

① Energie des elektrostatischen Feldes

Jackson Gl. (1.53) : $W = \frac{1}{2} \int d^3x g(\vec{x}) \bar{\Phi}(\vec{x})$

$$= \frac{1}{8\pi\epsilon_0} \int d^3x \int d^3x' \frac{g(x)g(x')}{|\vec{x} - \vec{x}'|^2}$$

(a) homogen geladene Kugel mit Radius R :

$$g(\vec{x}) = \Theta(R - r) \frac{1}{4\pi} \int d^3r' g(r') = Q = \int_0^R dr \int_0^{2\pi} d\phi \int_0^\pi r^2 \sin\theta g(\vec{r})$$

$$\Rightarrow k = \frac{3Q}{4\pi R^3}$$

$$= 4\pi \int_0^R dr \int_0^{2\pi} d\theta \int_0^\pi r^2 \sin\theta g(\vec{r}) = \frac{4\pi}{3} R^3 k$$

$$W = \frac{1}{8\pi\epsilon_0} \left(\frac{3Q}{4\pi R^3} \right)^2 (8\pi)^2 \int_0^R dr \int_0^{2\pi} d\theta \int_0^\pi r^2 \sin\theta \int_0^R dr' \int_0^{2\pi} d\theta' \int_0^\pi r'^2 \sin\theta' \frac{r^2 r'^2}{\sqrt{r^2 + r'^2 - 2rr' \cos(\theta - \theta')}}$$

alternativ (Schneller): $W = \frac{\epsilon_0}{2} \int d^3r |\vec{E}(\vec{r})|^2$

homogene Kugel: $\vec{E} = \frac{\vec{r}}{4\pi\epsilon_0} \left\{ \begin{array}{ll} \frac{1}{R^2} & r \leq R \\ \frac{1}{r^2} & r > R \end{array} \right. \quad (\text{Vorlesung?})$

$$W = \frac{\epsilon_0 Q^2}{2 (4\pi\epsilon_0)^2} \cdot 4\pi \left[\underbrace{\int_0^R dr r^2 \left(\frac{1}{R^3} \right)^2}_{\frac{r^4}{R^8}} + \underbrace{\int_R^\infty dr \frac{r^2}{r^4}}_{\frac{1}{r^2}} \right]$$

$$= \frac{1}{JR^8} \int_0^\infty r^5 dr - \frac{1}{2} \int_R^\infty \frac{1}{r^2} dr$$

$$= \frac{Q^2}{8\pi\epsilon_0} \left(\frac{R^5}{JR^8} - \underbrace{\frac{1}{\infty}}_0 + \frac{1}{R} \right) = \underline{\underline{\frac{Q^2}{4\pi\epsilon_0} \frac{3}{5} \frac{1}{R}}}$$

$$\frac{1}{J} + 1 = \frac{6}{5}$$

(b) homogen geladene Kugelschale

②

$$\vec{E}(r) = \frac{Q}{4\pi\epsilon_0} \vec{e}_r \frac{1}{r^2} \begin{cases} 0 & r \leq R_i \\ \frac{r^3 - R_i^3}{R_a^3 - R_i^3}, R_i < r < R_a \\ 1 & r \geq R_a \end{cases}$$

$$W = \frac{\epsilon_0}{2} \int d\Omega \int d\sigma |\vec{E}(r)|^2 = \frac{\epsilon_0 Q^2}{2 (4\pi\epsilon_0)^2} \cdot 4\pi \left[\int_{R_i}^{R_a} d\sigma r^2 \cdot 0 \right]$$

$$+ \int_{R_i}^{R_a} d\sigma \frac{r^2 (r^3 - R_i^3)^2}{2^4 (R_a^3 - R_i^3)} + \int_{R_a}^{\infty} d\sigma \frac{r^2}{2^4}$$

$$= \frac{Q^2}{8\pi\epsilon_0} \left[\frac{1}{(R_a^3 - R_i^3)^2} \int_{R_i}^{R_a} d\sigma \left(r^4 - 3r^2 R_i^2 + \frac{R_i^6}{r^2} \right) + \int_{R_a}^{\infty} d\sigma \frac{1}{r^2} \right]$$

$$\underbrace{\frac{1}{5} r^5 - r^2 R_i^3 - \frac{R_i^6}{r^2}}_{\text{R}_i} \Big|_{R_i}^{R_a}$$

$$- \frac{1}{r} \Big|_{R_a}^{\infty}$$

$$= \frac{Q^2}{8\pi\epsilon_0} \left[\frac{\frac{R_a^5 - R_i^5}{5} - R_i^3(R_a^2 - R_i^2) + R_i^6 \frac{2aR_i}{R_a^2 R_i}}{(R_a^3 - R_i^3)^2} + \frac{1}{R_a} \right]$$

