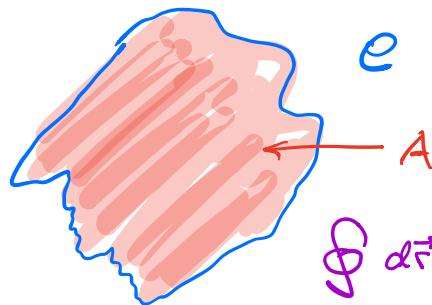


$$\begin{aligned}
 \vec{F} &= m \frac{d^2 \vec{r}}{dt^2} \\
 &\vdots \\
 \frac{m}{2} \vec{v}(t)^2 - \frac{m}{2} \vec{v}(t_0)^2 &= \int_{t_0}^t \vec{F} \cdot \underline{\frac{d\vec{r}}{dt'}} dt' = - \int \frac{du}{d\vec{r}} \frac{d\vec{r}}{dt} dt \\
 &= - \int du \\
 &= -u(t) + u(t_0)
 \end{aligned}$$

$$\vec{F}_m = - \vec{\nabla} u(\vec{r})$$

$$\vec{\nabla} u(\vec{r}) = \left(\frac{\partial u(x,y,z)}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z} \right)$$

Stokes'sches Satz

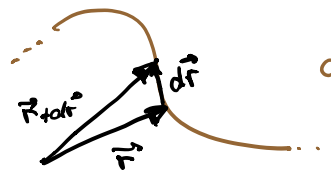


$$\oint_C d\vec{r} \cdot \vec{f}(\vec{r}) = \int_S d\vec{s} \cdot (\vec{\nabla} \times \vec{f})$$

i) $\oint$ Integration entlang einer geschlossenen Kurve



ii) $d\vec{r} \cdot \vec{f}(\vec{r})$

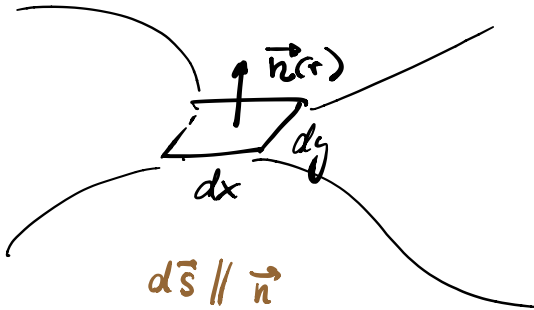


$$d\vec{r} \cdot \vec{f}(\vec{r}) = dr |\vec{f}(\vec{r})| \cos \theta(r)$$

$$dr = |d\vec{r}|$$



iii) $\int_A d\vec{s} \cdot \vec{\sigma}(\vec{r})$

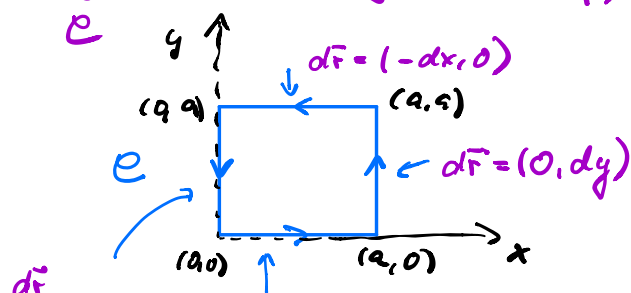


$d\vec{s} \parallel \vec{n}$

- bestimme an jedem Punkt $\vec{n}(\vec{r})$ den ($|\vec{n}|=1$) lokalen Normalenvektor der Fläche

$$\int_A d\vec{s} \cdot \vec{\sigma} = \int dx dy \vec{n} \cdot \vec{\sigma}$$

$$\oint_C d\vec{r} \cdot \vec{f}(\vec{r}) = \int d\vec{s} \cdot (\vec{\nabla} \times \vec{f})$$



$$d\vec{r} = (dx, 0) : d\vec{r} \cdot \vec{f}(\vec{r}) = (dx, 0) \begin{pmatrix} f_x(x, 0) \\ f_y(x, 0) \end{pmatrix} = f_x(x, 0) dx$$

$$\oint_C d\vec{r} \cdot \vec{f}(\vec{r}) = \int_0^a dx (f_x(x, 0) - f_x(x, a)) + \int_0^a dy (f_y(a, y) - f_y(0, y))$$

$$\oint_C d\vec{r} \cdot \vec{f}(\vec{r}) = - \oint_C d\vec{r} \cdot \vec{f}(\vec{r})$$

$$\int_A d\vec{s} = \vec{e}_2 \int_0^a dx \int_0^a dy \quad \vec{e}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

