

Theoretische Teilchenphysik II

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Exercise Sheet 10

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Renormalization group equations (RGEs) and the resummation of large logarithms (100 Points)

Exercise 10.1: (30 points) Recall from lecture 12 (c.f. Eq. (12.33)) that the QCD coupling constant α_s satisfies the following differential equation

$$\frac{d\alpha_s}{d \ln \mu} = \frac{-2\epsilon \alpha_s}{1 + \alpha_s \frac{1}{Z_{\alpha_s}} \frac{\partial Z_{\alpha_s}}{\partial \alpha_s}}, \quad (1)$$

where

$$Z_{\alpha_s} = 1 + \sum_{n=1}^{\infty} \frac{A_n}{\epsilon^n}, \quad (2)$$

and A_n are functions of α_s .

(a) Explain steps that are needed to rewrite Eq. (1) in the following way

$$\begin{aligned} \mu \frac{d\alpha_s}{d\mu} &= -2\epsilon\alpha_s + 2\alpha_s^2 A'_1, \\ A'_2 &= \alpha_s (A'_1)^2 + A_1 A'_1, \\ A'_3 &= \alpha_s A'_1 A'_2 + A_2 A'_1, \\ A'_4 &= \alpha_s A'_3 A'_1 + A_3 A'_1. \\ &\dots \end{aligned} \quad (3)$$

Assume that

$$A_1 = \sum_{n=1}^{\infty} a_{1n} \alpha_s^n. \quad (4)$$

Taking into account of a_{11} *only* write explicitly the solution to α_s and determine $A_{2,3,4,\dots}$ in Eq. (3). Repeat the same by accounting for a_{11} and a_{12} .

(b) We defined the anomalous dimension in the lecture 13 as (c.f. Eq. (13.6), the subscript is dropped)

$$\gamma = \frac{1}{Z} \frac{d}{d \ln \mu} Z, \quad (5)$$

where Z is also the same form as Z_{α_s} ,

$$Z = 1 + \sum_{n=1}^{\infty} \frac{B_n}{\epsilon^n}. \quad (6)$$

Find a similar structure as in Eq. (3) in this case.

Exercise 10.2: (40 points) Renormalization group equations provide a powerful way to study the structure of perturbation theory and, especially, large logarithms of kinematic variables that appear in arbitrary Green's functions. Such logarithms typically arise in high-order perturbative calculation.

To be more specific, consider renormalized four-point Green's function at one-loop in ϕ^4 theory,¹

$$G^{(4)}(\{p_i\}) = -i\lambda(m) \left(1 + \frac{\lambda(m)}{32\pi^2} [3L + f_1] \right) + \mathcal{O}(\lambda(m)^3), \quad E = \frac{\sqrt{s}}{2}, \quad (7)$$

where $L = \ln E^2/m^2$. This equation is valid for $s \sim |t| \sim |u| \gg m^2$ and the function f_1 does not contain the large logarithm L .

It is clear from the above equation that the perturbative expansion parameter for $G^{(4)}(\{p_i\})$ is not $\lambda(m)$ but rather $\lambda(m)L \gg \lambda(m)$ provided that $E \gg m$. Generally, it has the following form,

$$G^{(4)}(\{p_i\}) = -i\lambda(m) \sum_{n=0}^{\infty} \left(\frac{\lambda(m)}{32\pi^2} \right)^n \sum_{k=0}^n c_{n,k} L^k. \quad (8)$$

In each order of the perturbative expansion, the leading logarithm (LL) has the largest power followed by the next-to-leading (or sub-leading) one (NLL), etc (N^k LL). It is possible to use renormalization group analysis to re-sum these large logarithms to all orders in the perturbative expansion. Below we discuss how to do this.

- (a) Compute the β -function and the field anomalous dimensions in ϕ^4 theory at one loop in the $\overline{\text{MS}}$ scheme. Write an equation for the running coupling constant $\lambda(\mu)$ and solve it using $\lambda(\mu) = \lambda(m)$ at $\mu = m$ as the boundary condition.
- (b) Use the above result to write $\lambda(m)$ as a series expansion in $\lambda(E)$. Substitute it back to Eq. (7) and show that the Green's function assumes the following form

$$G^{(4)}(\{p_i\}) = -i\lambda(E) \left(1 + \frac{\lambda(E)}{32\pi^2} f_1 \right) + \mathcal{O}(\lambda(m)^3). \quad (9)$$

Why it is advantageous to use this expression for the Green's function as compared to the expression in Eq. (7)?

- (c) Use the results that you derived in item (a) to set up a RGE for the Green's function and explain how to solve it using the one-loop expressions for the β -function and the field anomalous dimension. Use the solution to explain how the LLs are re-summed to all orders.
- (d) To re-sum the NLLs, what ingredients do you need? Explain it.

Exercise 10.3: (30 points) In the QCD and ϕ^4 theory, there is only one coupling constant. Let's generalize to the case with multiple coupling constants. Consider a theory defined by the Lagrangian,

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi_1 \partial^\mu \phi_1 + \frac{1}{2} \partial_\mu \phi_2 \partial^\mu \phi_2 - \frac{\lambda_1}{4!} \phi_1^4 - \frac{\lambda_2}{4!} \phi_2^4 - \frac{\lambda_3}{4} \phi_1^2 \phi_2^2. \quad (10)$$

- (a) Draw all the Feynman diagrams that are needed to compute the three β -functions of the coupling constants, $\lambda_{1,2,3}$ at leading order.
- (b) Compute these β -functions in the $\overline{\text{MS}}$ scheme.
- (c) Try to solve the RGEs of $\lambda_{1,2,3}$ numerically and make stream plots to illustrate the solutions.

¹Note that the physical coupling constant $\lambda(m)$ is defined by the scattering amplitude of four particles at rest

$$\lambda(m) = i\mathcal{M}_{\phi\phi \rightarrow \phi\phi}(s = 4m^2, t = u = 0).$$