d\tau}$$





coherent light: $|g(\tau)| = 1 \quad \leadsto \quad \tau_c \rightarrow \infty$

general example: $g(\tau) = e^{-\frac{|\tau|}{A}}$

$$\longrightarrow \tau_c = 2 \int_0^{\infty} e^{-2\frac{|\tau|}{A}} d\tau = A(1 - 0) = A$$

equivalent coherence length: $l_c = c\tau_c$

all length scales of optical system $<$ than coherence length

$\longrightarrow$ system coherent

Theorem of Wiener-Khinchin

dependency between spectrum and temporal coherence

$$G(\tau) = \int_{-\infty}^{\infty} S(\omega) e^{-i\omega\tau} d\omega$$

signal: $S(\omega) = |U(\omega)|^2$

Fourier transformation of the temporal coherence function
corresponds to signal spectrum

proof: $G(\tau) = \langle u^*(t)u(t+\tau) \rangle = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T u^*(t)u(t+\tau) dt$

finite FT

$$U_T(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} u_T(t) \exp(\mathbf{i}\omega t) dt \quad u_T(t) = \begin{cases} u(t) & \text{for } |t| \leq T \\ 0 & \text{otherwise} \end{cases}$$

$$u_{\mathbf{T}}(t) = \int_{-\infty}^{\infty} U_{\mathbf{T}}(\omega) \exp(-\mathbf{i}\omega t) d\omega$$

$$G(\tau) = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-T}^T U_T^*(\omega') e^{i\omega' t} U_T(\omega'') e^{-i\omega''(t+\tau)} d\omega' d\omega'' dt$$

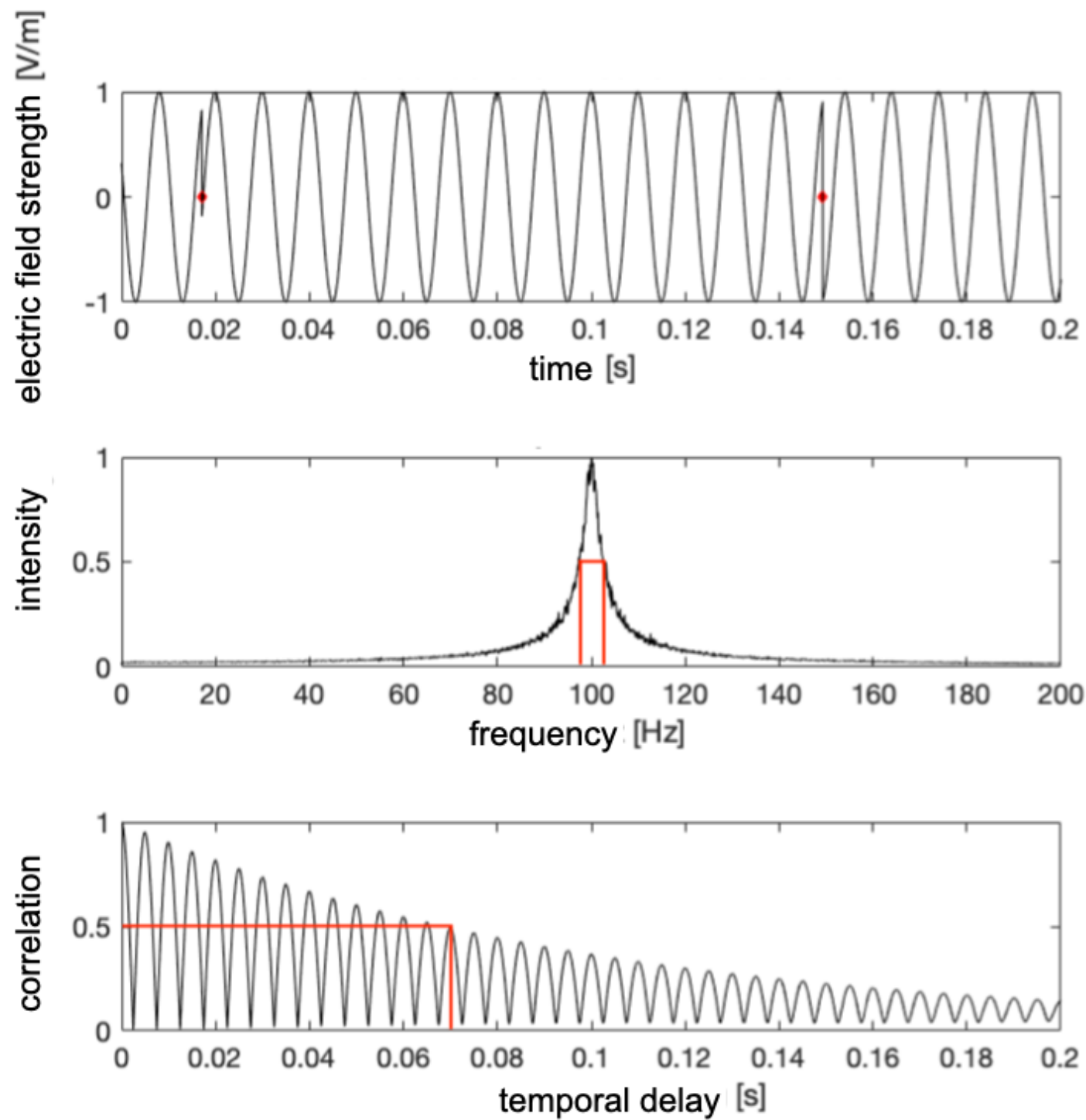
$$= \lim_{T \rightarrow \infty} \frac{1}{2T} 2\pi \int_{-\infty}^{\infty} |U_T(\omega')|^2 e^{-i\omega' \tau} d\omega'$$