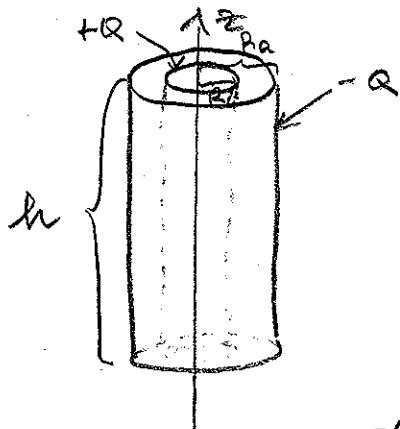
Grenzwerte: $R_i \rightarrow 0$: $W = \frac{Q^2}{8\pi\epsilon_0} \left[\frac{\frac{R_a^5}{5} - R_a^2}{R_a^6} + \frac{1}{R_a} \right]$

$$= \frac{Q^2}{4\pi\epsilon_0} \frac{3}{5} \frac{1}{R_a}$$

wie homogen geladene Kugel

$$R_i = R_a : W = \frac{Q^2}{8\pi\epsilon_0} \left[\frac{0 - 0 - 0}{0} + \frac{1}{R_a} \right] = \frac{Q^2}{4\pi\epsilon_0} \frac{1}{2R_a}$$

② Zylindrischer Kondensator



Ladungsverteilung

$$g(r) = \frac{Q}{2\pi h} [\epsilon_0 (\delta(r - R_i) - \delta(r - R_o))] \Theta\left(\frac{h}{2} - |r|\right)$$

$R_{i/o} \ll h$

$$\int d\sigma \int d\sigma' \int d\sigma'' \delta(r - R_i) \delta(r - R_o) = S \pi \epsilon_0 R_i R_o = Q$$

$$g(r) = \frac{Q}{2\pi h} \left[\frac{\delta(r - R_i)}{R_i} - \frac{\delta(r - R_o)}{R_o} \right] \Theta\left(\frac{h}{2} - |r|\right)$$

$$\Phi(r) = \int d\sigma \int d\sigma' \frac{g(r')}{|r - r'|} = \frac{Q}{2\pi h} \cdot 2\pi \dots \text{(hässliches Integral)}$$

Symmetrieverlegungen: $\vec{E}(r) = E_p \hat{e}_r$, 0 für $r > R_o$, 0 für $|r| > \frac{h}{2}$

Gaußscher Satz:

$$\underbrace{\int_Q d^3r g(r)}_{Q} = \underbrace{\int_{\partial V} \vec{E} \cdot d\vec{l}}_{2\pi h \epsilon_0 E(r)}, \quad dl = \int d\sigma \int d\sigma' \hat{e}_r$$

$R_i \leq r \leq R_o$

$$\vec{E}(r) = \begin{cases} 0, & r < R_i \\ \frac{Q}{2\pi \epsilon_0 r} \hat{e}_r, & R_i \leq r < R_o \\ 0, & r > R_o \end{cases}, \quad \vec{E} = -\nabla \Phi$$

$E_p = -\frac{\partial \Phi}{\partial r}$

$$\Phi(r) = - \int_{\rho_0}^r ds' E(s') + \Phi(\rho_0) = - \frac{Q}{2\pi \epsilon_0 h} \underbrace{\int_{\rho_0}^r \frac{ds'}{s'}}_{\ln s - \ln \rho_0} + \Phi(\rho_0)$$

$$\Phi(r) = \Phi(\rho_0) + \frac{Q}{2\pi \epsilon_0 h} \ln\left(\frac{s}{\rho_0}\right)$$

für $R_i \leq r < R_o$, $\Phi(r < R_i)$, $\Phi(r > R_o)$

$$\Phi(r \rightarrow \infty) \xrightarrow{!} 0 \Rightarrow \Phi(r > R_o) = 0 \Rightarrow \Phi(r = R_o) = 0 \xrightarrow{s_0 = R_o} \Phi(R_i \leq r < R_o) = \frac{Q}{2\pi \epsilon_0 h} \ln\left(\frac{R_o}{R_i}\right)$$

const.

$$\Phi(r < R_i) = \frac{Q}{2\pi \epsilon_0 h} \ln\left(\frac{R_i}{\rho_0}\right)$$

Φ steigt!

