

Bestimmung der Green'schen Funktion

$$\frac{d^2x(t)}{dt^2} + \gamma \frac{dx(t)}{dt} + w_0^2 x(t) = \delta(t-t')$$

$$x(t < t') = 0$$

$$\left. \frac{dx}{dt} \right|_{t < t'} = 0$$

$x(t)$ sei endlich und stetig

$$x_\epsilon(t) = G(t-t') \iff$$

$$\int_{t'-\epsilon}^{t'+\epsilon} dt \left(\frac{d^2}{dt^2} + \gamma \frac{d}{dt} + w_0^2 \right) G(t-t') = \int_{t'-\epsilon}^{t'+\epsilon} \delta(t-t') dt$$

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die linke Seite:

$$\begin{aligned} i) \lim_{\epsilon \rightarrow 0} \int_{t'-\epsilon}^{t'+\epsilon} w_0^2 G(t-t') dt &= w_0^2 G(0) \lim_{\epsilon \rightarrow 0} \int_{t'-\epsilon}^{t'+\epsilon} dt \\ &= w_0^2 G(0) \lim_{\epsilon \rightarrow 0} 2\epsilon = 0 \end{aligned}$$

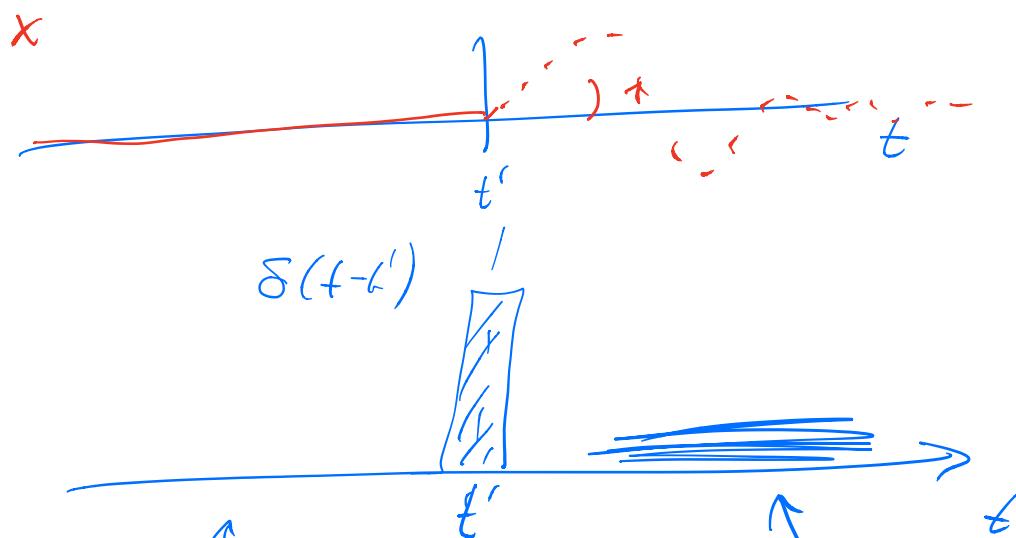
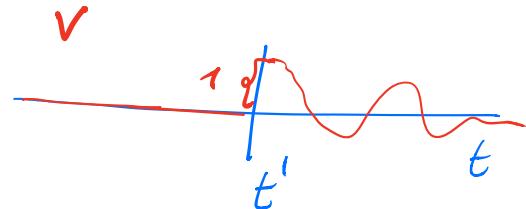
$$\begin{aligned} ii) \lim_{\epsilon \rightarrow 0} \gamma \int_{t'-\epsilon}^{t'+\epsilon} \frac{d}{dt} G(t-t') dt &= \gamma \lim_{\epsilon \rightarrow 0} [G(\epsilon) - G(-\epsilon)] \\ &= 0 \end{aligned}$$

$$\text{iii) } \lim_{\varepsilon \rightarrow 0} \int_{t'-\varepsilon}^{t+\varepsilon} dt \frac{d^2 G(t-t')}{dt^2} = \lim_{\varepsilon \rightarrow 0} \left(\frac{dG(t)}{dt} \Big|_{t=\varepsilon} - \frac{dG(t)}{dt} \Big|_{t=-\varepsilon} \right)$$

$$\Rightarrow \lim_{\varepsilon \rightarrow 0} \left(\frac{dG(t)}{dt} \Big|_{t=\varepsilon} - \frac{dG(t)}{dt} \Big|_{t=-\varepsilon} \right) = 1$$

||
0 Am Anfang in
Ruhe!