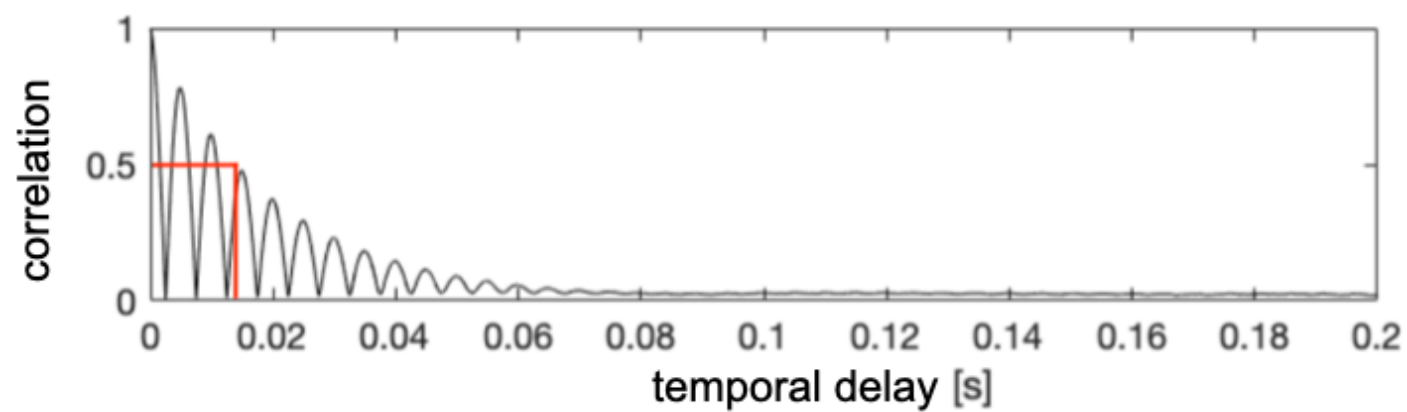
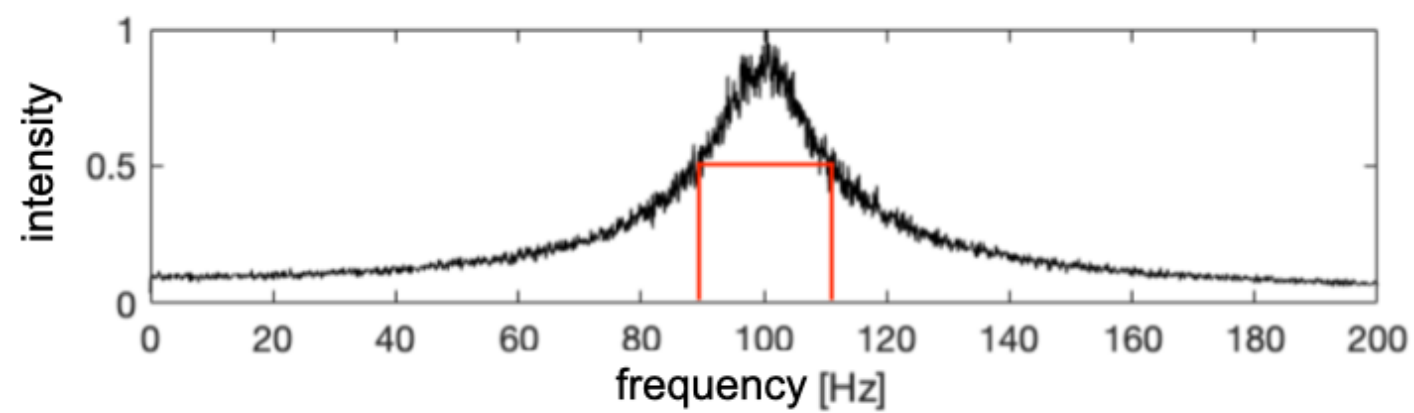
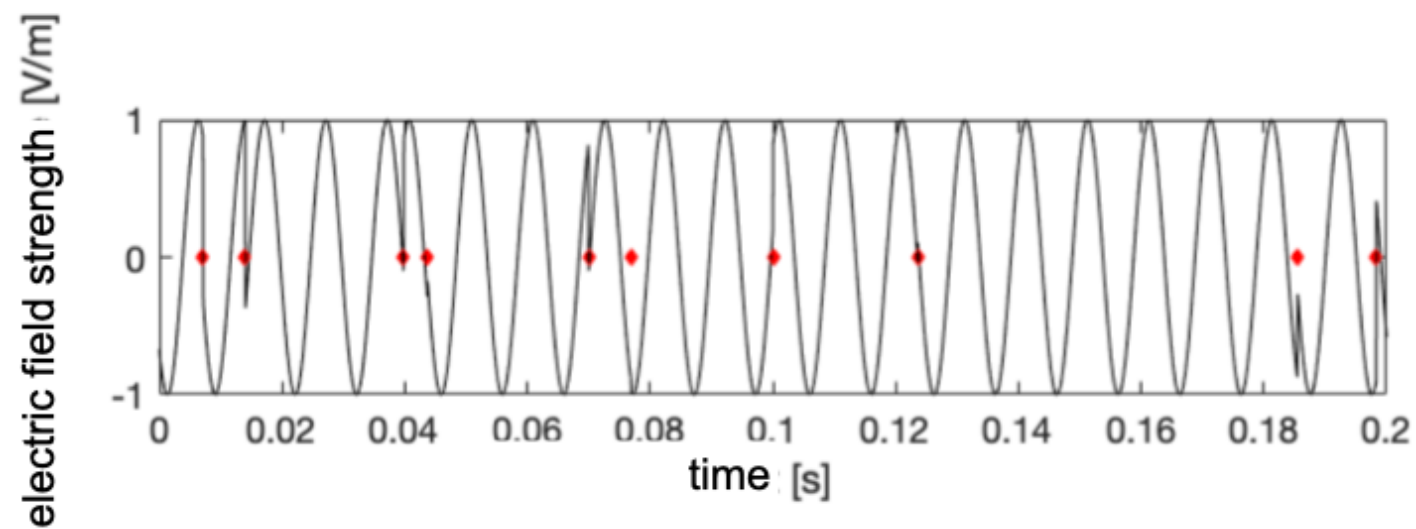
$$\curvearrowright S(\omega) = 2\pi \lim_{T \rightarrow \infty} \frac{1}{2T} |U_{\mathbf{T}}(\omega)|^2$$