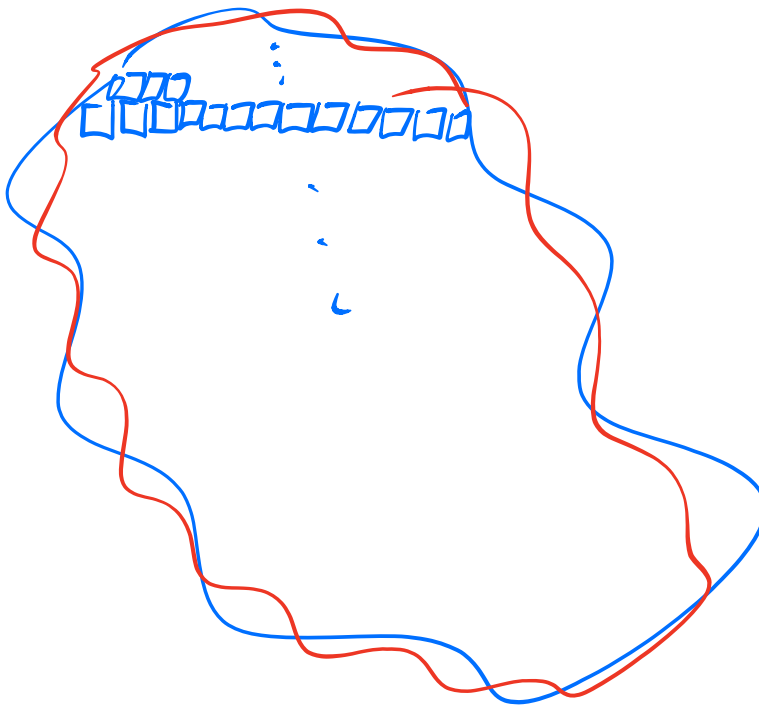
$$\begin{aligned} \int d\vec{s} \cdot (\vec{\nabla} \times \vec{f}) &= \int_0^a dx \int_0^a dy (\vec{\nabla} \times \vec{f})_z \\ &= \int_0^a dx \int_0^a dy \left(\frac{\partial f_y}{\partial x} - \frac{\partial f_x}{\partial y} \right) \end{aligned}$$

$$\int_0^a dx \frac{\partial f_y(x, y)}{\partial x} = f_y(a, y) - f_y(0, y)$$

$$\int_0^a dy \frac{\partial f_x(x, y)}{\partial y} = f_x(x, a) - f_x(x, 0)$$

$$\begin{aligned} &= \int_0^a dy (\underline{f_y(a, y)} - \underline{f_y(0, y)}) \\ &\quad - \int_0^a dx (\underline{f_x(x, a)} - \underline{f_x(x, 0)}) \end{aligned}$$

$$\oint_C d\vec{r} \cdot \vec{f}(\vec{r}) = \int_0^a dx \left(\underbrace{f_x(x, 0)}_{\text{red}} - \underbrace{f_x(x, a)}_{\text{green}} \right) + \int_0^a dy \left(\underbrace{f_y(a, y)}_{\text{magenta}} - \underbrace{f_y(0, y)}_{\text{blue}} \right)$$

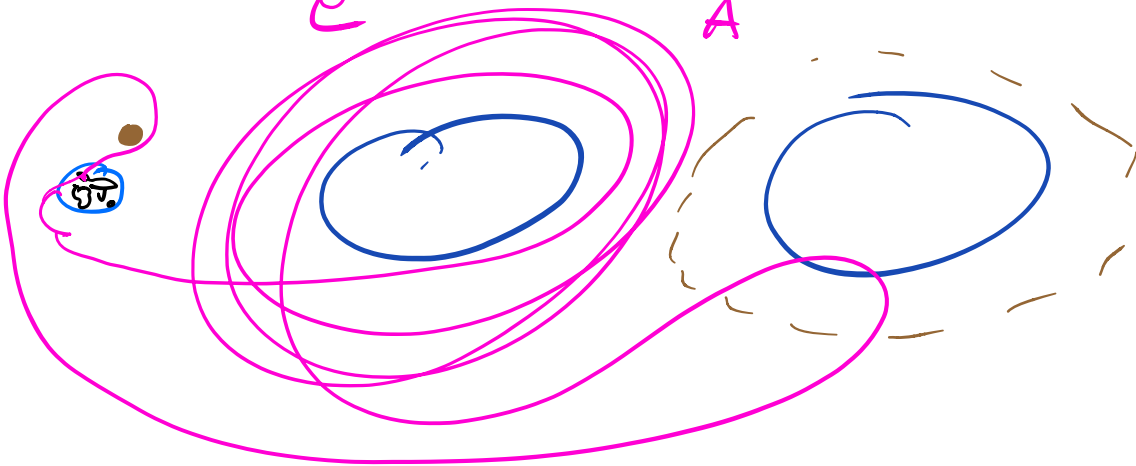


$$\begin{aligned} \frac{m}{2} \dot{\vec{r}}(t)^2 - \frac{m}{2} \dot{\vec{r}}(t_0)^2 &= \int_{t_0}^t \vec{F} \cdot \frac{d\vec{r}}{dt'} dt' \\ &= \int_{\vec{r}(t_0)}^{\vec{r}(t)} \vec{F} \cdot d\vec{r} = \text{Arbeit} \end{aligned}$$