$$\boxed{\frac{dG(t)}{dt} \Big|_{t=0^+} = 1}$$



✓ wir nennen die
Lösung

$$G(t < t') = 0, \quad \frac{dG(t < t')}{dt} = 0$$

wir nennen die
Lösung
(homogenes Problem)

$$x_{\text{hom}}(t) = e^{-\frac{\Omega}{2}t} (a \cos \omega t + b \sin \omega t)$$

$$\Omega = \sqrt{\omega_0^2 - \frac{\delta^2}{4}}$$

\Rightarrow

$$G(t-t') = e^{-\frac{\Omega}{2}(t-t')} \left(a \cos[\Omega(t-t')] + b \sin[\Omega(t-t')] \right)$$

$$G(0) = 0 \quad \Rightarrow \quad a = 0$$

$$\left. \frac{dG(t)}{dt} \right|_{t=0} = 1 = b \Omega$$

$$\begin{aligned} \frac{dX_{\text{hom}}}{dt} &= -\frac{\Omega}{2} e^{-\frac{\Omega}{2}t} b \sin(\omega t) \\ &\quad + b \Omega e^{-\frac{\Omega}{2}t} \cos(\omega t) \end{aligned}$$

$$b \Omega = b \Omega$$

$$G(t-t') = \theta(t-t') \frac{1}{\Omega} e^{-\frac{\Omega}{2}(t-t')} \sin(\omega(t-t'))$$

$$f(t) = f_0 \cos(\omega t) = \frac{f_0}{2} (e^{i\omega t} + e^{-i\omega t})$$

$$\bar{x}(t) = \int_{-\infty}^{\infty} dt' G(t-t') f(t')$$

$$= \frac{f_0}{2} \int_{-\infty}^{\infty} dt' G(t-t') e^{i\omega t'} \quad t-t' > 0$$

$$= \frac{f_0}{2} \int_{-\infty}^t dt' \frac{1}{\pi} e^{-\frac{i}{2}(t-t')} \sin(2\pi(t-t')) e^{i\omega t'}$$

$$= \frac{f_0}{2} \int_{-\infty}^t dt' \frac{1}{2\pi i} e^{-\frac{i}{2}(t-t')} \left(e^{i\omega(t-t')} - e^{-i\omega(t-t')} \right) e^{i\omega t'}$$

$$s = t' - t$$

$$= \frac{f_0}{2} \int_{-\infty}^0 ds \frac{1}{2\pi i} e^{+\frac{i}{2}s} \left(e^{-i\omega s} - e^{+i\omega s} \right) e^{i\omega s} e^{i\omega t}$$

$$\int_{-\infty}^0 ds e^{+\alpha s + i\beta s} = \frac{1}{\alpha + i\beta}$$

$$= \frac{f_0}{2} \frac{1}{\alpha i \omega} e^{i\omega t} \left(\frac{1}{\frac{1}{2} + i(\omega - \alpha)} - \frac{1}{\frac{1}{2} + i(\omega + \alpha)} \right)$$

$$= \frac{f_0}{2} \frac{1}{\omega^2 - \omega^2} e^{i\omega t} \frac{2i\omega}{\frac{\omega^2}{4} + i\gamma\omega - \omega^2 + \omega^2}$$