$S(\omega)d\omega \rightarrow$ spectral intensity in the interval between ω and $\omega + d\omega$

$$G(0) = I \longrightarrow I = \int_0^{\infty} S(\omega) d\omega$$

coherence of light source is given by spectral intensity





product between coherence time and bandwidth is constant

$$\tau_c \Delta\omega = \text{const.}$$

spectral width:
$$\Delta\omega = \frac{\left[\int_0^\infty S(\omega) d\omega \right]^2}{\int_0^\infty S^2(\omega) d\omega}$$

spectral filter can improve coherence

loss of intensity

TABLE 10.1-2 Spectral Widths of a Number of Light Sources Together with Their Coherence Times and Coherence Lengths in Free Space

Source	$\Delta\nu_c$ (Hz)	$\tau_c = 1/\Delta\nu_c$	$l_c = c\tau_c$
Filtered sunlight ($\lambda_o = 0.4\text{--}0.8\ \mu\text{m}$)	3.75×10^{14}	2.67 fs	800 nm
Light-emitting diode ($\lambda_o = 1\ \mu\text{m}$, $\Delta\lambda_o = 50\ \text{nm}$)	1.5×10^{13}	67 fs	20 μm
Low-pressure sodium lamp	5×10^{11}	2 ps	600 μm
Multimode He-Ne laser ($\lambda_o = 633\ \text{nm}$)	1.5×10^9	0.67 ns	20 cm
Single-mode He-Ne laser ($\lambda_o = 633\ \text{nm}$)	1×10^6	1 μs	300 m

Mutual coherence function

(spatial and temporal coherence)

cross-correlation at 2 spatial and 2 temporal positions

$u(\mathbf{r}_1, t_1)$ and $u(\mathbf{r}_2, t_2)$ with $\tau = t_2 - t_1$

$$G(\mathbf{r}_1, \mathbf{r}_2, \tau) = \langle u^*(\mathbf{r}_1, t) u(\mathbf{r}_2, t + \tau) \rangle$$

complex degree of mutual coherence

$$g(\mathbf{r}_1, \mathbf{r}_2, \tau) = \frac{G(\mathbf{r}_1, \mathbf{r}_2, \tau)}{\sqrt{I(\mathbf{r}_1)I(\mathbf{r}_2)}}$$

$$0 \leq |g(\mathbf{r}_1, \mathbf{r}_2, \tau)| \leq 1$$

example: plane monochromatic wave

$$|g(\mathbf{r}_1, \mathbf{r}_2, \tau)| = |e^{i\mathbf{k}(\mathbf{r}_1 - \mathbf{r}_2) - i\omega\tau}| = 1$$

special: temporal coherence $\mathbf{r}_1 = \mathbf{r}_2 = \mathbf{r}$

spatial coherence $\tau = 0$

Spatial coherence function

spatial case of mutual coherence function while assuming $\tau = 0$

$$G(\mathbf{r}_1, \mathbf{r}_2, 0) = \langle u^*(\mathbf{r}_1, t)u(\mathbf{r}_2, t) \rangle = G(\mathbf{r}_1, \mathbf{r}_2)$$

→ normalised mutual intensity

$$g(\mathbf{r}_1, \mathbf{r}_2) = \frac{G(\mathbf{r}_1, \mathbf{r}_2)}{\sqrt{I(\mathbf{r}_1)I(\mathbf{r}_2)}}$$

Van Cittert–Zernike theorem: equivalent to Wiener-Khinchin theorem in spatial domain

relation between spatial degree of coherence $|g(\mathbf{r}_1, \mathbf{r}_2)|$
and intensity distribution of an aperture $I(x', y')$

