

# Classical Theoretical Physics II

Lecture: Prof. Dr. K. Melnikov – Exercises: Dr. A. Behring

## Exercise Sheet 11

Issue: 02.07. – Submission: 09.07. @ 10:00 Uhr – Discussion: 12./13.07.

### Exercise 1: Hamiltonians of various systems

5 points

Find the canonical momenta, the Hamiltonian and Hamilton's equations of motion for each of the following cases.

- (a) A free particle with mass  $m$  in one dimension.
- (b) A harmonic oscillator with mass  $m$  and angular frequency  $\omega$  in one dimension.
- (c) A particle with mass  $m$  in the one-dimensional potential  $U(x) = \alpha x^n$ .
- (d) A particle with mass  $m$  in a three-dimensional potential  $U(r) = -k/r$ . Use polar coordinates  $r, \theta, \phi$ .
- (e) Two particles with masses  $m$  and  $M$  moving in a two-dimensional plane and interacting gravitationally. Use Cartesian coordinates  $x_i, y_i$  for each particle. (What could have been a better set of coordinates for this problem?)

### Exercise 2: Runge-Lenz vector and Poisson brackets

5 points

Consider a particle with mass  $m$ , position  $\vec{r}$  and momentum  $\vec{p}$  in three-dimensional space. The angular momentum of the particle is given by  $\vec{M} = \vec{r} \times \vec{p}$ .

- (a) Calculate the Poisson brackets

$$\{M_i, r_j\}, \quad \{M_i, p_j\}, \quad \{M_i, M_j\}, \quad \{M_i, M^2\}. \quad (2.1)$$

Make use of the properties of the Poisson brackets that were introduced in the lectures in Chapter 9 in Eqs. (20) to (22).

- (b) The Runge-Lenz vector is given by  $\vec{A} = \vec{p} \times \vec{M} - mk\hat{r}$ , where  $\hat{r} = \vec{r}/r$  is the unit vector in radial direction. Prove using Poisson brackets that the Runge-Lenz vector is conserved in the Kepler problem. That is, show that  $\{H, A_i\} = 0$  holds for the Hamiltonian  $H = \frac{p^2}{2m} - \frac{k}{r}$ .
- (c) We have already shown in a previous problem that the Runge-Lenz vector is no longer conserved if we add a  $r^{-2}$  in addition to the  $r^{-1}$  term in the potential. Consider the Hamiltonian

$$H = \frac{p^2}{2m} - \frac{k}{r} + \frac{C}{r^2} \quad (2.2)$$

and calculate the time derivative of the Runge-Lenz vector using Poisson brackets and compare your result to the result of problem 8.2 (d).

**Exercise 3: Particle in magnetic field****9 points**

The motion of a particle with mass  $m$  and charge  $e$  is described by the Lagrangian

$$L = \frac{m}{2} \dot{\vec{r}}^2 - \left( e\phi(\vec{r}, t) - \frac{e}{c} \vec{A}(\vec{r}, t) \cdot \dot{\vec{r}} \right), \quad (3.1)$$

where  $\phi$  and  $\vec{A}$  are the electromagnetic potentials.

- (a) Determine the canonical momentum  $\vec{p}$ , construct the Hamiltonian  $H$  and show that Hamilton's equations are given by

$$\dot{r}_i = \frac{1}{m} \left( p_i - \frac{e}{c} A_i \right), \quad \dot{p}_i = \frac{e}{mc} \left( p_j - \frac{e}{c} A_j \right) \frac{\partial A_j}{\partial r_i} - e \frac{\partial \phi}{\partial r_i}, \quad (3.2)$$

- (b) Derive the Lorentz force law from Hamilton's equations. *Hint:* use index notation and the following expressions for the electric and magnetic fields in terms of the electromagnetic potentials

$$E_i = -\frac{\partial \phi}{\partial r_i} - \frac{1}{c} \frac{\partial A_i}{\partial t}, \quad B_i = \epsilon_{ijk} \frac{\partial A_k}{\partial r_j}. \quad (3.3)$$

- (c) In the presence of a magnetic field we have that  $\vec{p} \neq m\vec{v}$ , where  $\vec{v} = \dot{\vec{x}}$ . This is evident from the first of Hamilton's equations in Eq. (3.2). Their Poisson brackets also reflect this fact: while  $\{p_i, p_j\} = 0$ , show that

$$\{v_i, v_j\} = -\frac{e}{m^2 c} \epsilon_{ijk} B_k. \quad (3.4)$$

- (d) Consider the specific case of a particle restricted to the  $x, y$ -plane, with potentials  $\phi(\vec{r}, t) = 0$  and  $\vec{A}(\vec{r}, t) = (-By, 0, 0)^T$ . What are the associated  $\vec{E}$  and  $\vec{B}$  fields? Give Hamilton's equations for this case. Show that they imply

$$m\dot{y} + \frac{eB}{c}x = K_1, \quad m\dot{x} - \frac{eB}{c}y = K_2, \quad (3.5)$$

with unknown constants  $K_{1,2}$ .

- (e) Show that the general solution to Eq. (3.5) is given by

$$x(t) = \frac{K_1}{m\omega} + R \sin(\omega t + \phi), \quad (3.6)$$

$$y(t) = -\frac{K_2}{m\omega} + R \cos(\omega t + \phi), \quad (3.7)$$

where  $\omega = \frac{eB}{mc}$  and  $R$  and  $\phi$  are (for now) unknown integration constants.

- (f) Impose the boundary conditions  $x(0) = x_0$ ,  $y(0) = y_0$ ,  $\dot{x}(0) = v_0$  and  $\dot{y}(0) = 0$ , and sketch the motion of the particle.