$$\vec{F} = -\vec{\nabla} u$$

Kraft entlang einer geschlossenen
Kurve

$$W = \oint_C \vec{F} \cdot d\vec{r} = \int_A d\vec{r} \cdot \vec{\nabla} \times \vec{F} = 0$$



$$\begin{aligned} \vec{\nabla} \times \vec{F} &= \sum_{ijk} \vec{e}_i \frac{\partial}{\partial x_j} F_k \varepsilon_{ijk} & F_k &= -\frac{\partial u}{\partial x_k} \\ &= - \sum_{ijk} \vec{e}_i \frac{\partial}{\partial x_j} \frac{\partial}{\partial x_k} \varepsilon_{ijk} u(\vec{r}) \\ &= - \sum_{ijk} \vec{e}_i \frac{\partial}{\partial x_k} \frac{\partial}{\partial x_j} \varepsilon_{ikj} u(\vec{r}) \\ &= \sum_{ijk} \vec{e}_i \frac{\partial}{\partial x_j} \frac{\partial}{\partial x_k} \varepsilon_{ijk} u(\vec{r}) = -\vec{\nabla} \times \vec{F} \\ &\Rightarrow \vec{0} = \vec{\nabla} \times \vec{F} \text{ wenn } \vec{F} = -\vec{\nabla} u \end{aligned}$$

Energieerhaltung $\Leftrightarrow$ Homogenität d.

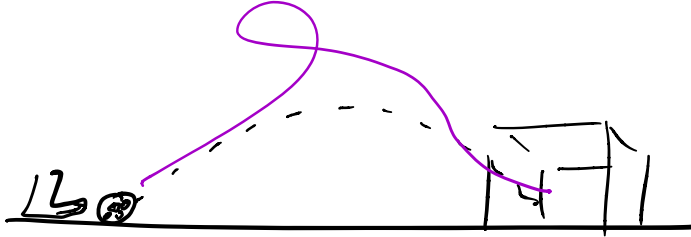
Zeit

$$U(\vec{r}, t) = U(\vec{r})$$



Arbeit ist Weg unabhängig

Das Variationsprinzip der Mechanik



Welches Prinzip entscheidet über die richtige Trajektorie?

Annahme:

$$S[\vec{r}(t)] = \int_{t_0}^t dt' L(\vec{r}(t'), \dot{\vec{r}}(t'), t')$$

physikalische Trajektorien minimieren S

betrachten eine räumliche Dimension

$$S[x(t)] = \int_{t_0}^t L(x(t'), \dot{x}(t'), t') dt'$$

Wir kennen $x(t_0) = x_0$ $x(t) = x_1$

$\tilde{x}(t')$ sei die physikalisch realisierte Trajektorie

$$\tilde{x}(t_0) = x_0 \quad \tilde{x}(t) = x_1$$

$$x(t') = \tilde{x}(t') + \delta x(t')$$



$$\delta x(t_0) = \delta x(t) = 0$$

$$\Delta S = \int_{t_0}^t L(\tilde{x} + \delta x, \dot{\tilde{x}} + \delta \dot{x}, t') dt' - \int_{t_0}^t L(\tilde{x}, \dot{\tilde{x}}, t') dt'$$

$$\begin{aligned} & L(\tilde{x}(t') + \delta x(t'), \dot{\tilde{x}}(t') + \delta \dot{x}(t'), t') \\ & \approx L(\tilde{x}, \dot{\tilde{x}}, t') + \frac{\partial L}{\partial \tilde{x}} \delta x + \frac{\partial L}{\partial \dot{\tilde{x}}} \delta \dot{x} \end{aligned}$$

$$\begin{aligned} & \int_{t_0}^t \frac{\partial L}{\partial \dot{\tilde{x}}} \frac{d}{dt'} \delta x(t') dt' \\ & = - \int_{t_0}^t \left(\frac{d}{dt'} \frac{\partial L}{\partial \dot{\tilde{x}}} \right) \delta x(t') dt' \end{aligned}$$

$$\Delta S = \int_{t_0}^t \left(\frac{\partial L}{\partial \tilde{x}} - \frac{d}{dt'} \frac{\partial L}{\partial \dot{\tilde{x}}} \right) \delta x(t') dt'$$

= 0

Euler-Lagrange
Gleichungen!