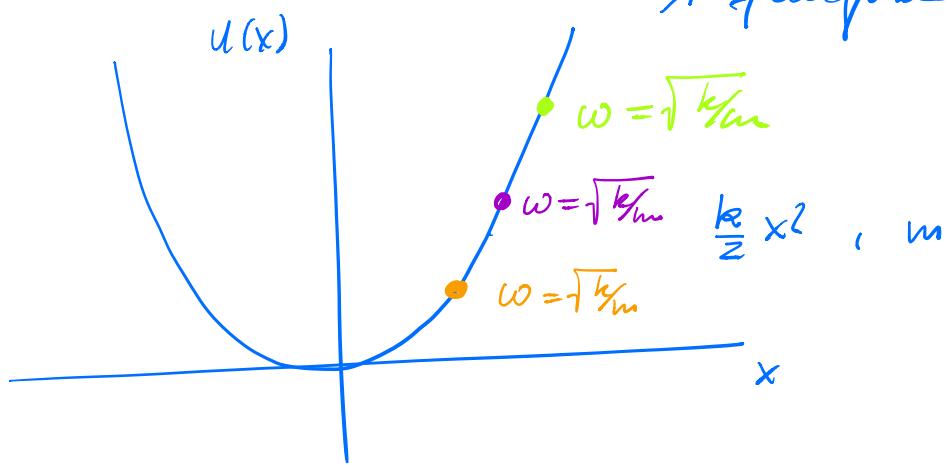
$$\omega^2 = \omega_0^2 - \gamma^2/4$$

$$\bar{x}_{inh}(t) = \frac{f_0}{2} e^{i\omega t} \frac{1}{\omega_0^2 - \omega^2 + i\gamma\omega}$$

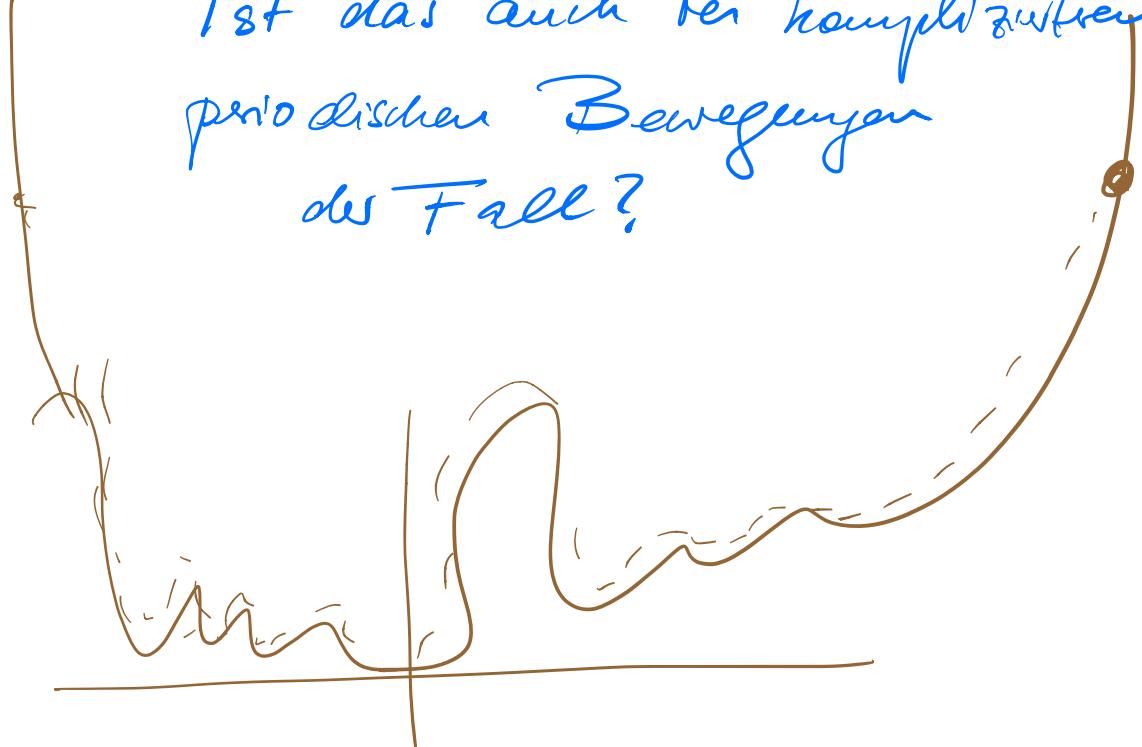
$$= f_0 e^{i\omega t} A(\omega)$$

✓

die Frequenz des harmonischen Oszillators ist
unabh. von den Anfangsbedingungen!



Ist das auch bei komplizierteren
periodischen Bewegungen
der Fall?