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$$(B) \text{ Kapazität: } C = \frac{Q}{U}$$

Anliegende Spannung: $U = \Phi(2i) - \Phi(Ra)$

$$= \frac{\alpha}{2\pi\epsilon_0 h} \ln\left(\frac{R_a}{R_i}\right) = 0$$

dann: $C = \frac{Q}{U} = \frac{2\pi\epsilon_0 h}{\ln(R_a/R_i)}$

(C) Gesamtenergie des Kondensators:

$$W = \frac{\epsilon_0}{2} \int d^3r |\vec{E}|^2 = \frac{\epsilon_0}{2} \left(\frac{\alpha}{2\pi\epsilon_0 h}\right)^2 \int d\varphi \int dz \int dr \frac{dS}{R_i} \frac{1}{r^2}$$

$$= \frac{Q^2}{4\pi\epsilon_0 h} \ln\left(\frac{R_a}{R_i}\right)$$

③ Multipolentwicklung

$$\Phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \left[\frac{Q}{r} + \frac{\vec{p} \cdot \vec{r}}{r^3} + \frac{1}{2} \sum_{i,j} Q_{ij} \frac{x_i x_j}{r^5} + \dots \right]$$

Merksatz: $Q = \int d^3x' g(\vec{r}')$ Monopolelement

$$p_i = \int d^3x' g(\vec{r}') x_i \quad \text{Dipolelement}$$

$$Q_{ij} = \int d^3x' g(\vec{r}') (3x'_i x'_j - r'^2 \delta_{ij}) \quad \text{Quadrupol-Element}$$

$$(a) (i) g(\vec{r}) = q \delta(x) \delta(y) \delta(z-d) = q \delta(x) \delta(y) \delta(z)$$

$$Q = \int d^3x' g(\vec{r}') = q - q = 0$$

$$p_i = \int d^3x' g(\vec{r}') x_i = dq - 0 \cdot q = dq \quad \text{für } i=3, i=1,2: p_{i2}=0$$

$$Q_{ij} = \int d^3x' g(\vec{r}') (3x'_i x'_j - \underbrace{r'^2 \delta_{ij}}_{x'^2 + y'^2 + z'^2}) = 3 \underbrace{\int d^3x' g(\vec{r}') x'_i x'_j}_{x'_i, x'_j \in \mathbb{R} \rightarrow 0} - \underbrace{\int d^3x' g(\vec{r}') \delta_{ij}}_{= q d^2 \delta_{ij}}$$

$$\Rightarrow Q_{33} = 3qd^2 - qd^2 = 2qd^2$$

$$Q_1 = Q_{22} = -qd^2, Q_{ij} = 0 \quad \text{für } i \neq j$$

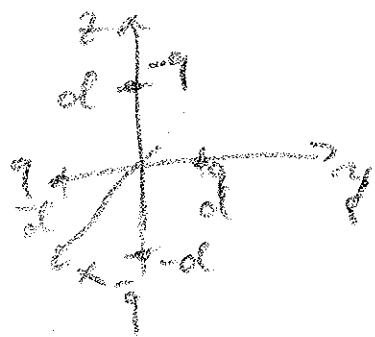
$$\text{Nebenbedingung: } \sum Q_{ij} = Q_1 + Q_{22} + Q_{33} = 0 \quad (\text{gilt, immer!})$$

$$(b) (i) \text{ Potential (s.o.): } \Phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \left[\frac{0}{r} + \frac{dq}{r^3} + \frac{qd^2}{2} \underbrace{\frac{-\vec{r}^2}{x^2 + y^2 + z^2}}_{+} \right]$$

$$\text{Zylinderkoord.: } \vec{r} = \hat{r} \hat{r} + \hat{\varphi} \hat{\varphi} + \hat{z} \hat{z} \quad |z| = \sqrt{x^2 + y^2}$$

$$\vec{E} = -\vec{\nabla} \Phi = \frac{qd}{4\pi\epsilon_0} \left[\left(\frac{3}{r^5} (z - \frac{d}{2}) + \frac{5}{r^7} \frac{3}{2} z^2 d \right) \hat{r} - \frac{1}{r^3} \left(\frac{3dz}{r^2} + 1 \right) \hat{z} \right]$$

$$(a)(ii) f(\theta) = q\delta(\theta)[\delta(y-d)\delta(\theta) + \delta(y+d)\delta_{\theta}] - \delta(y)\delta(\theta-d) - \delta(y)\delta(\theta+d)]$$



$$Q = \int d^3r' f(\theta') = q(1+1-1)=0$$

$$q_i = \int d^3r' f(\theta') x_i$$

$$= \begin{cases} 0 & i=1 \\ q(d-d)=0 & i=2 \\ q(-d-(d))=0 & i=3 \end{cases} = 0$$

$$\text{sign } \int d^3r' f(\theta') (\delta_{xx} \delta_{yy} - \delta_{xy}^2 \delta_{yz})$$