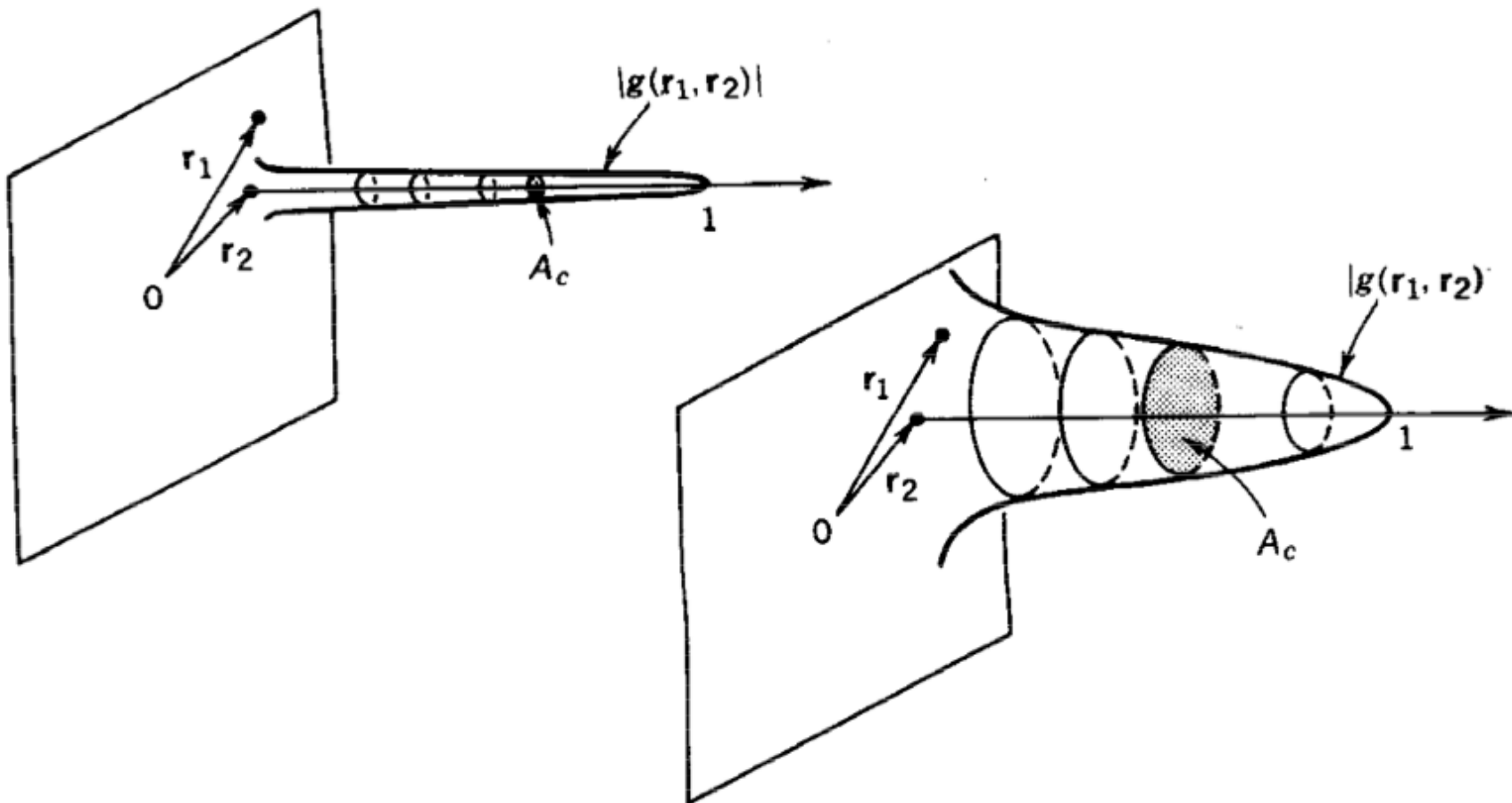
$$|g(\mathbf{r}_1, \mathbf{r}_2)| = \left| e^{ik \frac{(x_1^2 + y_1^2) - (x_2^2 + y_2^2)}{2z}} \cdot \frac{\iint_{\sigma} I(x', y') \cdot e^{-i2\pi(px' + qy')} dx' dy'}{\iint_{\sigma} I(x', y') dx' dy'} \right|$$

$$p = \frac{k}{2\pi} \frac{x_1 - x_2}{z} \quad q = \frac{k}{2\pi} \frac{y_1 - y_2}{z}$$

mutual intensity = prefactor • Fourier transform of intensity in aperture normalisation

**circular
aperture:**

$$|g(\mathbf{r}_1, \mathbf{r}_2)| \propto \left| 2J_1 \left(\frac{2\pi \cdot \rho \cdot |\mathbf{r}_1 - \mathbf{r}_2|}{\lambda_0 \cdot z} \right) / \left(\frac{2\pi \cdot \rho \cdot |\mathbf{r}_1 - \mathbf{r}_2|}{\lambda_0 \cdot z} \right) \right|$$



mean spatial extent
within which the field is
coherent

$$J_1 \left(\frac{2\pi \cdot \rho \cdot |\mathbf{r}_1 - \mathbf{r}_2|}{\lambda_0 \cdot z} \right) = 0 \longrightarrow |\mathbf{r}_1 - \mathbf{r}_2| = 1,22 \cdot \frac{\lambda_0 \cdot z}{2 \cdot \rho}$$

mean size of the
coherent area

$$\bar{x} = 1,22 \cdot \frac{\lambda_0 \cdot z}{2 \cdot \rho} \longrightarrow$$

$$\bar{x} = 1,22 \cdot \frac{\lambda_0}{2 \cdot \tan\left(\frac{\alpha}{2}\right)}$$

angular extent of light source

Theoretical Optics

Coherence Optics

Statistical properties of light

Prof. Carsten Rockstuhl