Anharmonische eindimensionale Oszillationen

$$\frac{m}{2} \left(\frac{dx}{dt} \right)^2 + U(x) = E$$

$$\frac{dx}{dt} = \pm \sqrt{\frac{2}{m} (E - U(x))}$$

es muß gelten, dass $U(x) \leq E$

$$\pm \int_{x_0}^x \frac{dx'}{\sqrt{\frac{2}{m} (E - U(x')}} = \int_{t_0}^t dt' = t - t_0$$

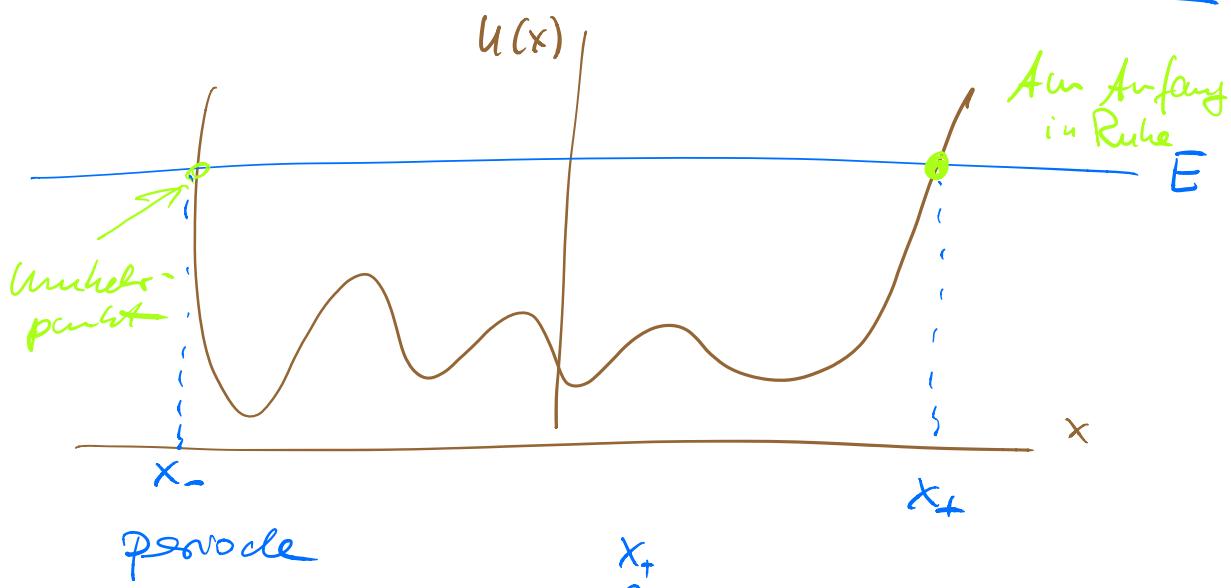
$$x_0 = x(t_0)$$

$$x = x(t)$$

$$t = t_0 \pm \sqrt{\frac{m}{2}} \int_{x_0}^{x(t)} \frac{dx}{\sqrt{E - U(x')}}$$

implizite Bestimmungsgleichung f. $x(t)$!

betrachten wir ein periodisches System



periode

$$T = 2 \sqrt{\frac{m}{2}} \int_{x_-}^{x_+} \frac{dx'}{\sqrt{E - U(x')}} \quad x_- < x_+$$

$$U(x) = r x^n \quad n \text{ gerade}$$

$$U(x_{\pm}) = E \Rightarrow x_{\pm} = \pm \left(\frac{E}{r} \right)^{1/n}$$

$$T(E) = \sqrt{2m} \int_{x_-}^{x_+} \frac{dx'}{\sqrt{E - r x^n}} \quad s = \frac{x}{x_+}$$

$$= \frac{E^{1/n}}{r^{1/n}} \sqrt{\frac{2m}{E}} \int_{-1}^1 \frac{ds}{\sqrt{1 - s^n}}$$

$$T(E) = \frac{\sqrt{2m}}{\delta^{\gamma_n}} E^{\gamma_n - 1/2} C_n$$

$$C_n = \int_{-1}^1 \frac{ds}{\Gamma_{1-s^n}} = \frac{\sqrt{\pi} \Gamma(\frac{n+1}{n})}{\Gamma(\frac{2n+1}{2n})}$$

$$\Gamma(x) = \int_0^\infty e^{-t} t^{x-1} dt$$

$$n=2 \quad C_2 = \pi$$

$$C_4 = 2.622 \dots$$

$$C_6 = 2.4285 \dots$$

$T(E)$ is nur für den harmonischen
Oszillatoren unabhängig von E !