$$Q_{xx} = -q d^2 \cdot 0, \quad Q_{yy} = q(3d^2 - 0d^2) = 6qd^2$$

$$Q_{xz} = q(3d^2 + 0d^2) = -6qd^2$$

$$Q_{xy} = 0, \quad Q_{yz} = q \cdot 0 = 0$$

$$(b)(ii) \Phi(x) = \frac{1}{4\pi\epsilon_0} \frac{1}{2} \cdot 6qd^2 \left(\frac{y^2 - z^2}{x^5} \right) = \frac{3qd^2}{4\pi\epsilon_0} \frac{y^2 - z^2}{x^5}$$

$$\vec{E}(x) = \nabla \Phi = \frac{3qd^2}{4\pi\epsilon_0} \left[\left(\frac{\partial y}{x^5} - \frac{3z^2(y^2-z^2)}{2x^7} \right) \hat{e}_y - \left(\frac{\partial z}{x^5} - \frac{3y^2(y^2-z^2)}{2x^7} \right) \hat{e}_x \right]$$

$$= \frac{3qd^2}{4\pi\epsilon_0} \left[\frac{5(y^2-z^2)}{x^7} \hat{x} - \frac{2y^2 \hat{e}_y + 2z^2 \hat{e}_x}{x^7} \right]$$

$$\frac{dy}{dx} = \frac{y^2 - d}{y}$$

?

$$Q = q \int dy \delta(y^2 - d)$$

$$2q \int_{-\infty}^{\infty} \delta(y^2 - d)$$

Nullstellen: $y^2 - d = 0 \Leftrightarrow y^2 = d^2$
 also $y_{\pm} = \pm d$

$$\frac{d}{dy} (y^2 - d) = \frac{1}{2} \frac{2y}{y^2}$$

$$= \frac{y}{y^2} = \text{sign}(y) \quad (\text{gibt nur das Vorzeichen des Argumenten } b)$$

$$\delta(f(x)) = \sum \frac{1}{|f'(x_0)|} \delta(x - x_0)$$

$$\rightarrow \delta(y^2 - d) = \frac{1}{\text{sign}(0)} \delta(y+d) + \frac{1}{\text{sign}(d)} \delta(y-d)$$

$$\Rightarrow \int_{-\infty}^{\infty} dy \delta(y^2 - d) = \int_{-\infty}^0 dy \delta(y+d) + \int_0^{\infty} dy \delta(y-d)$$

$$F(\vec{r}) = \frac{q\epsilon}{4\pi\epsilon_0} \left[\frac{q}{r^2} + \frac{\epsilon_0}{r^2} \left(\frac{3r^2 - 8^2}{r^2} \right) \right]$$

$$\frac{\partial F}{\partial r} = \frac{q\epsilon}{4\pi\epsilon_0} \left[-3 \frac{\partial q}{r^3} + \frac{\partial \epsilon_0}{r^3} + \frac{q\epsilon_0}{r^5} + \sqrt{\frac{q\epsilon}{3\pi\epsilon_0}} \left(2r^2 - 8^2 \right) \right]$$

$$\frac{\partial F}{\partial q} = \frac{q\epsilon}{4\pi\epsilon_0} \left[-3 \frac{\partial r}{r^3} - \frac{\partial \epsilon_0}{r^3} - \sqrt{\frac{q\epsilon}{3\pi\epsilon_0}} \left(2r^2 - 8^2 \right) \right]$$

$$E(\vec{r}) = -\nabla F = +\frac{q\epsilon}{4\pi\epsilon_0} \left[\left(\frac{3q}{r^2} + \frac{\partial q}{r^3} \right) \hat{r} - \left(\frac{3\epsilon_0}{r^2} + \frac{\partial \epsilon_0}{r^3} \right) \hat{r} - \left(\frac{3q^2 - 8^2}{r^2} + \frac{\partial q}{r^3} \right) \hat{\theta} \right]$$

$$= \frac{q\epsilon}{4\pi\epsilon_0} \left[-\frac{3q}{r^2} + \frac{\partial q}{r^3} - \frac{3\epsilon_0}{r^2} + \frac{\partial \epsilon_0}{r^3} + \frac{3q^2 - 8^2}{r^2} + \frac{\partial q}{r^3} \right] \hat{r} - \frac{3q^2 - 8^2}{r^2} \hat{\theta}$$

$$= \frac{q\epsilon}{4\pi\epsilon_0} \left[\left(\frac{3}{r^2} - \frac{\partial r}{r^3} + \frac{15q^2}{r^2} \right) \hat{r} - \left(\frac{3}{r^2} + \frac{1}{r^3} \right) \hat{\theta} \right]$$